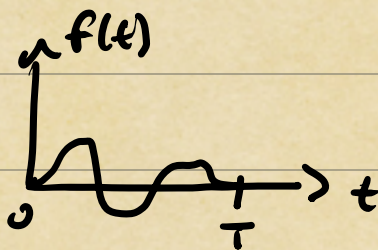


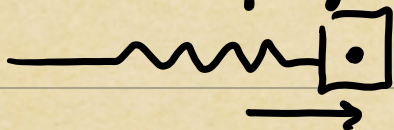
More on ODE Solution Ops

Consider the forced harmonic oscillator

$$(*) \quad u''(t) + k^2 u(t) = f(t)$$



$k = \text{spring const.}$



$\rightarrow k u(t) = \text{displacement from equilibrium}$

The solutions to the homogeneous eqn. are

$$(**) \quad v''(t) + k^2 v(t) = 0 \quad \Rightarrow \quad \begin{aligned} v_1(t) &= C_1 \cos kt \\ v_2(t) &= C_2 \sin kt \end{aligned}$$

So, for any solution $u(t)$ of $(*)$, we also have

$$\tilde{u}(t) = C_1 \cos kt + C_2 \sin kt + u(t)$$

is a solution of $(*)$ for any $C_1, C_2 \in \mathbb{C}$.

This leads to a whole **family** of solution operators with different behavior as $t \rightarrow \pm \infty$.

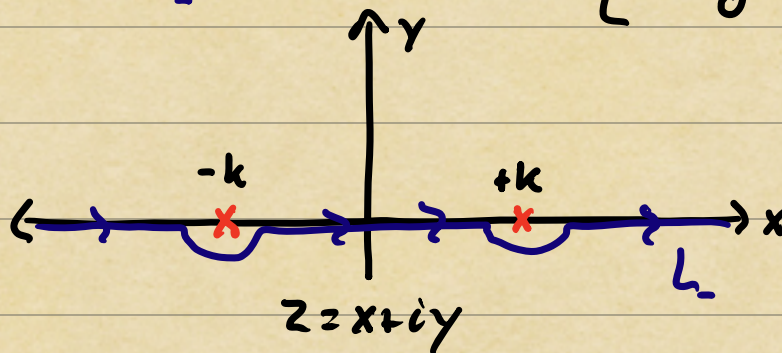
$$u(t) = \int_0^t G(t-\tau) f(\tau) d\tau$$

\uparrow
 Green's function
 = kernel of soln op.

Causal Green's Function

The *causal* Green's function is

$$G(t-\tau) = \frac{1}{2\pi} \int_{L_-} \frac{e^{iz(t-\tau)}}{k^2 - z^2} dz = \begin{cases} \frac{1}{k} \sin k(t-\tau) & \tau < t \\ 0 & \tau > t \end{cases}$$



The solution is "causal" because

$$u(t) = \frac{1}{k} \int_0^t \sin k(t-\tau) f(\tau) d\tau$$

only depends on $f(\tau)$ for $\tau < t$. In particular, $u(t) = 0$ for $t < 0$, which is before the forcing turns on.

How does $u(t)$ behave for $t > T$?

$$u(t) = \frac{1}{k} \int_0^T \underbrace{\sin k(t-\tau)}_{\sin kt \cos k\tau - \cos kt \sin k\tau} f(\tau) d\tau$$

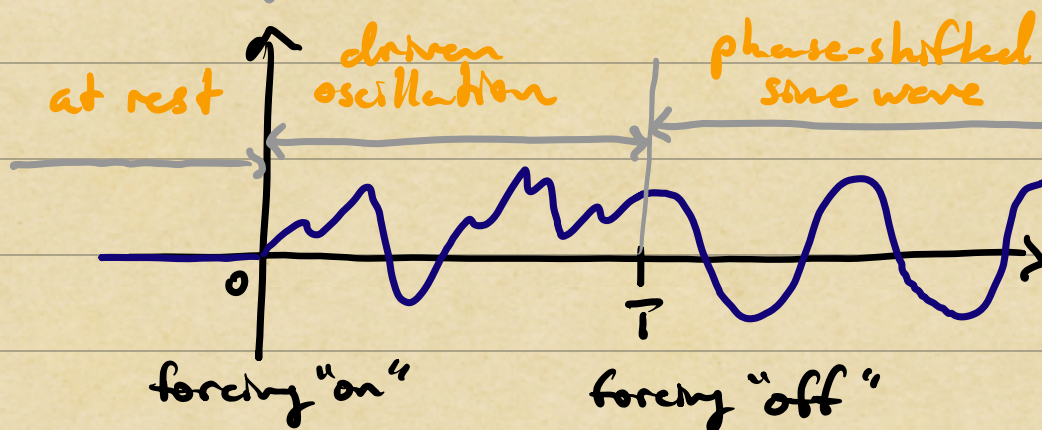
$$= \sin kt \left[\frac{1}{k} \int_0^T f(\tau) \cos k\tau d\tau \right] - \cos kt \left[\frac{1}{k} \int_0^T f(\tau) \sin k\tau d\tau \right]$$

$$= C_1 \sin kt + C_2 \cos kt, \quad t > T.$$

$$= A \sin(kt + \phi)$$

↙ phase shift = $\arctan(\frac{1}{2})$

$$\uparrow \text{amplitude} = \sqrt{C_1^2 + C_2^2}$$



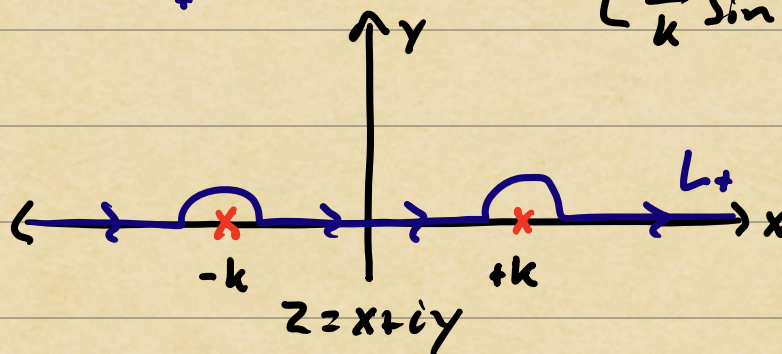
Q1: How are the causal & anti-causal Green's functions related?

Q2: How are the corresponding solutions related?

Anti-Causal Green's Function

The anti-causal Green's function is

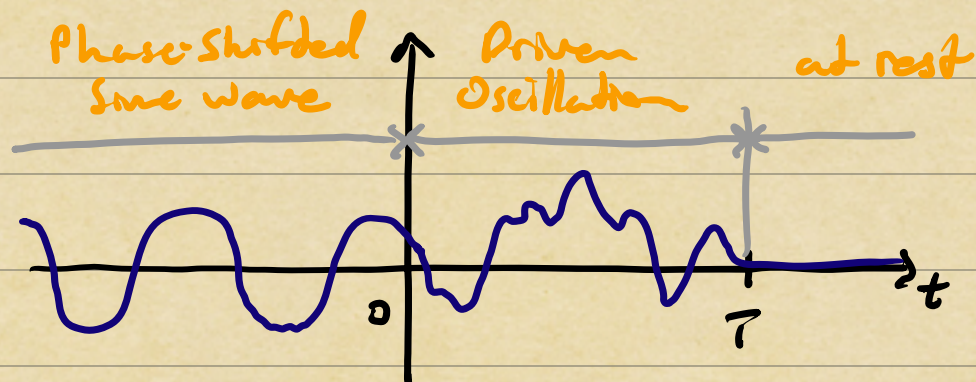
$$G_+(t-\tau) = \frac{1}{2\pi} \int_{L_+} \frac{e^{iz(t-\tau)}}{k^2 - z^2} dz = \begin{cases} 0 & \tau < t \\ \frac{1}{k} \sin k(t-\tau) & \tau > t \end{cases}$$



Here, the solution depends on $f(\tau)$ for $\tau > t$,

$$u(t) = \frac{1}{k} \int_t^{\infty} \sin k(t-\tau) f(\tau) d\tau.$$

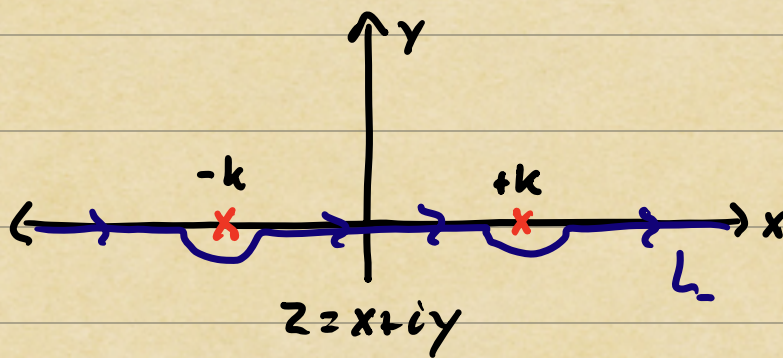
A similar calculation from above shows



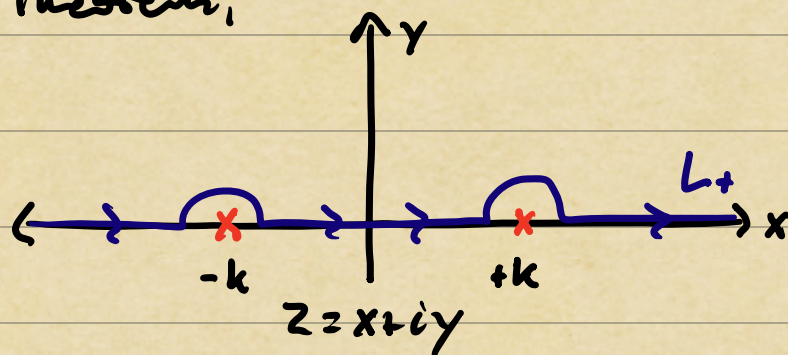
It's as if we set up an unforced oscillation

with just the right amplitude and phase
 so that the forcing between $t=0$ and $t=T$ has
 the net effect of eliminating all oscillation
 and leaving the oscillator at rest.

How are the two Green's functions related?



Residue
Theorem!



$$G_+(t-\tau) - G_-(t-\tau) = \frac{1}{2\pi} \int_{L_+} \frac{e^{iz(t-\tau)}}{k^2 - z^2} dz - \frac{1}{2\pi} \int_{L_-} \frac{e^{iz(t-\tau)}}{k^2 - z^2} dz$$

$$= i \operatorname{res}_{z=k} \frac{e^{iz(t-\tau)}}{k^2 - z^2} + i \operatorname{res}_{z=k} \frac{e^{iz(t-\tau)}}{k^2 - z^2}$$

$$= i \lim_{z \rightarrow -k} (z+k) \frac{e^{iz(t-\tau)}}{(k+z)(k-z)} + i \lim_{z \rightarrow k} (z-k) \frac{e^{iz(t-\tau)}}{(k+z)(k-z)}$$

$$= i \frac{e^{-ik(t-\tau)}}{2k} - i \frac{e^{ik(t-\tau)}}{2k}$$

$$= \frac{1}{k} \left[\frac{e^{ik(t-\tau)}}{2i} - \frac{e^{-ik(t-\tau)}}{2i} \right]$$

$$= \frac{1}{k} \sin k(t-\tau)$$

Much of this goes through for more general problems if we replace

$$1/p(z) \quad \text{by} \quad \hat{G}(z)$$