

# Disordered aragonite—the new high-pressure high-temperature phase of $\text{CaCO}_3$

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## Abstract

Phases of  $\text{CaCO}_3$  stabilised at high-pressures and temperatures are the potential agents of the global carbon cycle, transferring oxidised carbon in deep Earth's interiors and thus are of special interest for the Earth sciences. Here, we report finding of the new phase, named disarag, which is dynamically disordered aragonite with freely rotating  $\text{CO}_3$  groups, similar to that in  $\text{CaCO}_3$ -V phase with calcite-like structure. Disarag has stability field expanding from 3 to 10 GPa and from 1600 to 2000 K. Consideration of twinned structure enlarges this field, decreasing transition temperature from aragonite to disarag on 100–300 K. At P-T parameters corresponding to the transition from

aragonite to disarag, the marked disappearance of the diffraction peaks is observed in *in situ* experiments. We show that among known phases of  $\text{CaCO}_3$  disarag is the best candidate for the explanation of this reconstruction of diffraction pattern. Also for the first time, using *ab initio* molecular dynamics technique, we determine equilibrium curves between calcite and its disordered phases  $\text{CaCO}_3\text{-IV}$  and  $\text{CaCO}_3\text{-V}$ . We show that the transitions of alkaline-earth carbonates  $\text{CaCO}_3$ ,  $\text{SrCO}_3$ , and  $\text{BaCO}_3$  to the disordered states start, when the critical angle of librations of  $\text{CO}_3$  group about axis perpendicular to the molecular three-fold axis exceeds  $45^\circ$ . Calcite-like structure of  $\text{CaCO}_3$  is characterised by more intense librations, than aragonite-like structure of this compound and reaches the critical angle at lower temperatures. As a result, calcite transforms into the disordered state at lower temperatures than aragonite.

## Introduction

The recent decade was a revolutionary one in the investigation of the  $\text{CaCO}_3$  phase diagram. Lots of new phases were discovered with both experimental and theoretical methods. The P-T phase diagram of this simple compound is turned out to be far from simple. It is complicated by the presence of several metastable phases  $\text{CaCO}_3\text{-II}$ ,  $\text{III}$ ,  $\text{IIIb}$ , and  $\text{IV}$  (cc-II,  $\text{III}$ ,  $\text{IIIb}$ ,  $\text{VI}$ ), the joint appearance of several high-pressure phases ( $\text{CaCO}_3\text{-VII}$  with aragonite-II,<sup>1</sup> and cc-III with cc-IIIb<sup>2</sup>), the dynamical disordering of calcite (cc-I) structure with the formation of  $\text{CaCO}_3\text{-IV}$  (cc-IV) and  $\text{CaCO}_3\text{-V}$  (cc-V) phases,<sup>3,4</sup> and amorphisation of the sample at high-pressures and high-temperatures.<sup>5</sup> All set of experimental and theoretical techniques were used for the determination of high-pressure and high-temperature polymorphs, and the  $\text{CaCO}_3$  phase diagram gives several successful examples of cooperative use of *ab initio* theoretical methods and experimental techniques for the solution of crystal structures. Identification of high-pressure phases  $\text{CaCO}_3\text{-VII}$  and aragonite-II,<sup>1,6</sup> evidence of tetrahedral coordination of carbon atoms in  $\text{CaCO}_3\text{-}P2_1/c - h$ <sup>7</sup> and description of the disorder in high-temperature phases cc-IV and cc-V can be given as examples.<sup>4</sup>

Nevertheless, despite the numerous investigation have been performed, there is still an apparent *white spot* on the P-T phase diagram of  $\text{CaCO}_3$ , spreading from 3 to 11 GPa and from 800 to 1600 K. *In situ* X-ray diffraction experiments have revealed marked disappearance of diffraction peaks in this field,<sup>8,9</sup> with only three peaks left. With such a small amount of diffraction information, even simple indexing is problematic. The suggested hypothesis assuming the formation of cc-V<sup>8</sup> is inconsistent with available results of quenching experiments,<sup>10</sup> and no other candidates have been suggested yet.

In the present work, we perform *ab initio* molecular dynamic (MD) simulation of calcite and aragonite and show, that similar to calcite, aragonite adopts a disordered state at high temperatures. The disordered aragonite is a good candidate for the phase that appeared in the mentioned *white spot*.

**Aragonite and calcite crystal structures** Crystal structures of calcite and aragonite are similar in the sense, that both they can be presented as the close-packing of Ca-atoms, with  $\text{CO}_3$  molecular groups occupying octahedral voids.<sup>11</sup> In the case of calcite, the packing of Ca atoms is three-layered (*fcc*), and  $\text{CO}_3$  groups are in the centers of the octahedral voids, midway between close-packed (cp) Ca-layers (Figure 1a). In the case of aragonite, the packing of Ca-atoms is two-layered (*hcp*), with  $\text{CO}_3$  groups shifted from the centers of octahedral voids forming the double-layer.<sup>11</sup> The upper sublayer is at  $1/3$  of  $c'$  and the lower one—on  $2/3$  of  $c'$ , where  $c'$  is the distance between adjacent cp Ca-layers (Figure 1b). In adjacent sublayers,  $\text{CO}_3$  triangles have opposite orientations, being rotated on  $180^\circ$ .

Through the manuscript, we use the aragonite structure in a non-standard *Pmcn* (#162) setting. In this setting, the comparison of aragonite and calcite structures is more straightforward than in standard *Pnma* setting, as in both structures  $c$ -axes are perpendicular to the plane of cp-Ca layer.

Strontianite  $\text{SrCO}_3$  and witherite  $\text{BaCO}_3$  at 1 atm have the same structure as aragonite.<sup>12</sup> These compounds were also used as objects of MD simulations to compare aragonite-like

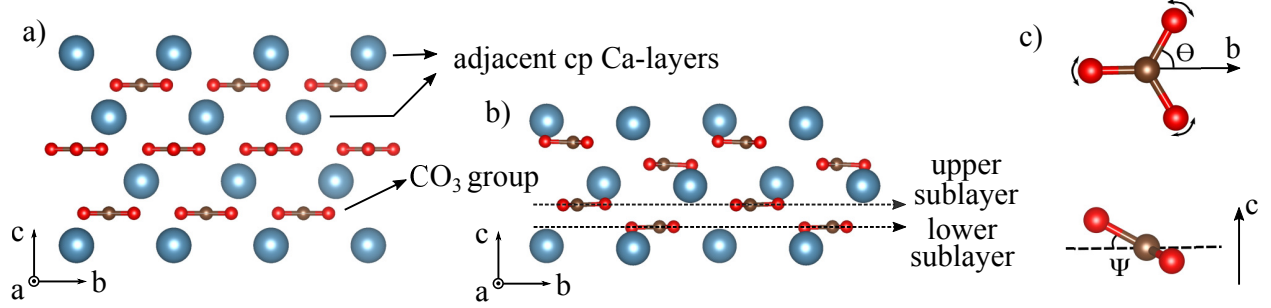


Figure 1: Calcite (a) and aragonite (b) crystal structures along  $a$ -axis and scheme, showing rotational angle  $\theta$  about three-fold axis of CO<sub>3</sub> molecular group and the libration angle  $\psi$  about the axis perpendicular to the three-fold axis (c)

disordered structures with different compositions.

**Microstructure of aragonite crystals** Twinning by  $\{110\}$  plane is the main microstructural feature of aragonite crystals.<sup>13–15</sup> This type of twinning can be produced during crystal growth, or as it was recently shown, mechanically, at crystal splitting.<sup>16</sup> In natural aragonite samples, we observed twinning down to the unit-cell size, with the formation of disordered and partially ordered sequences of twinning planes.<sup>15</sup> The ordered sequences of twinning planes are aragonite polytypes, which due to the orthorhombic symmetry, were designated as  $O2$ ,  $O3$ , *etc.*, according to Ramsdell notation.<sup>15</sup>

In the present work, to account for the effect of microstructure, we perform MD simulations of not only aragonite but also of  $O2$ -polytype, which is aragonite twinned by each second  $\{110\}$  plane. The consideration of the twinned structure is necessary, as powdered samples used in high-pressure experiments are enriched by the twin boundaries produced during grinding.

## Methods

The total energies and forces were calculated by solving the Schrödinger equation based on projector augmented plane-wave implementation of density functional theory, within the Vienna Ab Initio Simulation Package (VASP).<sup>17</sup> Exchange-correlation effects were treated

in the generalised gradient approximation (GGA) with the Perdew-Wang scheme.<sup>18</sup> Pseudo-potentials with  $3p^6 4s^2 3d^{0.1}$  (Ca),  $5s^2 4d^{0.01}$  (Sr),  $s^1$  (Ba),  $2s^2 2p^2$  (C), and  $2s^2 2p^4$  (O) electrons have been used.

We performed finite temperature *ab initio* MD simulations to verify the structures and calculate properties. All MD simulations were performed in the isothermal-isobaric *NPT* ensemble (N—the number of particles, P—pressure, and T—temperature) with Langevin thermostat. The integration of the classical Newton’s equations of motion uses the Verlet algorithm, and the ground state search is evaluated within an efficient iterative matrix diagonalisation scheme and a Pulay mixer for each step. A time step for the integration was set to 1 fs. Simulations were performed in the temperature range of 300-2200 K with the step of 100 K. Aragonite was simulated at pressures of 0, 1.5, 3, 4.5, 6, and 10 GPa, calcite—at 0, 1.5, 3, 4.5, and 6 GPa, while strontianite and witherite—only at 0 GPa. The integration of Brillouin zone was performed using  $\Gamma$ -point only. A plane-wave cutoff energy was set to 350 eV. The crystallographic properties were derived from time averages taken over at least 10 ps.

The aragonite, strontianite, and witherite were approximated by the supercells, containing 360 atoms, which are  $3 \times 3 \times 2$  supercells of the unit cells used. For the calcite, an ortho-hexagonal cell with 480 atoms was used. The supercells of aragonite, strontianite, and witherite were produced based on structural data of De Villeris,<sup>12</sup> and supercell of calcite—based on data of Graf.<sup>19</sup>

The MD simulations of calcite show that the formation of disordered phases cc-IV and cc-V takes place at 1150 and 1300 K, respectively. In the experiment, the formation of the cc-V phase is observed at 1240 K.<sup>3</sup> The close correspondence of experimental and theoretical temperatures shows the correctness of the used technique.

# Results

In order to determine the stability fields of calcite and aragonite and to estimate the effect of rotational disorder on Gibbs energies of these structures, we will use the results of our calculations of calcite $\leftrightarrow$ aragonite equilibrium curve based on quasi harmonic approximation (QHA) and GGA. At present, this work is at the second stage of review in the Crystal Growth & Design journal. Similar results were previously obtained by Ukita and co-authors based on the same QHA approximation but with local density approximation.<sup>20</sup>

Our results accordingly with the results of Ukita and co-authors,<sup>20</sup> show that in P-T coordinates the curve, characterising the equilibrium between calcite and aragonite, has a gentle slope at a pressure around 0.6 GPa and temperature around 700 K (Figure 2). At temperatures below this point, the curve is approximated by the subvertical line, and it became more gentle at higher temperatures. The presence of the slope on the aragonite $\leftrightarrow$ calcite equilibrium curve is in well agreement with the results of numerous experiments and thermodynamic modeling.<sup>21,22</sup>

The MD simulations of calcite show that at ambient pressure the disordered phases cc-IV and cc-V are formed at 1150 and 1300 K, respectively. Simulations in the pressure range of 0–3 GPa shows that the equilibrium curves corresponding to the formation of these phases are nearly parallel with a slightly positive Clapeyron slope in P-T coordinates (Figure 2). At pressures above 3 GPa, the slope of cc-I $\leftrightarrow$ cc-IV curve became steeper, and as the result at 4.5 GPa the equilibrium curves cc-I $\leftrightarrow$ cc-IV and cc-IV $\leftrightarrow$ cc-V intersect. At higher pressures, the transition from cc-I to cc-V proceeds directly, without the formation of intermediate phase cc-IV (Figure 2).

Performed MD simulations have shown, that aragonite, as well as calcite, adopts a disordered state on heating. The formation of disordered aragonite, which we have named *disarag*, is observed through the whole investigated pressure range up to 10 GPa. However, at pressures below 1.5 GPa, it is preceded by the formation of hexagonal aragonite, named hexarag. The hexarag is the polymorph of CaCO<sub>3</sub> metastable in all P-T range,<sup>23</sup> its crystal

structure and a possible appearance in the experiment were analysed in our subsequent work. On further heating, hexarag transforms into the disordered state. At 0 GPa it takes place at 1500 K.

No hysteresis has been observed for the transition of aragonite to the disordered state. The transition occurs at the same temperature, independently whether we increase or decrease the temperature in simulation, at least within the used temperature step of 100 K.

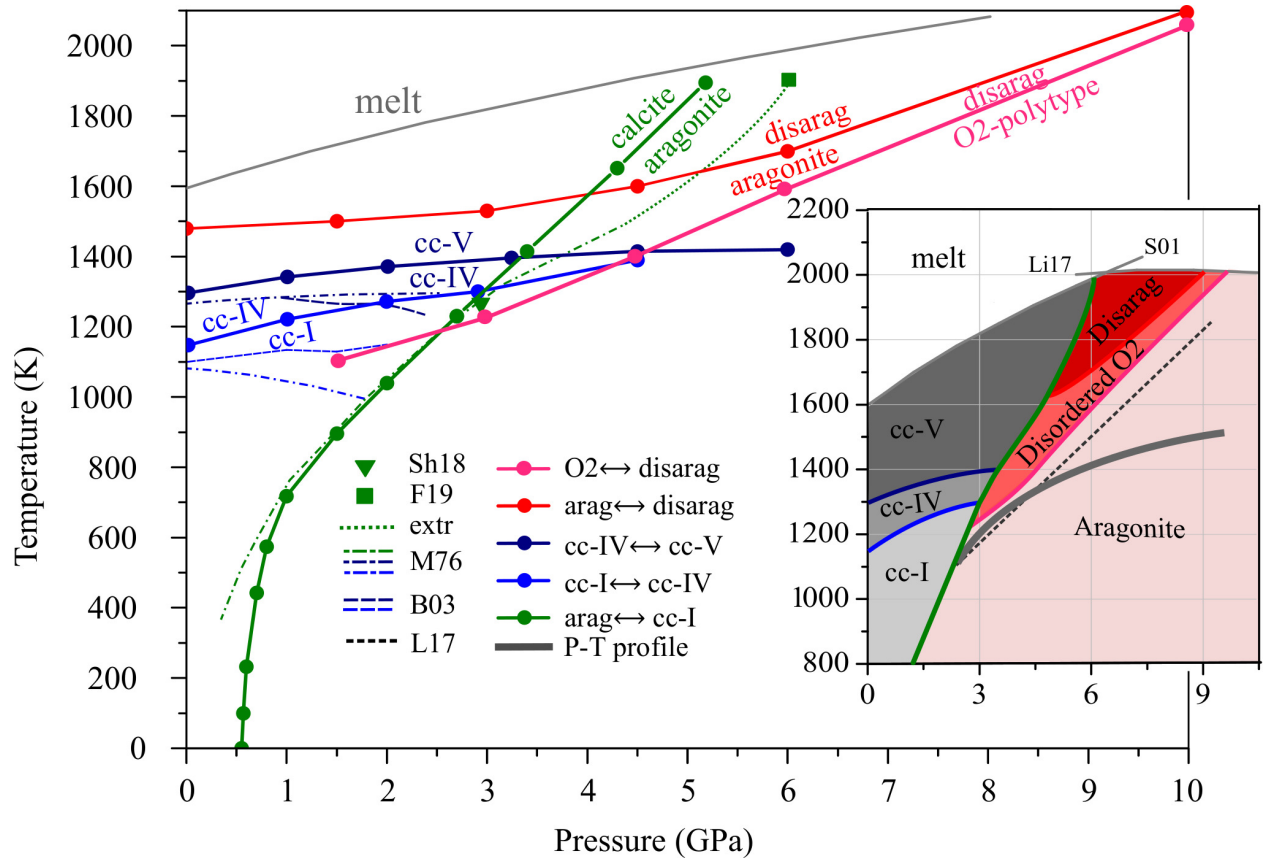


Figure 2: P-T diagram of  $\text{CaCO}_3$  with calculated equilibrium curves  $\text{cc-I} \leftrightarrow \text{cc-IV}$ ,  $\text{cc-IV} \leftrightarrow \text{cc-V}$ ,  $\text{calcite} \leftrightarrow \text{aragonite}$ ,  $\text{aragonite} \leftrightarrow \text{disarag}$ , and  $\text{O2} \leftrightarrow \text{disarag}$  shown as the solid lines and experimental curves of the phase transitions shown as the dashed and dotted lines; reference data are from: Sh18—, <sup>25</sup> F19—, <sup>26</sup> M76—, <sup>21</sup> B03—, <sup>27</sup> L17—, <sup>9</sup> S01—, <sup>8</sup> and Li17—, <sup>28</sup> extr—our extrapolation, P-T profile corresponds to the hot subducting slab<sup>29</sup>

At pressure above 1.5 GPa, hexarag is not observed in simulation, and aragonite directly transforms into the disordered state. The temperature of disordering is 1500–2100 K for the pressure range of 1.5–10 GPa (Figure 2). The performed simulations of strontianite and

witherite show, that these aragonite-like structures also adopt a disordered state on heating. At 0 GPa, strontianite became disordered at 1400 K, and witherite—at 1300 K. Thus, with increasing cation size from  $\text{Ca}^{2+}$  to  $\text{Ba}^{2+}$  the temperature of disordering decreases from 1500 K to 1300 K.

The transition of aragonite into the disordered state is accomplished by the abrupt increase of unit cell volume on  $\sim 2\%$  at a pressure of 4.5 GPa (Figure 3a). This value decreases with increasing pressure and at 10 GPa, the volume of disarag is only 0.5% higher than that of aragonite (Figure 3a). During transition, the  $b$  and  $c$  axes increase, while the  $a$ -axis decreases (Figure 3b). Our MD simulations show, that transition of calcite into the disordered state is also accomplished by the contraction of the  $a$ -axis, and expansion of the  $b$  and  $c$  axes, which is consistent with the experimental results<sup>30</sup> and modeling results obtained with empirical potentials.<sup>23</sup>

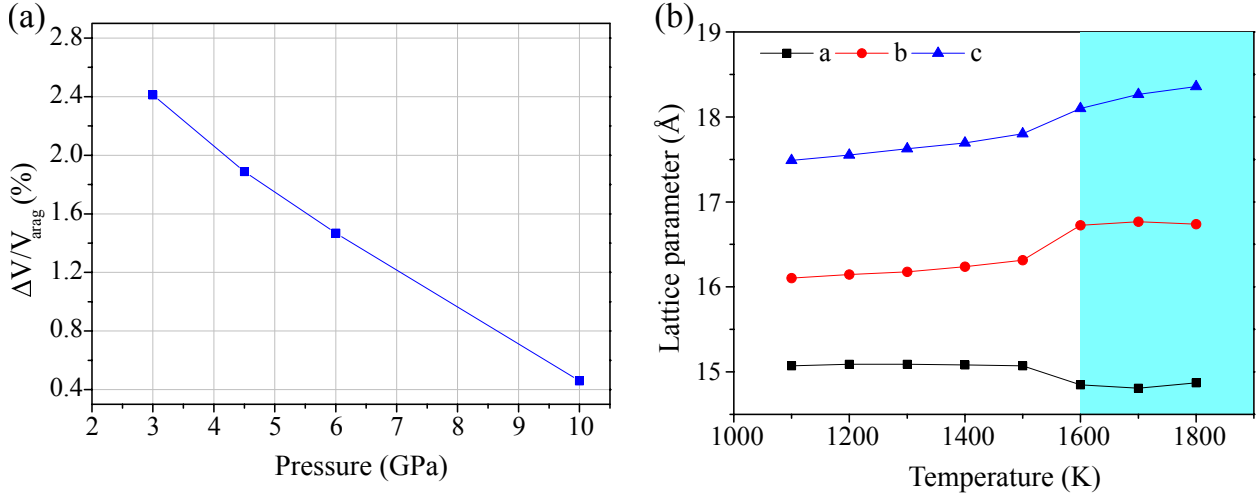


Figure 3: Ratio of unit cells volumes of aragonite and disarag at aragonite $\leftrightarrow$ disarag transformation in the range of 3–10 GPa (a) and dependencies of the cell parameters on temperature for aragonite and disarag (transition marked with the blue band) at 4.5 GPa (b)

At 300 K and 1.5 GPa, the amplitudes of rotations of  $\text{CO}_3$  groups about the molecular three-fold axis do not exceed  $5^\circ$  (Figure 4a), whereas with increasing temperature to 1100 K, this value increases to  $45^\circ$  (Figure 4b). At 1500 K, the amplitude reaches  $60^\circ$ , and disarag is formed. In disarag structure,  $\text{CO}_3$  triangles rotated on  $60^\circ$  and  $120^\circ$ , making the complete



circle around the molecular three-fold axis (Figure 4c). To reach the disordered state at 1500 K, the simulation run was increased to 30 000 MD-steps, as in a standard simulation with 15 000 steps, *disarag* was not formed (Figure 4c). At higher temperatures, 10 000 steps were enough to reach the disordered state. The distribution function of  $\text{CO}_3$  triangle orientations in *disarag* structure is characterised by the presence of six smooth maxima, three of which correspond to the initial equilibrium orientation, and three other—to the anti-orientation rotated on  $180^\circ$  (Figure 5a).

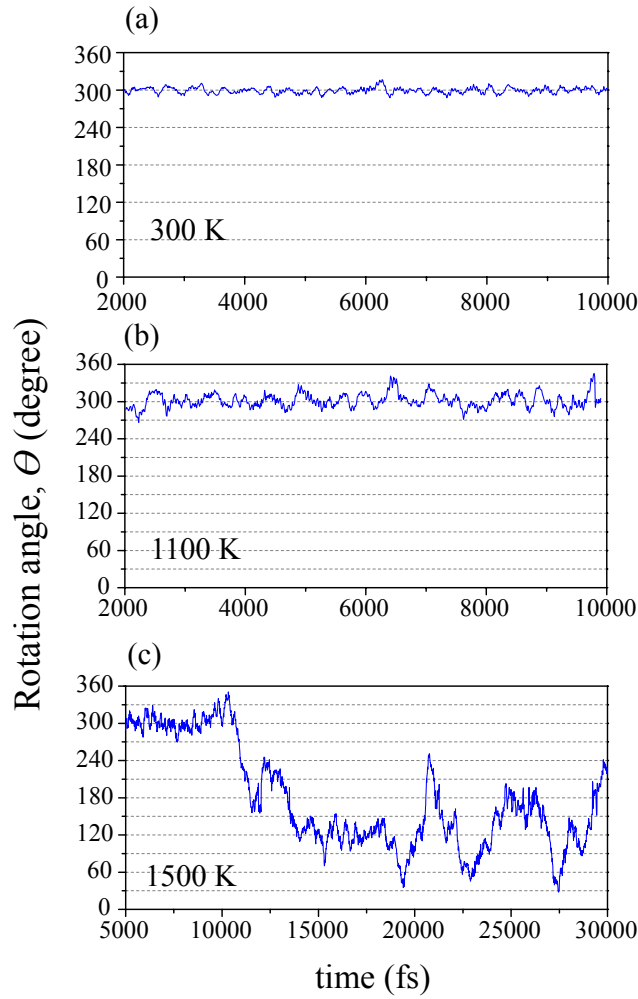


Figure 4: The dependence of the  $\theta$  angle (shown in Figure 1c) on time for aragonite, at 300, 1100, and 1500 K

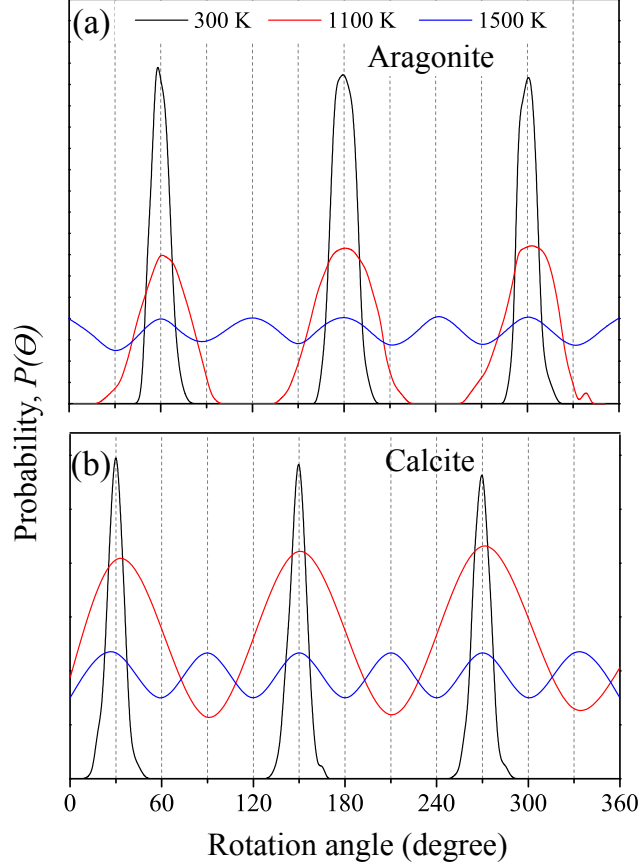


Figure 5: Probability distribution functions  $P(\theta)$  at 1.5 GPa and 300, 1100, and 1500 K for aragonite (a) and calcite (b)

The performed MD simulation of calcite at 1.5 GPa, shows that transitions to the disordered phases cc-IV and cc-V are observed at 1200 K and 1350 K, respectively. Thus, disordering of calcite is realised at lower temperatures than that of aragonite (Figure 2). The obtained distribution of  $\text{CO}_3$  triangles orientations for cc-IV and cc-V phases (Figure 5b) is consistent with available experimental<sup>3</sup> and theoretical results,<sup>23</sup> and it is similar to that of disarag (Figure 5a).

The performed simulations have shown, that the movement of the oxygen atoms about the three-fold axis of  $\text{CO}_3$  group is more complex than simple rotations in the plane of  $\text{CO}_3$  group. Such rotations are complicated by the intense librations about the axis orthogonal to the three-fold axis. Hereafter, the librations about the three-fold axis we will call *rotations*, and librations about perpendicular direction—*orthogonal librations*. As will be shown below,

orthogonal librations play a crucial role in the transition to the disordered state.

At 300 K and 1.5 GPa, the amplitude of orthogonal librations in aragonite structure reaches  $10^\circ$ , which is two times higher than the amplitudes of rotations at the same P-T conditions. Heating to 1500 K, increase the maximum amplitude of librations to  $45^\circ$  (Figure 6). Calcite is characterised by even higher amplitudes of orthogonal librations, with maximum value of  $15^\circ$  at 300 K and 1.5 GPa, and  $70^\circ$  at 1500 K, 1.5 GPa (Figure 6). The high values of amplitudes of orthogonal librations in calcite and aragonite are consistent with the almost text-book fact of the highly anisotropic thermal expansion of these structures with the high values of thermal expansion coefficients  $\alpha_c$ .<sup>31</sup> Noteworthy, that  $\alpha_a$  coefficient is negative for calcite<sup>31</sup> and positive for aragonite,<sup>30,32</sup> and performed simulations show more intense perpendicular librations for calcite structure.

At 1150 K, when cc-IV transforms to cc-V, the maximum value of amplitudes of orthogonal librations in calcite equals to  $45^\circ$ . The same value of amplitudes is observed for aragonite, when it transforms to disarag, at 1500 K and 1.5 GPa. Earlier it was suggested, that flipping of  $\text{CO}_3$  triangle in anti-orientation can be realised, once the amplitude of perpendicular librations reaches the critical angle.<sup>4</sup> To check this hypothesis, we determined the values of amplitudes of orthogonal librations at the transition to the disordered states for calcite and aragonite at pressures of 0, 1.5, 3, 4.5, and 6 GPa, and also for strontianite and witherite at ambient pressure. In all cases, we have obtained the same value of  $45^\circ$ . This confirms the crucial role of orthogonal librations in the transitions of carbonates to the disordered states. Noteworthy, that the value of critical angle does not depend no on the initial structure, no on the chemistry, and it is the same for both calcite and aragonite, and for  $\text{CaCO}_3$ ,  $\text{SrCO}_3$ , and  $\text{BaCO}_3$ .

At 1700 K and 1.5 GPa, the standard deviation of the amplitude of orthogonal librations is  $20^\circ$  for aragonite, and it is  $33^\circ$  for calcite (Figure 7), with maximum values of  $60^\circ$  and  $75^\circ$ , respectively (Figure 6). At such high values, the average distribution of oxygen atoms of  $\text{CO}_3$  atomic group actually loses the flat form and became elliptical. This sufficient deviation from

the flat distribution of C–O bond orientations, was earlier highlighted by G.Cai et al, based on their results of MD simulations with empirical potentials.<sup>33</sup> Meanwhile, according to the refinements performed based on X-ray diffractions experiments with calcite, the inclination of the orbital of oxygen atoms from the flat form is within  $\pm 7.7^\circ$ .<sup>3</sup> The obtained results show, that suggested from the experimental data umbrella form of the orbital of oxygen atoms does not characterise the movement of the separate  $\text{CO}_3$  group, and likely it is the artifact of the initial model used for the refinement.

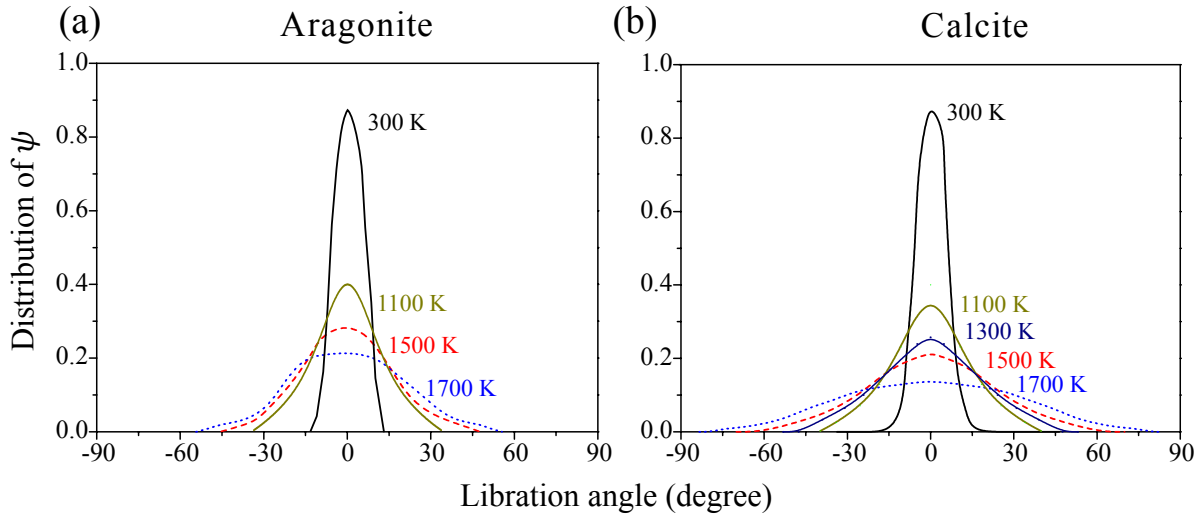


Figure 6: Probability function for the distribution functions of the amplitudes of orthognoal librations ( $\psi$  angle shown on Figure 1c) of  $\text{CO}_3$  triangle for aragonite (a) and calcite (b) at 1.5 GPa and 300, 1100, 1300, 1500, 1700 K

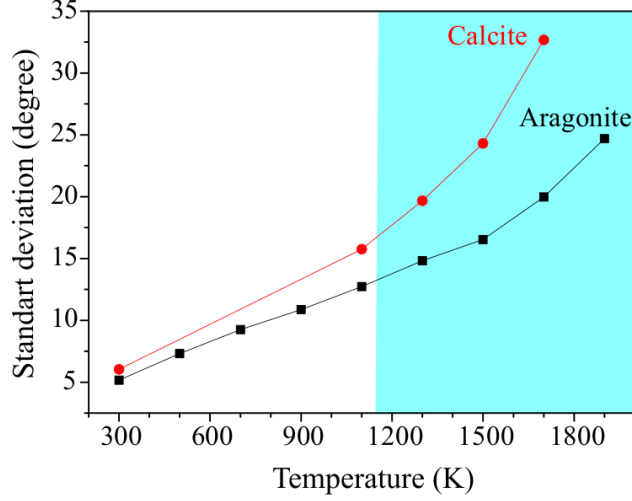


Figure 7: Dependence of the standard deviation of probability function shown on Figure 6 on temperature for aragonite and calcite at 1.5 GPa

Described librations of the  $\text{CO}_3$  group are also accompanied by the intensive movement of the whole  $\text{CO}_3$  triangle along the  $c$ -axis. In disarag, the amplitude of such a movement reaches  $1/3$  distance between cp Ca-layers, which facilitates hoppings of  $\text{CO}_3$  triangles between adjacent sublayers. As  $\text{CO}_3$  groups of adjacent sublayers have opposite orientations, the change of the sublayer is bounded to the change of orientation. This is illustrated in Figure 8. The dependence of  $z$ -coordinate on time (Figure 8a) shows that in the range of 5000–15000 MD steps, the  $\text{CO}_3$  triangle stays within the same sublayer. The average coordinate of *triangle1* is nearly  $2/3$  and that of *triangle2*—nearly  $1/3$ . After 15 000 MD-steps, the  $z$ -coordinate of *triangle1* abruptly drops to  $1/3$ , while that of *triangle2* abruptly increases to  $2/3$ . This mean, that *triangle1* have hopped from the upper sublayer to the lower sublayer, and *triangle2*—in the opposite direction. The change of the sublayer is accompanied by the rotation of the  $\text{CO}_3$  group on  $60^\circ$ , as can be seen from the dependence of rotation angle on time (Figure 8b). Within the new sublayer, the  $\text{CO}_3$  triangle stays in the range of 15000–25000 steps, and then the next hop between sublayers takes place (Figure 8a). This hop is also accompanied by the rotation on  $60^\circ$  (Figure 8a). Thus, disarag structure presents the new type of disorder related to the position of  $\text{CO}_3$  group along the  $c$ -axis. This type of disorder is strictly bound to the disorder of the orientation of the  $\text{CO}_3$  group. In the MD

simulation of witherite  $\text{BaCO}_3$ , the transition to the hexagonal symmetry has been noticed by Cai and co-authors.<sup>33</sup> Assuming rotational disorder of the carbonate group and its average position midway between cp-layers, the symmetry of the disarag structure will be also hexagonal the same as that of NiAs.

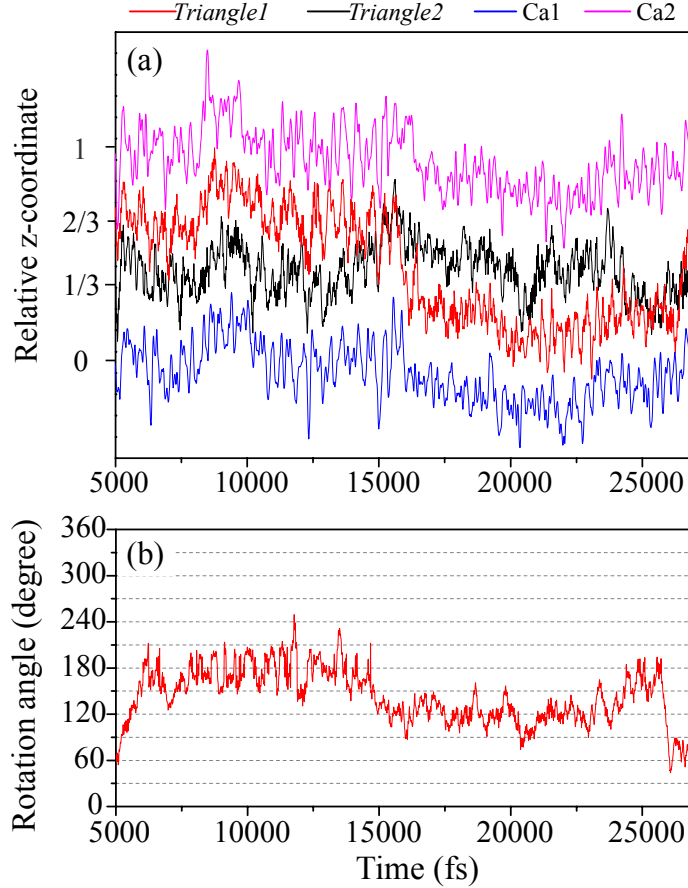


Figure 8: The dependence of relative  $z$ -coordinate on time for two Ca atoms (Ca1 and Ca2) lying in the adjacent cp-layers and for two carbon atoms, belonging to  $\text{CO}_3$  groups called *triangle1* and *triangle2* (a) and the dependence of *triangle1* orientation on time (b); pressure equals to 4.5 GPa, temperature to 1600 K

To account for the effect of defects on the temperature of aragonite disordering, the MD simulations of the aragonite structure twinned by each second (110) plane have been performed. As it was mentioned in the Introduction, aragonite structure twinned in such a manner corresponds to the so-called *O2-polytype*.<sup>15</sup> In this case, *O2-polytype* is the theoretical approximation of the real structure with the disordered sequence of twinned boundaries,

observed in experiments.<sup>15</sup> The simulation of the polytypes with the less density of twin boundaries and even more the structure with the disordered sequences of twin boundaries is too expansive in sense of computer resources.

Performed simulations have shown, that as well as aragonite, O2-polytype became disordered at high temperatures. The temperature of disordering is lower than that of aragonite. At pressures below 3 GPa, the difference in transition temperatures for twinned (O2-polytype) and normal (aragonite) structures is around 400 K (Figure 2). The difference decreases with pressure and at 9 GPa it is just 100 K.

## Discussion & Conclusions

Our paper is concerned with the physical chemistry of aragonite phases. We compute their thermodynamics using *ab initio* methods and derive the phase relations between the various structural modifications of  $\text{CaCO}_3$ . The thermodynamics is a natural part of physical chemistry, especially when the emphasise is on the bulk properties rather than on the microscopic characteristics. While it is obvious that the structure and thermodynamics are intimately connected, the major concern in this paper are the bulk properties of the new predicted phase.

The obtained results, describing equilibrium curves between calcite and its disordered forms cc-IV and cc-V, show that the curves corresponding to the transitions  $\text{cc-I} \leftrightarrow \text{cc-IV}$  and  $\text{cc-IV} \leftrightarrow \text{cc-V}$  are almost parallel within the P-T stability field of calcite. This is in full agreement with experimental data on  $\text{cc-IV} \rightarrow \text{cc-V}$  transition,<sup>21,27,34</sup> shown on Figure 2. A comparison of theoretical and experimental results on  $\text{cc-I} \rightarrow \text{cc-IV}$  transition is not so straightforward, as the appearance of the intermediate phase cc-IV phase on the heating history of the sample and maybe on some other factors. Based on differential thermal analysis experiments, Mirwald suggests the negative Clapeyron slope for the equilibrium curve  $\text{cc-I} \leftrightarrow \text{cc-IV}$ .<sup>21</sup> This assumes an expansion of the cc-IV stability field with pressure

increasing (Figure 2). However, the results of Mirwald depend on the direction, in which the phase boundary is crossed over. In case, when the boundary is crossed in the direction of cooling, the Clapeyron slope is close to zero, which is consistent with our results. Similar difficulties were revealed by Bagdassarov and co-authors in their measurements of electrical impedance.<sup>27</sup> According to these results, the phase boundary is subhorizontal with a slightly positive slope (Figure 2). Thus, the stability field of cc-IV not expands but contracts with pressure, and at a pressure above 3 GPa, this phase disappears, which is reproduced in our simulations. According to our results, the triple point of cc-I, cc-IV, and cc-V co-existence is located at 5 GPa and 1400 K. At pressures above this value, cc-IV does not form, and cc-I directly transforms into cc-V on heating. However, this can not be observed in experiments, as calcite is metastable at such high pressures.

For *ab initio* determination of the stability field of disarag the comparison of Gibbs energies of disordered phases of calcite (in ordered and disordered forms) and disarag is necessary. As *ab initio* calculation of the entropy of dynamically disordered phase is the complex task, we will use the results of quenching experiments for determination of equilibrium curve between calcite and aragonite. In this case, we assume that disordered aragonite is quenched in aragonite and disordered calcite—in calcite. The results of our unpublished calculations of calcite $\leftrightarrow$  equilibrium curve, which we have mentioned above, will be used for comparison.

The calculated equilibrium curve of calcite  $\leftrightarrow$  aragonite transition obtained within QHA is in remarkable agreement with the results of quenching laboratory experiments,<sup>25,26</sup> at temperatures where disordering does not take place, i.e. before the formation of cc-IV (Figure 2). After the formation of cc-IV, the theoretical results based on QHA, which do not account for the rotation of CO<sub>3</sub> triangles, start to deviate from the experimental ones. The disordering of calcite structure decreases its Gibbs energy, thus stabilising calcite against aragonite. This results in the expansion of the calcite stability field in comparison with the stability field determined with QHA. The aragonite disordering partially or completely compensates this energy difference, and above 1600 K experimental curve of calcite  $\leftrightarrow$  aragonite equilibrium



is parallel to the theoretical one, as it is shown in Figure 2.

The found disordered phase of aragonite has a distinct stability field on P-T diagram (Figure 2). In the pressure range of 3–10 GPa, formation of disarag precedes melting on 0–500 K. The transition from aragonite to disarag occurs at temperatures at least 300 K higher than the temperatures of the slab subducting in the upper mantle. The curve corresponding to the transition from O2-polytype to disarag almost touches the hot geotherm of the subducting slab at 3.5 GPa, which corresponds to the depth of nearly 110 km.<sup>29</sup> This assumes the possibility of disarag formation at these depths, although in the narrow pressure range.

Assuming that disarag is quenched in aragonite, as well as disordered calcite is quenched in calcite,<sup>3</sup> it can be concluded, that formation of disarag is consistent with available results of quenching experiments. However, only *in situ* experiments can approve the formation of disarag. Results of these experiments consistently show transition to the new phase at P-T parameters, where aragonite transforms to disarag in our simulations (Figure 2). The temperature difference between transition observed in experiments<sup>8,9</sup> and in our simulations is just 200 K (above 6 GPa) for normal aragonite, and 100 K—for twinned aragonite. However, the structure of the phase obtained in the experiment is debated. This is partly, due to the mentioned lack of diffraction information. The hypotheses, suggested cc-V,<sup>8</sup> some disordered phase<sup>9</sup> and amorphous phase<sup>5</sup> have been proposed. The appearance of cc-V phase at these P-T parameters is inconsistent with the results of quenching experiments (Figure 2). The unique character of amorphisation observed at high pressures and temperatures and absence of amorphisation in the other experiments performed at the same P-T parameters<sup>8,9</sup> states the question about the thermodynamic stability of the amorphous phase and reproducibility of this experiment. Summarising, we conclude that among known phases of  $\text{CaCO}_3$ , disarag is the best candidate for the explanation of structural changes observed in the experiment. However, the appearance of the other structure in ordered or disordered form can not be excluded.

In conclusion, we would like to highlight, that earlier only calcite-type disordered struc-

tures were known. These are  $\text{CaCO}_3\text{-IV}$ ,  $\text{CaCO}_3\text{-V}$ , rhombohedral  $\text{SrCO}_3$ , rhombohedral  $\text{BaCO}_3$ , and cubic  $\text{BaCO}_3$ . Here, we have shown that aragonite also became disordered, although at high pressures. Even at ambient pressure, the solution of disordered structures is a nontrivial task, as the example of disordered calcite has shown. This task became sufficiently more complex at conditions of high-pressures, where diffraction data are of intrinsically low quality. In this case, the application of the MD simulation technique became absolutely necessary step of the successful determination and characterisation of the structure.

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# Graphical TOC Entry

