

The Art of Computer Simulation

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C3-101 9:40 - 10:35

"The Art of Doing Science and Engineering" by Richard W. Hamming

Hamming code,
Hamming window,
Hamming distance,
etc.



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"The Art of Doing Science and Engineering" by Richard W. Hamming

► Contents

...

Chap 18 Simulation I

Chap 19 Simulation II

Chap 20 Simulation III

...



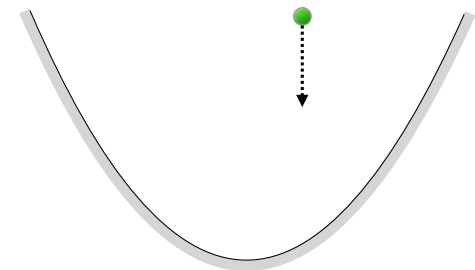
"A simulation is the answer to the question, "What if...?"

R. Hamming

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Let's make a simple simulation

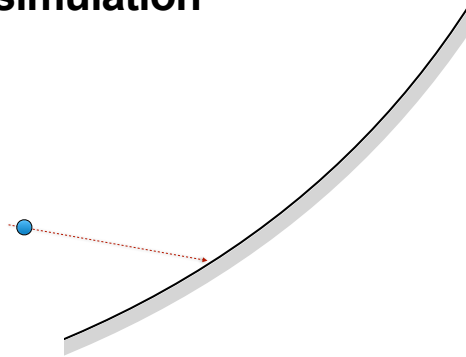
- What if we drop an elastic ball onto a parabola-curved floor?



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Let's make a simple simulation

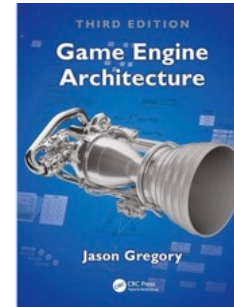
- Collision detection is a key.
 - Video games



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Let's make a simple simulation

- Collision detection is a key.
 - Video games



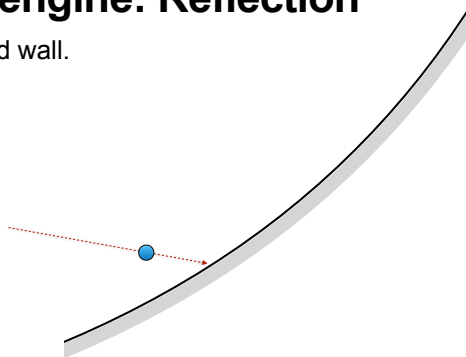
13. Collision and Rigid Body Dynamics

- 13.1 Do You Want Physics in Your Game?
- 13.2 Collision/Physics Middleware
- 13.3 The Collision Detection System
- 13.4 Rigid Body Dynamics
- 13.5 Integrating a Physics Engine into Your Game
- 13.6 Advanced Physics Features

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An element of game engine: Reflection

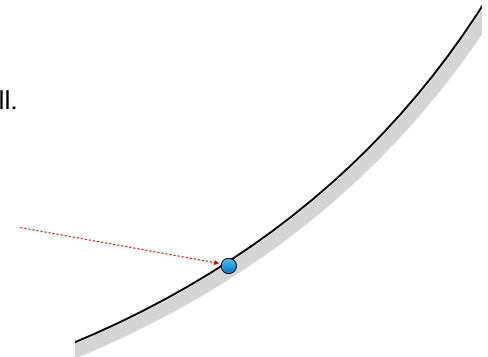
- A particle located outside a curved wall.



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Reflection

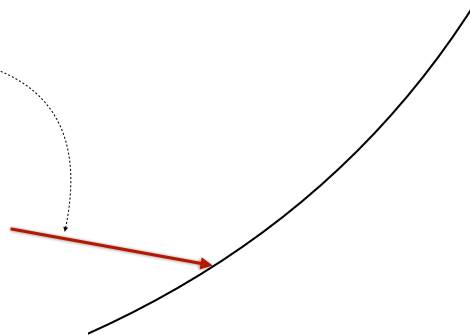
- It impinges on the wall.
 - It (slightly) invades into the wall.
 - Collision



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Reflection

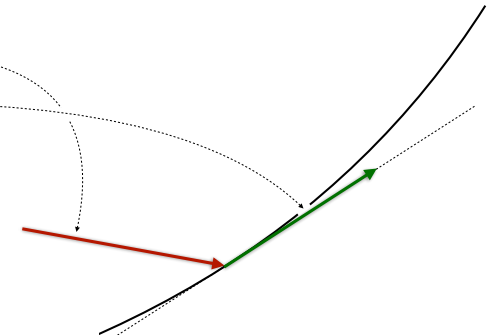
- Incoming velocity vector



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Reflection

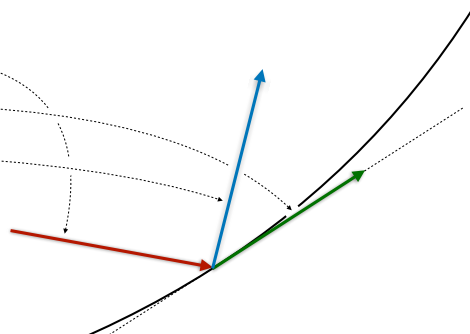
- Incoming velocity vector
- Tangential vector



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Reflection

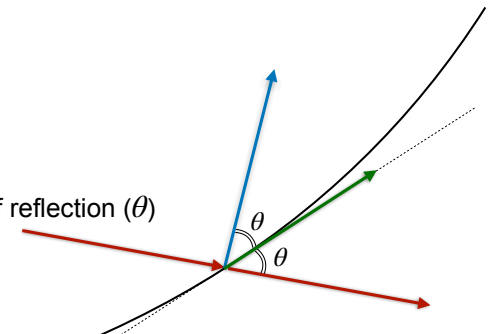
- Incoming velocity vector
- Tangential vector
- Reflected velocity vector



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Reflection

- Incoming velocity vector
- Tangential vector
- Reflected velocity vector
- Angle of incidence (θ) = angle of reflection (θ)

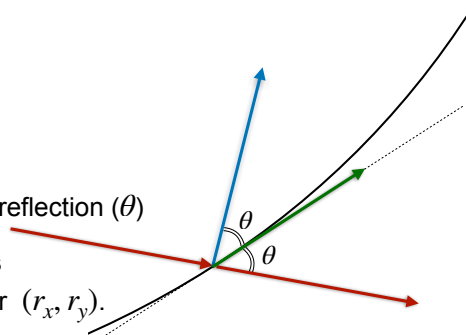


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Reflection

- Incoming velocity vector
- Tangential vector
- Reflected velocity vector
- Angle of incidence (θ) = angle of reflection (θ)
- Reflected (blue) vector (b_x, b_y) is 2θ rotation of incident (red) vector (r_x, r_y).

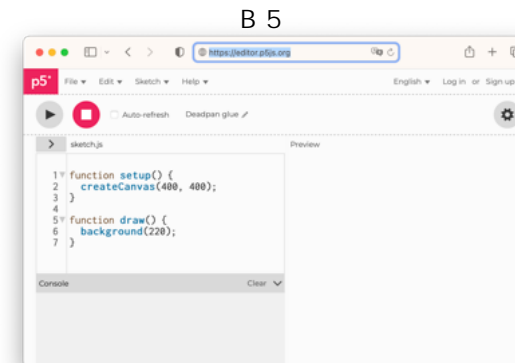
$$\begin{pmatrix} b_x \\ b_y \end{pmatrix} = \begin{pmatrix} \cos 2\theta & -\sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} \begin{pmatrix} r_x \\ r_y \end{pmatrix}$$



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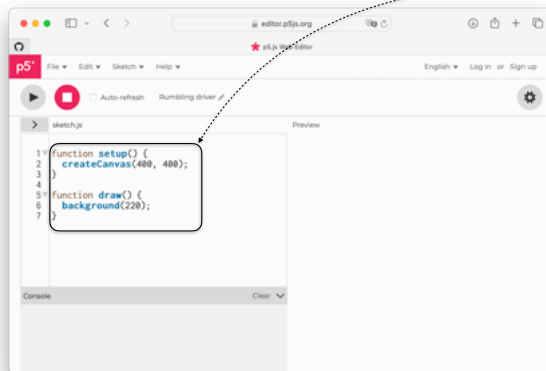
Let's try with P5.js on Web Editor

<https://editor.p5js.org/>



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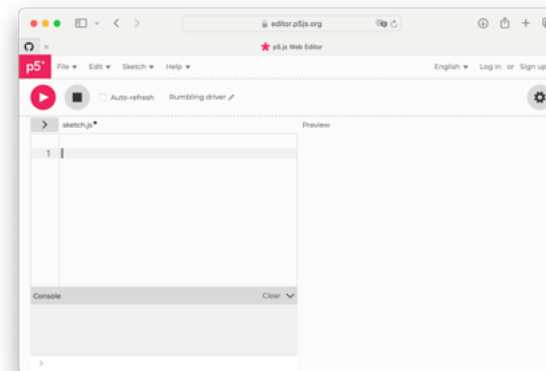
Let's try with P5.js Web Editor



- Clear the default code and leave this area blank by select all & delete.

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Let's try with P5.js Web Editor



- Preparation is done.

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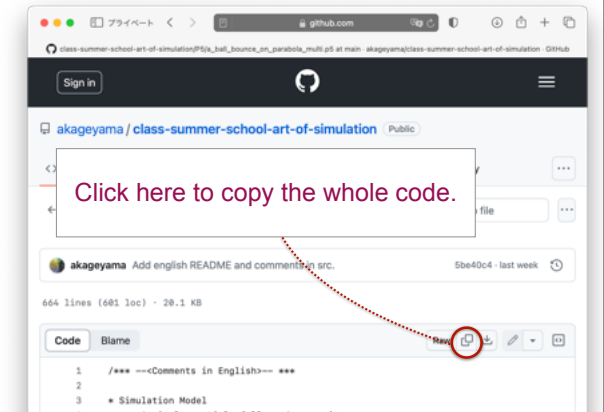
Get a sample code

<https://tinyurl.com/kobe-sc01>

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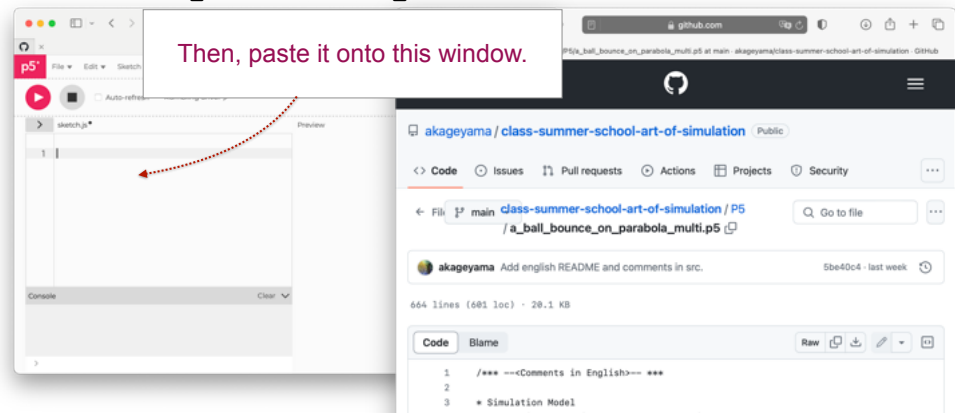
Get a sample code

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Get a sample code

<https://tinyurl.com/kobe-sc01>



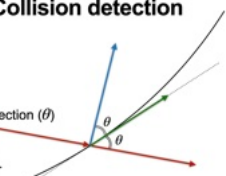
Reflection part

```
547 let dot_product
548 vecIx_norme
549 vecIy_norme
550 let angle_betwe
551 dot_product
552 );
553 let angle_for_reflection =
554   2 * angle_between_vecI_normed_and_vecT_normed;
555 let cos_angle = cos(angle_for_reflection);
556 let sin_angle = sin(angle_for_reflection);
557 let vecRx = cos_angle * vecIx - sin_angle * vecIy;
558 let vecRy = sin_angle * vecIx + cos_angle * vecIy;
559 ball.vx = vecRx;
560 ball.vy = vecRy;
```

Simple game engine: Collision detection

- Incoming velocity vector
- Tangential vector
- Reflected velocity vector
- Angle of incidence (θ) = angle of reflection (θ)
- Reflected vector (b_x, b_y) is 2θ rotation of incident vector (r_x, r_y).

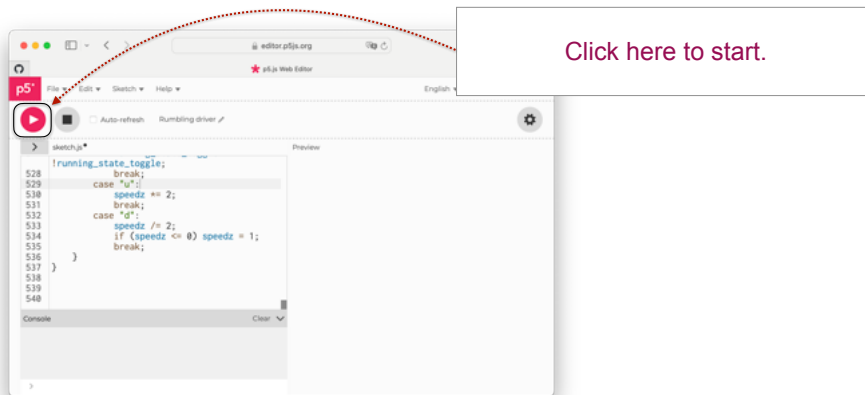
$$\begin{pmatrix} b_x \\ b_y \end{pmatrix} = \begin{pmatrix} \cos 2\theta & -\sin 2\theta \\ \sin 2\theta & \cos 2\theta \end{pmatrix} \begin{pmatrix} r_x \\ r_y \end{pmatrix}$$



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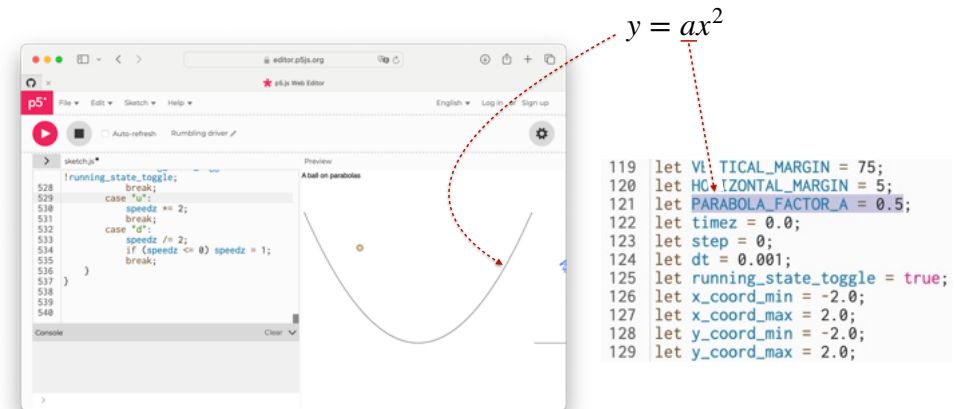
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Sample code



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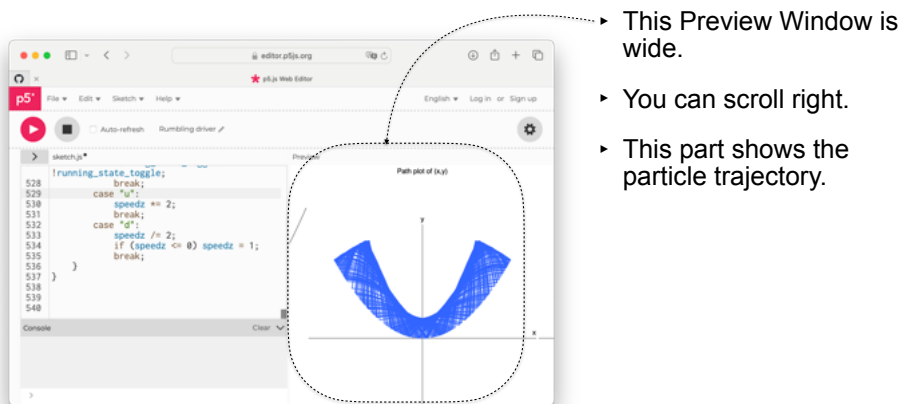
Sample code



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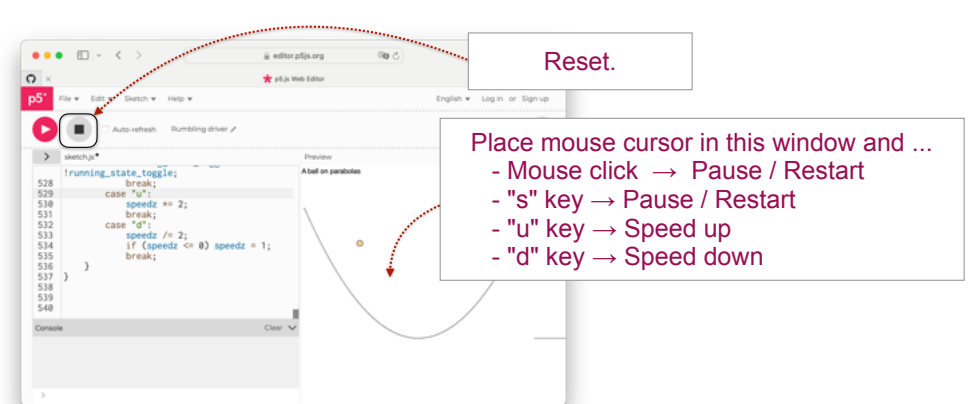
Let's try with P5.js Web Editor

<https://editor.p5js.org/>



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Usage



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Code explanation

```

1  /** --<Comments in English>-- **
2
3  * Simulation Model
4    - A single particle falls under gravity.
5    - There is a lower parabola given by:
6       $y = a \cdot x^2$ 
7    - The particle reflects upon hitting the
      parabola.
8
9  * Definition of Variables
10   - The coefficient of the parabola, a, is
11     PARABOLA_FACTOR_A.

```

```

381 windowz[1].label_y_axis("y");
382 windowz[2].draw_axes_x1_x2();
383 windowz[2].label_x_axis("x");
384 windowz[2].label_y_axis("vx");
385
386 header.title(" A ball on parabolas", LEFT);
387 header.title(" Path plot of (x,y)", CENTER);
388 header.title("Poincare map of (x,vx) on vy=0",
389             ", RIGHT); // 自由落下、下に凸な包絡線
390
391 ball = new GeneralCoords(
392     1.0, // Initial position x
393     0.0, // Initial velocity vx
394     1.5, // Initial position y
395     0.0 // Initial velocity vy
396 );
397 //

```

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Numerical experiments

```

381 windowz[1].label_y_axis("y");
382 windowz[2].draw_axes_x1_x2();
383 windowz[2].label_x_axis("x");
384 windowz[2].label_y_axis("vx");
385
386 header.title(" A ball on parabolas", LEFT);
387 header.title(" Path plot of (x,y)", CENTER);
388 header.title("Poincare map of (x,vx) on vy=0",
389             ", RIGHT); // 自由落下、下に凸な包絡線
390
391 ball = new GeneralCoords(
392     0.0, // Initial position x
393     2.0, // Initial velocity vx
394     1.5, // Initial position y
395     0.0 // Initial velocity vy
396 );
397 //

```

Try changing the initial conditions.

For example,

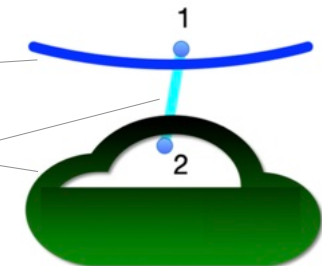
Note: Do not place the ball blow the parabola.

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Simulation of chaotic motion

Problem setting

- ▶ Bent metal rail
- ▶ A heavy object with a handle
- ▶ Connected by a rubber-rope
- ▶ Both ends (1 and 2) are pulleys



- The pulleys will show a kind of oscillating motion driven by the tension force.
- **Problem:** How do they move?

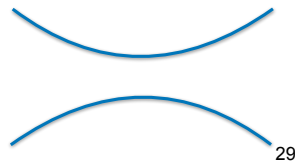
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Simplification

"Everything should be made as simple as possible, but not simpler."

A. Einstein

- ▶ No twisting motion (motion in a plane, i.e., 2-D simulation)
- ▶ Same pulleys.
- ▶ No friction.
- ▶ Rubber rope is a linear spring with no mass.
- ▶ The cable and handle are fixed.
- ▶ Symmetric parabolas.

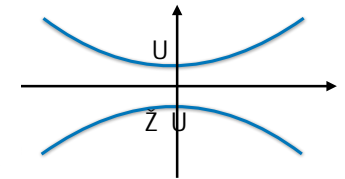


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Mathematical model

- ▶ m : Mass of Pulley 1 and pulley 2
- ▶ k : Spring constant of the rubber-rope. (Natural length is zero.)
- ▶ $2c$: Closest distance between the cable and handle.

$$y(x) = \begin{cases} y_1 = +x^2/c + c \\ y_2 = -x^2/c - c \end{cases}$$



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Equations of motion

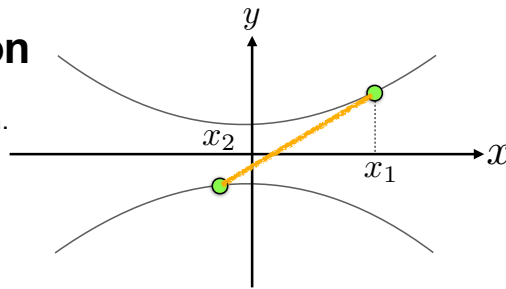
- ▶ See Appendix for the derivation.

$$\frac{dx_1}{dt} = v_1,$$

$$\frac{dx_2}{dt} = v_2,$$

$$\frac{dv_1}{dt} = \frac{-1}{4x_1^2 + c^2} \left\{ 4x_1 v_1^2 + \frac{kc^2}{m} \left(x_1 - x_2 + 2x_1 \frac{x_1^2 + x_2^2 + 2c^2}{c^2} \right) \right\},$$

$$\frac{dv_2}{dt} = \frac{-1}{4x_2^2 + c^2} \left\{ 4x_2 v_2^2 + \frac{kc^2}{m} \left(x_2 - x_1 + 2x_2 \frac{x_1^2 + x_2^2 + 2c^2}{c^2} \right) \right\}.$$



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Before you compute...

- ▶ Dimensional analysis
- ▶ Linear analysis
- ▶ Symmetry analysis
- ▶ Conservation quantities

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Dimensional analysis

- ▶ This system contains three parameters, c , k , and m .
- ▶ k and m appear always in the form k/m .
- ▶ The time scale of the system is determined by $\omega = \sqrt{\frac{k}{m}}$.
 - No other quantity with the dimension of time (measured by second).
- ▶ If you can find an oscillation solution in this problem under a specific circumstance, its frequency would be $\mathcal{O}(\omega)$.

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Variable normalization

- ▶ We re-scale positions by c .

$$x_i = c \tilde{x}_i$$

$$v_i = c \tilde{v}_i$$

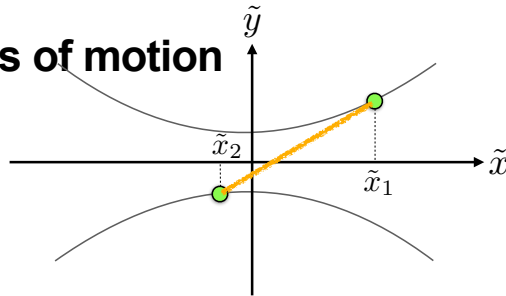
\tilde{x} : Non-dimensional position

\tilde{v}_i : Re-scaled velocity

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Re-scaled equations of motion

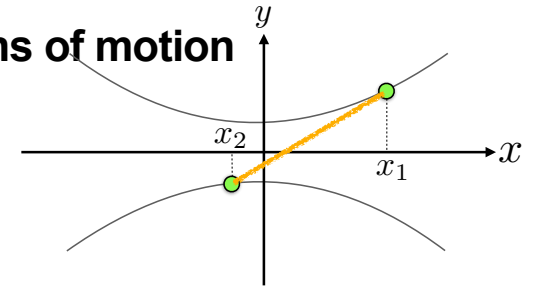
$$\begin{aligned} \frac{d\tilde{x}_1}{dt} &= \tilde{v}_1, \\ \frac{d\tilde{x}_2}{dt} &= \tilde{v}_2, \\ \frac{d\tilde{v}_1}{dt} &= \frac{-1}{1+4\tilde{x}_1^2} \{4\tilde{x}_1 \tilde{v}_1^2 + \omega^2(\tilde{x}_1 - \tilde{x}_2 + 2\tilde{x}_1(\tilde{x}_1^2 + \tilde{x}_2^2 + 2))\}, \\ \frac{d\tilde{v}_2}{dt} &= \frac{-1}{1+4\tilde{x}_2^2} \{4\tilde{x}_2 \tilde{v}_2^2 + \omega^2(\tilde{x}_2 - \tilde{x}_1 + 2\tilde{x}_2(\tilde{x}_1^2 + \tilde{x}_2^2 + 2))\}. \end{aligned}$$



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Re-scaled equations of motion

$$\begin{aligned} \frac{dx_1}{dt} &= v_1, \\ \frac{dx_2}{dt} &= v_2, \\ \frac{dv_1}{dt} &= \frac{-1}{1+4x_1^2} \{4x_1 v_1^2 + \omega^2(x_1 - x_2 + 2x_1(x_1^2 + x_2^2 + 2))\}, \\ \frac{dv_2}{dt} &= \frac{-1}{1+4x_2^2} \{4x_2 v_2^2 + \omega^2(x_2 - x_1 + 2x_2(x_1^2 + x_2^2 + 2))\}. \end{aligned}$$



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Linear analysis

- If x_i is small, then we can ignore nonlinear terms of x_i .

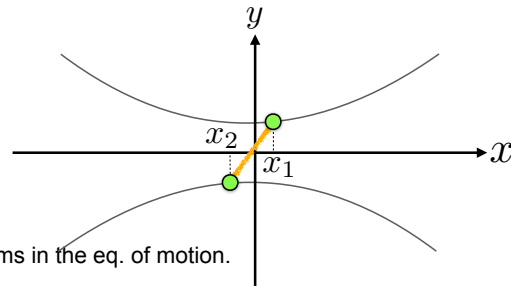
Example:

When $(x_1, x_2) = (0.1, -0.05)$

$$|x_1|^2 = 0.01$$

$$|x_2|^2 = 0.0025$$

$$|x_1 x_2| = 0.005$$



- Linearization Leave only linear terms in the eq. of motion.

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Linear analysis

$$\frac{dx_1}{dt} = v_1,$$

$$\frac{dx_2}{dt} = v_2,$$

$$\frac{dv_1}{dt} = \frac{-1}{1 + 4x_1^2} \{ 4x_1 v_1^2 + \omega^2 (x_1 - x_2 + 2x_1(x_1^2 + x_2^2 + 2)) \},$$

$$\frac{dv_2}{dt} = \frac{-1}{1 + 4x_2^2} \{ 4x_2 v_2^2 + \omega^2 (x_2 - x_1 + 2x_2(x_1^2 + x_2^2 + 2)) \}.$$

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Linear analysis

$$\frac{dx_1}{dt} = v_1$$

$$\frac{dx_2}{dt} = v_2$$

$$\frac{dv_1}{dt} = -\omega^2 (5x_1 - x_2)$$

$$\frac{dv_2}{dt} = -\omega^2 (5x_2 - x_1)$$

39

Linear analysis

$$\frac{dx_1}{dt} = v_1$$

$$\frac{dx_2}{dt} = v_2$$

$$\frac{dv_1}{dt} = -\omega^2 (5x_1 - x_2)$$

$$\frac{dv_2}{dt} = -\omega^2 (5x_2 - x_1)$$



$$\frac{d^2x_1}{dt^2} = -\omega^2 (5x_1 - x_2)$$

$$\frac{d^2x_2}{dt^2} = -\omega^2 (5x_2 - x_1)$$

40

Linear analysis

$$\frac{d^2 x_1}{dt^2} = -\omega^2 (5x_1 - x_2)$$

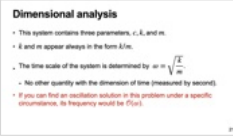
$$\frac{d^2 x_2}{dt^2} = -\omega^2 (5x_2 - x_1)$$

► Solutions

$$x_1(t) = c_1 \cos(2\omega t + c_2) + c_3 \cos(\sqrt{6}\omega t + c_4)$$

$$x_2(t) = c_1 \cos(2\omega t + c_2) - c_3 \cos(\sqrt{6}\omega t + c_4)$$

► Harmonic oscillations with frequencies 2ω and $\sqrt{6}\omega$



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Linear solutions

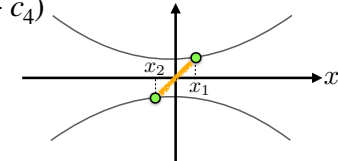
$$x_1(t) = c_1 \cos(2\omega t + c_2) + c_3 \cos(\sqrt{6}\omega t + c_4)$$

► Case $c_1 = 0$

$$x_2(t) = c_1 \cos(2\omega t + c_2) - c_3 \cos(\sqrt{6}\omega t + c_4)$$

$$x_1(t) = -x_2(t) = c_3 \cos(\sqrt{6}\omega t + c_4)$$

Harmonic oscillation of $\sqrt{6}\omega$



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Linear solutions

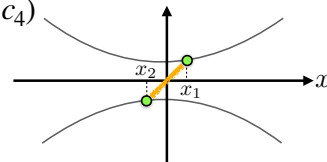
$$x_1(t) = c_1 \cos(2\omega t + c_2) + c_3 \cos(\sqrt{6}\omega t + c_4)$$

► Case $c_1 = 0$

$$x_2(t) = c_1 \cos(2\omega t + c_2) - c_3 \cos(\sqrt{6}\omega t + c_4)$$

$$x_1(t) = -x_2(t) = c_3 \cos(\sqrt{6}\omega t + c_4)$$

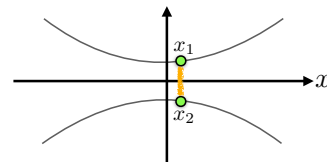
Harmonic oscillation of $\sqrt{6}\omega$



► Case $c_3 = 0$

$$x_1(t) = x_2(t) = c_1 \cos(2\omega t + c_2)$$

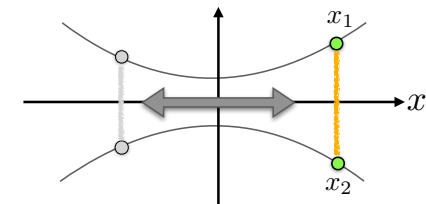
Harmonic oscillation of 2ω



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Nonlinear (large amplitude) symmetric solution 1

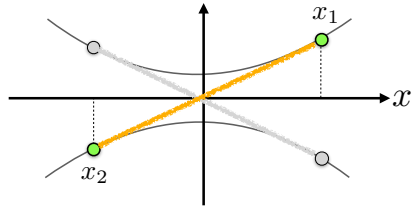
► If $(x_1, v_1) = (x_2, v_2)$ at $t = 0$



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Nonlinear (large amplitude) symmetric solution 2

- if $(x_1, v_1) = -(x_2, v_2)$ at $t = 0$,



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Conservation quantity

- Total energy $H = K + U$
- K : Kinetic energy of the two pulleys
- U : Potential energy of the rubber rope

$$H = \frac{m}{2} \{ \dot{x}_1^2 (1 + 4x_1^2) + \dot{x}_2^2 (1 + 4x_2^2) \} + \frac{k}{2} \{ (x_1 - x_2)^2 + (x_1^2 + x_2^2 + 2)^2 \}$$

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Simulation model

- State of the system
- Discretization
- Visualization

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State of the system

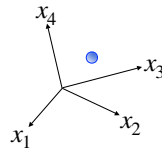
- (x_1, x_2, v_1, v_2)
 - x_1, x_2 : Position of point mass 1 and 2
 - v_1, v_2 : Velocity of point mass 1 and 2

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State of the system

- ▶ (x_1, x_2, v_1, v_2)
 - x_1, x_2 : Position of point mass 1 and 2
 - v_1, v_2 : Velocity of point mass 1 and 2
- ▶ Combine them → General coordinates
- ▶ 4-dimensional space ("phase space")
- ▶ "State" \Leftrightarrow A point in the 4-D space

$$\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix} = \begin{pmatrix} x_1 \\ x_2 \\ v_1 \\ v_2 \end{pmatrix}$$

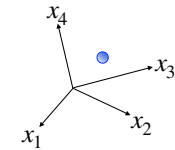


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State of the system

- ▶ (x_1, x_2, v_1, v_2)
 - x_1, x_2 : Position of point mass 1 and 2
 - v_1, v_2 : Velocity of point mass 1 and 2
- ▶ Combine them → General coordinates
- ▶ 4-dimensional space ("phase space")
- ▶ "State" \Leftrightarrow A point in the 4-D space
- ▶ "Time development of state"
 - \Leftrightarrow Directed curve in the 4-D space

$$\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix} = \begin{pmatrix} x_1 \\ x_2 \\ v_1 \\ v_2 \end{pmatrix}$$

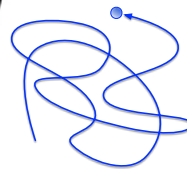
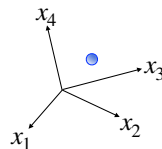


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State of the system

- ▶ (x_1, x_2, v_1, v_2)
 - x_1, x_2 : Position of point mass 1 and 2
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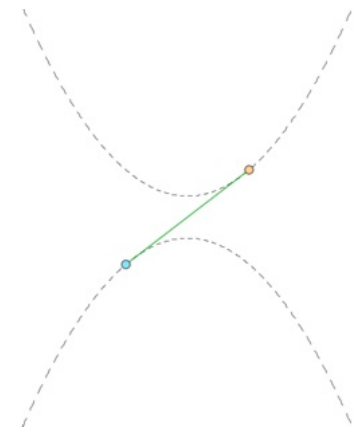
$$\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix} = \begin{pmatrix} x_1 \\ x_2 \\ v_1 \\ v_2 \end{pmatrix}$$



51

Visualization 1

- ▶ Parabola y_1 and parabola y_2
 - Dashed curve
- ▶ Point mass 1 and 2 $(x, y) = (x_i, y_i)$
 - Small circles
- ▶ Rubber rope
 - Green line



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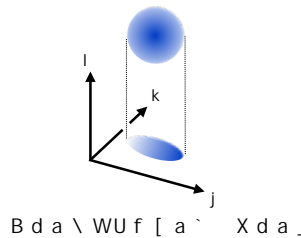
Visualization 2

- ▶ We want to "see" the time development of the system state, i.e., a curve in 4-D phase space.
- ▶ Trick: Project 4-D to 2-D (screen).

"Seeing" is a projection (from 3-D to 2-D).

- An example of projection:

$$(x_1(t), x_2(t), x_3(t), x_4(t)) \Rightarrow (x_1(t), x_2(t))$$



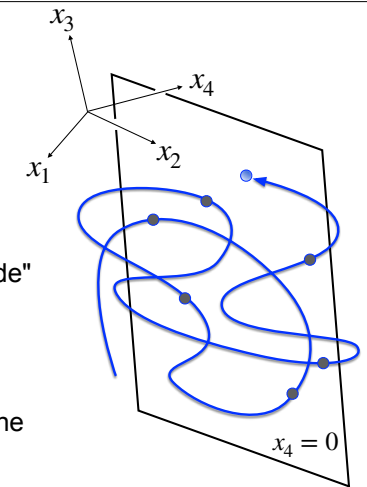
53

Visualization 3

- ▶ A cross section of a 4-D directed curve
 - For example, take $x_4 = 0$ plane.
- ▶ Focus on cross points when the curve crosses the plane ($x_4 = 0$), from "back side" ($x_4 < 0$) to "front side" ($x_4 > 0$).

$$(x_1, x_2, x_3, x_4 = 0)$$

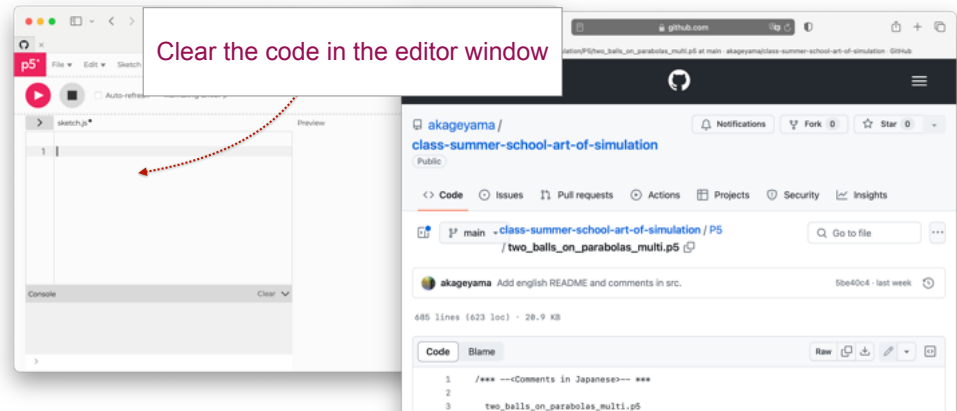
- ▶ Place a point on, say (x_1, x_3) , on x_1 - x_3 plane
- ▶ This is called "Poincare map."



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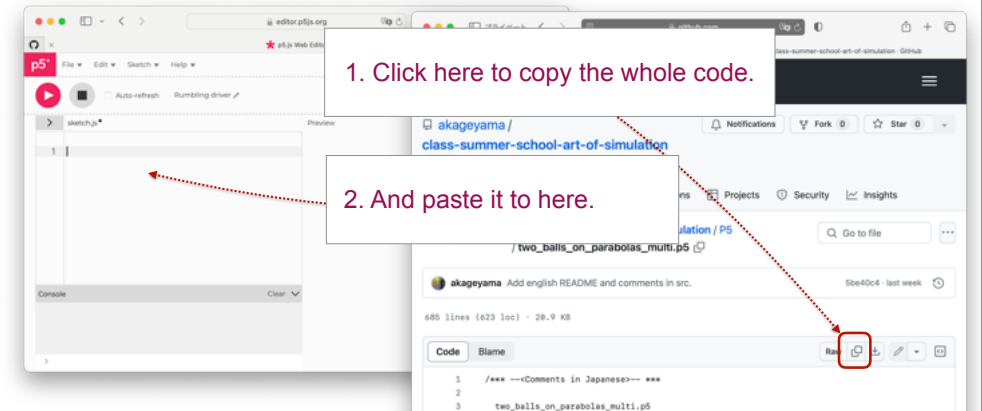
Sample code

<https://tinyurl.com/kobe-sc01b>



Sample code

<https://tinyurl.com/kobe-sc01b>



Wc g S f [a ` Q a X Q _ a f [a `

▶ > [` W ' " # S ` V

$$\frac{dx_1}{dt} = v_1, \rightarrow dx_1 = v_1 dt$$

$$\frac{dx_2}{dt} = v_2,$$

$$\frac{dv_1}{dt} = \frac{-1}{1 + 4x_1^2} \{ 4x_1 v_1^2 + \omega^2 (x_1 - x_2 + 2x_1(x_1^2 + x_2^2 + 2)) \},$$

$$\frac{dv_2}{dt} = \frac{-1}{1 + 4x_2^2} \{ 4x_2 v_2^2 + \omega^2 (x_2 - x_1 + 2x_2(x_1^2 + x_2^2 + 2)) \}.$$

```

510 let x1 = b.x1;
511 let v1 = b.v1;
512 let x2 = b.x2;
513 let v2 = b.v2;
514 let dx = x1 - x2;
515 let x1sq = x1 * x1;
516 let v1sq = v1 * v1;
517 let x2sq = x2 * x2;
518 let v2sq = v2 * v2;
519 let dy = x1sq + x2sq + 2;
520 let f1 = OMEGA_SQ * (dx + 2 * x1 * dy);
521 let f2 = OMEGA_SQ * (-dx + 2 * x2 * dy);
522 db.x1 = v1 * dt;
523 db.v1 = (-1.0 / (1 + 4 * x1sq)) * (4 * x1 * v1sq + f1) * dt;
524 db.x2 = v2 * dt;
525 db.v2 = (-1.0 / (1 + 4 * x2sq)) * (4 * x2 * v2sq + f2) * dt;

```

Initial condition

```

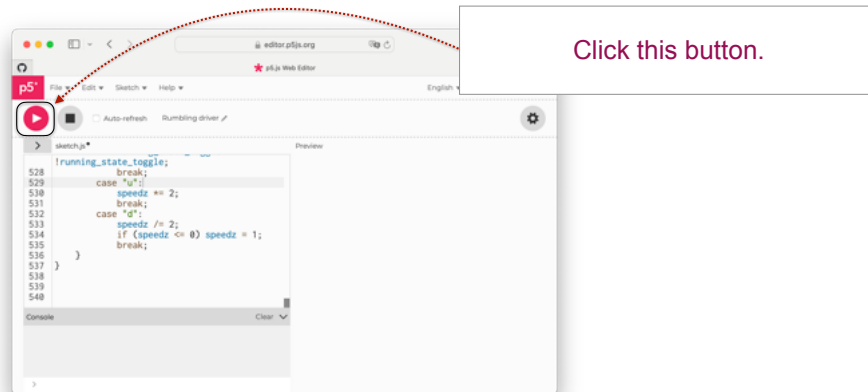
450 header.title(" Path plot of (x1,x2)", CENTER);
451 header.title("Poincare map of (x1,v1) on v2=0 ", RIGHT); //// 線形 (微小) 単振動
452 // balls = new GeneralCoords( x_coord_max*0.05,0, // x1 & v1
453 // x_coord_max*0.05,0); // x2 & v2
454 //// 線形 (微小) 振動
455 // balls = new GeneralCoords( x_coord_max*0.1,0, // x1 & v1
456 // -x_coord_max*0.05,0); // x2 & v2
457 //// 非線形単一周運動
458 // balls = new GeneralCoords( x_coord_max*0.4,0, // x1 & v1
459 // -x_coord_max*0.4,0); // x2 & v2
460 //// 比較的単純な非線形運動
461 // balls = new GeneralCoords( x_coord_max*0.4,0, // x1 & v1
462 // -x_coord_max*0.2,0); // x2 & v2
463 // 複雑な運動
464 balls = new GeneralCoords(
465 x_coord_max * 0.4,
466 0, // x1 & v1
467 -x_coord_max * 0.1,
468 0 // x2 & v2
469 );

```

x_1 : x-coord of particle 1
 v_1 : velocity of particle 1
 x_2 : x-coord of particle 2
 v_2 : velocity of particle 2

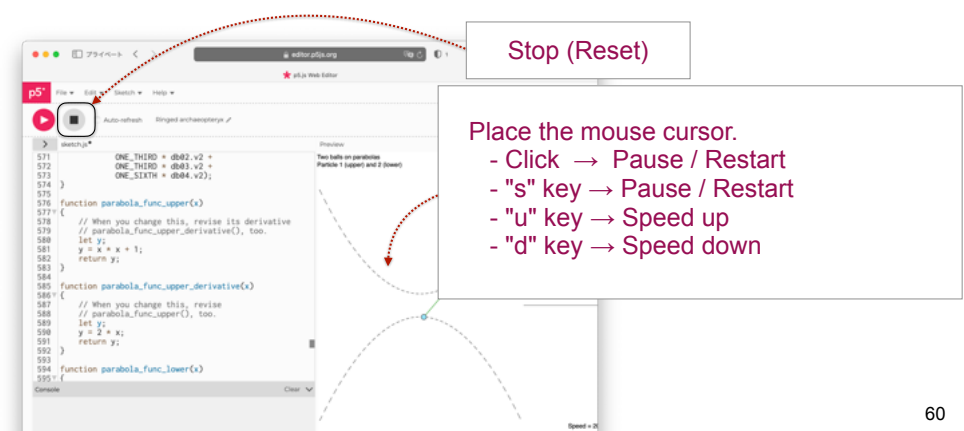
58

Execution of the simulation



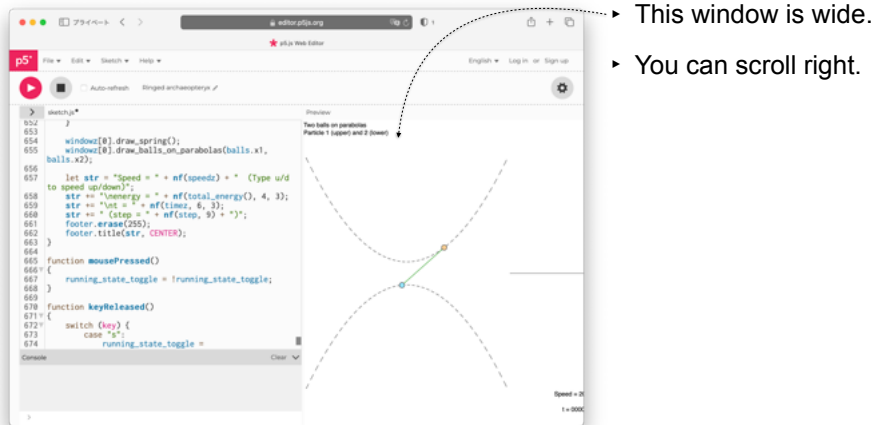
59

Usage (same as before)



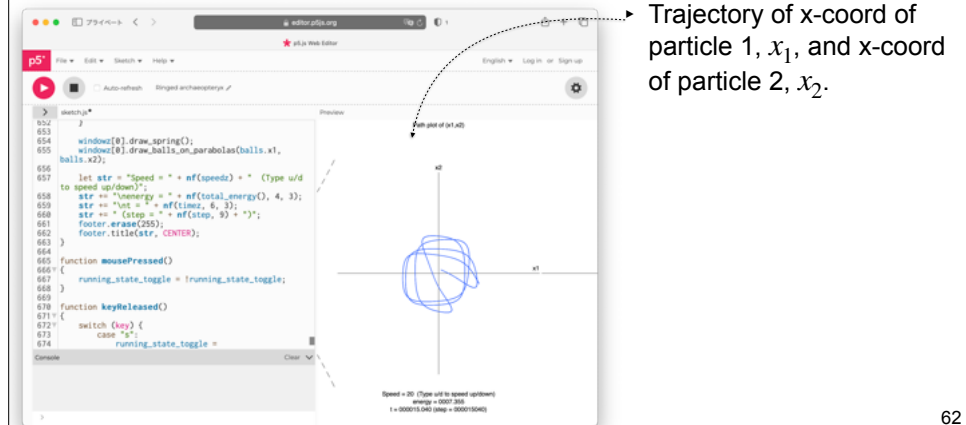
60

Visualization 1



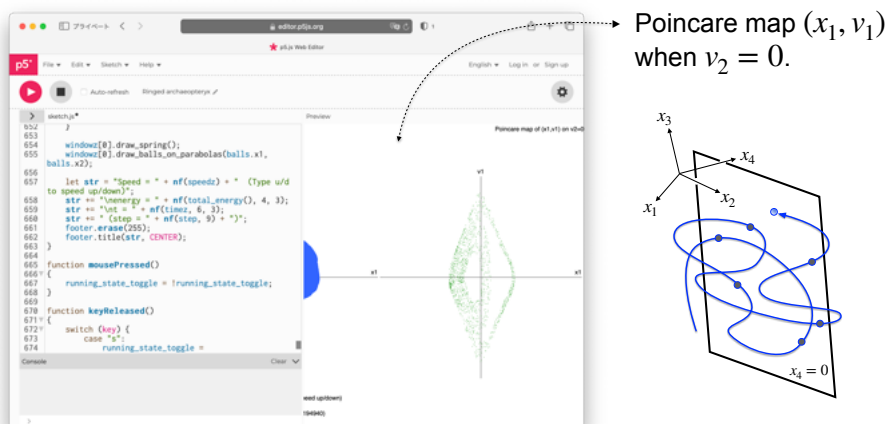
61

Visualization 2



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Visualization 3



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Let's try with different initial condition

```

450 header.title(" Path plot of (x1,x2)", CENTER);
451 header.title("Poincare map of (x1,v1) on v2=0 ", RIGHT); //// 線形 (微小) 単振動
452 // balls = new GeneralCoords( x1, v1, x2, v2 );
453 // balls = new GeneralCoords( x1, v1, x2, v2 );
454 //// 線形 (微小) 振動
455 // balls = new GeneralCoords( x1, v1, x2, v2 );
456 // balls = new GeneralCoords( x1, v1, x2, v2 );
457 // 非線形単一周期運動
458 balls = new GeneralCoords( x1, v1, x2, v2 );
459 // balls = new GeneralCoords( x1, v1, x2, v2 );
460 //// 比較的単純な非線形運動
461 // balls = new GeneralCoords( x1, v1, x2, v2 );
462 // balls = new GeneralCoords( x1, v1, x2, v2 );
463 // 複雑な運動
464 // balls = new GeneralCoords( x1, v1, x2, v2 );
465 // balls = new GeneralCoords( x1, v1, x2, v2 );
466 // balls = new GeneralCoords( x1, v1, x2, v2 );
467 // balls = new GeneralCoords( x1, v1, x2, v2 );
468 // balls = new GeneralCoords( x1, v1, x2, v2 );
469 // balls = new GeneralCoords( x1, v1, x2, v2 );

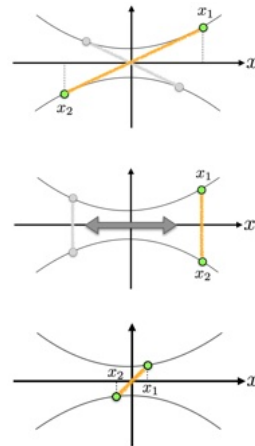
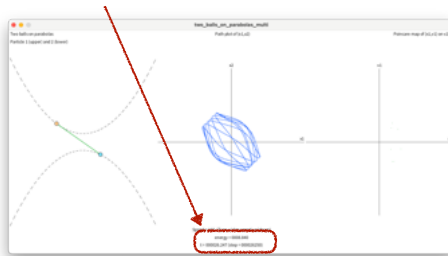
```

Comment out

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You can try various oscillations

- Symmetric nonlinear oscillation
- Small amplitude oscillation of $\sqrt{6} \omega$
- You can check the conservation of total energy.



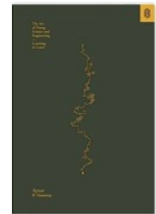
65

Summary

- Start with a model as simple as possible, but not simpler.
- Think before compute.
 - Linearized solutions, symmetry, conservation law, etc.
- Cultivate your intuition about the subject by quick experiments.
 - Then, move on to more complex models.
 - Visualization is a key.

"The purpose of computing is insight, not numbers."

R. Hamming



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Appendix: Derivation of the Equations of Motion

In the followings, $\mathbf{a} \cdot \mathbf{b}$ denotes the dot product of two vectors \mathbf{a} and \mathbf{b} , and $\mathbf{a} \cdot \mathbf{a}$ is written as \mathbf{a}^2 .

Let \mathbf{v}_1 and \mathbf{v}_2 represent the velocities of pulleys 1 and 2, respectively. The total kinetic energy of the system is

$$K = \frac{m}{2}(\mathbf{v}_1^2 + \mathbf{v}_2^2), \quad (1)$$

Neglecting the gravitation, the potential energy U of the system comes solely from the tension of the rubber band connecting the two pulleys. Assuming the natural length is 0 and the elastic coefficient (spring constant) k , the potential energy is given by

$$U = \frac{k}{2}(\mathbf{x}_1 - \mathbf{x}_2)^2. \quad (2)$$

Because we are assuming that the two curves near the pulleys of the cable and hanger are specified by one parameter c as

$$y(x) = \begin{cases} y_1 = +x^2/c + c & (\text{cable}), \\ y_2 = -x^2/c - c & (\text{hanger}), \end{cases} \quad (3)$$

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Appendix: Derivation of the Equations of Motion

$$(\mathbf{x}_1 - \mathbf{x}_2)^2 = (x_1 - x_2)^2 + (y_1 - y_2)^2 = (x_1 - x_2)^2 + \left(\frac{x_1^2}{c} + \frac{x_2^2}{c} + 2c\right)^2, \quad (4)$$

$$\mathbf{v}_i^2 = \dot{\mathbf{x}}_i^2 = \left\{1 + \left(\frac{dy_i}{dx_i}\right)^2\right\} \left(\frac{dx_i}{dt}\right)^2 = \left(1 + \frac{4x_i^2}{c^2}\right) \left(\frac{dx_i}{dt}\right)^2 \quad (i = 1, 2). \quad (5)$$

The Lagrangian $L(x_1, x_2, \dot{x}_1, \dot{x}_2) = K - U$ of this system is

$$L = \frac{m}{2} \left\{ \dot{x}_1^2 \left(1 + \frac{4x_1^2}{c^2}\right) + \dot{x}_2^2 \left(1 + \frac{4x_2^2}{c^2}\right) \right\} - \frac{k}{2} \left\{ (x_1 - x_2)^2 + \left(\frac{x_1^2}{c} + \frac{x_2^2}{c} + 2c\right)^2 \right\}. \quad (6)$$

The Lagrange's equations of motion,

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_i} \right) = \frac{\partial L}{\partial x_i} \quad (i = 1, 2), \quad (7)$$

is written as follows:

Appendix: Derivation of the Equations of Motion

$$m \left(1 + \frac{4x_1^2}{c^2}\right) \frac{d\dot{x}_1}{dt} + 8m \frac{x_1 \dot{x}_1^2}{c^2} = 4m \frac{x_1 \dot{x}_1^2}{c^2} - k \left\{ x_1 - x_2 + \frac{2x_1}{c} \left(\frac{x_1^2}{c} + \frac{x_2^2}{c} + 2c\right) \right\}, \quad (8)$$

$$m \left(1 + \frac{4x_2^2}{c^2}\right) \frac{d\dot{x}_2}{dt} + 8m \frac{x_2 \dot{x}_2^2}{c^2} = 4m \frac{x_2 \dot{x}_2^2}{c^2} - k \left\{ x_2 - x_1 + \frac{2x_2}{c} \left(\frac{x_1^2}{c} + \frac{x_2^2}{c} + 2c\right) \right\}. \quad (9)$$

They are coupled second-order differential equations of x_1 and x_2 . We split them into coupled first-order differential equations for the variables x_1, v_1, x_2, v_2 as

$$\frac{dx_1}{dt} = v_1, \quad (10)$$

$$\frac{dx_2}{dt} = v_2, \quad (11)$$

$$\frac{dv_1}{dt} = \frac{-1}{4x_1^2 + c^2} \left\{ 4x_1 v_1^2 + \frac{kc^2}{m} \left(x_1 - x_2 + 2x_1 \frac{x_1^2 + x_2^2 + 2c^2}{c^2} \right) \right\}, \quad (12)$$

$$\frac{dv_2}{dt} = \frac{-1}{4x_2^2 + c^2} \left\{ 4x_2 v_2^2 + \frac{kc^2}{m} \left(x_2 - x_1 + 2x_2 \frac{x_1^2 + x_2^2 + 2c^2}{c^2} \right) \right\}. \quad (13)$$