CL 716 - Modelling Chemical and Biological Patterns

Bifurcation analysis of nonlinear reaction-diffusion equations

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Abstract

The theoretical expressions are limited to the neighborhood of the marginal stability point. Computer simulations allow not only the verification of their predictions but also the investigation of the behavior of the system for larger deviations from the instability point.

Introduction

The general ideas underlying the theory of dissipative structures have been illustrated on a simple model system involving the following set of coupled chemical reactions:

$$A \rightleftharpoons X$$
 (a)
 $2X + Y \rightleftharpoons 3X$ (b)
 $B + X \rightleftharpoons Y + D$ (c)
 $X \rightleftharpoons E$ (d)

the system is open to the initial and final chemicals A, B, D and E, whose concentrations are imposed throughout the system; nonlinearity is introduced by the auto- and cross-catalytic steps (b) and (c);

We analyze some properties of the dissipative structures arising in nonlinear reaction-diffusion systems, within the framework of this model. Assuming a bounded, one-dimensional medium, the rate equations describing (1.1) are:

$$\frac{\partial X}{\partial t} = A + X^{2}Y - (B+1)X + D_{X} \cdot \frac{\partial^{2}X}{\partial r^{2}}$$

$$\frac{\partial Y}{\partial t} = BX - X^{2}Y + D_{Y} \cdot \frac{\partial^{2}Y}{\partial r^{2}} \quad (0 \le r \le L).$$
(1.2)

where Dx and Dr are the diffusion coefficients of X and Y assuming that Fick's law is valid. Two types of boundary conditions will be considered:

1. Zero flux boundary conditions (Neumann conditions):

$$\frac{\partial}{\partial r}X(0,t) = \frac{\partial}{\partial r}X(L,t) = \frac{\partial}{\partial r}Y(0,t) = \frac{\partial}{\partial r}Y(L,t) = 0 \quad (t \ge 0). \quad (1.3)$$

2. Fixed boundary conditions (Dirichlet conditions):

$$X(0,t) = X(L,t) = A$$

$$Y(0,t) = Y(L,t) = B/A \quad (t \ge 0).$$
(1.4)

MATLAB Code

The Matlab code consists of four files:-

The main script file

```
function pdex4
m=0; %slab
r=linspace(0,pi,100);
t=linspace(0,200,100);
sol=pdepe(m,@pdex4pde,@pdex4ic,@bc2fn,r,t);
disp(sol);
u1 = sol(:,:,1);
u2 = sol(:,:,2);
figure
surf(r,t,u1)
title('X(r,t)')
xlabel('Distance r')
ylabel('Time t')
figure
surf(r,t,u2)
title('Y(r,t)')
xlabel('Distance r')
ylabel('Time t')
```

Boundary condition definition file

Note: Two types of boundary conditions will be considered:1. Zero Flux Boundary Conditions (Neumann conditions)

2. Fixed Boundary Conditions (Dirichlet conditions)

```
function [pl,ql,pr,qr]=bc2fn(xl,ul,xr,ur,t)

% Constants
A = 2;
B = 0.4; %NOT GIVEN | TO BE CHANGED

% Case 1:- Zero Flux Boundary Conditions (Neumann conditions)
pl= [0;0];
ql=[1;1];
pr =[0;0];
qr =[1;1];

% Case 2:- Fixed Boundary Conditions (Dirichlet conditions)
% pl = [A; B/A];
% ql = [0; 0];
% pr = [A; B/A];
% qr = [0; 0];
```

Initial condition definition file

```
function u0 = pdex4ic(r);

% Constants
A = 2;
B = 3.7; %NOT GIVEN | TO BE CHANGED

c1 = 10;
c2 = 10;
L = 1;

u0 = [A;B/A];
```

The PDE solver file

```
function [c,f,s] = pdex4pde(r,t,u,DuDr)
% Diffusion Coefficients
Dy = 1.6*10^{(-3)};
Dx = 8.0*10^{(-3)};
% Constants
A = 2;
L = 1;
u1 = u(1);
u2 = u(2);
B = 3.7; %NOT GIVEN | TO BE CHANGED
c = [1; 1];
f = [Dx; Dy] .* DuDr;
% Rate equations describing the phenomenon
s1 = A + u1^2 + u2 - (B+1) + u1;
s2 = B*u1 - u1^2*u2;
\mbox{\ensuremath{\$}} Linearized equations for the perturbation x and y
% s1 = (B-1)*u1 + A^2*u2;
% s2 = -B*u1 - A^2*u2;
s = [s1; s2];
```

Numerical Analysis

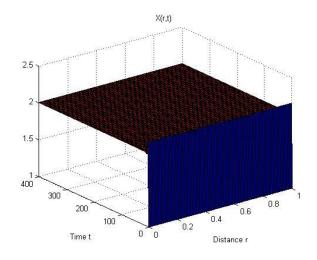
We use the simulation using MATLAB to verify the numerical results with the analytical ones. Following are the major numerical simulations:-

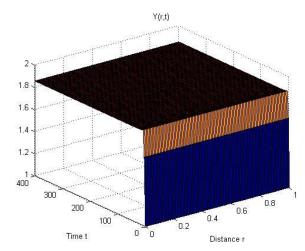
Case 1

All the equations are restricted to the constraint $0 \le r \le L$ and L is taken to be 1. Also diffusion coefficients are taken to be:-

- 1. $Dx = 1.6 \times 10^{-3}$
- 2. Dy = 8.0×10^{-3}

A = 2 and B = 3.7 for zero flux boundary condition





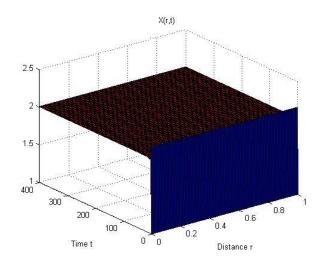
A spatio-temporal curve is plotted between X and Y vs distance 'r' and time 't' resulting from solving the system numerically by simulating on MATLAB. Figure 1A: X vs distance 'r' and time 't' for A=2, B=3.7, L=1, Dx = 1.6×10^{-3} , Dy = 8.0×10^{-3} . Figure 1B: Y vs distance 'r' and time 't' for A=2, B=3.7, L=1, Dx = 1.6×10^{-3} .

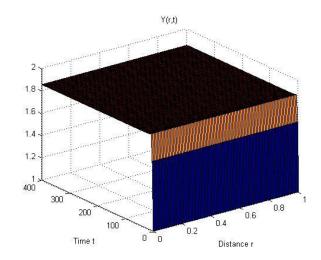
Case 2

All the equations are restricted to the constraint $0 \le r \le L$ and L is taken to be 1. Also diffusion coefficients are taken to be:-

- 1. $Dx = 8.0 \times 10^{-3}$
- 2. Dy = 1.6×10^{-3}

A = 2 and B = 3.7 for zero flux boundary condition





A spatio-temporal curve is plotted between X and Y vs distance 'r' and time 't' resulting from solving the system numerically by simulating on MATLAB. Figure 2A: X vs distance 'r' and time 't' for A=2, B=3.7, L=1, Dx = 8.0×10^{-3} , Dy = 1.6×10^{-3} . Figure 2B: Y vs distance 'r' and time 't' for A=2, B=3.7, L=1, Dx = 8.0×10^{-3} .

Analytical Solution

For zero – flux boundary conditions,

$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} e^{w_n t} \cos \frac{n \pi r}{L} \quad (n = 0, 1, 2, 3...).$$
 (2.5b)

Inserting this into the rate equations, we get secular equation relating w_n to the wavenumber n and the system's parameters:

$$w_n^2 - \operatorname{Tr} w_n + \Delta = 0 (2.6)$$

where

$$Tr = B - (A^2 + 1) - \beta(D_X + D_Y)$$

$$\Delta = A^2 + \beta[A^2D_X + (1 - B)D_Y] + \beta^2D_XD_Y$$

and

$$\beta = \left(\frac{n\pi}{L}\right)^2$$

Instability of the thermodynamic branch will occur for some value of n, if at least one of the roots of (2.6) has a positive Re w_n part. The main point is thus to establish the conditions for marginal stability, Re $w^{\sim} = 0$, corresponding either to 'exchange of stability', Im $w_n \neq 0$, or to 'overstability' Im $w_n \neq 0$. A close analysis of (2.6) shows that:

(a) the values of w_n are complex if:

$$(A - \delta^{1/2})^2 < B < (A + \delta^{1/2})^2 \tag{2.7}$$

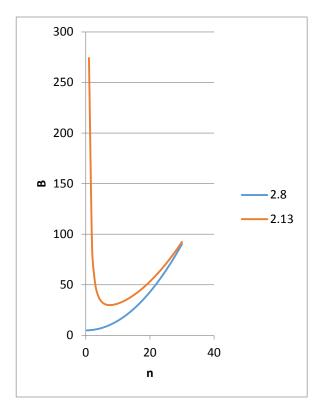
where $\delta = 1 - \beta(D_X - D_Y)$ must be a positive quantity. In this case marginal stability occurs at the critical point:

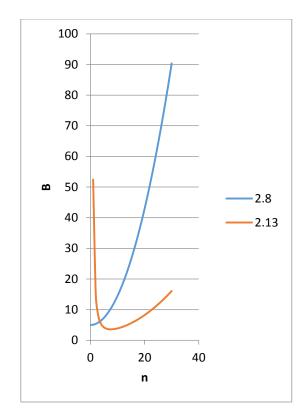
$$B = B_{1c}(n) = 1 + A^2 + \beta(D_X + D_Y). \tag{2.8}$$

For real w_n's the instability conditions reads:-

$$B \geqslant B(n_c) = \min_{\substack{n \geq 1 \\ \text{integer}}} \left\{ 1 + \frac{D_X}{D_Y} A^2 + \frac{A^2}{D_Y \beta} + \beta D_X \right\}$$
 (2.13)

Graphs are plotted for these in Microsoft Excel match with the ones shown in the paper:-





Linear stability diagrams resulting from above equations. Figure 3A: A=2, B=3.7, L=1, Dx = 8.0×10^{-3} , Dy = 1.6×10^{-3} . Figure 3B: A=2, B=3.7, L=1, Dx = 1.6×10^{-3} , Dy = 8.0×10^{-3} .

Critical Wave Number

The critical wave number corresponding to the onset of stability is given by n_{min} if it is an integer or by one of the two closest integers.

$$n_{\min} = \frac{L}{\pi} \cdot \frac{A^{1/2}}{(D_X D_Y)^{1/4}} \tag{2.11}$$

Case I: A=2, B=3.7, L=1, Dx = 8.0×10^{-3} , Dy = 1.6×10^{-3}

Critical Wave number = $[1/\Pi * 2^{1/2}/(8*10^{-3}*1.6*10^{-3})^{1/4}] = 8$

Case II: A=2, B=3.7, L=1, Dx = 1.6×10^{-3} , Dy = 8.0×10^{-3}

Critical Wave number = $[1/\Pi * 2^{1/2}/(1.6*10^{-3}*8*10^{-3})^{1/4}] = 8$ (same)

References

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