

Removal of pseudo-convergence in near-coplanar  
Riemann problems of ideal  
magnetohydrodynamics  
and  
parallel fluid simulations on shared memory  
processors

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George Mason University

28 August 2014

# Outline

## Introduction

Shocks in space plasma

Overview

Riemann problems of ideal MHD

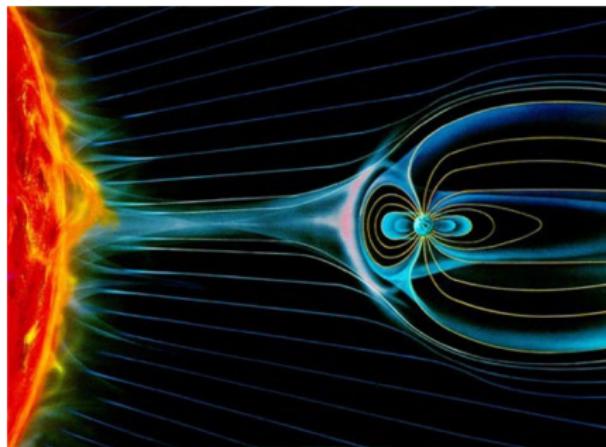
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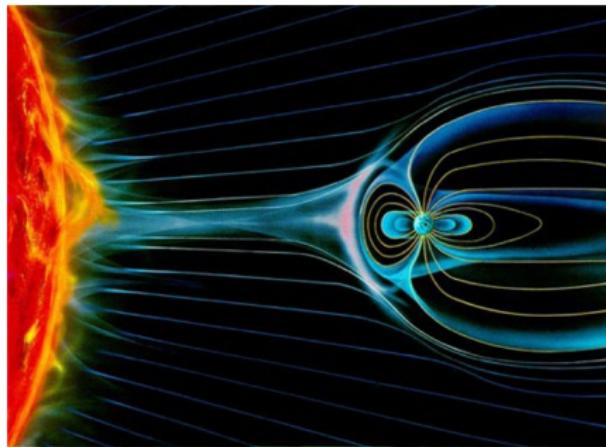
# Shocks in space plasma



Source [7]

- Sun-earth interaction:
  - ▶ Super-sonic plasma originating at sun travels toward earth.

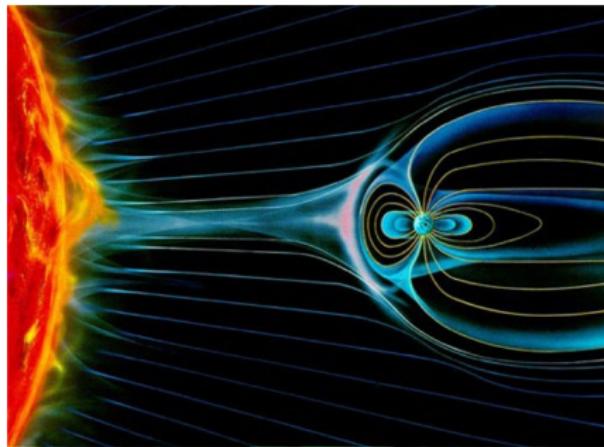
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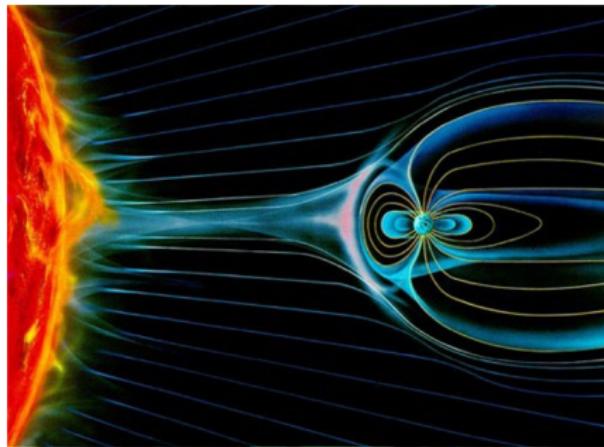
# Shocks in space plasma



Source [7]

- Sun-earth interaction:
  - ▶ Super-sonic plasma originating at sun travels toward earth.
  - ▶ Slowed down by magnetic field of Earth.
  - ▶ Plasma flow becomes subsonic forming a bow shock.

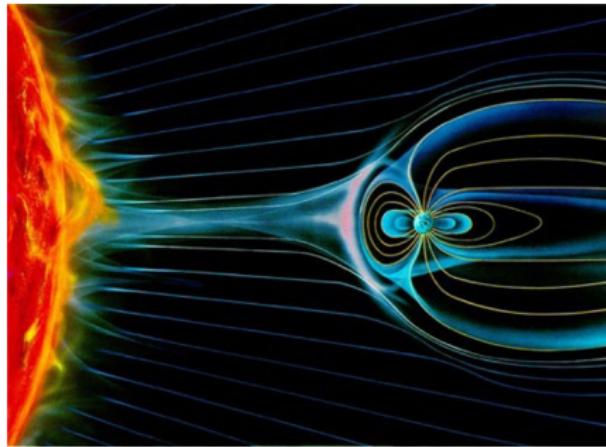
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Source [7]

- 2D ideal MHD bow shock simulation. [1]
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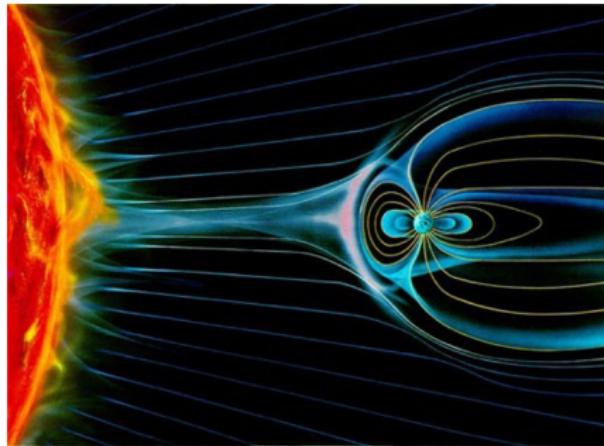
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  - ▶ Modification to the HLLD approximate Riemann solver of Athena [8].
  - ▶ Removes compound wave, correctly computes rotational discontinuity.
  - ▶ Converges at all grid resolutions.
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- Description multi-dimensional fluid solver capable of shared memory parallelism.
  - ▶ Hydrodynamics and ideal MHD.
  - ▶ Algorithms implemented for unstructured grids.

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## Ideal magnetohydrodynamics

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 ,$$

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot \left[ \rho \mathbf{v} \otimes \mathbf{v} + \left( p_g + \frac{B^2}{2} \right) \mathbf{I} - \mathbf{B} \otimes \mathbf{B} \right] = 0 ,$$

$$\frac{\partial E}{\partial t} + \nabla \cdot \left[ \left( E + p_g + \frac{B^2}{2} \right) \mathbf{v} - \mathbf{v} \cdot \mathbf{B} \otimes \mathbf{B} \right] = 0 , \text{ and}$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot [\mathbf{v} \otimes \mathbf{B} - \mathbf{B} \otimes \mathbf{v}] = 0 ,$$

where the energy density is defined as

$$E = \frac{p_g}{\gamma - 1} + \frac{\rho v^2}{2} + \frac{B^2}{2} ,$$

# Ideal magnetohydrodynamics

**Non-strictly** hyperbolic.

$$\frac{\partial \mathbf{U}}{\partial t} + \frac{\partial \mathbf{F}}{\partial x} = 0.$$

Real, but **not** necessarily distinct eigenvalues:

$\nu_n$  : contact or tangential discontinuity (entropy),

$\nu_n \pm c_s$  : slow rarefaction or shock,

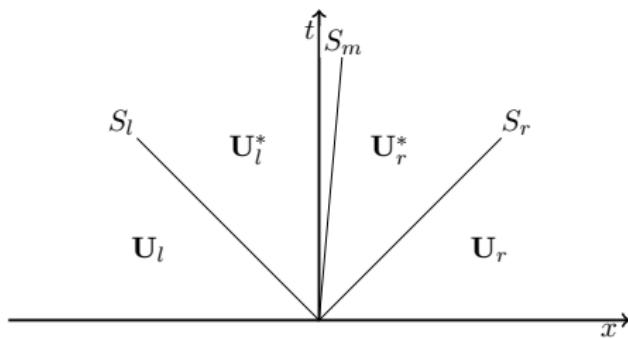
$\nu_n \pm c_a$  : rotational discontinuity (Alfvén), and

$\nu_n \pm c_f$  : fast rarefaction or shock,

$$c_{f,s}^2 = \frac{1}{2} \left[ a^2 + c_a^2 + c_t^2 \pm \sqrt{(a^2 + c_a^2 + c_t^2)^2 - 4a^2 c_a^2} \right],$$

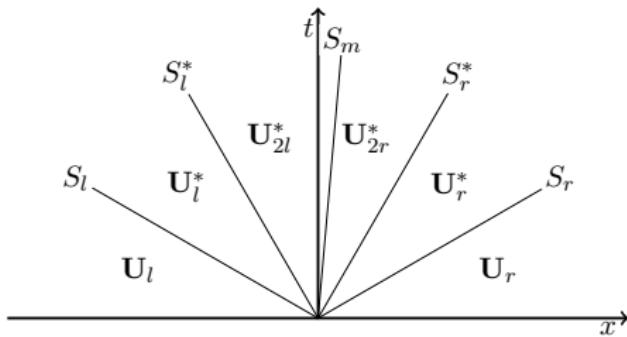
$$c_a^2 = \frac{B_n^2}{\rho}, \text{ and } c_t^2 = \frac{B_t^2}{\rho}.$$

# HLLD approximate Riemann solver



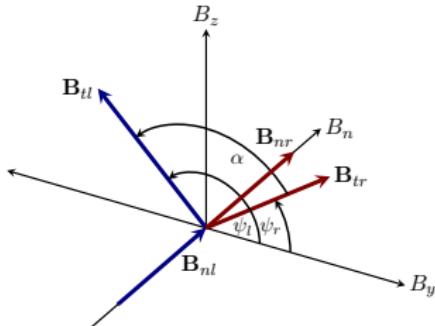
- HLLD [5] (D for *discontinuities*) extention of HLLC to MHD.

# HLLD approximate Riemann solver



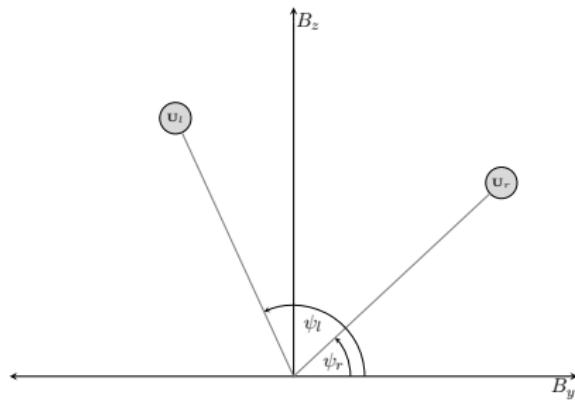
- HLLD [5] (D for *discontinuities*) extention of HLLC to MHD.
- Restores rotational discontinuity.

# Riemann problems of ideal MHD



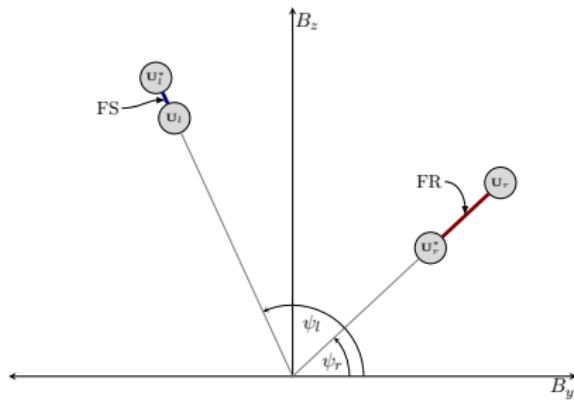
- Initial discontinuity separates two states.
- Rotation angle:  $\psi = \arctan(B_z/B_y)$ .
- Twist angle:  $\alpha = \psi_r - \psi_l$ .

# Riemann problems of ideal MHD



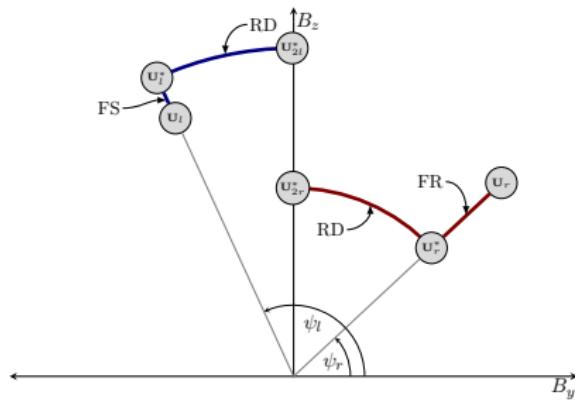
- Regular waves alter magnitude or orientation of  $B_t$ .

# Riemann problems of ideal MHD



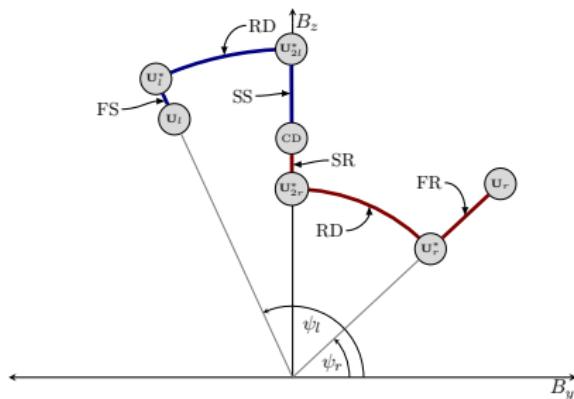
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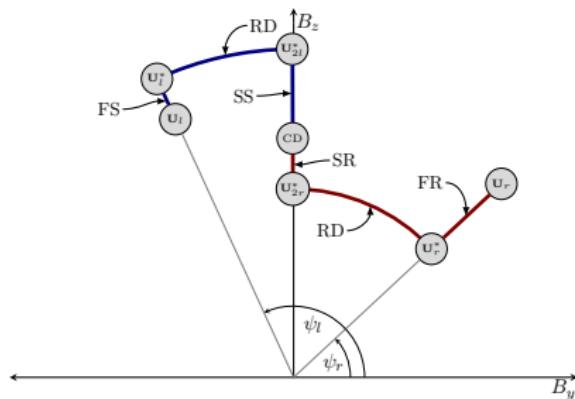
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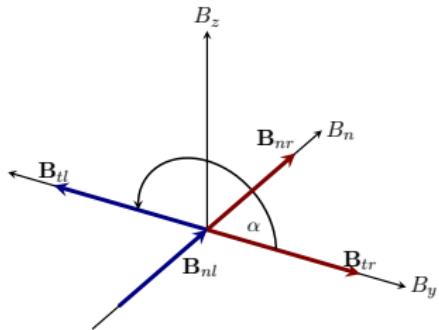
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- Slow shock:  $B_t \downarrow$ , slow rarefaction:  $B_t \uparrow$ .
- Contact discontinuity: no change.

# Riemann problems of ideal MHD



- Coplanar magnetic field

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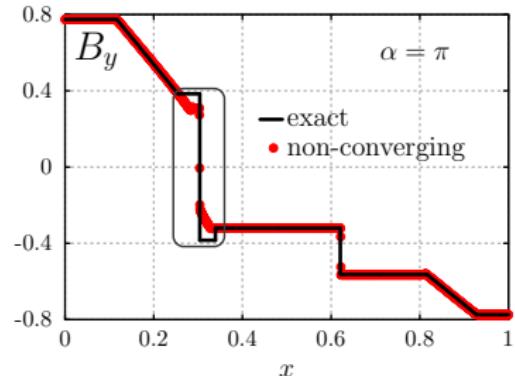
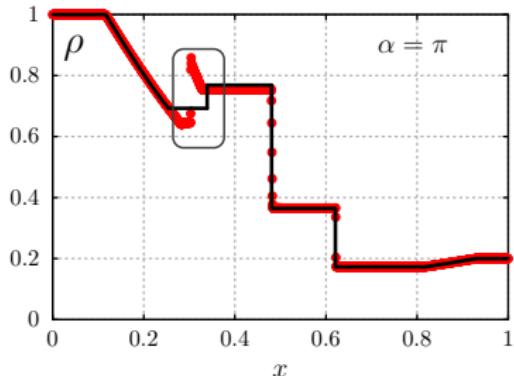
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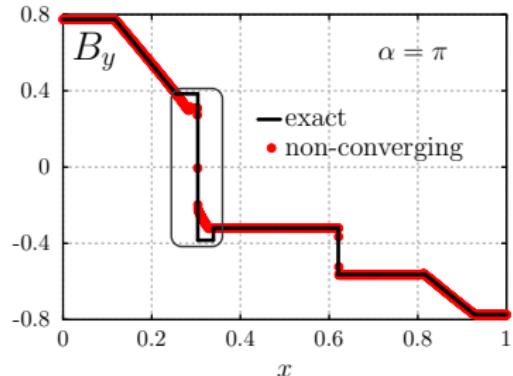
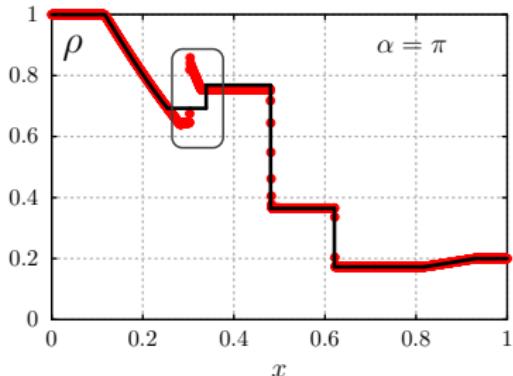
Shared memory parallelism

# Non-unique solutions



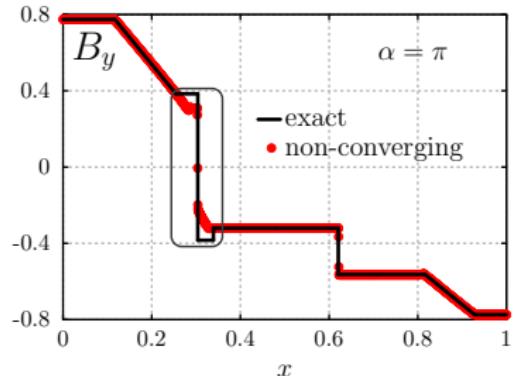
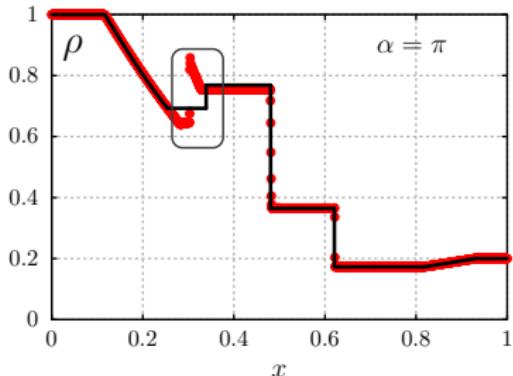
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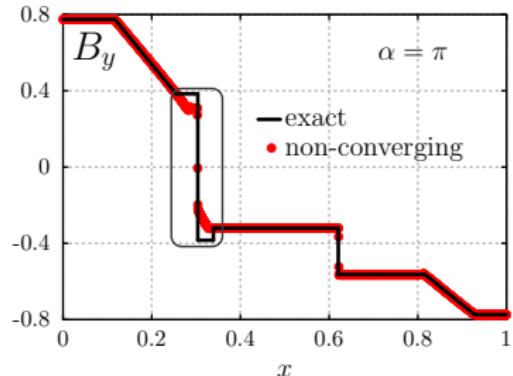
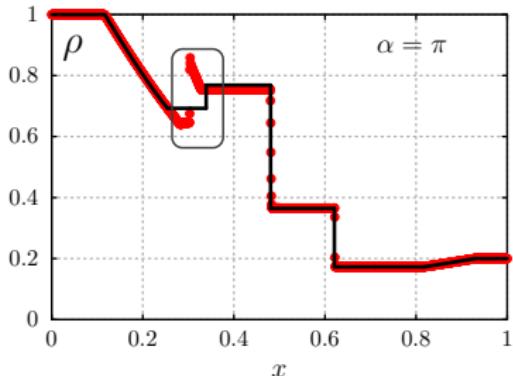
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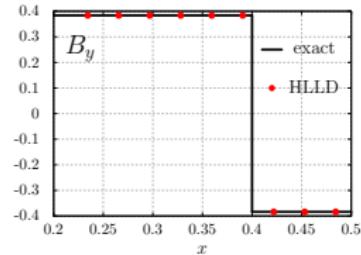
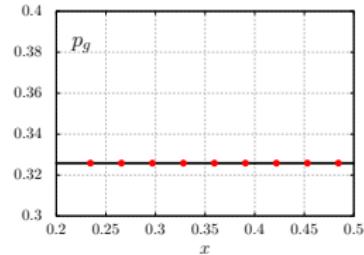
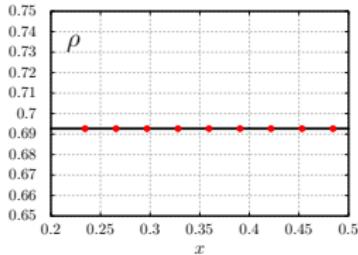
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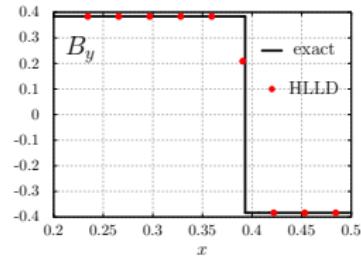
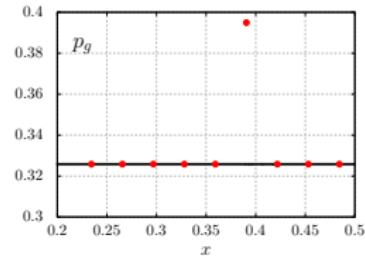
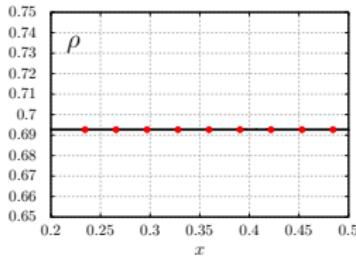
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  - ▶ Unstable for under small perturbations [4].
  - ▶ Satisfy jump conditions (coplanar case).

# Compound wave formation



- Compound waves are a product of numerical diffusion.

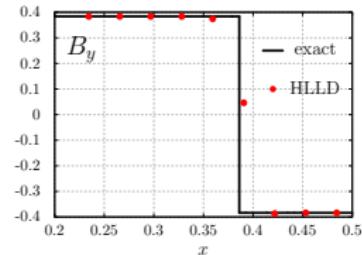
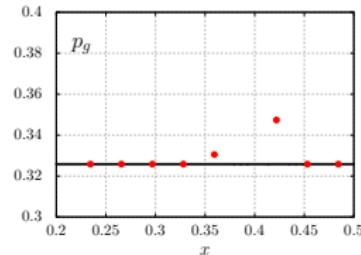
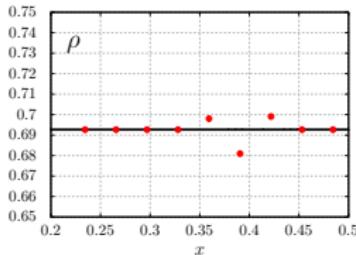
# Compound wave formation



- Compound waves are a product of numerical diffusion.
- Decrease  $B_t \rightarrow$  increase  $p_g$

$$p_g = (\gamma - 1) \left( E - \frac{\rho v^2}{2} - \frac{B^2}{2} \right)$$

# Compound wave formation

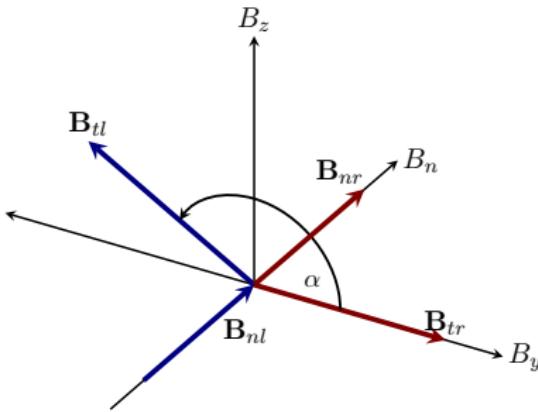


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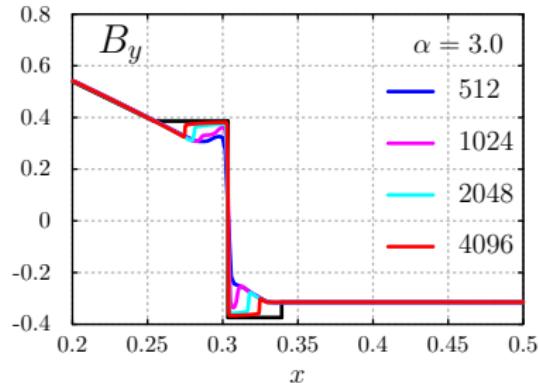
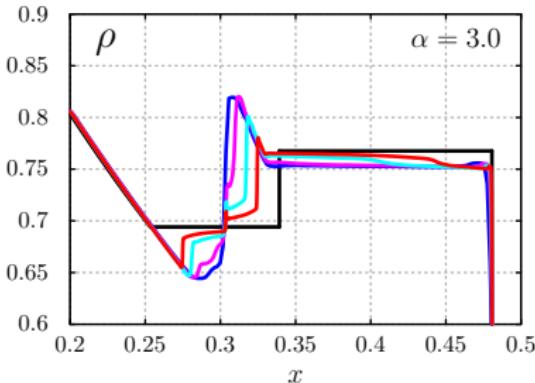
- Increase  $p_g \rightarrow$  compound wave.

# Pseudo-convergence



- $\alpha < \pi$ , compound waves at lower resolutions.

# Pseudo-convergence



- $\alpha < \pi$ , compound waves at lower resolutions.
- Pseudo-convergence: rotational discontinuity at higher resolutions [9].
- Transition between 1024 and 2048 grid points.
- $\alpha \rightarrow \pi$ , grid resolution increases.
- $\pi = \alpha$ , compound wave at all resolutions.

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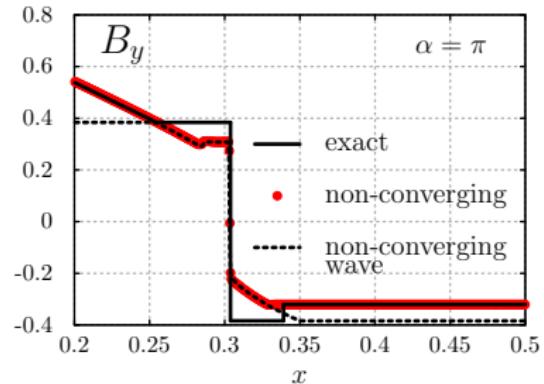
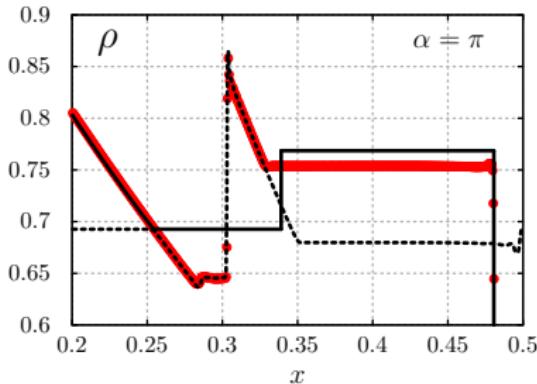
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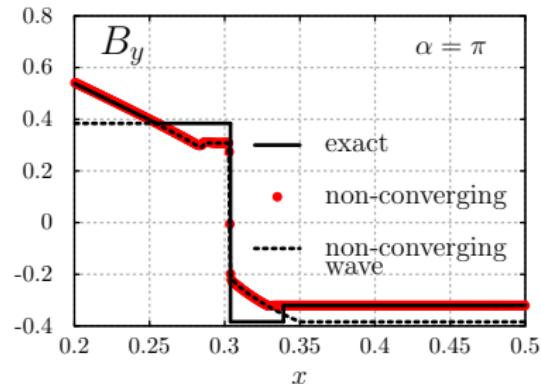
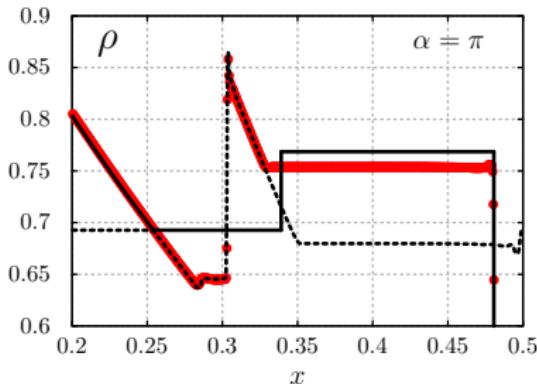
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# Compound wave modification



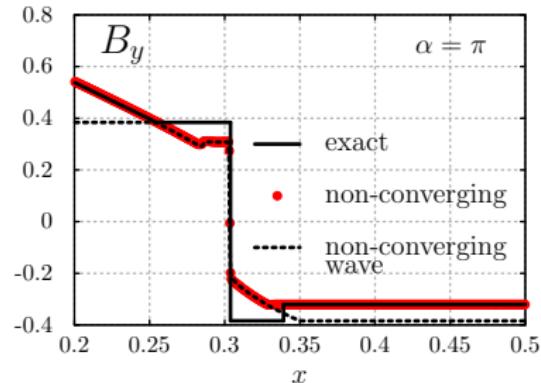
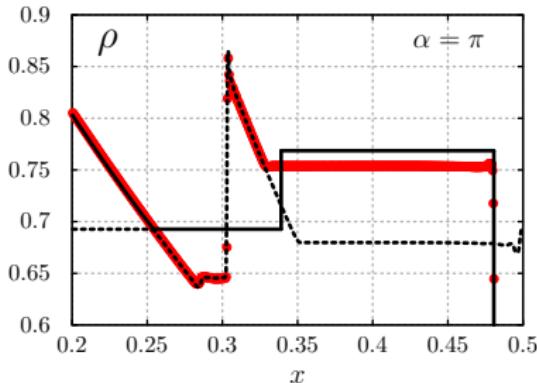
- Need to reduce numerical diffusion.

# Compound wave modification



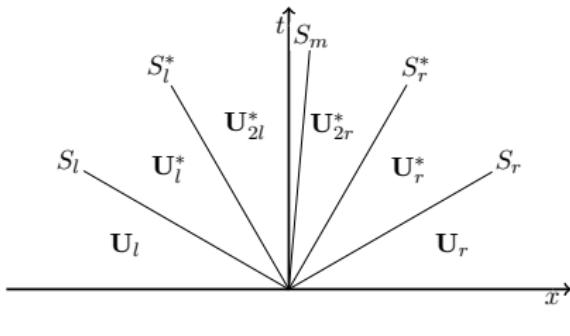
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- $\rho$  and  $B_t$  should remain constant across the rotational discontinuity  $x = 0.303$ .

# Compound wave modification

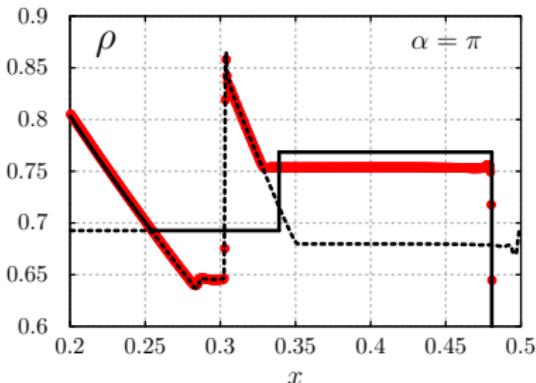


- Need to reduce numerical diffusion.
- $\rho$  and  $B_t$  should remain constant across the rotational discontinuity  $x = 0.303$ .
- Limit flux associated with the compound wave.

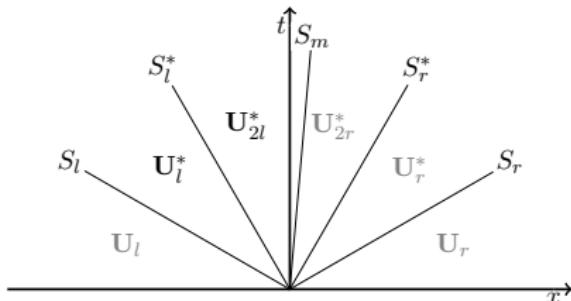
# Compound Wave modification



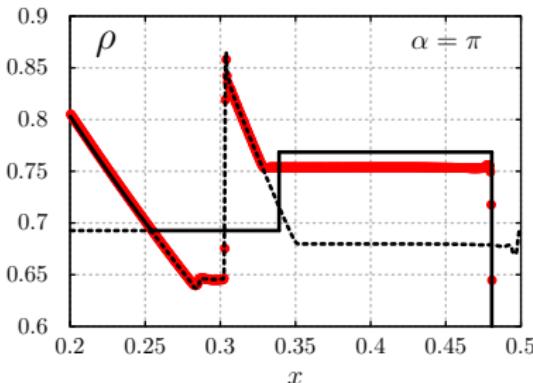
- Use HLLD to approximate intermediate states upstream and downstream of rotational discontinuity, i.e.,  $\mathbf{U}_l^*$  and  $\mathbf{U}_{2l}^*$ .



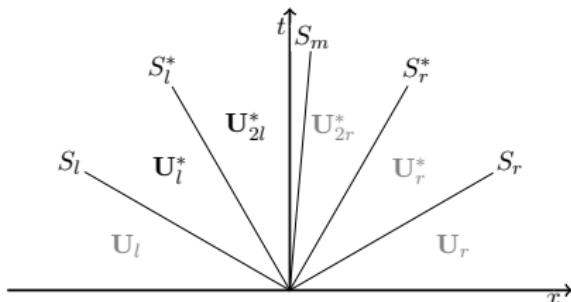
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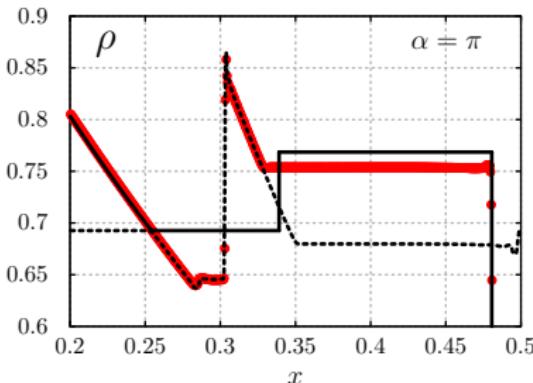
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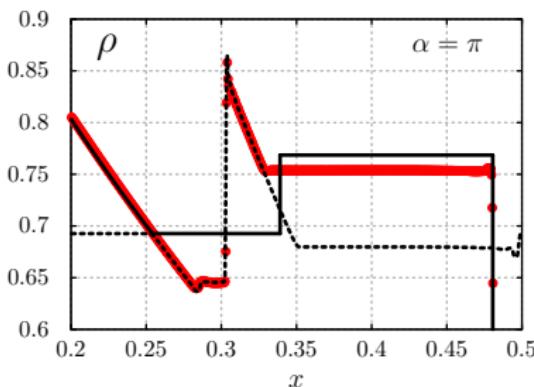
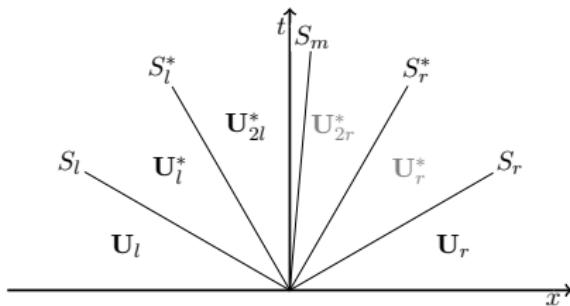
# Compound Wave modification



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- Reduce the contribution of  $\mathbf{F}^c = \mathbf{F}_{hlld}(\mathbf{U}_l^*, \mathbf{U}_{l2}^*)$  to the total flux.

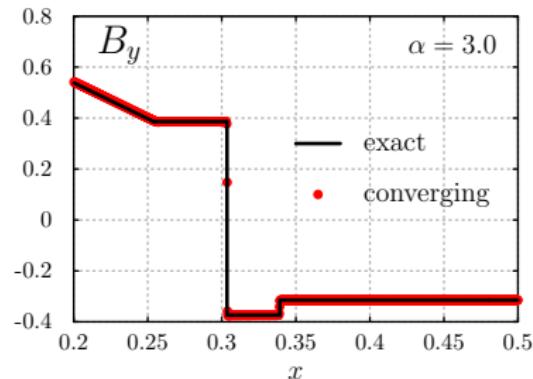
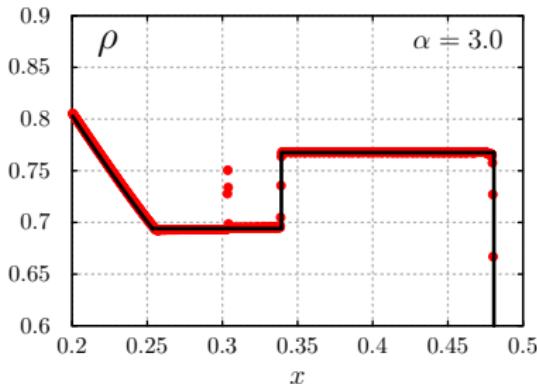


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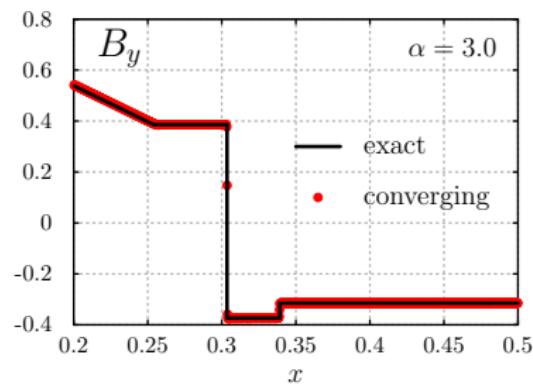
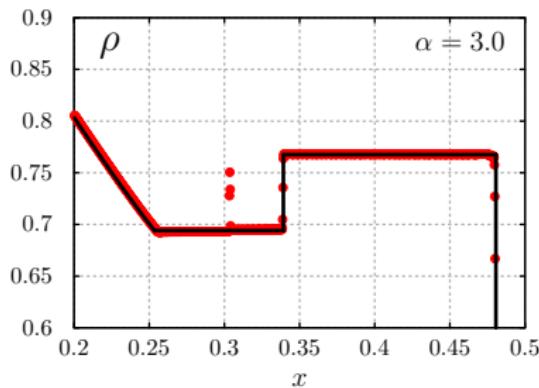
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- Reduce the contribution of  $\mathbf{F}^c = \mathbf{F}_{hlld}(\mathbf{U}_l^*, \mathbf{U}_{l2}^*)$  to the total flux.
- $\mathbf{F}_{cwm} = \mathbf{F}_{hlld}(\mathbf{U}_l, \mathbf{U}_r) - A\mathbf{F}^c$ , where  $A < 0.5$ .

# CWM solution



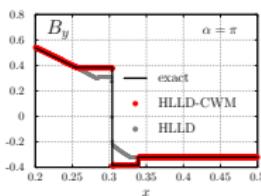
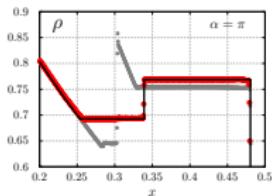
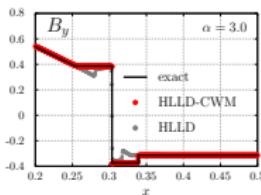
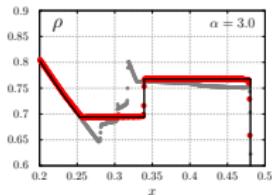
- CWM recovers the correct solution upstream and downstream of rotational discontinuity.
- Transition across rotational discontinuity is unresolved.
- Resolve it?

## CWM solution



- CWM recovers the correct solution upstream and downstream of rotational discontinuity.
- Transition across rotational discontinuity is unresolved.
- Resolve it?
- Simple but non-conservative.

# CWM Resolving the transition



- Jump conditions in Lagrangian mass coordinates,  $V = 1/\rho$ ,  $W$  is wave speed.
- Brackets denote difference across discontinuity.

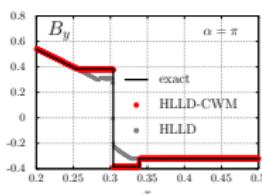
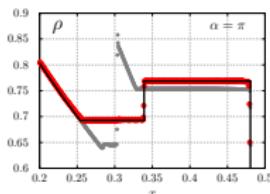
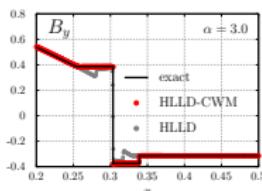
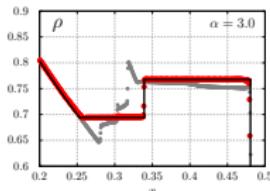
$$W[V] = -[v_n],$$

$$W[v_n] = -[P - B_n^2],$$

$$W[\mathbf{v}_t] = -B_n[\mathbf{B}_t],$$

$$W[V\mathbf{B}_t] = -B_n[\mathbf{v}_t],$$

# CWM Resolving the transition



- Across rotational discontinuity

$$[V] = 0,$$

$$[v_n] = 0,$$

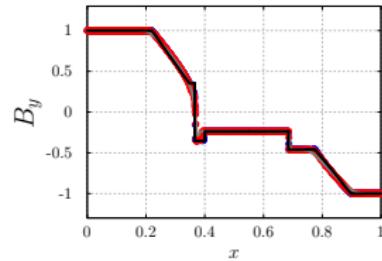
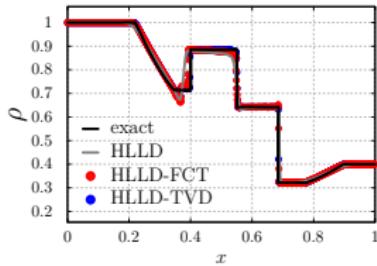
$$[p_g] = 0,$$

$$[B_t] = 0,$$

$$W[\mathbf{v}_t] = -B_n[\mathbf{B}_t],$$

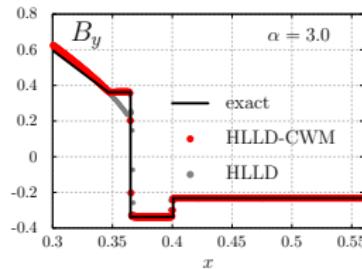
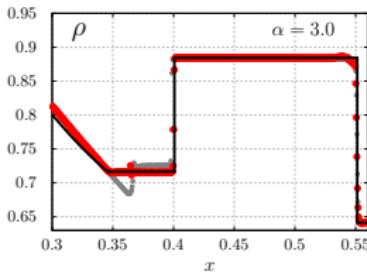
- $W = \sqrt{\rho}B_n$  is Lagrangian speed of linear wave.

# Examples



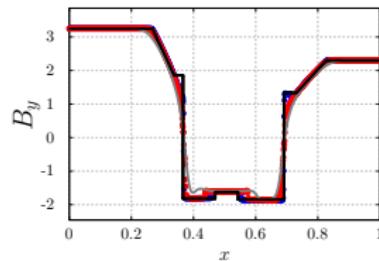
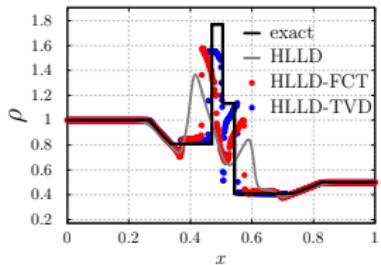
- Near coplanar initial conditions.
- Fast compound wave at  $x = 0.365$

# Examples



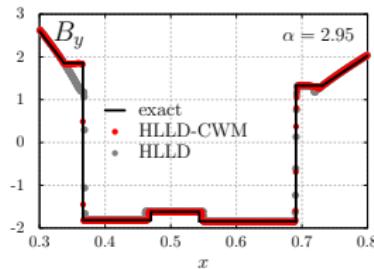
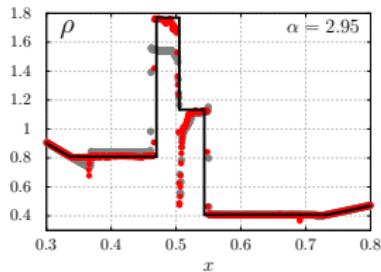
- Near coplanar initial conditions.
- Fast compound wave at  $x = 0.365$
- Small deviation for weak intermediate shocks.

# Examples



- Non-planar initial conditions.
- Fast compound waves at  $x = 0.366$  and  $x = 0.691$

# Examples



- Non-planar initial conditions.
- Fast compound waves at  $x = 0.366$  and  $x = 0.691$
- Small deviation for weak intermediate shocks.

# Outline

Introduction

Shocks in space plasma

Overview

Riemann problems of ideal MHD

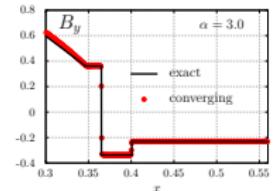
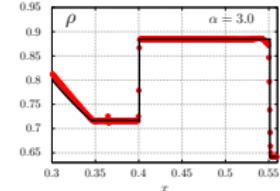
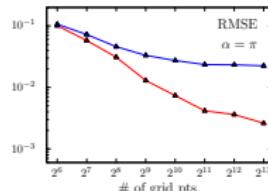
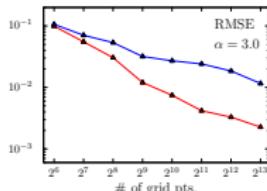
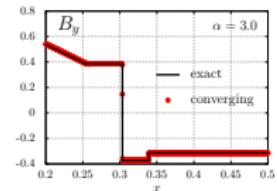
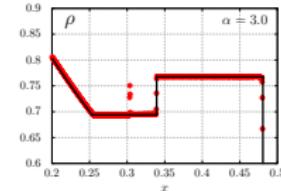
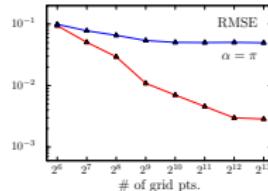
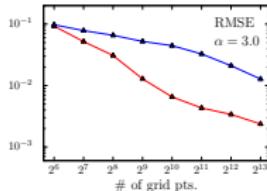
Compound wave formation in ideal MHD

Removing compound wave formation in ideal MHD

Convergence results using new modification

Shared memory parallelism

# Error analysis



- Error calculation without applying the correction at the rotational discontinuity.
- Without CWM, the compound wave starts to break apart between  $2^{10}$  and  $2^{11}$ .
- CWM produces convergence at low grid resolutions.

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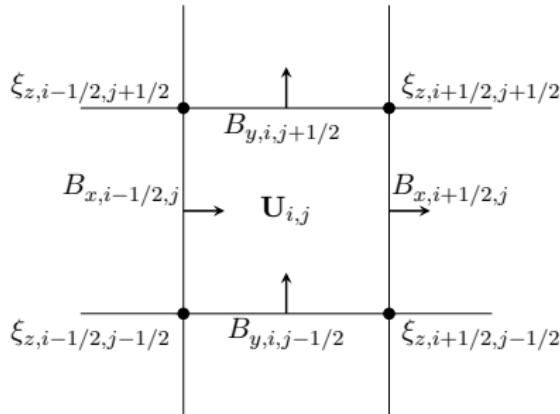
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# Higher dimensions

Maintaining  $\nabla \cdot \mathbf{B} = 0$  with constrained transport [2].

- Staggered grid.
- Hydrodynamical variables at cell centers.
- Magnetic field at interface.
- $\mathbf{E} = -\mathbf{v} \times \mathbf{B}$  at corners.
- Denote z-component of the emf as  $\xi_z$ .



Finite area integration of interface  $\mathbf{B}$

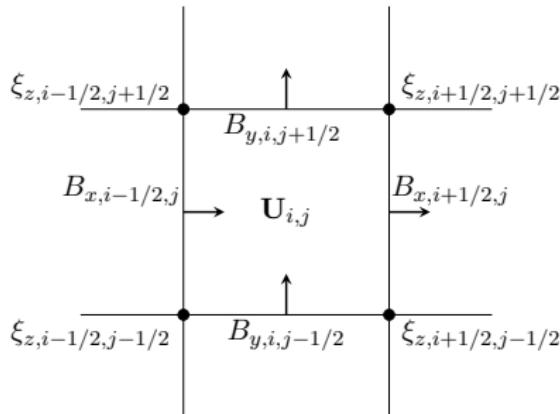
$$B_{x,i+1/2,j}^{n+1} = B_{x,i+1/2,j}^n - \frac{\delta t}{\delta y} (\xi_{z,i+1/2,j+1/2} - \xi_{z,i+1/2,j-1/2})$$

$$B_{y,i,j+1/2}^{n+1} = B_{y,i,j+1/2}^n + \frac{\delta t}{\delta x} (\xi_{z,i+1/2,j+1/2} - \xi_{z,i-1/2,j-1/2})$$

# Higher dimensions

Maintaining  $\nabla \cdot \mathbf{B} = 0$  with constrained transport [2].

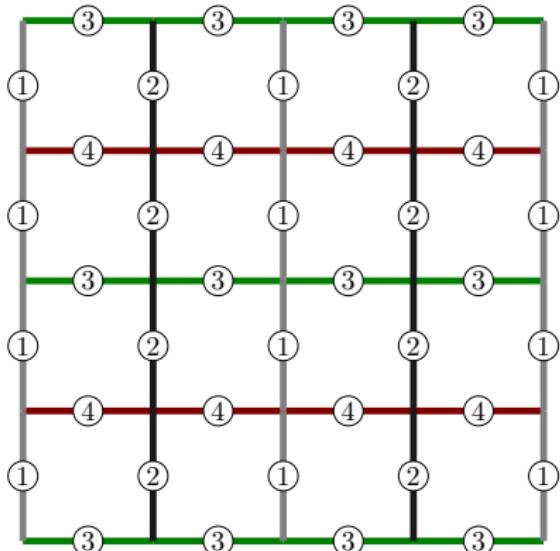
- Staggered grid.
- Hydrodynamical variables at cell centers.
- Magnetic field at interface.
- $\mathbf{E} = -\mathbf{v} \times \mathbf{B}$  at corners.
- Denote z-component of the emf as  $\xi_z$ .



Due to perfect cancellation, the numerical divergence in the cell remains zero to machine precision.

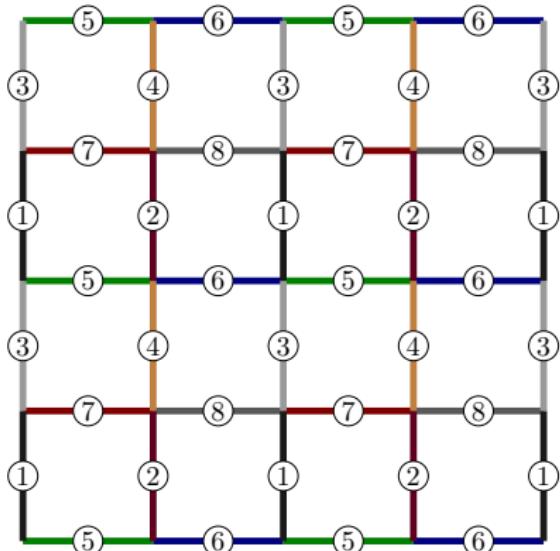
$$(\nabla \cdot \mathbf{B})_{i,j} = \frac{1}{\delta x} (B_{x,i+1/2,j} - B_{x,i-1/2,j}) + \frac{1}{\delta y} (B_{y,i,j+1/2} - B_{y,i,j-1/2})$$

# Shared memory parallelism



- Faces must be grouped by color to avoid memory contention.
- Loop over the faces becomes loop over the colors.

# Shared memory parallelism



- Faces must be grouped by color to avoid memory contention.
- Loop over the faces becomes loop over the colors.
- Constrained transport, faces and edges must be colored.

## Efficient algorithms

```
thrust::device_vector<float> x(n);    // independent  
                                         variable  
thrust::device_vector<float> y(n);    // y = f(x)  
thrust::device_vector<float> z(n);    // z = g(y)  
  
// compute y = f(x)  
thrust::transform(x.begin(),x.end(),y.begin(),f());  
  
// compute z = g(y)  
thrust::transform(y.begin(),y.end(),z.begin(),g());
```

- Function composition [3].
- $3n$  floats,  $2n$  reads,  $2n$  writes, and uses  $n$  temporary floats.

# Efficient algorithms

```
thrust::device_vector<float> x(n);    // independent
                                         variable
thrust::device_vector<float> z(n);    // z = g(y) = g(
                                         f(x))

// compute z = g(f(x))
thrust::transform(make_transform_iterator(x.begin(),
                                         f()),
                  make_transform_iterator(x.end(),
                                         () ),
                  z.begin(),
                  g());
```

- Function composition [3].
- $2n$  floats,  $n$  reads,  $n$  writes, and no temporary storage.

# Memory access

```
struct
    conservative_variables{
        float density;
        float momentum_x;
        float energy;
    }

    conservative_variables
    *state; // AoS

    state[i].density =
        some_number;
    state[i].momentum_x =
        another_number;
    state[i].energy =
        one_more_number;
```

```
struct
    conservative_variables{
        float *density;
        float *momentum_x;
        float *energy;
    }

    conservative_variables
    state; // SoA

    state.density[i] =
        some_number;
    state.momentum_x[i] =
        another_number;
    state.energy[i] =
        one_more_number;
```

- Memory coalescing occurs when multiple memory addresses are accessed with a single transaction.
- Memory does not coalesce with AoS (left).
- Memory does coalesce with SoA (right).

## Memory access

```
thrust::device_vector<primitive_variables>
    primitive_state(n); // AoS
thrust::device_vector<conservative_variables>
    conservative_state(n); // AoS
thrust::transform_n(primitive_state.begin(),
                   primitive_state.size(),
                   conservative_state.begin(),
                   convert_primitive_to_conservative(
                       gamma));
```

- Converting from primitive variables  $(\rho, v_x, p_g)$  to conservative variables  $(\rho, \rho v_x, en)$ .
- No coalescing.

## Memory access

```
thrust::device_vector<float> d(n), vx(n), pg(n);
thrust::device_vector<float> mx(n), en(n);
thrust::transform_n(
    thrust::make_zip_operator(
        make_tuple(d.begin(),
                   vx.begin(),
                   pg.begin())),
    n,
    thrust::make_zip_operator(
        make_tuple(d.begin(),
                   mx.begin(),
                   en.begin())),
    convert_primitive_to_conservative(gamma));
```

- Converting from primitive variables ( $\rho, v_x, p_g$ ) to conservative variables ( $\rho, \rho v_x, en$ ).
- Arrays can be combined on the fly with a `zip_operator` to achieve coalescing.

# Performance Comparison

Performance comparison for Orszag-Tang [6] test.

grid size	cells/second (GPU)	cells/second (CPU)	ratio
$64 \times 64$	$4.1894 \times 10^6$	$7.3955 \times 10^5$	5
$128 \times 128$	$1.6540 \times 10^7$	$7.5060 \times 10^5$	22
$256 \times 256$	$4.4497 \times 10^7$	$7.3155 \times 10^5$	60
$512 \times 512$	$6.3286 \times 10^7$	$7.9437 \times 10^5$	79
$1024 \times 1024$	$7.2134 \times 10^7$	$8.3354 \times 10^5$	86

- Dell Precision 7500 workstation with a (Dual CPU)
- CPU Intel Xeon E5645 @ 2.40 Ghz.
- GPU GeForce GTX TITAN with a memory bandwidth of 288.4 GB/sec and 2688 CUDA cores.
- Almost a factor of three increase of speed ratio from  $128 \times 128$  to  $256 \times 256$ .

# Conclusion

- Compound wave modification:
  - ▶ Produces rotational discontinuity for coplanar problems.
  - ▶ Removes pseudo-convergence for near coplanar problems.
  - ▶ Does not require exact solver.
  - ▶ FAST! HLLD intermediate states are already calculated.
  - ▶ Produces the correct result when other numerical inaccuracies are present.
- Parallel fluid solver.
  - ▶ General geometry on unstructured grids.
  - ▶ Capable of shared memory parallelism on GPU and CPU.

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