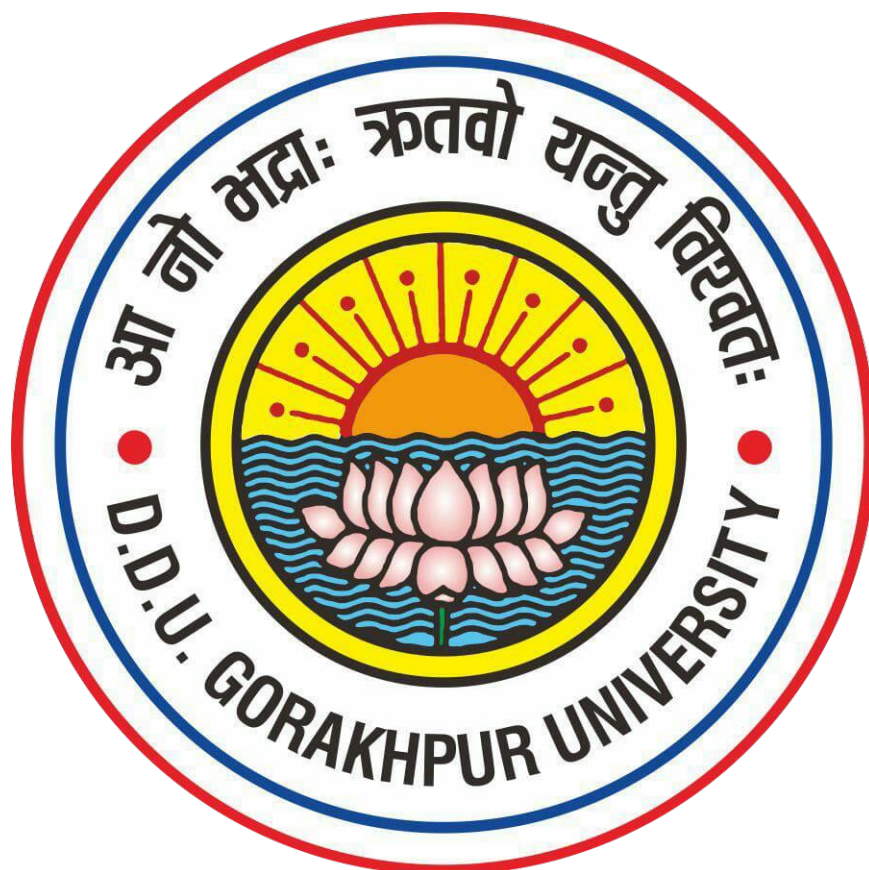


**Assignment**  
**Wavelet Analysis**

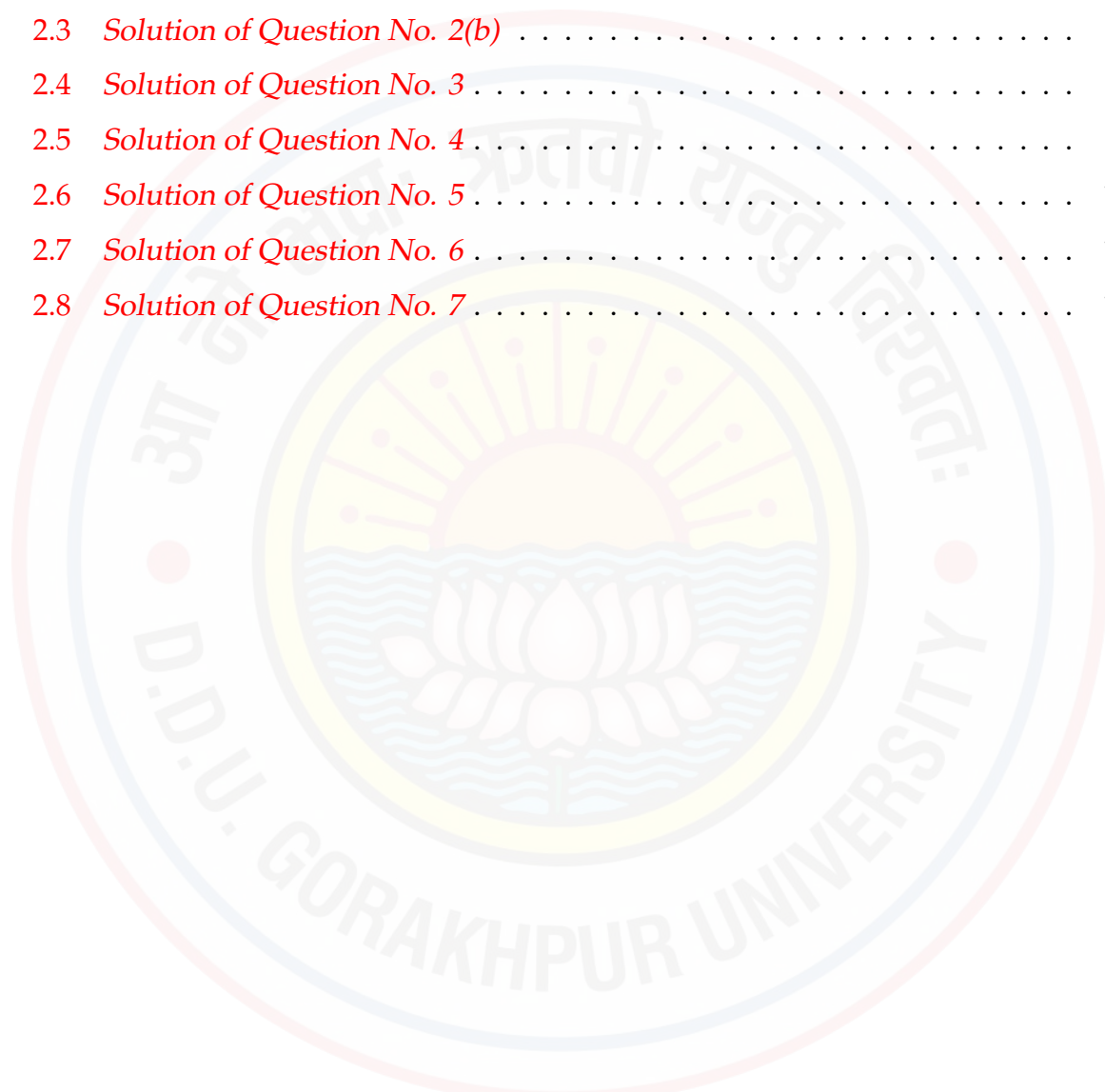


**Department of Mathematics and Statistics**  
**DDU Gorakhpur University, Gorakhpur (India)**

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# 1 Assignment Questions

1. Define a function  $f : \mathbb{R} \rightarrow \mathbb{R}$  by

$$f(t) = \begin{cases} \frac{1}{50\sqrt{t}} & , \quad t \in (0, 1); \\ 0 & , \quad \text{Otherwise} \end{cases}$$

Then show that  $f \in L^1(\mathbb{R})$  but  $f \notin L^2(\mathbb{R})$ .

2. Find the Fourier transform of the following functions:

$$(a) f(t) = \begin{cases} -20e^{-t} & , \quad t \geq 0; \\ 20e^t & , \quad t < 0 \end{cases}$$

$$(b) f(t) = \begin{cases} 10t & , \quad 0 \leq t < 1; \\ 0 & , \quad \text{Otherwise} \end{cases}$$

3. Define a function  $\psi(t)$  by

$$\psi(t) = \begin{cases} 3 & , \quad 0 \leq t < \frac{1}{2}; \\ -3 & , \quad \frac{1}{2} \leq t < 1; \\ 0 & , \quad \text{Otherwise} \end{cases}$$

Then show that (i)  $\|\psi\|_2 = 3$ . (ii)  $\int_{-\infty}^{\infty} \psi(t)dt = 0$ .

4. Let  $\psi \in L^2(\mathbb{R})$  be a basic wavelet and  $\phi$  be a bounded, integrable function, then  $\psi * \phi$  is also a basic wavelet.

5. Define Mexican Hat wavelet  $\psi_M(t)$  and prove the following

$$(i) \int_{-\infty}^{\infty} \psi_M(t)dt = 0.$$

$$(ii) C_{\psi_M} < \infty.$$

6. Let  $\delta(x)$  denotes the Dirac delta function. Then find the value of the following integrals:

$$(i) \int_{-\infty}^{\infty} \delta(x-1) \{ \cos^2(\pi x) + \sin^2(\pi x) + e^{\pi x} \} dx.$$

$$(ii) \int_{-\infty}^{\infty} \delta(x) \{ x^2 + x + 1 \} (e^x) dx.$$



7. Let us define a function  $g_\alpha(x)$  by

$$g_\alpha(x) = \frac{e^{-\frac{x^2}{4\alpha}}}{2\sqrt{\pi\alpha}}, \alpha > 0.$$

Then prove the following:

- (i)  $g \in L^1(\mathbb{R})$ .
- (ii)  $g \in L^2(\mathbb{R})$ .
- (iii)  $g \in L^1(\mathbb{R}) \cap L^2(\mathbb{R})$ .
- (iv)  $\|g_\alpha\|_1 = 1 \quad \forall \alpha > 0$ .

Does  $g \in L^4(\mathbb{R})$ ? Explain.



## 2 Solution of Assignment Questions

### 2.1 Solution of Question No. 1

For inclusivity of a function  $f(t)$  in  $L^1(\mathbb{R})$ , we have to show that

$$\|f\|_1^2 = \int_{-\infty}^{\infty} |f(t)| dt < \infty. \quad \forall t \in \mathbb{R}$$

Since in our problem

$$f(t) = \begin{cases} \frac{1}{50\sqrt{t}} & , \quad t \in (0, 1); \\ 0 & , \quad \text{Otherwise} \end{cases}$$

$$\begin{aligned} \text{So, } \|f\|_1^2 &= \int_{-\infty}^{\infty} |f(t)| dt = \int_{-\infty}^0 |f(t)| dt + \int_0^1 |f(t)| dt + \int_1^{\infty} |f(t)| dt \\ &= \int_{-\infty}^0 |0| dt + \int_0^1 \left| \frac{1}{50\sqrt{t}} \right| dt + \int_1^{\infty} |0| dt \\ &= \int_0^1 \frac{1}{50\sqrt{t}} dt = \frac{1}{50} \int_0^1 \frac{1}{\sqrt{t}} dt \\ &= \frac{1}{50} \lim_{\epsilon \rightarrow 0^+} \int_{\epsilon}^1 \frac{1}{\sqrt{t}} dt \\ &= \frac{1}{50} \lim_{\epsilon \rightarrow 0^+} \left[ 2\sqrt{t} \right]_{\epsilon}^1 \\ &= \frac{2}{50} \lim_{\epsilon \rightarrow 0^+} [1 - \sqrt{\epsilon}] \\ &= \frac{1}{25} \left[ 1 - \lim_{\epsilon \rightarrow 0^+} \sqrt{\epsilon} \right] \\ &= \frac{1}{25} [1 - 0] \\ &= \frac{1}{25} < \infty \end{aligned}$$

$$\text{i.e, } \|f\|_1^2 = \int_{-\infty}^{\infty} |f(t)| dt = \frac{1}{25} < \infty$$

thus,  $f \in L^1(\mathbb{R})$ .



Again, For inclusivity of a function  $f(t)$  in  $L^2(\mathbb{R})$ , we have to show that

$$\|f\|_2^2 = \int_{-\infty}^{\infty} |f(t)|^2 dt < \infty. \quad \forall t \in \mathbb{R}$$

$$\begin{aligned} \text{So, } \|f\|_2^2 &= \int_{-\infty}^{\infty} |f(t)|^2 dt = \int_{-\infty}^0 |f(t)|^2 dt + \int_0^1 |f(t)|^2 dt + \int_1^{\infty} |f(t)|^2 dt \\ &= \int_{-\infty}^0 |0|^2 dt + \int_0^1 \left| \frac{1}{50\sqrt{t}} \right|^2 dt + \int_1^{\infty} |0|^2 dt \\ &= \frac{1}{2500} \int_0^1 \frac{1}{|\sqrt{t}|^2} dt \\ &= \frac{1}{2500} \int_0^1 \frac{1}{t} dt \\ &= \frac{1}{2500} \lim_{\epsilon \rightarrow 0^+} \int_{\epsilon}^1 \frac{1}{t} dt \\ &= \frac{1}{2500} \lim_{\epsilon \rightarrow 0^+} [\ln t]_{\epsilon}^1 \\ &= \frac{1}{2500} \lim_{\epsilon \rightarrow 0^+} [\ln(1) - \ln(\epsilon)] \\ &= \frac{1}{2500} \left[ 0 - \lim_{\epsilon \rightarrow 0^+} \ln(\epsilon) \right] \\ &= \frac{1}{2500} [ -(-\infty) ] \\ &= \frac{1}{2500} [\infty] = \infty \\ \text{i.e, } \|f\|_2^2 &= \int_{-\infty}^{\infty} |f(t)|^2 dt = \frac{1}{2500} [\infty] = \infty \end{aligned}$$

thus,  $f \notin L^2(\mathbb{R})$ .

Thus we have shown that  $f \in L^1(\mathbb{R})$  but  $f \notin L^2(\mathbb{R})$ .



**2.2 Solution of Question No. 2(a)**

To find the Fourier transform of the function  $f(t)$ . Let's first check is it in  $L^1(\mathbb{R})$  or not ?

So,

$$\begin{aligned}
 \|f\|_1^2 &= \int_{-\infty}^{\infty} |f(t)| dt = \int_{-\infty}^0 |f(t)| dt + \int_0^{\infty} |f(t)| dt \\
 &= \int_{-\infty}^0 |20e^t| dt + \int_0^{\infty} |-20e^{-t}| dt \\
 &= 20 \left[ \int_{-\infty}^0 |e^t| dt + \int_0^{\infty} |e^{-t}| dt \right] \\
 &= 20 \left[ \int_{-\infty}^0 e^t dt + \int_0^{\infty} e^{-t} dt \right] \\
 &= 20 \left[ \lim_{M_1 \rightarrow -\infty} \int_{M_1}^0 e^t dt + \lim_{M_2 \rightarrow \infty} \int_0^{M_2} e^{-t} dt \right] \\
 &= 20 \left[ \lim_{M_1 \rightarrow -\infty} (e^t) \Big|_{M_1}^0 + \lim_{M_2 \rightarrow \infty} (-e^{-t}) \Big|_0^{M_2} \right] \\
 &= 20 \left[ \lim_{M_1 \rightarrow -\infty} (1 - e^{M_1}) + \lim_{M_2 \rightarrow \infty} (1 - e^{-M_2}) \right] \\
 &= 20 \left[ 1 - \lim_{M_1 \rightarrow -\infty} e^{M_1} + 1 - \lim_{M_2 \rightarrow \infty} e^{-M_2} \right] \\
 &= 20 [1 - 0 + 1 - 0]
 \end{aligned}$$

$$\int_{-\infty}^{\infty} |f(t)| dt = 40 < \infty$$

i.e,  $\|f\|_1^2 = \int_{-\infty}^{\infty} |f(t)| dt = 40 < \infty$

thus,  $f(t) \in L^1(\mathbb{R})$ .

Since the function  $f(t) \in L^1(\mathbb{R})$ , thus the Fourier transform exists and given as

$$\begin{aligned}
 \mathcal{F}[f(t)](\omega) &= \hat{f}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t} f(t) dt \quad \forall \omega \in \mathbb{R}. \\
 &= \int_{-\infty}^0 e^{-i\omega t} f(t) dt + \int_0^{\infty} e^{-i\omega t} f(t) dt
 \end{aligned}$$

$$\begin{aligned}
 \hat{f}(\omega) &= \int_{-\infty}^0 e^{-i\omega t} (20e^t) dt + \int_0^{\infty} e^{-i\omega t} (-20e^{-t}) dt \\
 &= 20 \left[ \int_{-\infty}^0 e^{-i\omega t} (e^t) dt - \int_0^{\infty} e^{-i\omega t} (e^{-t}) dt \right] \\
 &= 20 \left[ \int_{-\infty}^0 e^{-(1+i\omega)t} dt - \int_0^{\infty} e^{-(1+i\omega)t} dt \right] \\
 &= 20 \left[ \lim_{M_1 \rightarrow -\infty} \int_{M_1}^0 e^{-(1+i\omega)t} dt - \lim_{M_2 \rightarrow \infty} \int_0^{M_2} e^{-(1+i\omega)t} dt \right] \\
 &= 20 \left[ \left( -\frac{1}{(-1+i\omega)} \right) \lim_{M_1 \rightarrow -\infty} \left\{ e^{-(1+i\omega)t} \right\}_{M_1}^0 - \frac{1}{1+i\omega} \lim_{M_2 \rightarrow \infty} \left\{ e^{-(1+i\omega)t} \right\}_0^{M_2} \right] \\
 &= 20 \left[ \frac{1}{1-i\omega} \lim_{M_1 \rightarrow -\infty} \left\{ e^{-(1+i\omega)t} \right\}_{M_1}^0 - \frac{1}{1+i\omega} \lim_{M_2 \rightarrow \infty} \left\{ e^{-(1+i\omega)t} \right\}_0^{M_2} \right] \\
 &= 20 \left[ \frac{1}{1-i\omega} \lim_{M_1 \rightarrow -\infty} \left\{ 1 - e^{-(1+i\omega)M_1} \right\} - \frac{1}{1+i\omega} \lim_{M_2 \rightarrow \infty} \left\{ e^{-(1+i\omega)M_2} - 1 \right\} \right] \\
 &= 20 \left[ \frac{1}{1-i\omega} \left\{ 1 - \lim_{M_1 \rightarrow -\infty} e^{-(1+i\omega)M_1} \right\} - \frac{1}{1+i\omega} \left\{ \lim_{M_2 \rightarrow \infty} e^{-(1+i\omega)M_2} - 1 \right\} \right] \quad (1)
 \end{aligned}$$

Since,  $\lim_{M_1 \rightarrow -\infty} \left| e^{-(1+i\omega)M_1} \right|$

$$= \lim_{M_1 \rightarrow -\infty} \left\{ |e^{M_1}| \cdot |e^{-i\omega M_1}| \right\}$$

$$= \lim_{M_1 \rightarrow -\infty} e^{M_1} = 0$$

Therefore,  $\lim_{M_1 \rightarrow -\infty} e^{-(1+i\omega)M_1} = 0$

Similarly,  $\lim_{M_2 \rightarrow \infty} \left| e^{-(1+i\omega)M_2} \right|$

$$= \lim_{M_2 \rightarrow \infty} \left\{ |e^{-M_2}| \cdot |e^{-i\omega M_2}| \right\}$$

$$= \lim_{M_2 \rightarrow \infty} e^{-M_2} = 0$$

Therefore,  $\lim_{M_2 \rightarrow \infty} e^{-(1+i\omega)M_2} = 0$

Using these values in equation (1) we get as





$$\begin{aligned}
\hat{f}(\omega) &= 20 \left[ \frac{1}{1-i\omega} \{1-0\} - \frac{1}{1+i\omega} \{0-1\} \right] \\
&= 20 \left[ \frac{1}{1-i\omega} + \frac{1}{1+i\omega} \right] \\
&= 20 \left[ \frac{2}{1+\omega^2} \right]
\end{aligned}$$

$$\hat{f}(\omega) = \frac{40}{1+\omega^2}$$

### 2.3 Solution of Question No. 2(b)

Since,  $f(t) = \begin{cases} 10t & , \quad 0 \leq t < 1; \\ 0 & , \quad \text{Otherwise} \end{cases}$

To find the Fourier transform of the function  $f(t)$ , let's first check wheather is it in  $L^1(\mathbb{R})$  or not ?

So,

$$\begin{aligned}
\|f\|_1^2 &= \int_{-\infty}^{\infty} |f(t)| dt = \int_{-\infty}^0 |f(t)| dt + \int_0^1 |f(t)| dt + \int_1^{\infty} |f(t)| dt \\
&= \int_0^1 |10t| dt \\
&= 10 \int_0^1 t dt = 10 \left[ \frac{t^2}{2} \right]_0^1 = 5t \Big|_0^1 = 5 < \infty \\
\int_{-\infty}^{\infty} |f(t)| dt &= 5 < \infty
\end{aligned}$$

So  $f \in L^1(\mathbb{R})$ .Thats why Fourier transform exists and given by

$$\begin{aligned}
\mathcal{F}[f(t)](\omega) &= \hat{f}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t} f(t) dt \quad \forall \omega \in \mathbb{R}. \\
&= \int_{-\infty}^0 e^{-i\omega t} f(t) dt + \int_0^1 e^{-i\omega t} f(t) dt + \int_1^{\infty} e^{-i\omega t} f(t) dt \\
&= 10 \int_0^1 t e^{-i\omega t} dt \\
&= 10 \left[ \left. \frac{t e^{-i\omega t}}{-i\omega} \right|_0^1 - \int_0^1 \frac{e^{-i\omega t}}{-i\omega} dt \right]
\end{aligned}$$



$$\begin{aligned}
\hat{f}(\omega) &= 10 \left[ \frac{e^{-i\omega}}{-i\omega} - \frac{e^0}{-i\omega} - \int_0^1 \frac{e^{-i\omega t}}{-i\omega} dt \right] \\
&= 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) - \int_0^1 \frac{e^{-i\omega t}}{-i\omega} dt \right] \\
&= 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) - \frac{e^{-i\omega t}}{(-i\omega)^2} \Big|_0^1 \right] \\
&= 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) - \frac{e^{-i\omega}}{(-i\omega)^2} + \frac{1}{(-i\omega)^2} \right] \\
&= 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) + \frac{e^{-i\omega}}{\omega^2} + \frac{1}{\omega^2} \right] \\
&= 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) + \frac{e^{-i\omega}}{\omega^2} + \frac{1}{\omega^2} \right]
\end{aligned}$$

So, the Fourier transform of the function,  $f(t)$  is

$$\hat{f}(\omega) = 10 \left[ \frac{1}{-i\omega} (e^{-i\omega} - 1) + \frac{e^{-i\omega}}{\omega^2} + \frac{1}{\omega^2} \right]$$

## 2.4 Solution of Question No. 3

Given function is,

$$\psi(t) = \begin{cases} 3 & , \quad 0 \leq t < \frac{1}{2}; \\ -3 & , \quad \frac{1}{2} \leq t < 1; \\ 0 & , \quad \text{Otherwise} \end{cases}$$

(i)  $\|\psi\|_2$  is given by

$$\begin{aligned}
\|\psi\|_2^2 &= \int_{-\infty}^{\infty} |\psi(t)|^2 dt \\
&= \int_{-\infty}^0 |\psi(t)|^2 dt + \int_0^{\frac{1}{2}} |\psi(t)|^2 dt + \int_{\frac{1}{2}}^1 |\psi(t)|^2 dt + \int_1^{\infty} |\psi(t)|^2 dt \\
&= \int_0^{\frac{1}{2}} |3|^2 dt + \int_{\frac{1}{2}}^1 |-3|^2 dt \\
\|\psi\|_2^2 &= \int_0^1 3^2 dt = 9t \Big|_0^1 = 9 \implies \|\psi\|_2 = 3
\end{aligned}$$



(ii)  $\int_{-\infty}^{\infty} \psi(t)dt$  is evaluated as

$$\begin{aligned}\int_{-\infty}^{\infty} \psi(t)dt &= \int_{-\infty}^0 \psi(t)dt + \int_0^{\frac{1}{2}} \psi(t)dt + \int_{\frac{1}{2}}^1 \psi(t)dt + \int_1^{\infty} \psi(t)dt \\ &= \int_0^{\frac{1}{2}} 3dt + \int_{\frac{1}{2}}^1 (-3)dt \\ &= \int_0^1 (3 - 3)dt \\ \int_{-\infty}^{\infty} \psi(t)dt &= 0\end{aligned}$$

## 2.5 Solution of Question No. 4

It is provided in our problem that  $\psi \in L^2(\mathbb{R})$  be a basic wavelet. i.e

$$C_{\psi} = \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega < \infty$$

where  $\hat{\psi}(\omega)$  is the Fourier transform of the function  $\psi \in L^2(\mathbb{R})$ .

Again, it is also provided that  $\phi$  is a bounded, integrable function.

Then we have to show that  $C_{\psi * \phi} < \infty$ . i.e  $\psi * \phi$  is also a basic wavelet.

To do this, since we have

$$C_{\psi} = \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega)|^2}{|\omega|} d\omega < \infty$$

and we also know that,

$$(\psi * \phi)(t) = \int_{-\infty}^{\infty} \psi(t - u)\phi(u)du$$

thus,

$$(\hat{\psi * \phi})(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t} (\psi * \phi)(t)dt = \hat{\psi}(\omega) \cdot \hat{\phi}(\omega)$$

then,

$$\begin{aligned}C_{\psi * \phi} &= \int_{-\infty}^{\infty} \frac{|(\hat{\psi * \phi})(\omega)|^2}{|\omega|} d\omega \\ &= \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega) \cdot \hat{\phi}(\omega)|^2}{|\omega|} d\omega\end{aligned}$$



$$C_{\psi * \phi} = \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega)|^2 \cdot |\hat{\phi}(\omega)|^2}{|\omega|} d\omega$$

Since  $\psi \in L^2(\mathbb{R})$  is a basic wavelet and the  $\phi$  is a bounded, integrable function then the overall integration should be less than  $\infty$ .i.e.

$$C_{\psi * \phi} = \int_{-\infty}^{\infty} \frac{|\hat{\psi}(\omega)|^2 \cdot |\hat{\phi}(\omega)|^2}{|\omega|} d\omega < \infty$$

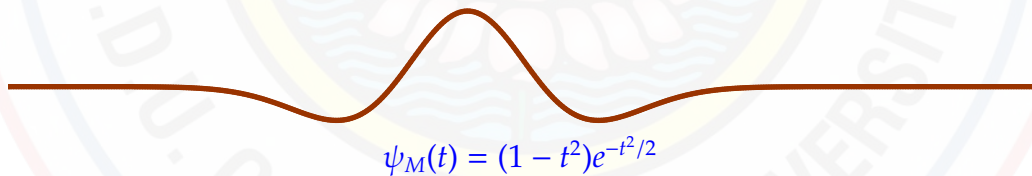
Hence,  $C_{\psi * \phi} < \infty$ . So,  $(\psi * \phi)$  is a basic wavelet.

## 2.6 Solution of Question No. 5

### Mexican Hat wavelet(Definition)

The Mexican Hat wavelet/Ricker wavelet/Second derivative wavelet is denoted by  $\psi_M(t)$  and defined by

$$\psi_M(t) = -\frac{d^2}{dt^2} (e^{-t^2/2}) = (1 - t^2) e^{-t^2/2}$$



(I) We have to show  $\int_{-\infty}^{\infty} \psi_M(t) dt = 0$ . To do so

$$\int_{-\infty}^{\infty} \psi_M(t) dt = \int_{-\infty}^{\infty} (1 - t^2) e^{-t^2/2} dt$$

$$\int_{-\infty}^{\infty} \psi_M(t) dt = \int_{-\infty}^{\infty} e^{-t^2/2} dt - \int_{-\infty}^{\infty} t^2 e^{-t^2/2} dt$$



$$\int_{-\infty}^{\infty} \psi_M(t) dt = I_1 + I_2 \quad (1)$$

$$\text{where, } I_1 = \int_{-\infty}^{\infty} e^{-t^2/2} dt \text{ and } I_2 = \int_{-\infty}^{\infty} t^2 e^{-t^2/2} dt.$$

Let's evaluate  $I_1$ ,

$$I_1 = \int_{-\infty}^{\infty} e^{-t^2/2} dt = 2 \int_0^{\infty} e^{-t^2/2} dt$$

$$\text{Let, } \frac{t^2}{2} = u \implies t = \sqrt{2u} \implies dt = \frac{1}{\sqrt{2}} u^{-\frac{1}{2}} du$$

$$I_1 = 2 \int_0^{\infty} e^{-u} \frac{1}{\sqrt{2}} u^{-\frac{1}{2}} du$$

$$= \frac{2}{\sqrt{2}} \int_0^{\infty} e^{-u} u^{-\frac{1}{2}} du$$

$$= \sqrt{2} \int_0^{\infty} e^{-u} u^{\frac{1}{2}-1} du$$

$$= \sqrt{2} \Gamma \frac{1}{2}$$

$$= \sqrt{2} \cdot \sqrt{\pi}$$

$$I_1 = \sqrt{2\pi} \quad (2)$$

Next, Let's evaluate  $I_2$ ,

$$I_2 = \int_{-\infty}^{\infty} t^2 e^{-t^2/2} dt = 2 \int_0^{\infty} t^2 e^{-t^2/2} dt$$

$$\text{Let, } \frac{t^2}{2} = u \implies t^2 = 2u \implies dt = \frac{1}{\sqrt{2}} u^{-\frac{1}{2}} du$$

$$I_2 = 2 \int_0^{\infty} 2u e^{-u} \frac{1}{\sqrt{2}} u^{-\frac{1}{2}} du$$

$$= \frac{2 \cdot 2}{\sqrt{2}} \int_0^{\infty} u e^{-u} u^{-\frac{1}{2}} du$$



$$\begin{aligned} I_2 &= \frac{2 \cdot 2}{\sqrt{2}} \int_0^{\infty} e^{-u} u^{\frac{1}{2}} du \\ &= \frac{2 \cdot 2}{\sqrt{2}} \int_0^{\infty} e^{-u} u^{\frac{3}{2}-1} du \\ &= \frac{2 \cdot 2}{\sqrt{2}} \Gamma \frac{3}{2} \\ &= \frac{2 \cdot 2}{\sqrt{2}} \cdot \frac{1}{2} \Gamma \frac{1}{2} \\ &= \frac{2 \cdot 2}{\sqrt{2}} \cdot \frac{\sqrt{\pi}}{2} \\ I_2 &= \sqrt{2\pi} \end{aligned} \quad (3)$$

Using equation (2) and (3) in equation (1), we get as

$$\int_{-\infty}^{\infty} \psi_M(t) dt = \sqrt{2\pi} - \sqrt{2\pi} = 0$$

(II) We have to show that  $C_{\psi_M} < \infty$ . To do so

Since we know that

$$\hat{\psi}_M(\omega) = \sqrt{2\pi} \omega^2 e^{-\omega^2/2}$$

Next,

$$\begin{aligned} C_{\psi_M} &= \int_{-\infty}^{\infty} \frac{|\hat{\psi}_M(\omega)|^2}{|\omega|} d\omega \\ &= \int_{-\infty}^{\infty} \frac{|2\pi \omega^4 e^{-\omega^2}|}{|\omega|} d\omega \\ &= 4\pi \int_0^{\infty} \frac{|\omega^4 e^{-\omega^2}|}{|\omega|} d\omega \\ &= 4\pi \int_0^{\infty} \frac{\omega^4 e^{-\omega^2}}{\omega} d\omega \\ &= 4\pi \int_0^{\infty} \omega^3 e^{-\omega^2} d\omega \end{aligned}$$





$$\text{Let } \omega^2 = u \implies \omega = \sqrt{u} \text{ and } d\omega = \frac{1}{2}u^{-\frac{1}{2}}du$$

$$\begin{aligned} C_{\psi_M} &= 4\pi \int_0^\infty u^{3/2} \cdot e^{-u} \frac{1}{2}u^{-1/2} du \\ &= 2\pi \int_0^\infty e^{-u} u du \\ &= 2\pi \cdot 1 = 2\pi \end{aligned}$$

Hence,

$$C_{\psi_M} = 2\pi < \infty$$

Thus,  $\psi_M$  is a basic wavelet.

## 2.7 Solution of Question No. 6

The Dirac Delta function  $\delta(x)$  is very cool in the sense that

$$\delta(x) = \begin{cases} +\infty, & x = 0 \\ 0, & x \neq 0 \end{cases}$$

Its unique characteristics do not end there though, because when integrating the Dirac Delta function we would get

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

Or, if we have another function  $f(x)$  multiplied to the Dirac Delta function and integrating them we would get

$$\int_{-\infty}^{\infty} f(x)\delta(x) dx = f(x) \int_{-\infty}^{\infty} \delta(x) dx = f(0)$$

Since

$$\int_{-\infty}^{\infty} \delta(x) dx = \begin{cases} 0, & x \neq 0 \\ 1, & x = 0 \end{cases}$$

Therefore in the previous integral we would have



$$\int_{-\infty}^{\infty} f(x)\delta(x) dx = \int_{-\infty}^{\infty} f(0)\delta(0) dx = f(0)$$

What if we have  $\delta(x - a)$  ? It's the same thing! The only thing here is we need to satisfy the condition  $x - a = 0$  such that  $\int_{-\infty}^{\infty} \delta(x - a)dx = 1$ . If the bounds of our integral though is not infinity, we need to make sure that the if we let  $x = a$ ,  $a$  would be in the bounds of the integral or else the integral would evaluate into zero.

$$\int_{-b}^b f(x)\delta(x - a) dx = \begin{cases} 0, & \text{if } b < a \text{ or } -b > a \text{ such that we cannot let } x = a \\ f(0), & \text{if } -b < a < b \text{ such that we can let } x = a \end{cases}$$

(i) We have to evaluate,  $\int_{-\infty}^{\infty} \delta(x - 1) \{ \cos^2(\pi x) + \sin^2(\pi x) + e^{\pi x} \} dx$ .

In the view of above discussion about Dirac Delta function we can evaluate above integral easily, So

$$\int_{-\infty}^{\infty} \delta(x - 1) \{ \cos^2(\pi x) + \sin^2(\pi x) + e^{\pi x} \} dx = f(1)$$

where,  $f(x) = \{ \cos^2(\pi x) + \sin^2(\pi x) + e^{\pi x} \}$ .

Thus

$$\int_{-\infty}^{\infty} \delta(x - 1) \{ \cos^2(\pi x) + \sin^2(\pi x) + e^{\pi x} \} dx = \{ \cos^2(\pi) + \sin^2(\pi) + e^{\pi} \} = 1 + e^{\pi}$$

(ii) We have to evaluate,  $\int_{-\infty}^{\infty} \delta(x) \{ x^2 + x + 1 \} (e^x) dx$ .

As, similar to the above solution, we can found this one also like,

$$\int_{-\infty}^{\infty} \delta(x) \{ x^2 + x + 1 \} (e^x) dx = f(0)$$

where,  $f(x) = \{ x^2 + x + 1 \} (e^x)$ .

Thus

$$\int_{-\infty}^{\infty} \delta(x) \{ x^2 + x + 1 \} (e^x) dx = \{ 0^2 + 0 + 1 \} (e^0) = 1.$$

**2.8 Solution of Question No. 7**

We have given a function  $g_\alpha(x)$  defined as

$$g_\alpha(x) = \frac{e^{-\frac{x^2}{4\alpha}}}{2\sqrt{\pi\alpha}}, \alpha > 0.$$

(i) To show that  $g \in L^1(\mathbb{R})$ , we have to show  $\|g_\alpha\|_1^2 = \int_{-\infty}^{\infty} |g_\alpha(x)| dx < \infty$ . Thus

$$\begin{aligned} \|g_\alpha\|_1^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)| dx = \int_{-\infty}^{\infty} \left| \frac{e^{-\frac{x^2}{4\alpha}}}{2\sqrt{\pi\alpha}} \right| dx \\ &= \int_{-\infty}^{\infty} \frac{|e^{-\frac{x^2}{4\alpha}}|}{|2\sqrt{\pi\alpha}|} dx \\ &= \frac{1}{2\sqrt{\pi\alpha}} \int_{-\infty}^{\infty} e^{-\frac{x^2}{4\alpha}} dx \\ &= \frac{1}{\sqrt{\pi\alpha}} \int_0^{\infty} e^{-\frac{x^2}{4\alpha}} dx \end{aligned}$$

$$\text{Let, } u = \frac{x^2}{4\alpha} \implies dx = \sqrt{\alpha} u^{-\frac{1}{2}} du$$

$$\begin{aligned} \|g_\alpha\|_1^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)| dx = \frac{1}{\sqrt{\pi\alpha}} \int_0^{\infty} e^{-u} \sqrt{\alpha} u^{-\frac{1}{2}} du \\ &= \frac{\sqrt{\alpha}}{\sqrt{\pi\alpha}} \int_0^{\infty} e^{-u} u^{-\frac{1}{2}} du \\ &= \frac{1}{\sqrt{\pi}} \int_0^{\infty} e^{-u} u^{\frac{1}{2}-1} du \\ &= \frac{1}{\sqrt{\pi}} \Gamma\left(\frac{1}{2}\right) \\ &= \frac{\sqrt{\pi}}{\sqrt{\pi}} \end{aligned}$$

$$\|g_\alpha\|_1^2 = \int_{-\infty}^{\infty} |g_\alpha(x)| dx = 1 < \infty$$

Hence,  $g_\alpha \in L^1(\mathbb{R})$ .



(ii) To show that  $g \in L^2(\mathbb{R})$ , we have to show  $\|g_\alpha\|_2^2 = \int_{-\infty}^{\infty} |g_\alpha(x)|^2 dx < \infty$ . Thus

$$\begin{aligned}\|g_\alpha\|_2^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)|^2 dx = \int_{-\infty}^{\infty} \left| \frac{e^{-\frac{x^2}{4\alpha}}}{2\sqrt{\pi\alpha}} \right|^2 dx \\ &= \int_{-\infty}^{\infty} \frac{|e^{-\frac{x^2}{4\alpha}}|^2}{|2\sqrt{\pi\alpha}|^2} dx \\ &= \frac{1}{4\pi\alpha} \int_{-\infty}^{\infty} e^{-\frac{x^2}{2\alpha}} dx \\ &= \frac{1}{2\pi\alpha} \int_0^{\infty} e^{-\frac{x^2}{2\alpha}} dx\end{aligned}$$

$$\text{Let, } u = \frac{x^2}{2\alpha} \implies dx = \sqrt{\frac{\alpha}{2}} u^{-\frac{1}{2}} du$$

$$\begin{aligned}\|g_\alpha\|_2^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)|^2 dx = \frac{1}{2\pi\alpha} \int_0^{\infty} e^{-u} \sqrt{\frac{\alpha}{2}} u^{-\frac{1}{2}} du \\ &= \frac{1}{2\pi\alpha} \cdot \sqrt{\frac{\alpha}{2}} \int_0^{\infty} e^{-u} \cdot u^{\frac{1}{2}-1} du \\ &= \frac{1}{2\pi\alpha} \cdot \sqrt{\frac{\alpha}{2}} \Gamma\left(\frac{1}{2}\right) \\ &= \frac{\sqrt{\pi}}{2\pi\alpha} \cdot \sqrt{\frac{\alpha}{2}}\end{aligned}$$

$$\|g_\alpha\|_2^2 = \int_{-\infty}^{\infty} |g_\alpha(x)|^2 dx = \frac{1}{\sqrt{8\pi\alpha}} < \infty$$

Hence,  $g_\alpha \in L^2(\mathbb{R})$ .

(iii) We have to show that  $g_\alpha \in L^1(\mathbb{R}) \cap L^2(\mathbb{R})$ .

Since from above discussion we see that  $g_\alpha \in L^1(\mathbb{R})$  as well as  $g_\alpha \in L^2(\mathbb{R})$ .

Thus,  $g_\alpha \in L^1(\mathbb{R}) \cap L^2(\mathbb{R})$ .

(iv) We have to show that  $\|g_\alpha\|_1 = 1$ .

Since from (i), it is obvious that  $\|g_\alpha\|_1 = 1$ . Hence, there is nothing to do more.



Yes,  $g_\alpha \in L^4(\mathbb{R})$ , Here is the explanation,

**Explanation that  $g_\alpha \in L^4(\mathbb{R})$**

To show the inclusivity of  $g_\alpha \in L^4(\mathbb{R})$ . We must show that  $\|g_\alpha\|_4^2 = \int_{-\infty}^{\infty} |g_\alpha(x)|^4 dx < \infty$ .

Thus,

$$\begin{aligned}\|g_\alpha\|_4^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)|^4 dx = \int_{-\infty}^{\infty} \left| \frac{e^{-\frac{x^2}{4\alpha}}}{2\sqrt{\pi\alpha}} \right|^4 dx \\ &= \int_{-\infty}^{\infty} \frac{|e^{-\frac{x^2}{4\alpha}}|^4}{|2\sqrt{\pi\alpha}|^4} dx \\ &= \frac{1}{16\pi^2\alpha^2} \int_{-\infty}^{\infty} e^{-\frac{x^2}{\alpha}} dx \\ &= \frac{1}{8\pi^2\alpha^2} \int_0^{\infty} e^{-\frac{x^2}{\alpha}} dx\end{aligned}$$

$$\boxed{\text{Let, } u = \frac{x^2}{\alpha} \implies dx = \sqrt{\alpha} \cdot \frac{1}{2} u^{-\frac{1}{2}} du}$$

$$\begin{aligned}\|g_\alpha\|_4^2 &= \int_{-\infty}^{\infty} |g_\alpha(x)|^4 dx = \frac{1}{8\pi^2\alpha^2} \int_0^{\infty} e^{-u} \sqrt{\alpha} \cdot \frac{1}{2} u^{-\frac{1}{2}} du \\ &= \frac{\sqrt{\alpha}}{16\pi^2\alpha^2} \int_0^{\infty} e^{-u} \cdot u^{\frac{1}{2}-1} du \\ &= \frac{\sqrt{\alpha}}{16\pi^2\alpha^2} \Gamma\left(\frac{1}{2}\right) \\ &= \frac{\sqrt{\pi\alpha}}{16\pi^2\alpha^2}\end{aligned}$$

$$\|g_\alpha\|_4^2 = \int_{-\infty}^{\infty} |g_\alpha(x)|^4 dx = \frac{\sqrt{\pi\alpha}}{16\pi^2\alpha^2} < \infty$$

Hence,  $g_\alpha \in L^4(\mathbb{R})$ . It completes the explanation.