

Излучение нестационарных полей и их распространение в нелинейном пространстве

В данной теоретической работе получены нелинейные свойства электромагнитного поля путем сравнения аналитических решений классических задач нестационарной электродинамики для линейных и нелинейных случаев распространения.

Решения модельных задач получены без перехода в частотную область с применением метода эволюционных уравнений. Нелинейная модель построена на гипотезе появления вторичных источников и предложен итеративный метод их учета.

Пространство распространения неограниченно, поэтому метод модового базиса не может быть применен напрямую, а упомянутый метод эволюционных уравнений, фактически, является его расширением на бесконечное или полубесконечное пространство. На рассмотрение предложена нестационарная слоисто-неоднородная среда.

©Ахмедов Ролан

1 Введение

Equation Section (Next)

1.1 Основные определения

1.2 Теоретическое обоснование

1.3 Область применения

2 Функция Бесселя первого рода

Equation Section (Next)

2.1 Определение и линейные свойства

$$J_{-n} z = -1^n J_n z \quad (2.1)$$

$$J_{n+1} z + J_{n-1} z = \frac{2n}{z} J_n z \quad (2.2)$$

2.2 Интегро-дифференциальные свойства

$$2 \frac{d}{dz} J_n z = J_{n-1} z - J_{n+1} z \quad (2.3)$$

$$\frac{d}{dz} J_n z = J_{n-1} z - \frac{n}{z} J_n z \quad (2.4)$$

$$\frac{d}{dz} J_n z = \frac{n}{z} J_n z - J_{n+1} z \quad (2.5)$$

$$\frac{d}{dz} \frac{J_n z}{z^n} = - \frac{J_{n+1} z}{z^n} \quad (2.6)$$

$$\frac{d}{dz} [z^n J_n z] = z^n J_{n-1} z \quad (2.7)$$

2.3 Интеграл 1

$$I_1 = \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} \quad (2.8)$$

Ватсон «Теория бесселевых функций» (п. 13.46, ст. 450)

$$\int_0^\infty \frac{dt}{t^{\lambda+\nu}} J_\mu at J_\nu bt J_\nu ct = \frac{bc/2^\nu}{\Gamma(\nu+1/2) \Gamma(1/2)} \int_0^\pi \int_0^\pi \frac{J_\mu at J_\nu \omega t}{\omega^\nu t^\lambda} \sin^{2\nu} \varphi d\varphi dt$$

$$\omega = \sqrt{b^2 + c^2 - 2bc \cdot \cos \varphi}$$

$$\int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} = \frac{\rho R}{2\Gamma(3/2) \Gamma(1/2)} \int_0^\pi \frac{\sin^2 \varphi}{\sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi}} \times$$

$$\times \int_0^\infty d\nu J_1 \omega \nu J_0 \nu \sqrt{c^2 t^2 - z^2} d\varphi$$

$$\omega = \sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi}$$

Прудников, Брычков, Маричев, 1983 (2.13.32, ст. 209)

$$\int_0^{\infty} dv J_n \left(\frac{a}{v} \right) J_{n-1} \left(\frac{b}{v} \right) = \begin{cases} b^{n-1}/a^n & , b < a \\ 0 & , b > a \end{cases}$$

$$\sin \frac{\varphi}{2} = \pm \sqrt{\frac{1 - \cos \varphi}{2}}$$

$$1 - \cos \varphi = \begin{cases} 2 \sin^2 \frac{\varphi}{2} & , 0 < \varphi < \pi \\ -2 \sin^2 \frac{\varphi}{2} & , \pi < \varphi < 2\pi \end{cases}$$

$$\int_0^{\infty} dv J_1 \left(\nu \sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi} \right) J_0 \left(\nu \sqrt{c^2 t^2 - z^2} \right) = \frac{1}{\sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi}}$$

$$\begin{aligned} \sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi} &> \sqrt{c^2 t^2 - z^2} \\ \rho^2 + R^2 - 2\rho R \cdot \cos \varphi - 2\rho R + 2\rho R &> c^2 t^2 - z^2 \\ \rho - R^2 + 2\rho R \cdot 1 - \cos \varphi &> c^2 t^2 - z^2 \\ \rho - R^2 + 4\rho R \sin^2 \frac{\varphi}{2} &> c^2 t^2 - z^2 \\ \varphi &> 2 \arcsin \sqrt{\frac{c^2 t^2 - z^2 - \rho - R^2}{4\rho R}} \end{aligned}$$

$$\psi = 2 \arcsin \sqrt{\frac{c^2 t^2 - z^2 - \rho - R^2}{4\pi\rho}}; \quad \psi \leq \varphi \leq \pi$$

$$\Gamma\left(\frac{3}{2}\right)\Gamma\left(\frac{1}{2}\right)=\left(\frac{1}{2}\sqrt{\pi}\right)\cdot\sqrt{\pi}=\frac{\pi}{2}$$

$$\begin{aligned} &\int_0^{\infty} \frac{dv}{v} J_1 \left(\nu R \right) J_1 \left(\nu \rho \right) J_0 \left(\nu \sqrt{c^2 t^2 - z^2} \right) = \frac{\rho R}{2\Gamma\left(\frac{3}{2}\right)\Gamma\left(\frac{1}{2}\right)} \times \\ &\times \int_0^{\pi} \frac{\sin^2 \varphi}{\sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi}} \int_0^{\infty} dv J_1 \left(\omega \nu \right) J_0 \left(\nu \sqrt{c^2 t^2 - z^2} \right) d\varphi = \\ &= \frac{\rho R}{\pi} \int_0^{\pi} \frac{\sin^2 \varphi}{\sqrt{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi}} \int_0^{\infty} dv J_1 \left(\omega \nu \right) J_0 \left(\nu \sqrt{c^2 t^2 - z^2} \right) d\varphi = \frac{\rho R}{\pi} \int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{\rho^2 + R^2 - 2\rho R \cdot \cos \varphi} = \\ &= \frac{\rho}{\pi R} \int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{\frac{\rho^2}{R^2} + 1 - \frac{2\rho}{R} \cdot \cos \varphi} \end{aligned}$$

$$\begin{aligned}
\int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{a+b \cdot \cos \varphi} &= \int_{\psi}^{\pi} \frac{1-\cos^2 \varphi}{a+b \cdot \cos \varphi} d\varphi = \int_{\psi}^{\pi} \frac{d\varphi}{a+b \cdot \cos \varphi} - \int_{\psi}^{\pi} \frac{\cos^2 \varphi + \frac{a}{b} \cos \varphi}{a+b \cdot \cos \varphi} d\varphi + \frac{a}{b} \int_{\psi}^{\pi} \frac{\cos \varphi d\varphi}{a+b \cdot \cos \varphi} = \\
&= \int_{\psi}^{\pi} \frac{d\varphi}{a+b \cdot \cos \varphi} - \frac{1}{b} \int_{\psi}^{\pi} \cos \varphi d\varphi + \frac{a}{b^2} \int_{\psi}^{\pi} \frac{a+b \cdot \cos \varphi}{a+b \cdot \cos \varphi} d\varphi - \frac{a^2}{b^2} \int_{\psi}^{\pi} \frac{d\varphi}{a+b \cdot \cos \varphi} = \\
&= \left(1 - \frac{a^2}{b^2}\right) \int_{\psi}^{\pi} \frac{d\varphi}{a+b \cdot \cos \varphi} + \frac{a}{b^2} \int_{\psi}^{\pi} d\varphi - \frac{1}{b} \int_{\psi}^{\pi} \cos \varphi d\varphi
\end{aligned}$$

Прудников, Бычков, Маричев «Интегралы и ряды: Элементарные функции» ст. 181,

$$\int \frac{d\varphi}{a+b \cdot \cos \varphi} = \frac{2}{\sqrt{a^2-b^2}} \arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\varphi}{2} \right)$$

$$\tan \frac{\pi}{2} = \infty \Rightarrow \arctan \left(\tan \frac{\pi}{2} \right) = \arctan \infty = \frac{\pi}{2}$$

$$\begin{aligned}
\int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{a+b \cdot \cos \varphi} &= \left(1 - \frac{a^2}{b^2}\right) \int_{\psi}^{\pi} \frac{d\varphi}{a+b \cdot \cos \varphi} + \frac{a}{b^2} \int_{\psi}^{\pi} d\varphi - \frac{1}{b} \int_{\psi}^{\pi} \cos \varphi d\varphi = \\
&= \left(1 - \frac{a^2}{b^2}\right) \frac{2}{\sqrt{a^2-b^2}} \arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\varphi}{2} \right) + \frac{a}{b^2} \varphi - \frac{\sin \varphi}{b} \Bigg|_{\psi}^{\pi} = \\
&= \left(1 - \frac{a^2}{b^2}\right) \frac{2}{\sqrt{a^2-b^2}} \left[\arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\pi}{2} \right) - \arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\psi}{2} \right) \right] + \frac{a}{b^2} \pi - \frac{a}{b^2} \psi - \\
&\quad - \frac{\sin \pi}{b} + \frac{\sin \psi}{b} = -\frac{2}{b^2} \frac{a^2-b^2}{\sqrt{a^2-b^2}} \left[\arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\pi}{2} \right) - \arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\psi}{2} \right) \right] + \frac{a}{b^2} \pi - \\
&\quad - \frac{a}{b^2} \psi - \frac{\sin \pi}{b} + \frac{\sin \psi}{b} = \frac{2\sqrt{a^2-b^2}}{b^2} \left[\arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\psi}{2} \right) - \frac{\pi}{2} \right] + \frac{a}{b^2} (\pi - \psi) \frac{\sin \psi}{b}
\end{aligned}$$

$$\tan \frac{\psi}{2} = \frac{\sin \frac{\psi}{2}}{\cos \frac{\psi}{2}} = \frac{\sin \frac{\psi}{2}}{\sqrt{1-\sin^2 \frac{\psi}{2}}} = \sqrt{\frac{\frac{c^2 t^2 - z^2 - \rho - R^2}{4\pi\rho}}{1 - \frac{c^2 t^2 - z^2 - \rho - R^2}{4\pi\rho}}} = \sqrt{\frac{c^2 t^2 - z^2 - \rho - R^2}{\rho + R^2 - c^2 t^2 + z^2}}$$

$$\int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{a+b \cdot \cos \varphi} = \frac{2\sqrt{a^2-b^2}}{b^2} \left[\arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \tan \frac{\psi}{2} \right) - \frac{\pi}{2} \right] + \frac{a}{b^2} \pi - \psi - \frac{\sin \psi}{b} =$$

$$= 2 \frac{\sqrt{a^2-b^2}}{b^2} \arctan \left(\frac{\sqrt{a^2-b^2}}{a+b} \sqrt{\frac{c^2 t^2 - z^2 - \rho - R^2}{\rho + R^2 - c^2 t^2 + z^2}} \right) - \pi \frac{\sqrt{a^2-b^2}}{b^2} + \frac{a}{b^2} \pi - \psi - \frac{\sin \psi}{b}$$

$$I_1 = \int_0^{\infty} \frac{d\nu}{\nu} J_1(\nu R) J_1(\nu \rho) J_0(\nu \sqrt{c^2 t^2 - z^2}) = \frac{\rho}{\pi R} \int_{\psi}^{\pi} \frac{\sin^2 \varphi d\varphi}{\frac{\rho^2}{R^2} + 1 - \frac{2\rho}{R} \cos \varphi}$$

$$I_1 = \frac{\rho^2 + R^2}{4\pi R \rho} \pi - 2\psi + \frac{\rho^2 - R^2}{4\pi R \rho} \left(\arctan \left(\frac{\rho + R}{\rho - R} \tan \frac{\psi}{2} \right) - \arctan \left(\frac{\rho - R}{\rho + R} \tan^{-1} \frac{\psi}{2} \right) \right) \quad (2.9)$$

2.4 Интеграл 2

$$I_2 = \int_0^{\infty} d\nu J_1(\nu R) J_0(\nu \rho) J_0(\nu \sqrt{c^2 t^2 - z^2}) \quad (2.10)$$

Прудников

$$\int_0^{\infty} J_0(ax) J_0(bx) J_1(cx) dx = \frac{1}{\pi c} \arccos \frac{a^2 + b^2 - c^2}{2ab}$$

$$|a-b| < c < a+b \quad a, b > 0 \quad (2.11)$$

$$I_2 = \frac{1}{\pi R} \arccos \frac{c^2 t^2 - z^2 + \rho^2 - R^2}{2\rho \sqrt{c^2 t^2 - z^2}}$$

$$I_2 = \quad (2.12)$$

3 Функция Ломмеля двух переменных

3.1 Определение и линейные свойства

3.2 Интегро-дифференциальные свойства

3.3 Интеграл 3

3.4 Интеграл 4

3.5 Интеграл 5

4 Метод эволюционных уравнений для цилиндрикой системы

$$\left\{ \begin{aligned} \partial_{ct} \varepsilon \partial_{ct} \mu h_m - \partial_z \mu^{-1} \partial_z \mu h_m + v^2 h_m &= \sqrt{\mu_0} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho [\vec{z}_0 \times \vec{J}] \cdot \nabla_\perp \Psi_m^* \nu - \\ &- \partial_{ct} \left\{ \sqrt{\varepsilon_0} \varepsilon \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Psi_m^* \nu I_z \right\} - \partial_z \left\{ \sqrt[3]{\mu_0} \mu^{-1} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Psi_m^* \nu g \right\} \\ \partial_{ct} \mu \partial_{ct} \varepsilon e_n - \partial_z \varepsilon^{-1} \partial_z \varepsilon e_n + \chi^2 e_n &= \sqrt{\varepsilon_0} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho [\vec{I} \times \vec{z}_0] \cdot \nabla_\perp \Phi_n^* \chi - \\ &- \partial_{ct} \left\{ \sqrt{\mu_0} \mu \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Phi_n^* \chi \right\} - \partial_z \left\{ \sqrt[3]{\varepsilon_0} \varepsilon^{-1} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Phi_n^* \chi \right\} \\ I_n^e &= -\partial_{ct} \left\{ \varepsilon e_n \right\} \sqrt{\mu_0} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Phi_n^* \chi \\ V_n^e &= \varepsilon^{-1} \partial_z \left\{ \varepsilon e_n \right\} \sqrt[3]{\varepsilon_0} \varepsilon^{-1} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Phi_n^* \chi \\ I_m^h &= \mu^{-1} \partial_z \left\{ \mu h_m \right\} \sqrt[3]{\mu_0} \mu^{-1} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Psi_m^* \nu \\ V_m^h &= -\partial_{ct} \left\{ \mu h_m \right\} \sqrt{\varepsilon_0} \frac{1}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \Psi_m^* \nu \end{aligned} \right.$$

$$\begin{aligned} \vec{E} &= \sqrt[3]{\varepsilon_0} \left\{ \sum_{m=-\infty}^{\infty} \int_0^\infty d\nu W_m^h \nabla_\perp \Psi_m \times \vec{z}_0 + \sum_{m=1}^{\infty} \int_0^\infty d\chi \mathcal{W}_n^e \nabla_\perp \Phi_n \right\} \\ \vec{H} &= \sqrt[3]{\mu_0} \left\{ \sum_{m=-\infty}^{\infty} \int_0^\infty d\nu I_m^h \nabla_\perp \Psi_m + \sum_{m=1}^{\infty} \int_0^\infty d\chi I_n^e \vec{z}_0 \times \nabla_\perp \Phi_n \right\} \end{aligned}$$

$$\begin{aligned} E_z(\rho, \phi, z, t) &= \sqrt[3]{\varepsilon_0} \sum_{n=0}^{\infty} \int_0^\infty \chi^2 d\chi e_n(z, t; \chi) \Phi_n(\rho, \phi; \chi) \\ H_z(\rho, \phi, z, t) &= \sqrt[3]{\mu_0} \sum_{m=0}^{\infty} \int_0^\infty \nu^2 d\nu h_m(z, t; \nu) \Psi_m(\rho, \phi; \nu) \end{aligned}$$

5 Нелинейное приложение эволюционного подхода

6 Методы учета нелинейные свойства слоистых сред

7 Излучение точечного источника

8 Импульсное излучение плоского диска с током

8.1 Постановка задачи

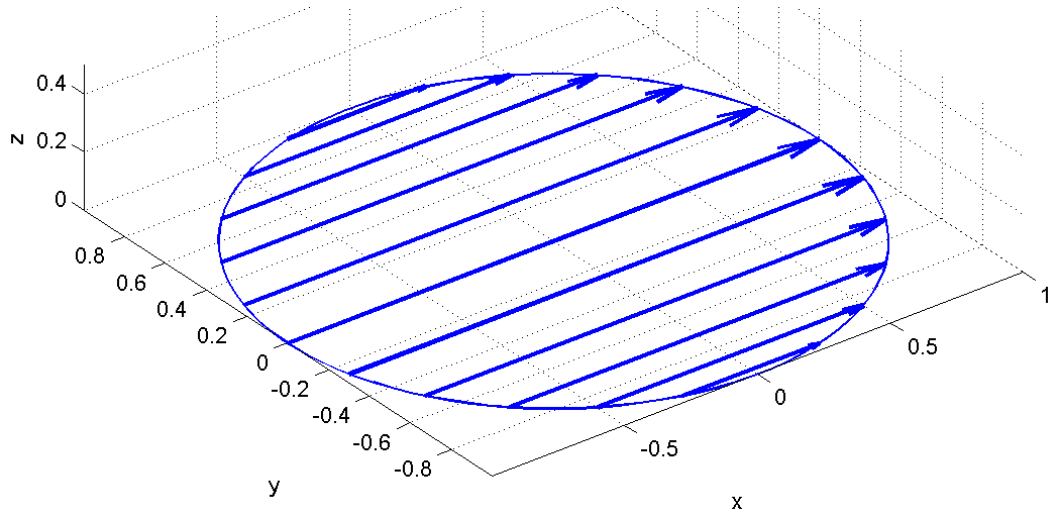


Рисунок 1. Распределение электрического тока

$$\vec{j}_0 = \vec{J} = \vec{x}_0 H(t) \delta z \quad H \rho - H \rho - R \quad (2.13)$$

8.2 Источник в базисной форме

$$\begin{cases} \vec{\rho}_0 = \vec{x}_0 \cos \varphi + \vec{y}_0 \sin \varphi \\ \vec{\varphi}_0 = -\vec{x}_0 \sin \varphi + \vec{y}_0 \cos \varphi \end{cases} \Rightarrow \vec{A} = \begin{pmatrix} \cos \varphi & \sin \varphi \\ -\sin \varphi & \cos \varphi \end{pmatrix}$$

$$\vec{j}_0 \vec{\rho}_0, \vec{\varphi}_0 = \vec{A} \times \vec{j}_0 \vec{x}_0, \vec{y}_0$$

$$\vec{j}_0 \vec{x}_0, \vec{y}_0 = \vec{A}^{-1} \times \vec{j}_0 \vec{\rho}_0, \vec{\varphi}_0$$

$$\vec{j}_0 = \begin{pmatrix} \cos \varphi & \sin \varphi \\ -\sin \varphi & \cos \varphi \end{pmatrix} \begin{pmatrix} H(t) \delta z \quad H \rho - H \rho - R \\ 0 \end{pmatrix}$$

$$= H(t) \delta z \quad H \rho - H \rho - R \quad \vec{\rho}_0 \cos \varphi - \vec{\varphi}_0 \sin \varphi$$

$$\vec{j}_0 = H(t) \delta z \quad H \rho - H \rho - R \quad \vec{\rho}_0 \cos \varphi - \vec{\varphi}_0 \sin \varphi$$

$$j_m(z, t, \nu) = \frac{\sqrt{\mu_0}}{2\pi} \int_0^{2\pi} d\varphi \int_0^\infty \rho d\rho \vec{j}_0 \left[\nabla_\perp \Psi_m^* \times \vec{z}_0 \right]$$

$$J_{m-1} \nu \rho + J_{m+1} \nu \rho = \frac{2m}{\nu \rho} J_m \nu \rho$$

$$im \vec{\rho}_0 \frac{J_m \nu \rho}{\rho \sqrt{\nu}} = i \vec{\rho}_0 \frac{\sqrt{\nu}}{2} \frac{2m}{\rho \nu} J_m \nu \rho = i \vec{\rho}_0 \frac{\sqrt{\nu}}{2} [J_{m-1} \nu \rho + J_{m+1} \nu \rho]$$

$$\begin{aligned} \left[\nabla_{\perp} \Psi_m^* \times \vec{z}_0 \right] &= -e^{-im\varphi} \left(\vec{\phi}_0 \sqrt{v} \frac{J_{m-1} v\rho - J_{m+1} v\rho}{2} + im\vec{\rho}_0 \frac{J_m v\rho}{\rho\sqrt{v}} \right) = \\ &= -e^{-im\varphi} \frac{\sqrt{v}}{2} \vec{\phi}_0 \left[J_{m-1} v\rho - J_{m+1} v\rho \right] + i\vec{\rho}_0 \left[J_{m-1} v\rho + J_{m+1} v\rho \right] \end{aligned}$$

$$\begin{aligned} \vec{j}_0 \left[\nabla_{\perp} \Psi_m^* \times \vec{z}_0 \right] &= -\sqrt{v} \frac{\cos m\varphi - i \sin m\varphi}{2} H(t) \delta_z \left[H\rho - H\rho - R \right] \times \\ &\times \vec{\phi}_0 \left[J_{m-1} v\rho - J_{m+1} v\rho \right] + i\vec{\rho}_0 \left[J_{m-1} v\rho + J_{m+1} v\rho \right] \vec{\rho}_0 \cos \varphi - \vec{\phi}_0 \sin \varphi = \\ &= -\sqrt{v} \frac{\cos m\varphi - i \sin m\varphi}{2} H(t) \delta_z \left[H\rho - H\rho - R \right] \times \\ &\times -\sin \varphi \left[J_{m-1} v\rho - J_{m+1} v\rho \right] + i \cos \varphi \left[J_{m-1} v\rho + J_{m+1} v\rho \right] \end{aligned}$$

$$\begin{aligned} j_m(z, t; v) &= \frac{\sqrt{\mu_0}}{2\pi} \int_0^{2\pi} d\varphi \int_0^{\infty} \rho d\rho \vec{j}_0 \left[\nabla_{\perp} \Psi_m^* \times \vec{z}_0 \right] = \\ &= \frac{\sqrt{\mu_0}}{2\pi} \frac{\sqrt{v}}{2} H(t) \delta_z \int_0^{2\pi} \sin \varphi \left[\cos m\varphi - i \sin m\varphi \right] d\varphi \int_0^R \rho d\rho \left[J_{m-1} v\rho - J_{m+1} v\rho \right] - \\ &- i \frac{\sqrt{\mu_0}}{2\pi} \frac{\sqrt{v}}{2} H(t) \delta_z \int_0^{2\pi} \cos \varphi \left[\cos m\varphi - i \sin m\varphi \right] d\varphi \int_0^R \rho d\rho \left[J_{m-1} v\rho + J_{m+1} v\rho \right] \end{aligned}$$

$$\begin{aligned} \cos x \cos y &= \frac{1}{2} \left[\cos x - y + \cos x + y \right] \\ \sin x \cos y &= \frac{1}{2} \left[\sin x - y + \sin x + y \right] \\ \sin x \sin y &= \frac{1}{2} \left[\cos x - y - \cos x + y \right] \end{aligned} \quad \begin{aligned} \delta_{m,n} &= \frac{1}{2\pi} \int_0^{2\pi} e^{i(m-n)} d\varphi \\ \delta_{m,-n} &= \frac{1}{2\pi} \int_0^{2\pi} e^{i(-m-n)} d\varphi \end{aligned}$$

$$\begin{aligned} \int_0^{2\pi} \left[\cos m\varphi - i \sin m\varphi \right] \sin \varphi d\varphi &= \frac{1}{2} \int_0^{2\pi} \sin \varphi - m\varphi d\varphi + \frac{1}{2} \int_0^{2\pi} \sin m\varphi + \varphi d\varphi - \\ &- \frac{i}{2} \int_0^{2\pi} \cos m\varphi - \varphi d\varphi + \frac{i}{2} \int_0^{2\pi} \cos m\varphi + \varphi d\varphi = -\frac{i}{2} \left(\int_0^{2\pi} \cos m\varphi - \varphi d\varphi - i \int_0^{2\pi} \sin m\varphi - \varphi d\varphi \right) + \\ &+ \frac{i}{2} \left(\int_0^{2\pi} \cos m\varphi + \varphi d\varphi - i \int_0^{2\pi} \sin m\varphi + \varphi d\varphi \right) = \frac{i}{2} \left(-\int_0^{2\pi} e^{-im-1}\varphi d\varphi + \int_0^{2\pi} e^{-im+1}\varphi d\varphi \right) = \\ &= i\pi \left(-\delta_{m,1} + \delta_{m,-1} \right) - i\pi\delta_{m,1} + i\pi\delta_{m,-1} \end{aligned}$$

$$\begin{aligned}
& \int_0^{2\pi} [\cos m\varphi - i \sin m\varphi] \cos \varphi d\varphi = \frac{1}{2} \int_0^{2\pi} \cos m\varphi - \varphi d\varphi + \frac{1}{2} \int_0^{2\pi} \cos m\varphi + \varphi d\varphi - \\
& - \frac{i}{2} \int_0^{2\pi} \sin m\varphi - \varphi d\varphi - \frac{i}{2} \int_0^{2\pi} \sin m\varphi + \varphi d\varphi = \frac{1}{2} \left[\int_0^{2\pi} \cos m\varphi - \varphi d\varphi - i \int_0^{2\pi} \sin m\varphi - \varphi d\varphi \right] + \\
& + \frac{1}{2} \left[\int_0^{2\pi} \cos m\varphi + \varphi d\varphi - i \int_0^{2\pi} \sin m\varphi + \varphi d\varphi \right] = \frac{1}{2} \left(\int_0^{2\pi} e^{-i m-1 \varphi} d\varphi + \int_0^{2\pi} e^{-i m+1 \varphi} d\varphi \right) = \pi \delta_{m,1} + \pi \delta_{m,-1}
\end{aligned}$$

$$\begin{aligned}
j_m(z, t; \nu) &= i \frac{\sqrt{\mu_0}}{2\pi} \frac{\sqrt{\nu}}{2} H(t) \delta z \left\{ \left[\delta_{m,-1} - \delta_{m,1} \right] \int_0^R \rho d\rho \left[J_{m-1} \nu \rho - J_{m+1} \nu \rho \right] - \right. \\
& \left. - \left[\delta_{m,-1} + \delta_{m,1} \right] \int_0^R \rho d\rho \left[J_{m-1} \nu \rho + J_{m+1} \nu \rho \right] \right\} = \\
&= i \frac{\sqrt{\mu_0}}{2} \frac{\sqrt{\nu}}{2} H(t) \delta z \left\{ \delta_{m,-1} \int_0^R \rho d\rho \left[J_2 \nu \rho - J_0 \nu \rho \right] - \delta_{m,1} \int_0^R \rho d\rho \left[J_0 \nu \rho - J_2 \nu \rho \right] - \right. \\
& \left. - \delta_{m,-1} \int_0^R \rho d\rho \left[J_2 \nu \rho + J_0 \nu \rho \right] - \delta_{m,1} \int_0^R \rho d\rho \left[J_0 \nu \rho + J_2 \nu \rho \right] \right\} = \\
&= i \frac{\sqrt{\mu_0}}{2} \frac{\sqrt{\nu}}{2} H(t) \delta z \left\{ \delta_{m,-1} \int_0^R \rho d\rho \left[J_2 \nu \rho - J_0 \nu \rho \right] - \delta_{m,1} \int_0^R \rho d\rho \left[J_0 \nu \rho - J_2 \nu \rho \right] - \right. \\
& \left. - \delta_{m,-1} \int_0^R \rho d\rho \left[J_2 \nu \rho + J_0 \nu \rho \right] - \delta_{m,1} \int_0^R \rho d\rho \left[J_0 \nu \rho + J_2 \nu \rho \right] \right\} = \\
&= -\sqrt{\mu_0} \frac{i}{2} \sqrt{\nu} H(t) \delta z \left\{ \left[\delta_{m,1} + \delta_{m,-1} \right] \int_0^R \rho d\rho J_0 \nu \rho \right\}
\end{aligned}$$

$$\begin{aligned}
& \int z J_0 z dx = z J_1 z + C \\
& \int_0^R \rho d\rho J_0 \nu \rho = \frac{1}{\nu^2} \int_0^R \nu \rho d\nu \rho J_0 \nu \rho = \frac{\nu R J_1 \nu R - \nu 0 J_1 \nu 0}{\nu^2} = \frac{R J_1 \nu R}{\nu}
\end{aligned}$$

$$\begin{aligned}
j_m(z, t; \nu) &= -\sqrt{\mu_0} \frac{i}{2} \sqrt{\nu} H(t) \delta z \left[\delta_{m,1} + \delta_{m,-1} \right] \int_0^R \rho d\rho J_0 \nu \rho = \\
&= -\sqrt{\mu_0} \frac{i}{2} \sqrt{\nu} H(t) \delta z \left[\delta_{m,1} + \delta_{m,-1} \right] \frac{R J_1 \nu R}{\nu} = -\sqrt{\mu_0} \frac{i R}{2 \sqrt{\nu}} H(t) J_1 \nu R \delta z \left[\delta_{m,1} + \delta_{m,-1} \right]
\end{aligned}$$

8.3 Продольные эволюционные коэффициенты

$$G(z', t', z, t) = \frac{c}{2} H(c t - t' - z - z') J_0 \nu \sqrt{c^2 t - t'^2 - z - z'^2}$$

$$h_m(z, t; \nu) = -\sqrt{\mu_0} \frac{icR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \int_0^{\frac{t-z}{c}} \delta(z') dz' \int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - (z-z')^2} =$$

$$= -\sqrt{\mu_0} \frac{icR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - (z-z')^2}$$

$$\int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - z^2} = \frac{1}{\nu c} \int_{\nu z}^{\nu ct} ds J_0 \sqrt{s^2 - \nu^2 z^2}$$

$$\int_{\nu z}^{\nu ct} ds e^{-i0s} J_0 \sqrt{s^2 - \nu^2 z^2} = 2iU_1 \left[-i\nu ct - z, \sqrt{\nu^2 c^2 t^2 - \nu^2 z^2} \right]$$

$$\int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - z^2} = \frac{2i}{\nu c} U_1 \left[-i\nu ct - z, \sqrt{\nu^2 c^2 t^2 - \nu^2 z^2} \right]$$

$$h_m(z, t; \nu) = -\sqrt{\mu_0} \frac{ic}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - (z-z')^2} =$$

$$= \sqrt{\mu_0} \frac{R}{2} \frac{\delta_{m,1} + \delta_{m,-1}}{\nu^{3/2}} J_1 \nu R U_1 \left[-i\nu ct - z, \sqrt{\nu^2 c^2 t^2 - \nu^2 z^2} \right] =$$

$$= -\sqrt{\mu_0} \frac{R}{2} \frac{\delta_{m,1} + \delta_{m,-1}}{\nu^{3/2}} J_1 \nu R U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right]$$

8.4 Поперечные электрические эволюционные коэффициенты

$$\frac{\partial}{\partial t} \int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - z^2} = J_0 \nu \sqrt{c^2 t^2 - z^2}$$

$$V_m^h = -\partial_{ct} \mu h_m \Big|_{\mu=const} = -\mu \frac{\partial h_m}{\partial t} =$$

$$= -\mu \frac{\partial}{\partial t} \left(-\sqrt{\mu_0} \frac{iR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \int_0^{\frac{t-z}{c}} dt' J_0 \nu \sqrt{c^2 (t-t')^2 - (z-z')^2} \right) =$$

$$= \mu \sqrt{\mu_0} \frac{iR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R J_0 \nu \sqrt{c^2 t^2 - z^2}$$

8.5 Поперечные магнитные эволюционные коэффициенты

$$I_m^h = \frac{1}{\mu} \frac{\partial h_m}{\partial z} = -\frac{\sqrt{\mu_0}}{\mu} \frac{R}{2} \frac{\delta_{m,1} + \delta_{m,-1}}{\nu^{3/2}} J_1 \nu R \frac{\partial}{\partial z} U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right]$$

$$\frac{\partial}{\partial z} U_1(W_{+,Z}) = -\frac{i\nu}{2} [U_0(W_{+,Z}) - U_2(W_{+,Z})]$$

$$I_m^h = \frac{1}{\mu} \frac{\partial h_m}{\partial z} = \frac{\sqrt{\mu_0}}{\mu} \frac{iR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \begin{bmatrix} U_0 & W_+, Z & -U_2 & W_+, Z \end{bmatrix}$$

8.6 Компоненты электрического поля

$$\nabla_{\perp} \Psi_m \times \vec{z}_0 = -e^{im\varphi} \left(\vec{\phi}_0 \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} - im \vec{\rho}_0 \frac{J_m \nu \rho}{\rho \sqrt{\nu}} \right)$$

$$\begin{aligned} \vec{E} &= -\sqrt{\varepsilon_0} \left\{ \sum_{m=-\infty}^{\infty} \int_0^{\infty} d\nu V_m^h \nabla_{\perp} \Psi_m \times \vec{z}_0 + \sum_{m=-\infty}^{\infty} \int_0^{\infty} d\chi V_n^e \nabla_{\perp} \Phi_n \right\} = \\ &= -\frac{1}{\sqrt{\varepsilon_0}} \sum_{m=-\infty}^{\infty} \int_0^{\infty} d\nu \mu \sqrt{\mu_0} \frac{iR}{4} \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R J_0 \nu \sqrt{c^2 t^2 - z^2} \times \\ &\times e^{im\varphi} \left(\vec{\phi}_0 \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} - im \vec{\rho}_0 \frac{J_m \nu \rho}{\rho \sqrt{\nu}} \right) = \\ &= -\mu \frac{iR}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \delta_{m,1} + \delta_{m,-1} \sum_{m=-\infty}^{\infty} e^{im\varphi} \int_0^{\infty} \frac{d\nu}{\sqrt{\nu}} J_1 \nu R J_0 \nu \sqrt{c^2 t^2 - z^2} \\ &\times \left(\vec{\phi}_0 \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} - im \vec{\rho}_0 \frac{J_m \nu \rho}{\rho \sqrt{\nu}} \right) \end{aligned}$$

$$\begin{aligned} E_{\varphi} &= -\mu \frac{iR}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \delta_{m,1} + \delta_{m,-1} \sum_{m=-\infty}^{\infty} e^{im\varphi} \int_0^{\infty} \frac{d\nu}{\sqrt{\nu}} J_1 \nu R J_0 \nu \sqrt{c^2 t^2 - z^2} \times \\ &\times \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} = -\mu \frac{iR}{8} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \left(e^{i\varphi} \int_0^{\infty} d\nu J_1 \nu R J_0 \nu \rho - J_2 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} - \right. \\ &\left. - e^{-i\varphi} \int_0^{\infty} d\nu J_1 \nu R J_0 \nu \rho - J_2 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} \right) = \\ &= \mu \frac{R}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{e^{i\varphi} - e^{-i\varphi}}{2i} \int_0^{\infty} d\nu J_1 \nu R J_0 \nu \rho - J_2 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} = \\ &= \mu \frac{R}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \sin \varphi \int_0^{\infty} d\nu J_1 \nu R J_0 \nu \rho - J_2 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} \end{aligned}$$

$$J_{m+1} z = \frac{2m}{z} J_m z - J_{m-1} z ; \quad J_2 \nu \rho = \frac{2}{\nu \rho} J_1 \nu \rho - J_0 \nu \rho$$

$$\begin{aligned} E_{\varphi} &= \mu \frac{R}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \sin \varphi \int_0^{\infty} d\nu J_1 \nu R \left(J_0 \nu \rho - \frac{2}{\nu \rho} J_1 \nu \rho + J_0 \nu \rho \right) J_0 \nu \sqrt{c^2 t^2 - z^2} = \\ &= \mu \frac{R}{2} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \sin \varphi \int_0^{\infty} d\nu J_1 \nu R \left(J_0 \nu \rho - \frac{J_1 \nu \rho}{\nu \rho} \right) J_0 \nu \sqrt{c^2 t^2 - z^2} = \mu \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{R \sin \varphi}{2} \left(I_2 - \frac{I_1}{\rho} \right) \end{aligned}$$

$$\begin{aligned}
E_\rho &= \mu \frac{iR}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \delta_{m,1} + \delta_{m,-1} \sum_{m=-\infty}^{\infty} e^{im\varphi} \int_0^\infty \frac{d\nu}{\sqrt{\nu}} J_1 \nu R J_0 \nu \sqrt{c^2 t^2 - z^2} im \frac{J_m \nu \rho}{\rho \sqrt{\nu}} = \\
&= -\mu \frac{R}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{\delta_{m,1} + \delta_{m,-1}}{\rho} \sum_{m=-\infty}^{\infty} m e^{im\varphi} \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_m \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} = \\
&= -\mu \frac{R}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{e^{i\varphi} + e^{-i\varphi}}{\rho} \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} = \\
&= -\frac{\mu}{4} \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{R \cos \varphi}{\rho} \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho J_0 \nu \sqrt{c^2 t^2 - z^2} = -\mu \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{R \cos \varphi}{2\rho} I_1
\end{aligned}$$

$$\vec{E}_{\rho, \phi, z, t} = \mu \frac{\sqrt{\mu_0}}{\sqrt{\varepsilon_0}} \frac{R}{2} \left[\vec{\phi}_0 \sin \varphi \left(I_2 - \frac{I_1}{\rho} \right) - \vec{\rho}_0 \cos \varphi \frac{I_1}{\rho} \right] \quad (2.14)$$

8.7 Компоненты магнитного поля

$$\vec{H} = \sqrt{2\mu_0} \left\{ \sum_{m=0}^{\infty} \int_0^\infty d\nu I_m^h \nabla_\perp \Psi_m + \sum_{m=1}^{\infty} \int_0^\infty d\nu \chi I_n^e \vec{z}_0 \times \nabla_\perp \Phi_n \right\}$$

$$\nabla_\perp \Psi_m = e^{im\varphi} \left(\vec{\rho}_0 \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} + \vec{\phi}_0 im \frac{J_m \nu \rho}{\sqrt{\nu} \rho} \right)$$

$$I_m^h = \frac{1}{\mu} \frac{\partial h_m}{\partial z} = \frac{\sqrt{\mu_0}}{\mu} \frac{iR}{4} \int_0^\infty d\nu \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix}$$

$$\begin{aligned}
\vec{H} &= \frac{1}{\sqrt{\mu_0}} \sum_{m=-\infty}^{\infty} \int_0^\infty d\nu I_m^h \nabla_\perp \Psi_m = \frac{1}{\mu} \frac{iR}{4} \int_0^\infty d\nu \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} \times \\
&\times e^{im\varphi} \left(\vec{\rho}_0 \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} + \vec{\phi}_0 im \frac{J_m \nu \rho}{\sqrt{\nu} \rho} \right)
\end{aligned}$$

$$\begin{aligned}
H_\varphi &= \frac{iR}{\mu} \frac{\delta_{m,1} + \delta_{m,-1}}{4} J_1 \nu R \int_0^\infty d\nu \frac{U_0 W_{+,Z} - U_2 W_{+,Z}}{\sqrt{\nu}} im \frac{J_m \nu \rho}{\sqrt{\nu} \rho} = \\
&= -\frac{R}{\mu} \frac{\delta_{m,1} + \delta_{m,-1}}{4} \int_0^\infty d\nu J_1 \nu R \frac{U_0 W_{+,Z} - U_2 W_{+,Z}}{\sqrt{\nu}} m \frac{J_m \nu \rho}{\sqrt{\nu} \rho} = \\
&= -\frac{R}{\mu} \frac{\delta_{m,1} + \delta_{m,-1}}{4\rho} e^{im\varphi} m \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_m \nu \rho \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} = \\
&= -\frac{R}{\mu} \frac{e^{i\varphi} + e^{-i\varphi}}{4\rho} \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} = \\
&= -\frac{R \cos \varphi}{\mu} \frac{1}{2\rho} \int_0^\infty \frac{d\nu}{\nu} J_1 \nu R J_1 \nu \rho \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix}
\end{aligned}$$

$$\begin{aligned}
H_\rho &= \frac{1}{\sqrt{\mu_0}} \sum_{m=-\infty}^{\infty} \int_0^\infty d\nu I_m^h \nabla_\perp \Psi_m = \frac{1}{\mu} \frac{iR}{4} \int_0^\infty d\nu \frac{\delta_{m,1} + \delta_{m,-1}}{\sqrt{\nu}} J_1 \nu R \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} \times \\
&\times e^{im\varphi} \sqrt{\nu} \frac{J_{m-1} \nu \rho - J_{m+1} \nu \rho}{2} = \frac{iR}{\mu} \frac{\delta_{m,1} + \delta_{m,-1}}{8} e^{im\varphi} \int_0^\infty d\nu J_1 \nu R \times \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \\ J_{m-1} \nu \rho & -J_{m+1} \nu \rho \end{bmatrix} = \\
&= \frac{iR}{\mu} \frac{e^{i\varphi} - e^{-i\varphi}}{8} \int_0^\infty d\nu J_1 \nu R \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} \begin{bmatrix} J_0 \nu \rho & -J_2 \nu \rho \end{bmatrix}
\end{aligned}$$

$$J_2 \nu \rho = \frac{2}{\nu \rho} J_1 \nu \rho - J_0 \nu \rho$$

$$H_\rho = -\frac{R \sin \varphi}{\mu} \frac{\varphi}{2} \int_0^\infty d\nu J_1 \nu R \begin{bmatrix} U_0 & W_{+,Z} & -U_2 & W_{+,Z} \end{bmatrix} \begin{bmatrix} J_0 \nu \rho & -\frac{J_1 \nu \rho}{\nu \rho} \end{bmatrix}$$

$$H_\rho = -\frac{R \sin \varphi}{\mu} \frac{\varphi}{2} \left\{ I_3 + \frac{I_4}{\rho} \right\}$$

$$\vec{H}_{\rho, \phi, z, t} = \vec{\varphi}_0 \frac{R \cos \varphi}{\mu} \frac{I_4}{2} \frac{1}{\rho} - \vec{\rho}_0 \frac{R \sin \varphi}{\mu} \frac{\varphi}{2} \left\{ I_3 + \frac{I_4}{\rho} \right\} = \frac{R}{2\mu} \left\{ \vec{\varphi}_0 \frac{I_4 \cos \varphi}{\rho} - \vec{\rho}_0 \sin \varphi \left(I_3 + \frac{I_4}{\rho} \right) \right\}$$

$$\vec{H}_{\rho, \phi, z, t} = \frac{R}{2\mu} \left\{ \vec{\varphi}_0 \frac{I_4 \cos \varphi}{\rho} - \vec{\rho}_0 \sin \varphi \left(I_3 + \frac{I_4}{\rho} \right) \right\} \quad (2.15)$$

$$h_m = -\sqrt{\mu_0} \frac{R}{2} \frac{\delta_{m,1} + \delta_{m,-1}}{\nu^{3/2}} J_1 \nu R U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right]$$

$$\Psi_m \nu = \frac{J_m \nu \rho}{\sqrt{\nu}} e^{im\varphi}$$

$$\begin{aligned}
H_z \rho, \phi, z, t &= \sqrt[2]{\mu_0} \sum_{m=0}^{\infty} \int_0^\infty \nu^2 d\nu h_m(z, t; \nu) \Psi_m \rho, \phi; \nu = \\
&= -\frac{R}{2} \int_0^\infty \nu^2 \frac{\delta_{m,1} + \delta_{m,-1}}{\nu^{3/2}} J_1 \nu R \frac{J_m \nu \rho}{\sqrt{\nu}} e^{im\varphi} U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right] d\nu = \\
&= -\frac{R}{2} \frac{\delta_{m,1} + \delta_{m,-1}}{e^{im\varphi}} \int_0^\infty d\nu J_1 \nu R J_m \nu \rho U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right] = \\
&= -iR \frac{e^{i\varphi} - e^{-i\varphi}}{2i} \int_0^\infty d\nu J_1 \nu R J_1 \nu \rho U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right] = \\
&= -iR \sin \varphi \int_0^\infty d\nu J_1 \nu R J_1 \nu \rho U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right]
\end{aligned}$$

$$H_z \rho, \phi, z, t = -iR \sin \varphi \int_0^\infty d\nu J_1 \nu R J_1 \nu \rho U_1 \left[i\nu ct - z, \nu \sqrt{c^2 t^2 - z^2} \right] \quad (2.16)$$

8.8 Декартовы координаты

9 Нелинейная поправка для задачи плоского диска

$$\vec{\mathbf{J}} = \frac{\partial}{\partial t} \vec{\mathbf{P}}' \vec{\mathbf{E}} + \vec{\mathbf{J}}^\sigma \vec{\mathbf{E}}, \vec{\mathbf{H}} + \vec{\mathbf{J}}^e; \quad \vec{\mathbf{I}} = \frac{\partial}{\partial t} \vec{\mathbf{M}}' \vec{\mathbf{H}} + \vec{\mathbf{J}}^h;$$

$$\vec{\mathbf{J}} = \frac{\partial}{\partial t} \vec{\mathbf{P}}' \vec{\mathbf{E}} + \vec{\mathbf{J}}^\sigma \vec{\mathbf{E}}, \vec{\mathbf{H}} + \vec{\mathbf{J}}^e = \frac{\partial}{\partial t} \vec{\mathbf{P}}' \vec{\mathbf{E}}$$

$$\vec{\mathbf{J}}' \vec{\mathbf{E}}, \vec{\mathbf{H}} = \begin{pmatrix} J'_\rho \\ J'_\varphi \\ J'_z \end{pmatrix} = \begin{pmatrix} \vec{J}' \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial t} \vec{P}' \vec{\mathbf{E}} \\ \frac{\partial}{\partial t} P'_z \vec{\mathbf{E}} \end{pmatrix}$$

$$P'_z \vec{\mathbf{E}} = ?$$

$$\vec{\mathbf{J}}^1 = \frac{\partial}{\partial t} \chi_3 \vec{\mathbf{E}}^3 + \sigma \vec{E} = \chi_3^E \vec{x}_0 \frac{\partial}{\partial t} E_x^3 + \sigma \vec{x}_0 E_x$$