



Evaluating the performance of microstructure generation algorithms for 2-d foam-like representative volume elements



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ABSTRACT

This investigation evaluates various numerical algorithms; each designed to generate periodic 2-D Representative Volume Elements (RVEs) containing foam-like microstructures suitable for direct import into commercial finite element software for mechanical evaluation. The operation of each algorithm is discussed and the resulting RVEs are examined from both a mechanical and a morphological perspective. A basic Voronoi-based algorithm is found to be simple to implement but the method is shown to produce inherently anisotropic microstructures. Increasing the degree of irregularity of the microstructure reduces the anisotropy but at the cost of creating unrealistic microstructures, containing highly angular cells. A method of modifying such unrealistic microstructures using a centroidal tessellation relaxation algorithm is demonstrated, ultimately producing RVEs with relatively realistic mono-disperse microstructures. An alternative algorithm is also investigated the advantage of this algorithm is its ability to generate poly-disperse microstructures, with a controllable degree of poly-dispersity and an almost fully isotropic mechanical response.

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1. Introduction

The microstructure of cellular foam is very important in determining its bulk mechanical response. The wide variety of cellular foam reflects the large diversity in the types of microstructure that can be found. Important parameters that can be used to describe the microstructure include the relative density (Gibson and Ashby, 1997), cellular area fraction, cellular morphology (open or closed) (Raj and Kerr, 2010), poly-dispersity of cell size (Kraynik et al. 2004; Jang et al. 2008; Jang et al. 2010; Kraynik, 2006; Mills, 2007) and statistical anisotropy of cellular dimensions (Jebur, 2013; Benouali et al. 2005). In some cases, macro-scale models able to accurately model a foam's mechanical response might not yet be available due to the complexity of the material behaviour (e.g. hyperelastic transversely isotropic compressible continuum models) and if they are, it can often be a time consuming, difficult and expensive process to perform all the lab-based experiments required to determine the model parameters. This traditional macro-scale approach also provides little information on the underlying deformation mechanisms occurring within the material's microstructure, and consequently provides limited scope for microstructural

optimisation. With the increase in computational power, multi-scale modelling is becoming a powerful and increasingly relevant computational approach (Liebscher et al. 2013; Verhoosel et al. 2015), useful in optimising a material's underlying microstructure for specific applications. For example, changing relative density and microstructural anisotropy as a function of position, can optimise the structural response of a foam-based component according to anticipated loading conditions. Such microstructural optimisation is commonplace in many naturally occurring structures such as bone and wood. Nevertheless, micro to macro simulation of entire parts and components is computationally expensive, usually prohibitively so, leading many modellers to employ the concept of a Representative Volume Element (RVE) (Hill, 1963) in combination with a Periodic Boundary Condition (PBC) (Guedes and Kikuchi, 1990; Anthoine, 1995). The aim here is to obtain a homogenised macro-response for a material's bulk behaviour from a small sample of the material.

The accuracy and practicality of computational homogenisation or any other micro-to-macro mapping technique, strongly depends on the choice of RVE. The latter should ideally be both realistic and computationally efficient, two criteria that are often at odds with one another (Kouznetsova et al. 2001; Smit et al. 1998; Swan, 1994). The optimum size of the RVE and its level of detail are two important considerations (Alsayednoor et al. 2013).

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When it comes to the level of detail within the RVE, one extreme strategy is to employ RVEs based on, for example, actual 3-D topologies, measured using techniques such as micro-CT imaging (Shan and Gokhale, 2001; Maire et al. 2003; Michailidis et al. 2011). In this case, issues associated with strict application of the PBC have to be resolved (Youssef et al. 2005) together with problems related to identification of the distinct phases within the material. An alternative is to use numerical algorithms to automatically generate RVEs with periodic microstructures. Particulate composites, or porous materials with high relative density can be modelled in this way (Guo et al. 2014). Nevertheless, long computation times resulting from use of the large number of continuum elements required in these simulations, remains a practical restriction to wider exploitation.

Modelling foams of low relative density, with microstructures containing narrow ribs or faces offers the opportunity to use structural elements, such as beam or shell, rather than continuum elements (Zhu et al. 2000; Körner et al. 2002; Roberts and Garboczi, 2002; Kraynik et al. 2003; Chen et al. 1999; Fazekas et al. 2002; Schmidt, 2004).

Despite the large number of prior numerical investigations that have employed digitally generated microstructures to model foams of low relative density, concerns regarding the realism and accuracy of these morphologies exist. Namely, can the numerical algorithms produce microstructures that accurately represent the diverse range of real morphologies observed in actual low-density cellular foams? A recognised method of generating representative randomised microstructures for cellular foams is through Voronoi tessellation (Voronoi, 1908) or alternatively through use of variants on this technique. Basic Voronoi tessellation (Aurenhammer, 1991; Zhu and Windle, 2002; Zhu et al. 2006) provides a fast and effective way to create beam-based finite element microstructures possessing morphologies that serve as a rough, first order approximation to those observed in several classes of polymeric and metallic, open and closed cell foams (Silva et al. 1995; Van der Burg et al. 1997; Shulmeister et al. 1998). However, changes to the basic Voronoi algorithm can substantially alter the predicted microstructural morphologies. Altering the fundamental operation of the basic Voronoi algorithm, using for example the Laguerre-Voronoi approach, can produce even greater statistical changes (Fazekas et al., 2002; Hardenacke and Hohe, 2009). The main goal of the current investigation is to evaluate the realism of microstructures generated using these different techniques by recreating isotropic microstructural cellular morphologies. Note that in this investigation, the constitutive behaviour of the beam elements is the same throughout all simulations and is restricted to a simple linear-elastic material model in order to focus attention on the effects of the RVE microstructure rather than the material behaviour. In so doing, the investigation aims to highlight the relative advantages and disadvantages of the different microstructure generation algorithms.

2. Description of numerical RVE generation algorithms

2.1. Basic Voronoi Tessellation

Voronoi tessellation is an algorithm designed to partition a given Euclidian surface into distinct regions. To do this individual seed points are positioned across the surface and the Voronoi algorithm is used to partition the surface into individual cells, separating each seed point from its neighbour with a cell wall positioned half way between the two seed points (Aurenhammer et al. 2000). Various methods have been implemented to generate foam-like RVEs based on Voronoi tessellation. The methods tend to differ mainly in how the position of seed-points required by the Voronoi tessellation algorithm, are generated (e.g. Van der Burg et al. 1997;

Grenestedt and Tanaka 1998; Silva and Gibson, 1997; Zhu et al. 2001b). In this investigation, a combination of methods proposed by Zhu et al. (2001b) and Grenestedt and Tanaka (1998) is utilised (Alsayednoor et al. 2013). This combined technique has been found to be computationally faster than the original method proposed by Zhu et al. (2001b), upon which it is based, nevertheless the resulting microstructures are very similar. In order to generate a periodic 2-D RVE using this method, first a unit square is populated with regular seed positions (the precursor seeding for a honeycomb pattern). Eq. (1) is then used to perturb these initial seed positions, by a distance, Δ , in an arbitrary direction, to produce a unit square containing an irregular distribution of seeds. Here Δ is given as,

$$\Delta = \frac{\alpha \times \delta_0}{100}, \alpha \in (0, 100) \quad (1)$$

where

$$\delta_0 = \sqrt{\frac{2A}{N\sqrt{3}}} \quad (2)$$

In Eqs. (1) and (2), the parameter, α , controls the degree of irregularity of the microstructure, similar to the same parameter used by Zhu et al. (2001a), δ_0 is the minimum specified distance between each adjacent seed position, which is calculated from Eq. (2). A and N represent the area of the RVE and the number of seeds, respectively. ' α ' can be set to values up to 100, though in practice setting $\alpha=50$ is found to produce an isotropic RVE (the main motivation for randomising the seed positions) and using higher values merely increases the angularity of the cells within the RVE, resulting in progressively less realistic microstructural morphologies. The 'perturbed' unit square seeding is copied nine times to create a larger square containing 3×3 identical perturbed unit square seedings. Applying a Voronoi tessellation algorithm to the resulting 3×3 structure produces a fully periodic cellular pattern within the central square, with pairs of counterpart nodes on opposing boundaries (see Fig. 1).

Later in this investigation it will be shown that the basic Voronoi method described here can introduce unwanted anisotropy in the microstructural morphology; anisotropy that is evident in the RVE's mechanical response. The same anisotropic behaviour was also observed in RVEs generated using the method proposed by Zhu et al. (2001a), particularly for RVEs with only low to moderate degrees of irregularity.

2.2. Centroidal Voronoi Tessellation

The Centroidal Voronoi Tessellation (CVT) algorithm can be described as an iterative 'relaxation process'. As will be shown, the CVT method is a convenient approach to removing unintended anisotropy from RVEs generated using the basic Voronoi algorithm, producing mono-disperse digital microstructures closely resembling those of various real polymer foams. The technique runs in a reasonable computational time (less than 30 seconds when using a 64 bit Windows OS with Intel (R) Xeon (R) CPU @ 2.66 GHz and 12 GB RAM) and can be used to construct both 2-D and 3-D periodic RVEs with a highly mono-disperse cellular size distribution. This technique is a modification of the basic Voronoi method (Burkardt et al. 2002). Liu et al. 2009 considered the Lloyd's algorithm (Lloyd, 1982) to be the best technique for computing a CVT. The algorithm can be described as an iterative 'relaxation process' where a basic Voronoi cell is generated in each iteration, though the seed position of the cell is updated at each iteration and moved towards the position of the Voronoi cell's centroid (hence the name of the algorithm), until the two eventually converge, resulting in more regular polygon containing less acute internal angles (see, for example, Fig. 2).

An example of the effect of the algorithm on the final microstructure is demonstrated in Fig. 3. The figure shows the

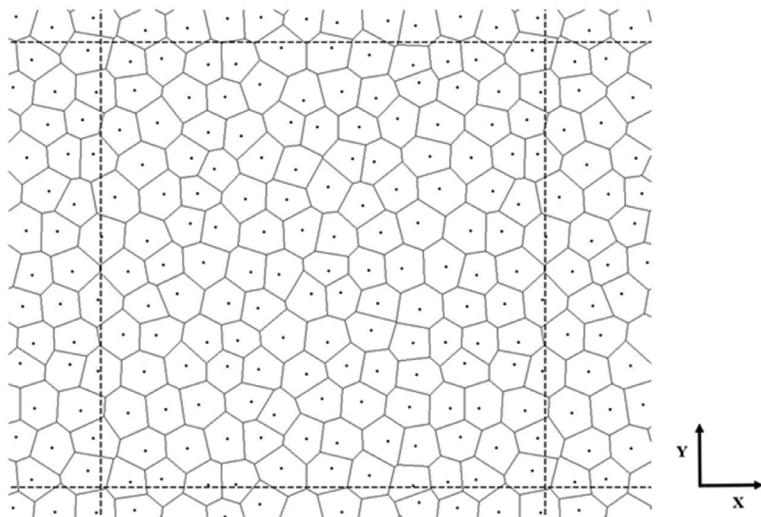


Fig. 1. An example of fully periodic classical Voronoi tessellation (central square).

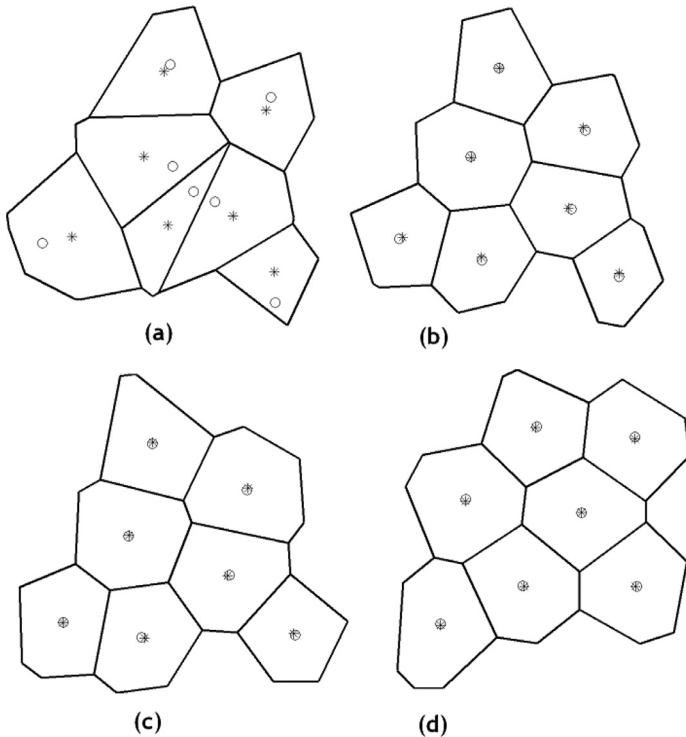


Fig. 2. Illustrative example of Lloyd's algorithm. At each iteration (a, b, c, d) a centroid for each Voronoi cell is calculated (black stars) and replaced by a seed point (circles).

microstructure before iterative application (about 500 iterations for 500 cells) of the Lloyds relaxation algorithm. The resulting cells (**Fig. 3b**) are of more realistic geometry and of lower polydispersity compared to those generated using the basic Voronoi algorithm when employing a fully randomised seeding configuration (**Fig. 3a**).

2.3. Laguerre-Voronoi Tessellation

One important limitation of both the basic and Centroidal Voronoi Tessellation algorithms, is their inability to model polydisperse microstructures. In general, all seed-based modelling techniques lead to symmetric Gaussian cell-size distributions ([Fazekas et al. 2002; Gervois et al. 2002; Kanaun and Tkachenko, 2006](#)).

In reality, the microstructural cell-size distribution of a polydisperse foam tends more towards a log-normal or Gamma distribution ([Kraynik, 2006; Montminy et al. 2004](#)); significantly different in form to a Gaussian distribution. To address this issue, instead of constructing Voronoi cells based on their seed point coordinates (i.e. the basic Voronoi tessellation method), each individual cell can instead be represented using a 'hard-disk' with a specified size and coordinate. Using such an approach means that cell-partitioning can be based on the location of the contact points between the hard-disks rather than the mid-way bi-sector point lying between seed positions (as is used in the basic Voronoi approach). In the Laguerre-Voronoi approach, bisectors are located at the contact point of adjacent disks rather than at the mid-way distance between their centres. The difference in the resulting microstructure is illustrated in **Fig. 4**.

In general, CVT are suitable for the generation of isotropic 2-D and 3-D periodic RVEs with different degrees of irregularity, for cells with low to moderate polydispersity. However, for cases of large polydispersity, it will be shown that the Laguerre-Voronoi method is more effective. One of the main challenges of the Laguerre-Voronoi method is in creating the initial packing arrangement of hard-discs (i.e. the position and size of the non-overlapping discs shown in **Fig. 4**). Generally, two main methods are used: (i) Random Sequential Adsorption (RSA) ([Cooper, 1988](#)) and (ii) classical hard-disk Molecular Dynamics (MD). In the RSA method, hard-disks are randomly generated with different sizes and coordinates. Random sized discs are then randomly, sequentially and irreversibly placed in the RVE, though if any overlap is detected then the disc placement is cancelled. The advantages of using RSA are its simplicity and relatively small computational resource requirement, an advantage that becomes more apparent for polydisperse cell size distributions. Different techniques for applying the RSA method are reported in the literature; most techniques are suitable for single mode and bimodal cell size generation ([Harder and Silbert, 1980; Adamczyk et al. 1997; Brilliantov et al. 1998; Rouault and Assouline, 1998; Gray et al. 2001; Richard et al. 2001; Wu et al. 2003; Gan et al. 2010](#)). [Fan et al. \(2004\)](#) applied a modified version of a collective rearrangement RSA algorithm, initially developed by [He and Ekere \(1998\)](#), to generate a polydisperse polycrystalline metallic grain microstructure. In another attempt, [Farr and Groot \(2009\)](#) developed a new method for packing 3-D polydisperse spheres based on a one-dimensional mapping. Although all these approaches are applicable to both 2-D

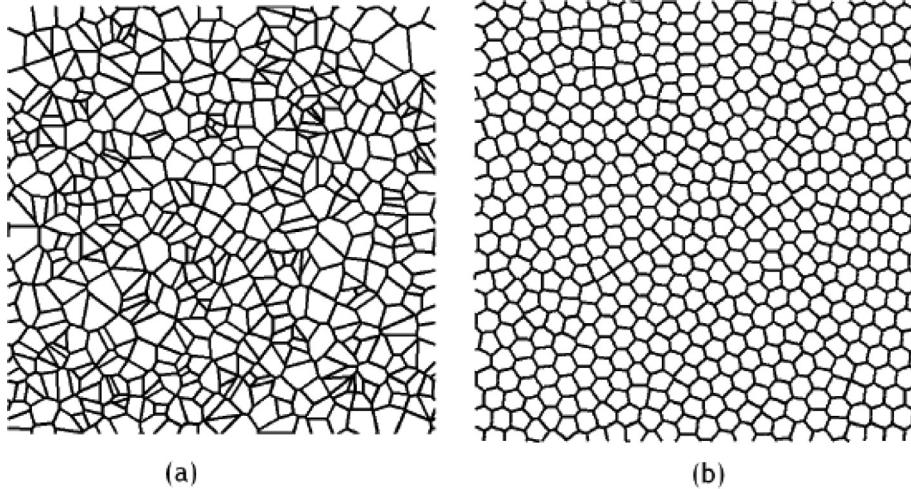


Fig. 3. (a) Microstructure generated by basic Voronoi technique (with $\alpha = 50$) before application of the relaxation algorithm and (b) the same seeding arrangement but after application of the Lloyd's algorithm.

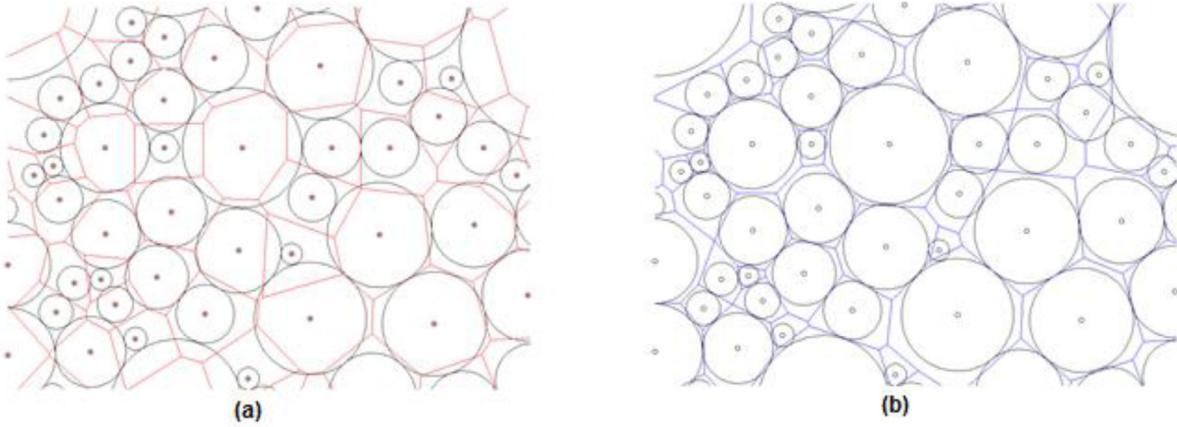


Fig. 4. (a) Voronoi tessellation generated from seed positions produced by an arrangement of hard-discs (b) Laguerre Voronoi Tessellation generated using the same arrangement of hard-disks. Note the latter results in a much more polydisperse microstructure.

and 3-D cases, they are best suited to single or bimodal cell size distributions, a higher degree of polydispersity tends to be more complicated, time consuming and requires alternative algorithms. After studying most of the RSA techniques published in the literature, the Drop and Roll method was chosen for this investigation (Visscher and Bolsterli, 1972; Okubo and Odagaki, 2004). Important advantages of the Drop and Roll method, compared to other packing algorithms, include increased contact between neighbouring disks, which results in a higher packing density, greater simplicity and a more efficient computational algorithm (Gibson, 2007). The implementation process for generating fully periodic and relatively isotropic RVEs using a novel, modified version of the Drop and Roll method is outlined in the supplementary material.

As with the basic Voronoi technique, described in Section 2.1, when using the Drop and Roll method, a control mechanism is implemented to define the geometrical degree of irregularity of the resulting RVE. The approach used here was originally suggested by Kraynik et al. (2004) and involves generating the disk size (radius) distribution using a log-normal function, i.e.

$$\mu = \ln\left(\frac{m}{\sqrt{1 + \frac{\gamma}{m^2}}}\right) \quad (3)$$

$$\sigma_{sd} = \sqrt{\ln\left(1 + \frac{\gamma}{m^2}\right)} \quad (4)$$

$$r = \exp(\mu + z\sigma_{sd}) \quad (5)$$

where m is the disk mean radius, γ is the variance of the disk radii, μ and σ_{sd} are the mean and standard deviation of the disk radii required to generate random data following a lognormal distribution and z is a pseudo-random variable. γ is the main parameter controlling the degree of polydispersity since it defines the range of disk sizes and plays a similar role to α in the basic Voronoi method, described in Section 2.1. The 'lognrnd.m' function available in MatlabTM, see Eq (5), directly calculates random values based on a log-normal frequency distribution function. Using Eqs. (3) and (4), with disk mean radius and variance inputs (m and γ), mean and standard deviations (μ and σ_{sd}) are calculated and used in Eq. (5). In order to avoid generation of very large or very small disks, a limit is set within a specified interval by choosing r_{min} and r_{max} for the minimum and maximum disk radii. The disk size-limits depend on the desired final target microstructure.

2.4. RVE boundary condition

The concepts of a Representative Volume Element (RVE) and a Periodic Boundary Condition (PBC) are intrinsically linked. A two dimensional PBC is required: (i) to constrain the motion of counterpart nodes on each pair of RVE boundaries to each other and (ii) to ensure stress continuity across boundaries. To fulfil these conditions,

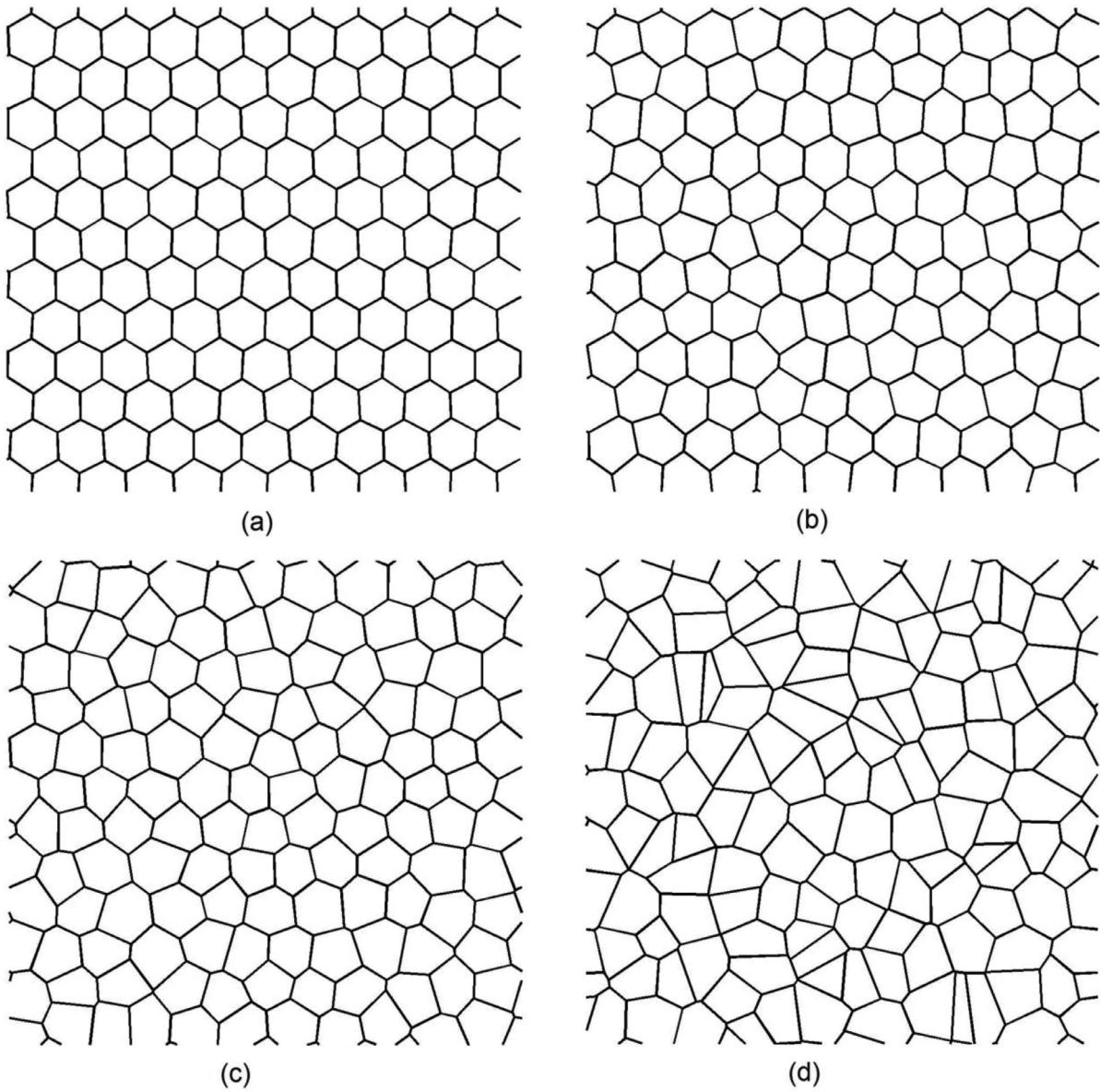


Fig. 5. RVEs with 6 different degree of irregularity (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$.

the following constraints are applied on the boundaries of the RVE using the 'EQUATION' keyword option available in Abaqus™:

$$U_i^R = U_i^L + U_i^{di}, U_i^T = U_i^B + U_i^{di}, i = 1 \text{ to } 2 \quad (6)$$

$$U_i^R = U_i^L, U_i^T = U_i^B, i = 3 \quad (7)$$

where the superscripts *L*, *R*, *T* and *B* indicate the left, right, top and bottom boundaries, respectively. *U* is the nodal displacement in the *i*-th direction, the subscript, *i*, represents the degree of freedom with *i*=1, 2 indicating displacements in the X (transverse) and Y (perpendicular) directions and *i*=3 indicating rotation about an axis lying perpendicular to the X-Y plane. Eqs. (6) and (7) automatically lead to perpendicular RVE compression. Hence, to apply compressive loading in any other arbitrary direction the plane of reference must be rotated accordingly (Schmidt, 2004).

3. Microstructure generation and Finite Element method

3.1. Finite element computation

Throughout this investigation a constant relative density of $\rho_R = 0.05$ is used. Hence, a constant element cross sectional area is maintained in all the finite element models of this investigation. The same procedure used in both Alsayednoor et al. (2013) and Zhu et al. (2001a) is applied to calculate the beam thickness, *t*, i.e.

$$t = \frac{\rho_R A}{\sum_{k=1}^n l_k} \quad (8)$$

here *n* is the number of elements, *A* is the area of the RVE and *l* is the length of each individual element. For all cases, the material model is linear elastic with a Young's modulus, E_S , of 1 GPa, a material density, ρ_m , of 2000 kg m⁻³ and a Poisson's ratio of 0.33 consequently it should be noted that all conclusions from this work are restricted to linear elastic behaviour of the constituent material. This is clearly a significant simplification compared to

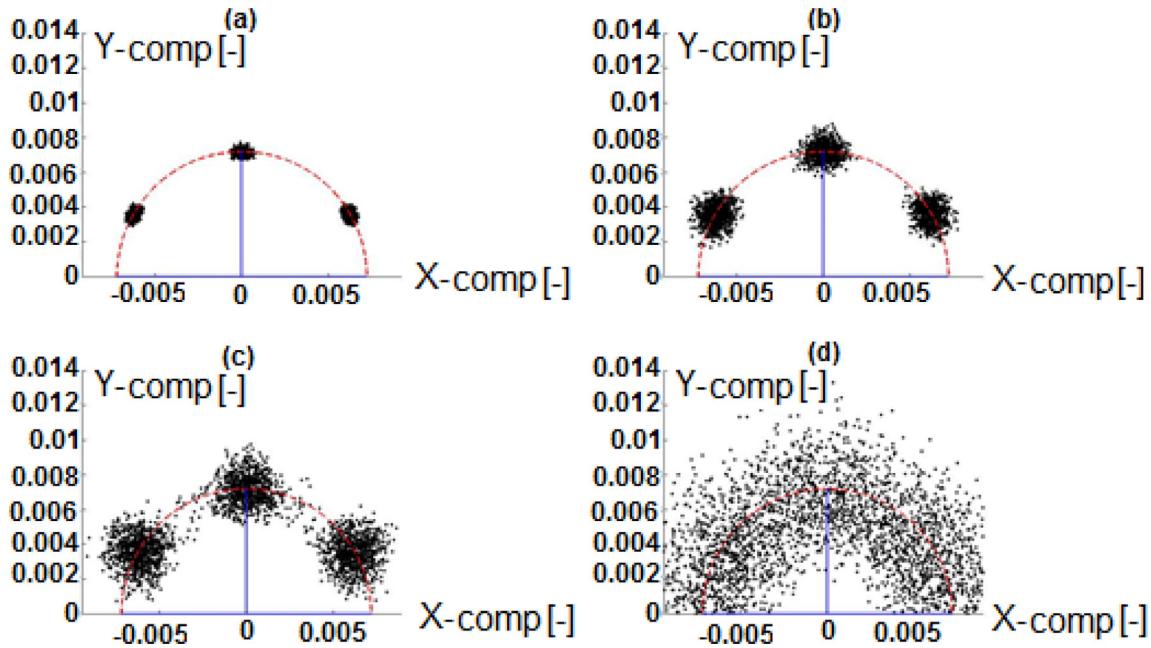


Fig. 6. Polar representation of cellular orientation. The X and Y axis represent the length of the normalised components of the semi-major axis of an ellipse, fitted within each individual cell of the RVE, multiplied by the cell's area (black dots). The results are for microstructures generated using a basic Voronoi tessellation algorithm. (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$. The radius of the red circle represents an average value of the data for a given RVE.

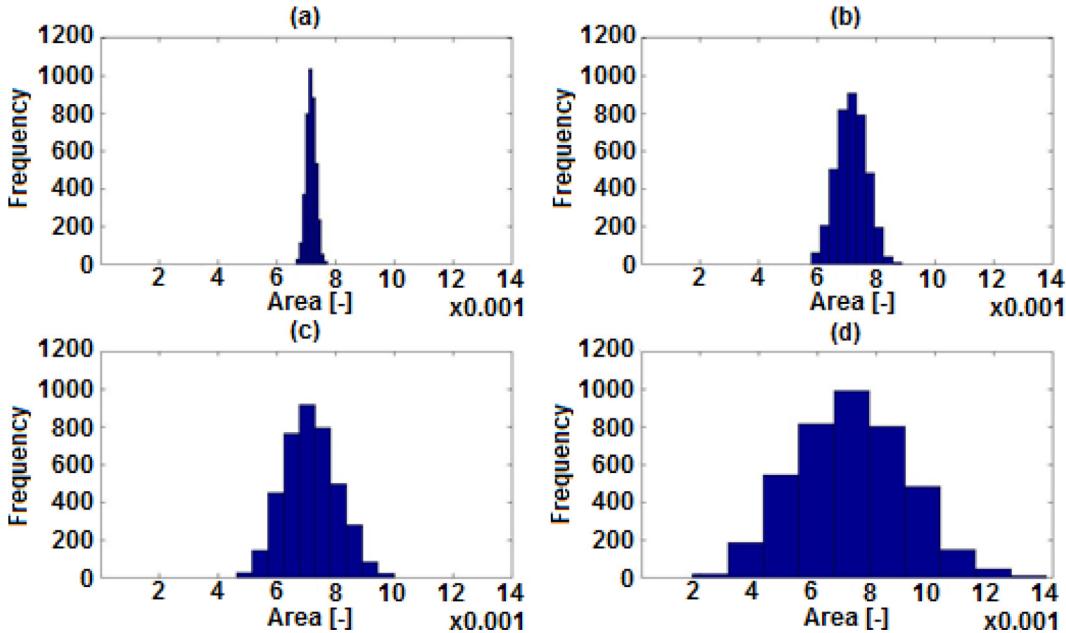


Fig. 7. histogram of cellular area for undeformed state of microstructure generated from Basic Voronoi Tessellation. (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$.

real foam materials undergoing large deformations, but one which allows the investigation to focus on the influence of microstructural morphology rather than material behaviour. In order to eliminate the effect of unrealistically short elements near the border of the RVE, resulting from the cropping operation, (see Fig. 4), elements with thickness-to-length ratio above 1/8 were deleted prior to using the RVE in the FE simulations. Periodicity between the two counterpart boundaries was preserved, by considering each small element as a point and adjusting counterpart boundary elements accordingly. A mesh sensitivity study was performed, varying the number of elements per strut from 1 to 10. Good convergence was found using 5 elements and so this number was adopted for all subsequent simulations reported in this investigation.

Zhu and Windle, 2002 also used 5 elements per strut in their investigation). Timoshenko beam elements (B22 in the Abaqus library) were employed and the Abaqus™ Standard implicit solver was used. Contact between elements was not modelled as the maximum compressive strain applied to the RVEs was no more than 10%; results from Alsayednoor et al. (2013) suggested contact between elements only becomes important at compressive strains of around 30%.

3.2. Evaluation of RVEs

When creating RVEs using numerical algorithms, it is clearly desirable to generate realistic microstructural morphologies. But what

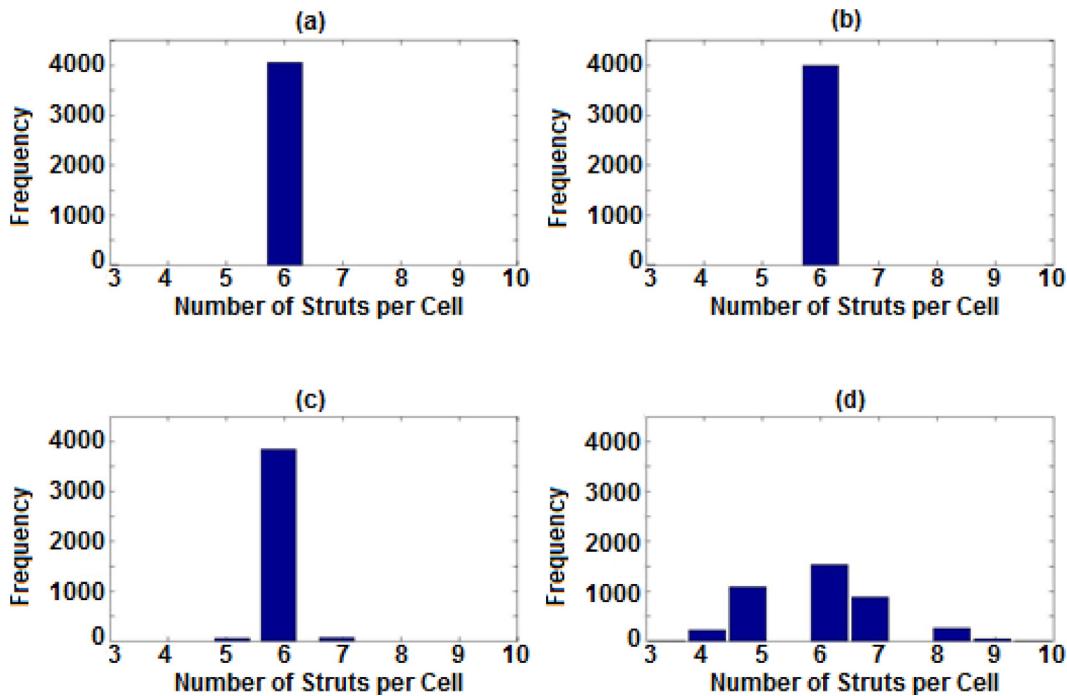


Fig. 8. histograms for number of struts within each cell for undeformed state of microstructure for different degrees of irregularity: (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$.

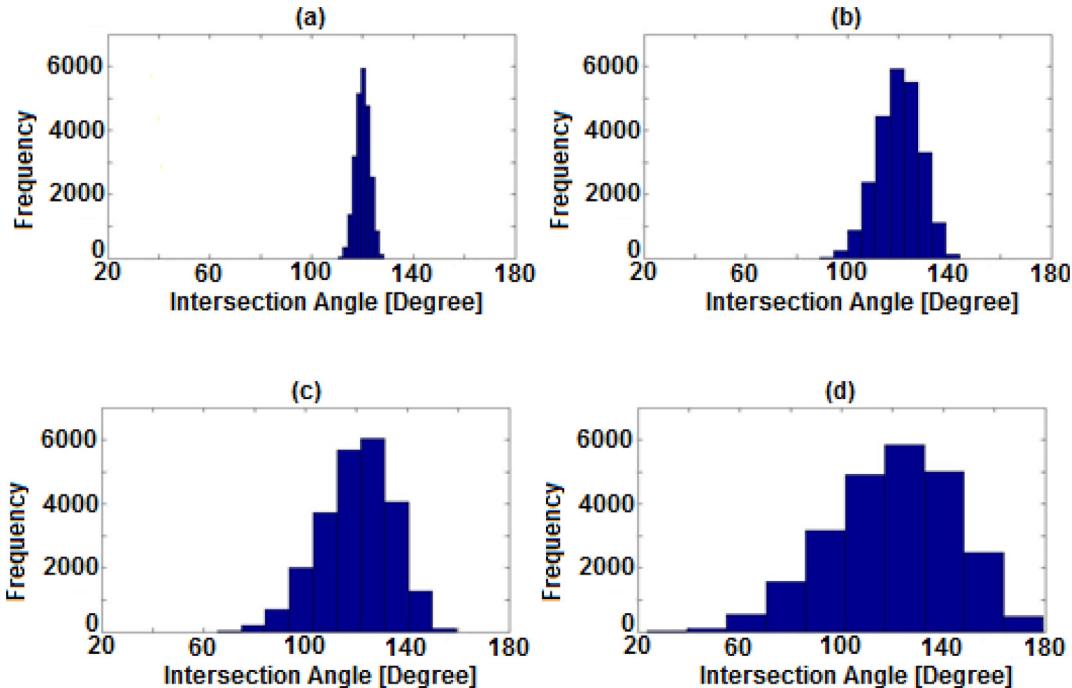


Fig. 9. histograms for strut intersection angle within each cell for undeformed microstructures using different degrees of irregularity: (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$.

constitutes a 'realistic microstructure'? Certain realistic microstructural properties are relatively straightforward to anticipate. For example, a foam showing an isotropic mechanical response is generally expected to have no statistical directional dependence within its microstructure, i.e. there should be no preferred direction of cell orientation (Benouali et al. 2005; Jebur et al. 2012). Likewise, polydispersity of real foams (or the lack of it) is evident when examining real micro-CT data (Montminy et al. 2004; Benouali et al. 2005) and is seen to vary from one type of foam to the next.

It is therefore desirable to develop numerical algorithms that can control the degree of polydispersity in the digital microstructure. Other measures of realism relating to a foam's microstructure are less obvious. Referring to Fig. 3, comparison with micro-CT images (Jebur, 2013) suggests that Fig. 3b is more realistic than Fig. 3a. Cells in Fig. 3a are rather angular; a morphology unlikely to occur in a real gas-expanded foam due to minimisation of the energy associated with free surfaces (related to the second law of thermodynamics) (Jang et al. 2010). Various statistical measures can be used

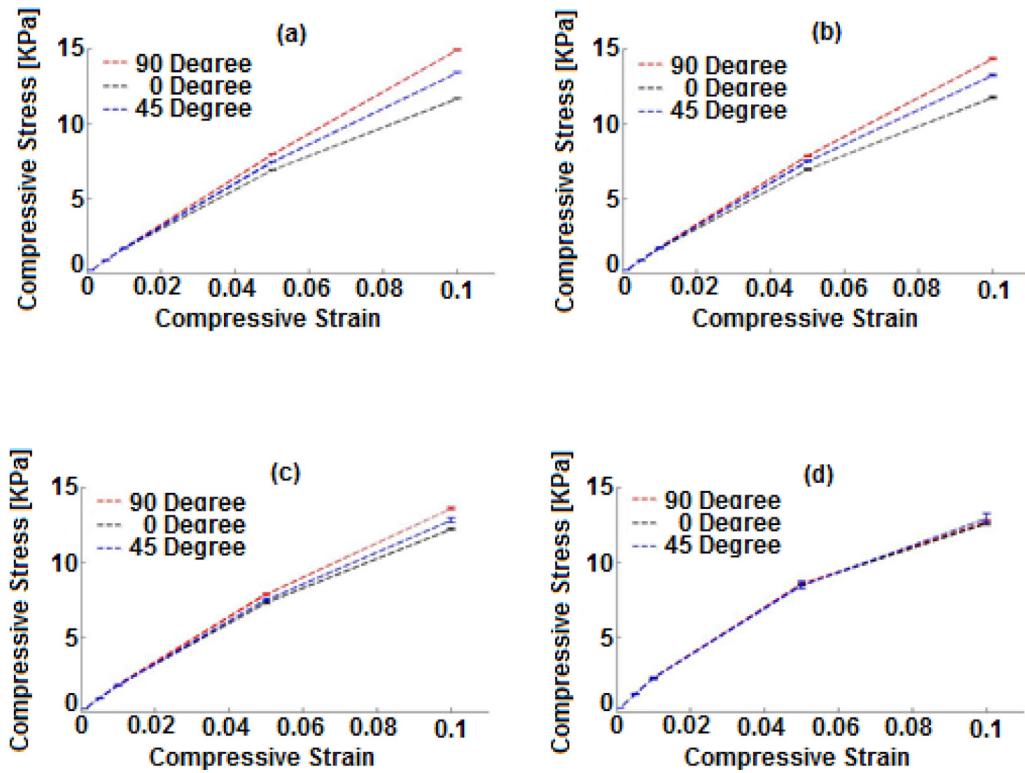


Fig. 10. Averaged stress-strain curves for compressed RVEs generated fusing basic Voronoi Tessellation and with degree of irregularity: (a) $\alpha = 5$, (b) $\alpha = 15$, (c) $\alpha = 25$, (d) $\alpha = 50$.

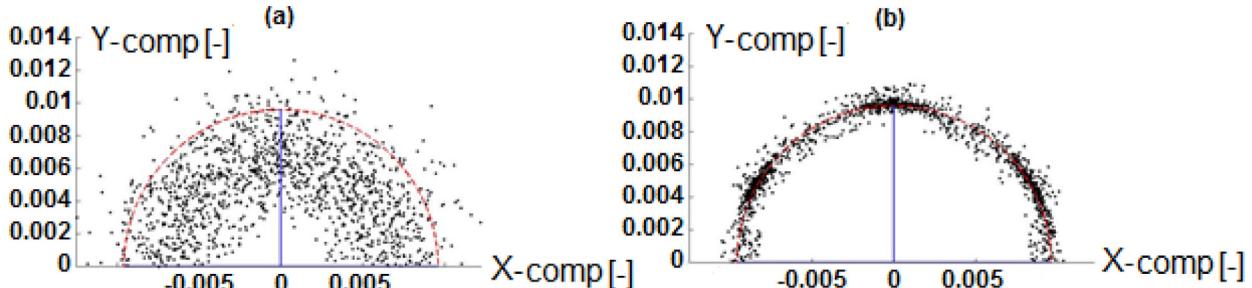


Fig. 11. Polar representation of cellular orientation. The X and Y axis represent the length of the normalised components of the semi-major axis of an ellipse, fitted within each individual cell of the RVE, multiplied by the cell's area (black dots). The results are for microstructures generated using: (a) before Lloyd's relaxation, $\alpha = 50$ (b) after Lloyd's relaxation. The radius of the red circle represents an average value of the data for a given RVE.

to analyse the realism of the cell morphology, including the number of struts per cell and the strut intersection angles (Montminy et al. 2004; Kraynik, 2006). In this section, the goal is to scrutinise each aspect of the digital microstructures and evaluate the performance of each of the three numerical methods discussed in Section 2. In each case, the mechanical response of the RVEs is examined and results of the morphological and mechanical analyses are compared. A simple and effective way to examine the mechanical anisotropy of the various RVEs is to apply uniaxial compressive loads along different directions with respect to the orientation of the RVE. Note that the response to alternative loading cases is also of interest, and indeed would be necessary if the RVEs were to be used to accurately determine the parameters of macro-scale constitutive models (e.g. hyperelastic models). Nevertheless, such loading cases are beyond the requirements of the current investigation and so are deferred to future work.

3.2.1. Evaluation of microstructure generated using the basic Voronoi Tessellation

To study the possible relationship between the degree of irregularity and morphological / mechanical anisotropy of microstruc-

tures generated using basic Voronoi tessellation (see Section 2.1), four different values of α have been employed ($\alpha = 5, 15, 25$ and 50). To obtain useful representative results using an RVE-based study, one must use either an RVE containing enough information (i.e. a large enough number of cells), or alternatively perform enough repeat simulations using smaller RVEs, to provide reliable representative results. Alsayednoor et al. (2013) demonstrated that the second strategy is the more computationally efficient approach and so for each degree of irregularity, 20 random RVEs with periodic boundary and boundary conditions, each containing 550 cells have been generated. Examples are shown in Fig. 5.

Algorithms were written to automatically generate statistical data from all 20 RVEs used for each case. Techniques used by these algorithms are described in (Alsayednoor, 2013). Typical statistical data relating to cell orientation and size, generated from RVEs similar to the examples shown in Fig. 5 are provided in Figs. 6–9. A variety of statistical measures can be analysed (Alsayednoor, 2013; Montminy et al. 2004; Kraynik, 2006), for brevity, this investigation is limited to cell orientation, normalised size distribution, the number of struts per cell and the strut internal intersection angle. All RVE simulations in this investigation have a side length of

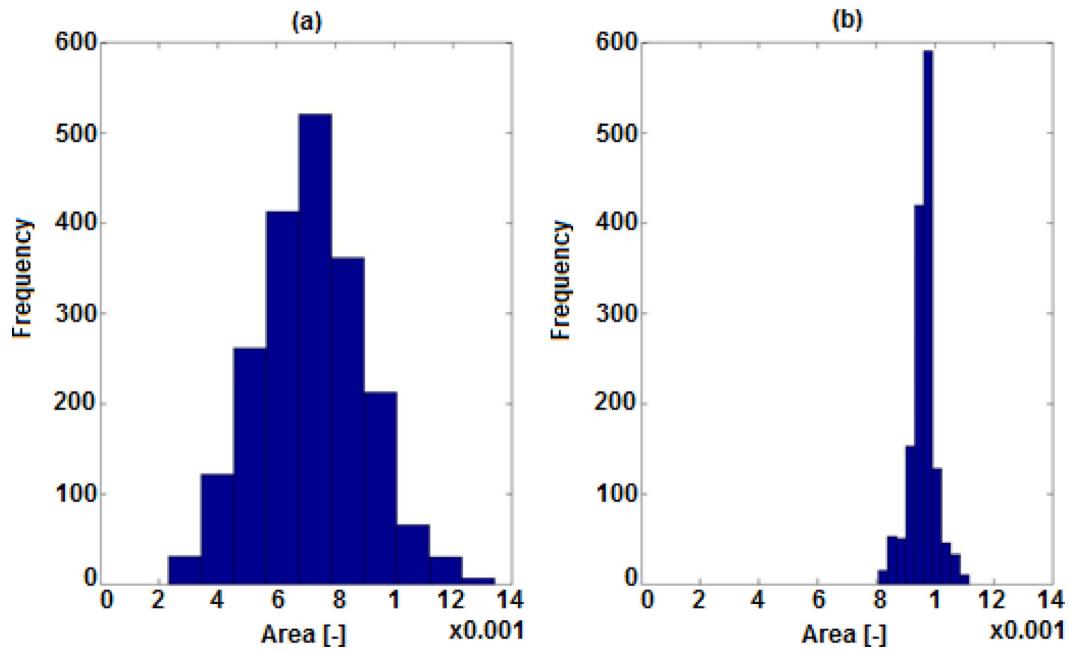


Fig. 12. Histograms of cellular area for microstructure: (a) before Lloyd's relaxation, $\alpha = 50$ and (b) after Lloyd's relaxation.

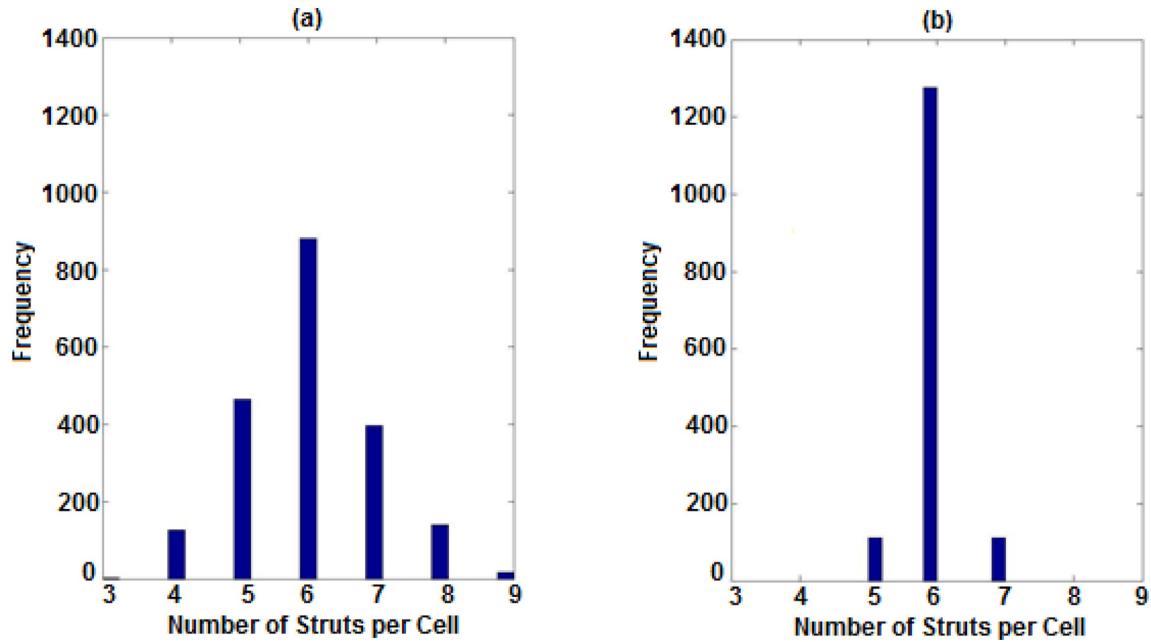


Fig. 13. Histograms for number of struts within each cell: (a) before Lloyd's relaxation, $\alpha = 50$ and (b) after Lloyd's relaxation.

1mm. All cellular areas were normalised based on the area of a single regular hexagonal cell.

In Fig. 6, each black point represents an individual cell. The location of the point in the figure indicates both the area of the individual cell and its orientation. Radial distance from the origin indicates the area of the cell while the polar location of the point indicates cell orientation. The radius of the red circle represents the average cell size within the RVE. Fig. 6, shows that for RVEs with very slight irregularity ($\alpha=5$) almost all the cells are orientated close to the angles 30° , 90° and 150° . This is to be expected as the structure is close to that of a regular honeycomb. The area of the individual cells is almost constant, i.e. mono-disperse, as indicated by the fact that most of the points lie close to the red semi-circle. As the degree of irregularity increases from 5 to 50 the data becomes

progressively more scattered around those angles and the points show increasing scatter away from the red semi-circle, indicating that the cell microstructure becomes increasingly poly-disperse as irregularity increases.

This can also be seen in Fig. 7, which shows a histogram of cell area for each degree of irregularity. The histogram distributions become broader as the degree of irregularity increases. A truly isotropic RVE should have no preferred directional orientation in the microstructure, i.e. in terms of Fig. 6, there should be no concentration of the points in any particular direction. By increasing the degree of irregularity, α , the cellular orientations tend to become less concentrated in any particular direction and the RVE morphology tends to become more random and isotropic. Fig. 5 clearly shows that the cost of this isotropy is a less realistic

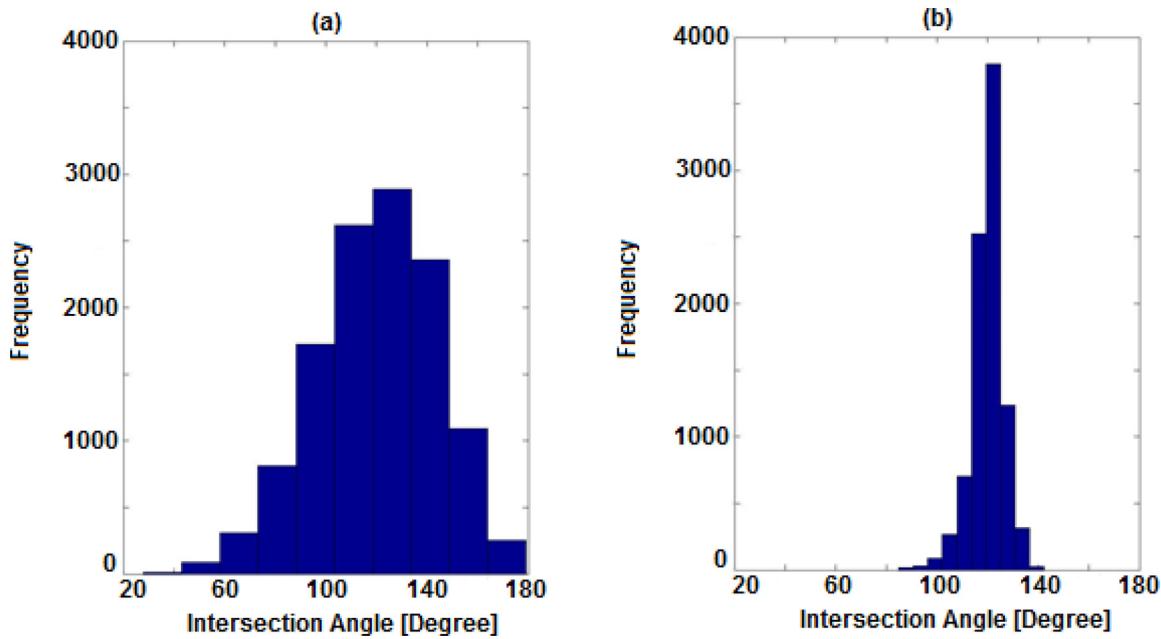


Fig. 14. Histograms for struts intersection angle within each cell for: (a) before Lloyd's relaxation, $\alpha=50$ (b) after Lloyd's relaxation.

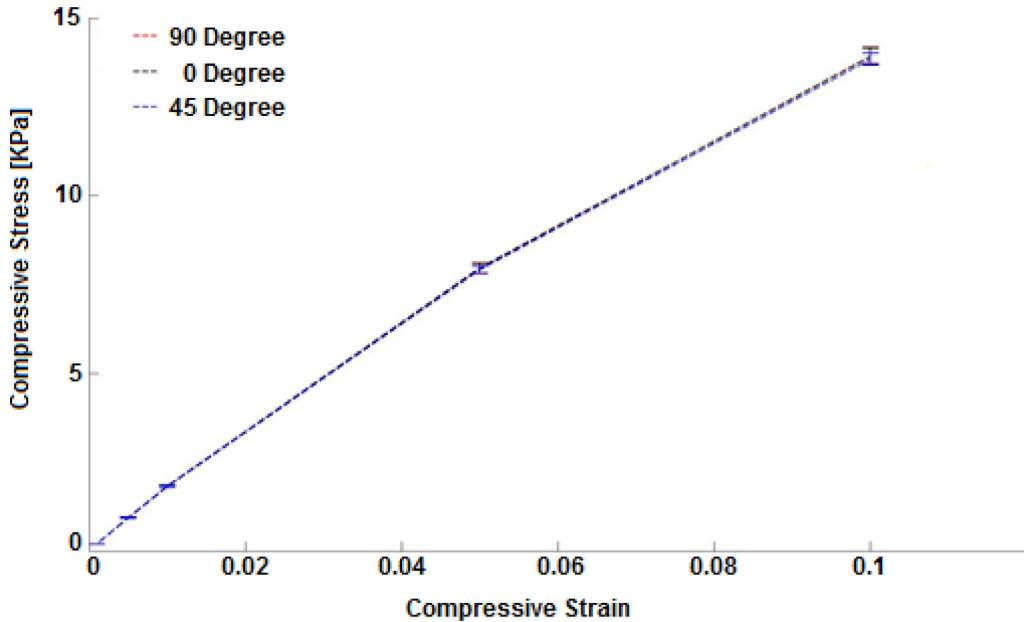


Fig. 15. Averaged stress-strain curves for compressed RVEs generated using Lloyd's algorithm.

microstructure with individual cells taking a less realistic, more angular form.

Histograms showing the number of struts per cell and the strut internal intersection angle are plotted in Figs. 8 and 9 and are an effective way to quantify the realism of average cell morphology. For a low degree of irregularity, the microstructure remains similar to that of a hexagonal honeycomb and all cells contain 6 struts (see Fig. 8a). As might be expected, the strut internal intersection angle remains close to 120 degrees. For high degrees of irregularity some cells containing 5 and 7 struts ($\alpha=25$) and even as few as 4 and as many as 9 struts ($\alpha=50$) are generated (see Fig. 8c and 8d). If these were regular cells, one may expect to see strut internal intersection angles as low as 108° and as high as 129° when $\alpha=25$ and as low as 90° and as high as 140° when $\alpha=50$. Examination of Fig. 9c and 9d, reveals a much wider range of strut internal inter-

section angles; lower than 65° and higher than 160° when $\alpha=25$ and lower than 28° and higher than 177° when $\alpha=50$, this exaggerated range in internal intersection angles is a good indicator of the highly irregular and angular cell shapes shown in Fig. 5c and 5d.

To examine the mechanical response, the RVEs were subject to 10% compression in the vertical, horizontal and 45° directions. The results are shown in Fig. 10. At very low strains (< 2% compression) the stiffness response in the three directions is practically identical for any of the three directions. Given that analytical theory predicts an isotropic response for honeycomb structures compressed under small strains (Gibson and Ashby, 1997), this result is to be expected. At larger strains (> 2% compression), the results diverge for low degrees of irregularity, reflecting an anisotropic stiffness response of a regular honeycomb structure

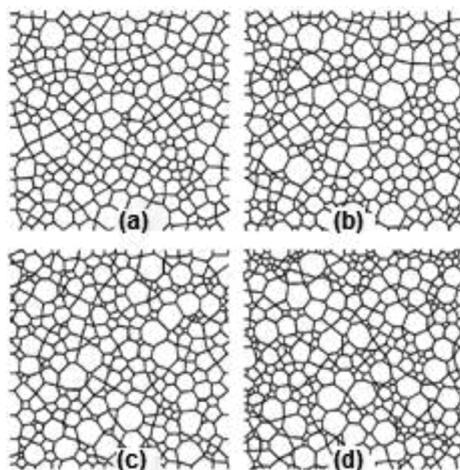


Fig. 16. RVEs with 4 different degree of poly-dispersity: (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$. The final beam-based structures are obtained by solving the Apollonius problem [Wolfram Mathematica®] for packed hard circular disks and then by applying a Laguerre-Voronoi tessellation.

under large strain compression (Zhu et al. 2006). Increasing the degree of irregularity clearly leads to a more isotropic response at large strains, reflecting the increasingly isotropic morphology observed in the randomised microstructure, as indicated in Fig. 6. It is worth noting that the small strain stiffness of the RVEs tends to increase with increasing degree of irregularity (for $\alpha = 5, 15, 25$ and 50 the moduli are: 171.66 ± 1.15 , 175.00 ± 2.65 , 186.00 ± 1 and 238.33 ± 6.35 KPa respectively), a result in agreement with that found previously by Zhu et al. 2001a.

3.2.2. Evaluation of microstructure generated using the Centroidal Voronoi Tessellation method

As discussed in Section 2.2, the Centroidal Voronoi Tessellation method used here employs the Lloyds algorithm to relax the microstructures originally generated using the basic Voronoi technique (see Section 2.1). To demonstrate the method, the maximum degree of irregularity, $\alpha = 50$ (also known as Poisson Voronoi

(Muche, 2005)), was used to create the RVEs. Once again 20 RVEs, each with periodic microstructure and boundary conditions, were generated. The microstructure was then ‘relaxed’ via application of the Lloyds algorithm (see, for example, Fig. 3); there is an obvious visual change in the morphology of the relaxed microstructure with cells clearly less angular and of lower poly-dispersity than the microstructure prior to application of the Lloyds algorithm. Fig. 11 is an equivalent polar plot to Fig. 6 and shows that the relaxed microstructure is isotropic and the cell size distribution is rather mono-disperse. Comparison of Fig. 12 and Fig. 7d confirms the reduction of poly-dispersity after application of the Lloyd’s algorithm. Fig. 13 shows that the relaxed microstructure contains almost no cells with less than 5 and more than 7 sides, indicating that the relaxation process does change the distribution for the number of struts within each cell. Nevertheless, Fig. 14 reveals a range of strut internal intersection angles of between 95° to 130° ; much closer to the range of 108° and 129° expected for regular pentagonal and heptagonal cells and much lower than the range of about 28° to 177° associated with the microstructure prior to relaxation.

The relaxed RVEs were also subject to 10% compression in the vertical, horizontal and at 45° directions. As expected, Fig. 15 shows an isotropic mechanical response. The stiffness is lower than that of the fully randomised microstructure from which it is generated (189 ± 1.08 KPa after relaxation compared to 238.33 ± 6.35 KPa prior to relaxation from a microstructure generated with $\alpha = 50$) and shows a more linear stress versus strain curve. It can be concluded that the CVT method can be applied to fully randomised RVEs (created using the basic Voronoi method) to generate isotropic RVEs with a realistic, predominantly mono-disperse microstructural morphology.

3.2.3. Evaluation of microstructure generated using the Laguerre-Voronoi Tessellation method

To generate a dense packing of circular hard disks using the Drop and Roll method (see Section 2.3), a square with size 80×80 is considered (see supplementary materials for detailed explanation of the Drop and Roll algorithm). The following parameter values are used in Eq. 5; $r_{min} = 0.5$, $r_{max} = 5$, $m = 2.5$, $\gamma = [0.05, 0.1, 0.15,$

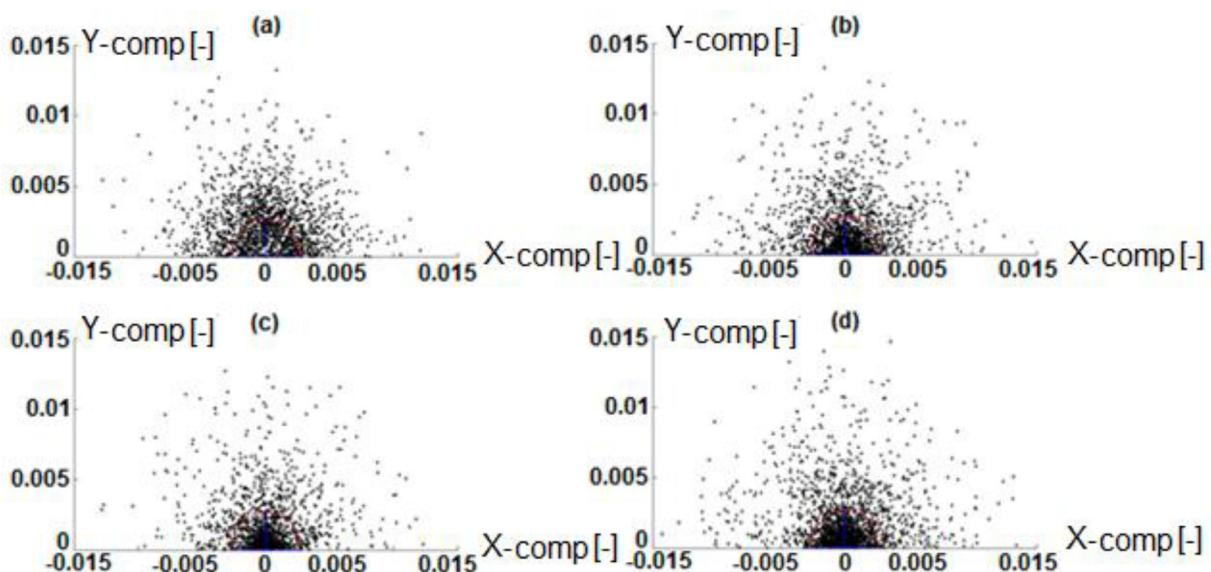


Fig. 17. Polar representation of cellular orientation. The X and Y axis represent the length of the normalised components of the semi-major axis of an ellipse, fitted within each individual cell of the RVE, multiplied by the cell’s area (black dots). The results are for microstructures generated using Laguerre-Voronoi Tessellation. (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$. The radius of the red circle represents an average value of the data for a given RVE.

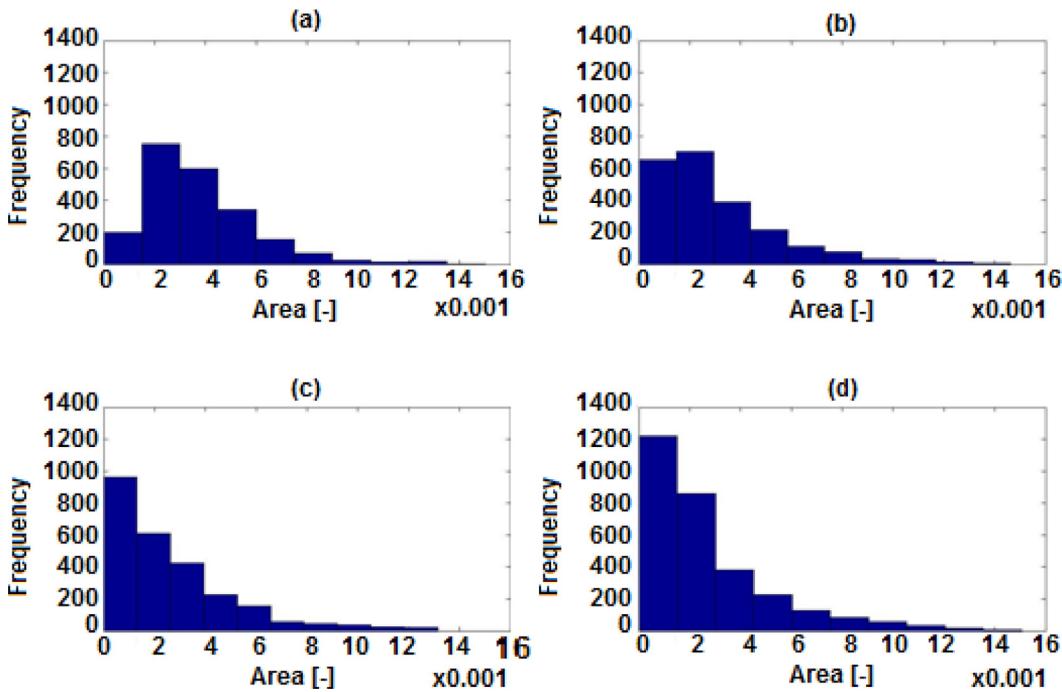


Fig. 18. histograms of cellular area for undeformed state of microstructure generated from Laguerre-Voronoi Tessellation. (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$.

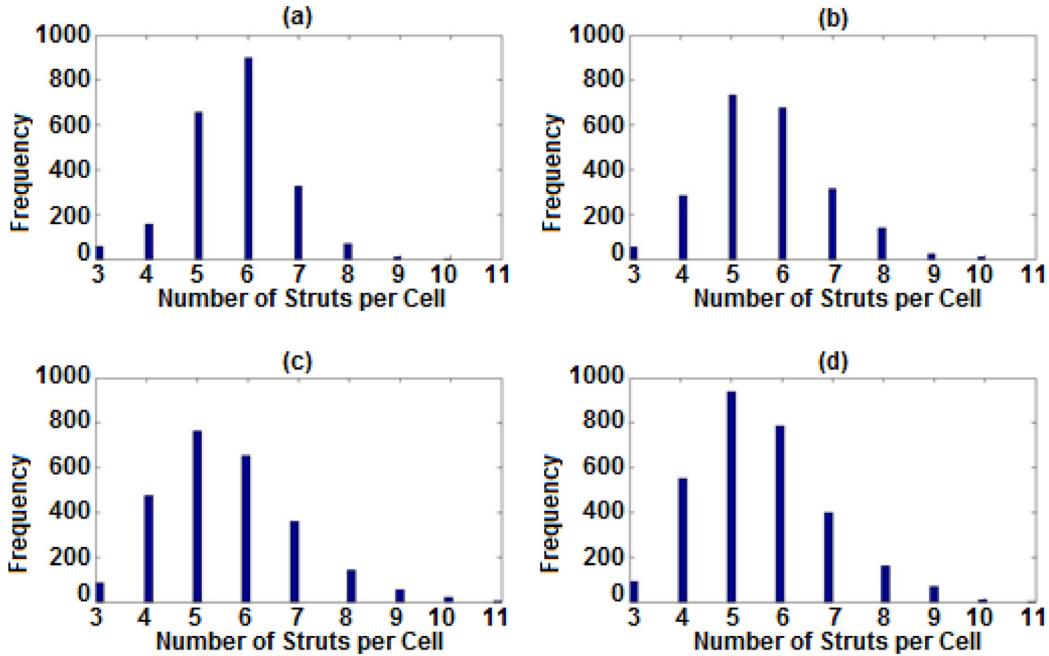


Fig. 19. histograms for number of struts within each cell for undeformed state of microstructure for different degrees of poly-dispersity: (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$.

0.20]. Increasing γ results in a higher poly-dispersity of the hard disc diameters, this is clearly reflected in the microstructural morphology shown in Fig. 16. To evaluate the average behaviour for each degree of poly-dispersity, 20 RVEs were generated for each value of γ . All RVEs generated in this part of the study have the same relative density as RVEs generated using the alternative generation algorithms discussed in Sections 3.2.1 & 3.2.2. This was done in order to facilitate comparison of the average small strain compressive modulus of the different RVEs.

Examining the microstructural morphology in more detail, Fig. 17 shows no obvious preferred direction for all cases and a high scatter of cell areas. The high poly-dispersity is confirmed in the histogram of Fig. 18, which reveals how the distribution range increases with increasing γ , i.e. the poly-dispersity increases with increasing γ . Fig. 18 shows that the shape of the distribution of cell areas is very different to that shown in Figs. 7 and 12, and can be well-fitted using a log-normal distribution. This is to be expected given that the same type of function was used to generate

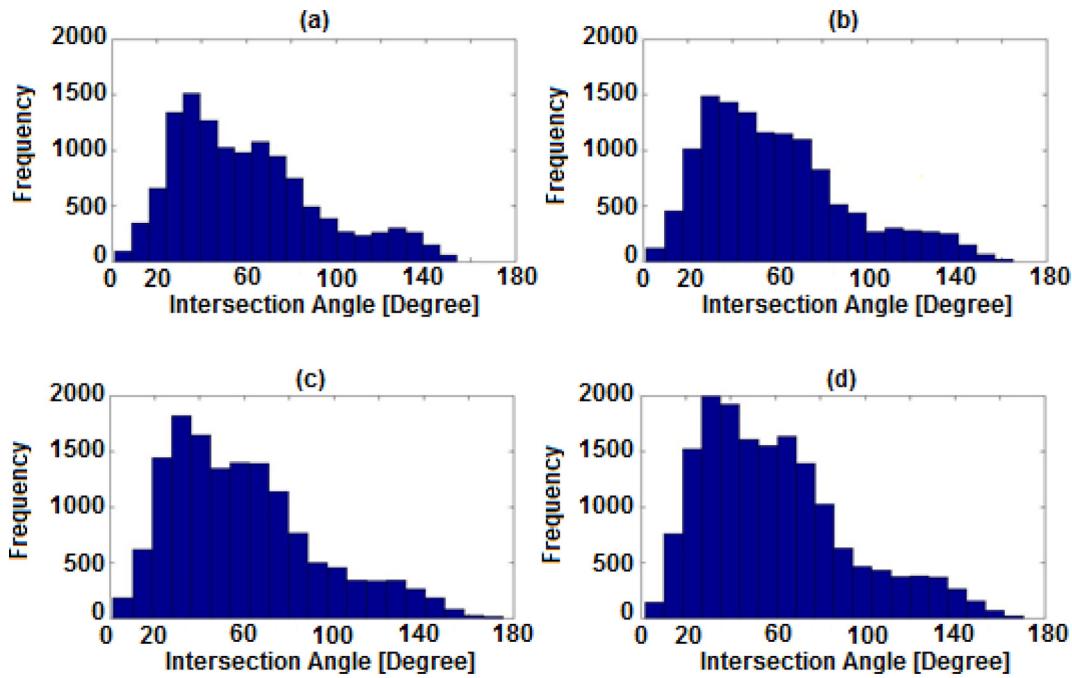


Fig. 20. histograms for struts intersection angle within each cell for the undeformed microstructure generated using different degrees of poly-dispersity: (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$.

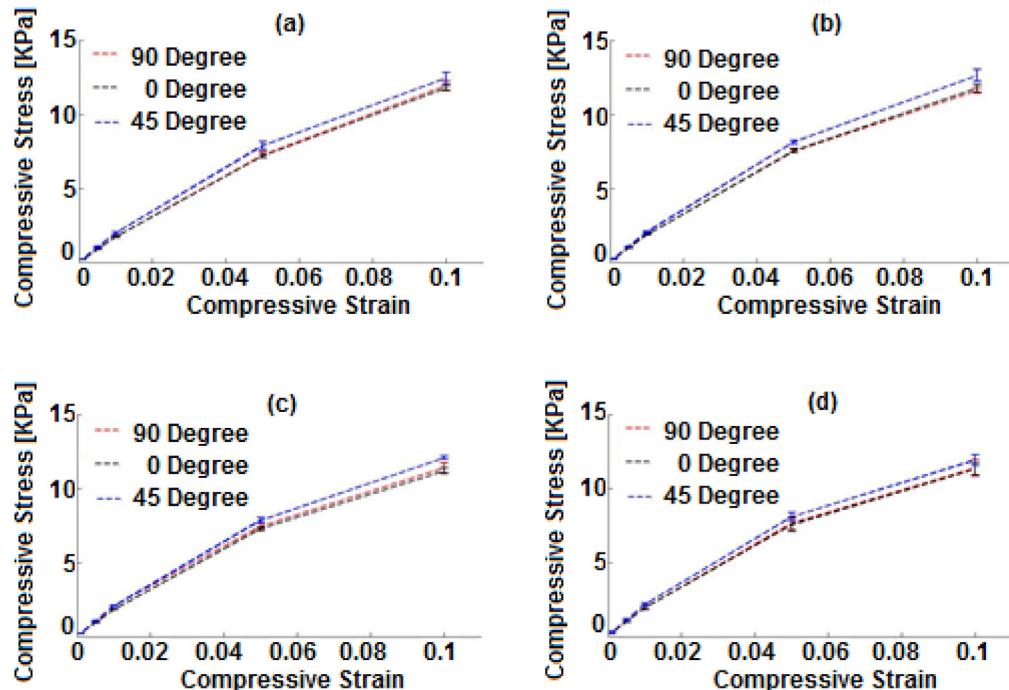


Fig. 21. Averaged stress-strain curves for compressed RVEs generated from Laguerre-Voronoi Tessellation with degree of poly-dispersity: (a) $\gamma = 0.05$, (b) $\gamma = 0.10$, (c) $\gamma = 0.15$, (d) $\gamma = 0.20$.

the seeding discs using the Drop and Roll method. As γ increases, the probability of generating smaller cells is seen to increase. This can also be seen in Fig. 17 which shows a denser clustering of points within the red circle at higher values of γ and a shrinking of the radius of the circle. Fig. 19 and Fig. 20 show that the distributions of both the number of sides per cell and the strut internal intersection angles are relatively insensitive to the degree of poly-dispersity.

The resulting RVEs are also subject to 10% compression in the perpendicular (along Y-axis), transverse (along X-axis) and

orthotropic directions (45° rotated from X-axis). Fig. 21 reveals that a small amount of anisotropy is evident in the mechanical response of the RVEs; the response in the 45° direction is slightly stiffer than in the X or Y directions. The origin of this anisotropy stems from anisotropy induced in the microstructure due to the basic functioning of the modified Drop and Roll algorithm implemented in this investigation. A full discussion of the latter is provided in the supplementary material. Nevertheless, the anisotropy is much smaller than that resulting from the basic Voronoi algorithm (see Section 3.2.1). There is a slight

trend towards increasing stiffness with increasing degree of polydispersity (for $\gamma = 0.05, 0.10, 0.15$ and 0.20 the modulus calculated at 0.1% strain were: 197.38 ± 25.9 , 219.72 ± 38.9 , 206.23 ± 18.43 and 233.78 ± 45.55 kPa, respectively), in agreement with results by previous researchers (Fazekas et al. 2002). This range of stiffness values lies within that found using the basic Voronoi algorithm, see Section 3.2.1.

4. Conclusions

Various generation algorithms for 2-D foam-like RVE microstructures have been examined. The aim of the investigation was to evaluate the capabilities of these algorithms in generating realistic microstructures suitable for modelling the mechanical response of 2-D foams with fully random microstructures. While actual foams may in some cases possess transversely isotropic mechanical properties, a pre-requisite of microstructure generation algorithms should be the ability to produce mechanically isotropic RVEs; introduction of any intended anisotropy is subsequently a simple matter that can be achieved by simple affine stretching of the RVE microstructure.

The basic Voronoi tessellation method has been shown to generate microstructures possessing anisotropy; an unintended artefact of the seeding method used by the algorithm. A novel method of examining the microstructural morphology has been introduced using polar plots; these diagrams are a simple technique to visually highlight the statistical orientation and size distribution of cells within the RVEs. Mechanical compression tests conducted in three different directions on the RVEs confirmed the anisotropy revealed by the polar plots. It was found that morphological and mechanical anisotropy could be erased in RVEs with very high degrees of irregularity, though this created highly angular cell morphologies that were considered to be unrealistic. This effect has been studied by examining the strut internal angle distribution; an overly wide distribution is linked to an unrealistic, highly angular cellular morphology. Nevertheless, it has been shown that isotropic RVEs with realistic cellular morphologies can be created from such RVEs by applying a Centroidal Voronoi Tessellation algorithm to relax the microstructure. This combined approach results in the successful creation of isotropic mono-disperse RVEs, though the method is restricted to the generation of relatively mono-disperse microstructures. In order to generate poly-disperse microstructures the so-called Laguerre-Voronoi tessellation algorithm has been employed and combined with a modified version of the Drop-and-Roll method in order to create the seeding for the cells. The method is shown to successfully produce almost fully isotropic RVEs with a controlled degree of polydispersity. Together, the numerical algorithms analysed in this investigation offer the possibility to generate isotropic 2-D RVEs with realistic microstructures ranging from mono-disperse to highly poly-disperse cellular morphologies. The small strain response of such RVEs is known to be reasonably well-predicted using analytical theory (Gibson and Ashby, 1997) and so the main value of type of RVEs presented in this paper would ultimately be exploited when simulating their large-strain mechanical response (e.g. Alsayednoor 2013). The main restriction of the current investigation is its focus on 2-D RVEs. A longer term goal is therefore to extend the techniques developed in this work, to the generation of equivalent 3-D RVEs, a more challenging task but ultimately more useful in modelling real foam materials.

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Supplementary materials

Supplementary material associated with this article can be found, in the online version, at doi:10.1016/j.mechmat.2016.04.001.

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