



Problem Sheet 1

Solve any three problems.

1. A Lorentz transformation $x^{\mu} \to x'^{\mu} = \Lambda^{\mu}{}_{\nu}x^{\nu}$ is such that $\eta_{\mu\nu}x^{\mu}x^{\nu} = \eta_{\mu\nu}x'^{\mu}x'^{\nu}$ where $\eta_{\mu\nu}$ is the Minkowski metric and x^{μ} is any 4-vector. Show that this implies that

$$\eta_{\mu\nu} = \eta_{\sigma\tau} \Lambda^{\sigma}{}_{\mu} \Lambda^{\tau}{}_{\nu} \,. \tag{1}$$

Use this result to show that an infinitesimal transformation around identity of the form

$$\Lambda^{\mu}{}_{\nu} = \delta^{\mu}{}_{\nu} + \omega^{\mu}{}_{\nu} \tag{2}$$

 (ω^{μ}_{ν}) is infinitesimally small) is a Lorentz tranformation when $\omega_{\mu\nu}$ is antisymmetric: i.e. $\omega_{\mu\nu} = -\omega_{\nu\mu}$.

2. The Lagrangian density for a complex scalar field $\Phi(x) = \phi_1(x) + i\phi_2(x)$ is given by

$$\mathcal{L} = \partial_{\mu} \Phi^* \partial^{\mu} \Phi - m^2 \Phi^* \Phi - \frac{\lambda}{2} (\Phi^* \Phi)^2 . \tag{3}$$

Write down the Euler-Lagrange equations of motion (EOM). Verify that the Lagrangian is invariant under the infinitesimal transformations

$$\delta \Phi = i\alpha \Phi \quad , \quad \delta \Phi^* = -i\alpha \Phi^* \,, \tag{4}$$

which is an infinitesimal version of $\Phi \to e^{i\alpha}\Phi$.

Derive the Noether current associated with this transformation and verify explicitly that it is conserved using the field equations satisfied by ϕ .

It is also instructive to write \mathcal{L} in terms of the two real scalar fields ϕ_1 , ϕ_2 . Can you justify the numerical factors appeared in the Lagrangian?

- 3. Verify Wick's theorem for
 - (a) The product of three scalars:

$$T(\phi(x_1)\phi(x_2)\phi(x_3)) = : \phi(x_1)\phi(x_2)\phi(x_3) : +\phi(x_1)\Delta_F(x_2 - x_3) + \phi(x_2)\Delta_F(x_3 - x_1) + \phi(x_3)\Delta_F(x_1 - x_2)$$
(5)

(b) The product of four scalars:

$$T\left(\phi\left(x_{1}\right)\phi\left(x_{2}\right)\phi\left(x_{3}\right)\phi\left(x_{4}\right)\right). \tag{6}$$

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4. The Fourier decompositions of a real scalar field and its conjugate momentum are (note that they are at t=0, as compared to the lectures)

$$\phi(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_{\vec{p}}}} \left[a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} + a_{\vec{p}}^{\dagger} e^{-i\vec{p}\cdot\vec{x}} \right]$$

$$\pi(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} (-i) \sqrt{\frac{E_{\vec{p}}}{2}} \left[a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} - a_{\vec{p}}^{\dagger} e^{-i\vec{p}\cdot\vec{x}} \right]$$
(7)

with (equal time at t = 0) commutation relationships

$$[\phi(\vec{x}), \phi(\vec{y})] = [\pi(\vec{x}), \pi(\vec{y})] = 0$$
 and $[\phi(\vec{x}), \pi(\vec{y})] = i\delta^{(3)}(\vec{x} - \vec{y})$ (8)

implying

$$[a_{\vec{p}}, a_{\vec{q}}] = [a_{\vec{p}}^{\dagger}, a_{\vec{q}}^{\dagger}] = 0$$
 and $[a_{\vec{p}}, a_{\vec{q}}^{\dagger}] = (2\pi)^3 \delta^{(3)}(\vec{p} - \vec{q})$ (9)

- (a) Generalize above analysis for the complex scalar field whose lagrangian is given in problem 2.
- (b) Compute the Feynman propagator for the complex scalar field defined by $\langle 0|T\left(\Phi(x)\Phi^{\dagger}(y)\right)|0\rangle$.
- (c) Use the mode expansion of the field and derive the propagator in momentum space.
- (d) Lastly, compute the time-ordered product $T(\Phi(x)\Phi^{\dagger}(y)\Phi(z))$.
- 5. Solve the exercises mentioned in the lectures 1 and 2.

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