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MA-473

MidSem-Exam

PAPER -①

① Given transformation

$$x = \ln S.$$

$$\tau = \frac{\sigma^2}{2}(T-t)$$

$$V(S,t) = e^{\alpha x + \beta \tau} u(x,\tau) \text{ where } \begin{aligned} \alpha &= -\left(\frac{2r}{\sigma^2} - 1\right)/2 \\ \beta &= -\left(\frac{2r}{\sigma^2} + 1\right)/4 \end{aligned}$$

We have

$$\frac{\partial V}{\partial t} = \frac{\partial}{\partial t} \left[e^{\alpha x + \beta \tau} u(x,\tau) \right] = \frac{\partial}{\partial \tau} \left[e^{\alpha x + \beta \tau} u(x,\tau) \right] \cdot \frac{d\tau}{dt}$$

$$\boxed{\frac{\partial V}{\partial t} = \left(\beta e^{\alpha x + \beta \tau} u(x,\tau) + e^{\alpha x + \beta \tau} \frac{\partial u}{\partial \tau} \right) \left(-\frac{\sigma^2}{2} \right)}$$

$$\frac{\partial V}{\partial S} = \frac{\partial}{\partial S} \left[e^{\alpha x + \beta \tau} u(x,\tau) \right] = \frac{\partial}{\partial x} \left[e^{\alpha x + \beta \tau} u(x,\tau) \right] \cdot \frac{\partial x}{\partial S}$$

$$\boxed{\frac{\partial V}{\partial S} = \left(\alpha e^{\alpha x + \beta \tau} u(x,\tau) + e^{\alpha x + \beta \tau} \frac{\partial u}{\partial x} \right) \frac{1}{S}}$$

$$\begin{aligned}
 \frac{\partial^2 v}{\partial s^2} &= \frac{\partial}{\partial s} \left(\left[\frac{\alpha}{S} e^{\alpha x + \beta T} u + e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right] \frac{1}{S} \right) \\
 &= -\frac{1}{S^2} \left[\alpha e^{\alpha x + \beta T} u + e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right] + \frac{1}{S} \frac{\partial}{\partial s} \left[\alpha e^{\alpha x + \beta T} u + e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right] \\
 &= -\frac{1}{S^2} \left[\alpha e^{\alpha x + \beta T} u + e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right] + \frac{1}{S} \left\{ \alpha \frac{\partial}{\partial x} \left(e^{\alpha x + \beta T} u \right) \cdot \frac{\partial x}{\partial s} \right. \\
 &\quad \left. + \frac{\partial}{\partial x} \left(e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right) \cdot \frac{\partial x}{\partial s} \right\} \\
 &= -\frac{e^{\alpha x + \beta T}}{S^2} \left[du + \frac{\partial u}{\partial x} \right] + \frac{1}{S} \left\{ \frac{\partial}{\partial s} \left[\alpha e^{\alpha x + \beta T} u + e^{\alpha x + \beta T} \frac{\partial u}{\partial x} \right] \right. \\
 &\quad \left. + \frac{1}{S} \left[\alpha e^{\alpha x + \beta T} \frac{\partial u}{\partial x} + e^{\alpha x + \beta T} \frac{\partial^2 u}{\partial x^2} \right] \right\} \\
 &= -\frac{e^{\alpha x + \beta T}}{S^2} \left[\alpha u + \frac{\partial u}{\partial x} \right] + \frac{\alpha}{S^2} e^{\alpha x + \beta T} \left[\alpha u + \frac{\partial u}{\partial x} \right] + \frac{\alpha}{S^2} e^{\alpha x + \beta T} \left[\alpha \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} \right] \\
 \boxed{\frac{\partial^2 v}{\partial s^2} = \frac{e^{\alpha x + \beta T}}{S^2} \left\{ -du - \frac{\partial u}{\partial x} + \alpha^2 u + \alpha \frac{\partial u}{\partial x} + \alpha \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} \right\}}
 \end{aligned}$$

Now putting the values in B.S.PDE i.e,

$$\frac{\partial v}{\partial t} + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2 v}{\partial s^2} + rS \frac{\partial v}{\partial s} - rv = 0$$

we get;

$$\left. \begin{aligned} -\frac{\sigma^2}{2} e^{\alpha x + \beta t} \left(\beta u + \frac{\partial u}{\partial t} \right) + \frac{1}{2} \sigma^2 s^2 \cdot \frac{1}{s^2} \left\{ -\alpha u - \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} + \alpha \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} \right\} e^{\alpha x + \beta t} \end{aligned} \right.$$

$$+ \gamma S \cdot \frac{e^{\alpha x + \beta t}}{S} \left(\alpha u + \frac{\partial u}{\partial x} \right) - \gamma e^{\alpha x + \beta t} u = 0$$

$$\Rightarrow -\frac{\sigma^2}{2} \left(\beta u + \frac{\partial u}{\partial t} \right) + \frac{\sigma^2}{2} \left(-\alpha u - \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} + 2\alpha \frac{\partial u}{\partial x} + \frac{\partial^2 u}{\partial x^2} \right) + \gamma \left(\alpha u + \frac{\partial u}{\partial x} \right) - \gamma u = 0$$

$$\Rightarrow -\frac{\sigma^2}{2} \left(\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} \right) + \left(-\frac{\sigma^2 \beta}{2} u - \frac{\sigma^2 \alpha}{2} u - \frac{\sigma^2}{2} \frac{\partial u}{\partial x} + \frac{\sigma^2 \alpha^2}{2} u + \frac{\sigma^2}{2} \frac{\partial u}{\partial x} + \gamma \alpha u + \gamma \frac{\partial u}{\partial x} - \gamma u \right) = 0$$

$$\Rightarrow \boxed{\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0} \quad \left. \begin{array}{l} -\infty < x < \infty \\ t \geq 0 \end{array} \right.$$

Now we will show:

$$-\frac{\sigma^2 \beta}{2} u - \frac{\sigma^2 \alpha}{2} u - \frac{\sigma^2}{2} \frac{\partial u}{\partial x} + \frac{\sigma^2 \alpha^2}{2} u + \frac{\sigma^2 \alpha}{2} \frac{\partial u}{\partial x} + \gamma \alpha u + \gamma \frac{\partial u}{\partial x} - \gamma u \geq 0$$

$$\Rightarrow \left(-\frac{\sigma^2 \beta}{2} - \frac{\sigma^2 \alpha}{2} + \frac{\sigma^2 \alpha^2}{2} + \gamma \alpha - \gamma \right) u + \left(\frac{-\sigma^2}{2} + \sigma^2 \alpha + \gamma \right) \frac{\partial u}{\partial x}$$

$$= \left(-\frac{\sigma^2 \beta}{2} + \frac{\sigma^2 \alpha(\alpha-1)}{2} + \gamma(\alpha-1) \right) u + \left(\frac{-\sigma^2}{2} + \frac{\sigma^2}{2} \left(\frac{2\alpha}{\sigma} + 1 \right) + \gamma \right) \frac{\partial u}{\partial x}$$

$$\Rightarrow -\frac{\sigma^2}{2}\beta + \frac{\sigma^2}{2}\alpha(\alpha-1) + \gamma(\alpha-1)$$

on putting values of $\alpha = -\frac{1}{2} \left(\frac{2\gamma}{\sigma^2} - 1 \right)$

$$\text{and } \beta = -\frac{1}{4} \left(\frac{2\gamma}{\sigma^2} + 1 \right)^2$$

we will get

$$\boxed{-\frac{\sigma^2}{2}\beta + \frac{\sigma^2}{2}\alpha(\alpha-1) + \gamma(\alpha-1) = 0}$$

Hence Proved

Now ~~the~~ final condition

$$V(S, T) = V_T(S)$$

$$U(x, 0) \cdot e^{\alpha x} = V(S, T) = V_T(S)$$

$$\Rightarrow \boxed{U(x, 0) = e^{-\alpha x} V_T(S)}$$

Hence B-S PDE is converted to

1-D heat conduction PDE.

(2) Given BVP:

$$\left\{ \begin{array}{l} \frac{\partial u}{\partial \tau} = \frac{\partial^2 u}{\partial x^2}, \quad -\infty < x < \infty, \tau > 0 \\ u(x, 0) = u_0(x) \end{array} \right. \rightarrow *$$

Consider solution of form:

$$u(x, \bar{\tau}) = \bar{\tau}^{-1/2} v(\eta) \text{ where } \eta = \frac{x-\xi}{\sqrt{\bar{\tau}}} \text{ and } \xi \rightarrow \text{parameter}$$

Then

$$\frac{\partial u}{\partial \tau} = \frac{\partial}{\partial \bar{\tau}} \left(\bar{\tau}^{-1/2} v(\eta) \right) = \frac{\partial}{\partial \bar{\tau}} (\bar{\tau}^{-1/2}) v(\eta) + \bar{\tau}^{-1/2} \frac{\partial}{\partial \bar{\tau}} v(\eta)$$

$$= -\frac{1}{2} \bar{\tau}^{-3/2} v(\eta) + \bar{\tau}^{-1/2} \frac{\partial u}{\partial \eta} \cdot \frac{d\eta}{d\bar{\tau}}$$

$$= -\frac{1}{2} \bar{\tau}^{-3/2} v(\eta) + \bar{\tau}^{-1/2} \frac{\partial u}{\partial \eta} (x-\xi) \bar{\tau}^{-3/2} / 2$$

$$= -\frac{1}{2} \bar{\tau}^{-3/2} \left(v + \eta \frac{\partial u}{\partial \eta} \right)$$

$$\Rightarrow \boxed{\frac{\partial u}{\partial \tau} = -\frac{1}{2} \bar{\tau}^{-3/2} \frac{d}{d\eta} \left(\eta v(\eta) \right)}$$

$$\frac{\partial u}{\partial x} = \frac{2}{\partial x} \left[\bar{\tau}^{-1/2} v(\eta) \right] = \bar{\tau}^{-1/2} \frac{du}{d\eta} \cdot \frac{d\eta}{dx}$$

$$\boxed{\frac{\partial u}{\partial x} = \bar{\tau}^{-1} \frac{du}{d\eta}}$$

$$\Rightarrow \frac{\partial^2 u}{\partial x^2} = \frac{2}{\partial x} \left[\bar{\tau}^{-1} \frac{du}{d\eta} \right] = \bar{\tau}^{-1} \frac{d^2 u}{d\eta^2} \cdot \frac{d\eta}{dx}$$

$$\Rightarrow \boxed{\frac{\partial^2 u}{\partial x^2} = \bar{\tau}^{-3/2} \frac{d^2 v}{d\eta^2}}$$

Putting the values in *, we get

$$-\frac{\bar{\tau}^{-3/2}}{2} \frac{d}{d\eta} [\eta v(\eta)] = \bar{\tau}^{-3/2} \frac{d^2 v}{d\eta^2}$$

$$\Rightarrow \boxed{\frac{d^2 v}{d\eta^2} + \frac{1}{2} \frac{d}{d\eta} [\eta v] = 0}$$

Integrating this \Leftrightarrow ODE we get

$$\boxed{\frac{du}{d\eta} + \frac{\eta}{2} v = C_1} \rightarrow \text{const of Integration}$$

choose $C=0$

Then we have linear homogeneous eqn

$$\frac{du}{d\eta} + \frac{\eta}{2} u = 0$$

$$\frac{du}{u} = -\frac{\eta}{2} d\eta \Rightarrow \ln\left(\frac{u}{C}\right) = -\frac{\eta^2}{4}$$

$$\Rightarrow \boxed{u = C e^{-\eta^2/4}} \rightarrow C \text{ is const.}$$

Hence for heat eqn we have solⁿ of form.

$$\boxed{u(x, \bar{t}) = C \bar{t}^{-1/2} e^{-(x-\xi)^2/4\bar{t}}}$$

Now we need

$$\int_{-\infty}^{\infty} C \bar{t}^{-1/2} e^{-(x-\xi)^2/4\bar{t}} d\xi = 1$$

$$\text{Putting } \xi = \frac{x-\eta}{\sqrt{2\bar{t}}} \Rightarrow \boxed{\xi = x - (2\bar{t})^{1/2} \eta}$$

$$d\xi = (2\bar{t})^{1/2} d\eta \Rightarrow \bar{t}^{-1/2} d\xi = \sqrt{2} d\eta$$

$$\Rightarrow \int_{-\infty}^{\infty} c\sqrt{2} e^{-\eta^2/2} d\eta = 1$$

$$\Rightarrow C\sqrt{2} \cdot \sqrt{2\pi} \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\eta^2/2} d\eta = 1$$

$\brace{= 1 \text{ by Standard } N(0,1) \text{ CDF}}$

$$\Rightarrow \boxed{C = \frac{1}{2\sqrt{\pi}}}$$

then we get special solution as.

$$\boxed{u(x, \bar{x}) = \frac{1}{2\sqrt{\pi \bar{x}}} e^{-(x-q)^2/4\bar{x}}} \rightarrow \begin{array}{l} \text{fundamental soln} \\ \text{or} \\ \text{Green function} \\ \text{for heat Eq} \end{array}$$

let $g(q; x, \bar{x})$ be class of green func. with q parameter
then

$$\frac{\partial g}{\partial \bar{x}} = \frac{\partial^2 g}{\partial x^2} \quad \text{for any } q$$

thus for $u_0(q)$, we have

$$\int_{-\infty}^{\infty} u_0(q) \cdot \frac{\partial g}{\partial \bar{x}} dq = \int_{-\infty}^{\infty} u_0(q) \frac{\partial^2 g}{\partial x^2} dq$$

$$\Rightarrow \frac{\partial}{\partial \bar{c}} \left[\int_{-\infty}^{\infty} u_0(q) g dq \right] = \frac{\delta^2}{\delta u^2} \left[\int_{-\infty}^{\infty} u_0(q) g dq \right]$$

$$\Rightarrow \boxed{u(x, \bar{c}) = \int_{-\infty}^{\infty} u_0(q) x \frac{1}{2\sqrt{\pi \bar{c}}} e^{-\frac{(x-q)^2}{4\bar{c}}} dq}$$

\downarrow
is also sol' of * with D.C. $u_0(x)$

Now

$$\lim_{\bar{c} \rightarrow 0} \frac{1}{2\sqrt{\pi \bar{c}}} e^{-\frac{(x-q)^2}{4\bar{c}}} dq = \begin{cases} 0 & , x-q \neq 0 \\ \infty & , x-q = 0 \end{cases}$$

and also

$$\int_{-\infty}^{\infty} \frac{1}{2\sqrt{\pi \bar{c}}} e^{-\frac{(x-q)^2}{4\bar{c}}} dq = 1 \text{ is true } \forall \bar{c}$$

we have:

$$\boxed{\lim_{\bar{c} \rightarrow 0} \frac{1}{2\sqrt{\pi \bar{c}}} e^{-\frac{(x-q)^2}{4\bar{c}}} = \delta(x-q)}$$

Dirac-Delta function

$$\boxed{\lim_{\bar{c} \rightarrow 0} \int_{-\infty}^{\infty} u_0(q) x \frac{1}{2\sqrt{\pi \bar{c}}} e^{-\frac{(x-q)^2}{4\bar{c}}} dq = u_0(x)}$$

 \downarrow
I.C.

Now for European call option

$$\text{Using } V(S,t) = e^{-r(T-t)} u(x, \bar{t}).$$

we have for B-S PDE

$$V(S,t) = e^{-r(T-t)} \int_{-\infty}^{\infty} u_0(q) \cdot \frac{1}{2\sqrt{\pi\bar{t}}} e^{-(x-q)^2/4\bar{t}} dq.$$

$$V(S,t) = e^{-r(T-t)} \int_{-\infty}^{\infty} V_T(e^q) \frac{1}{2\sqrt{\pi\bar{t}}} e^{-(q-x)^2/4\bar{t}} dq$$

$$V(S,t) = e^{-r(T-t)} \frac{1}{\sigma \sqrt{2\pi(T-t)}} X$$

$$\int_0^{\infty} V_T(s') e^{-\left\{ \ln s' - \left[\ln S + \left(r - \frac{\sigma^2}{2} \right)(T-t) \right] \right\}^2 / 2\sigma^2(T-t)} \frac{ds'}{s'}$$

where $G(s', T; s, t) = \frac{1}{\sigma \sqrt{2\pi(T-t)}} e^{-\left\{ \ln s' - \left[\ln S + \left(r - \frac{\sigma^2}{2} \right)(T-t) \right] \right\}^2 / 2\sigma^2(T-t)}$

is Green's function for B-S PDE.

for call option

$$V_T(s') = \max \{ s' - K, 0 \}$$