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MA-473: COMPUTATIONAL FINANCE

ENDSEM -> PAPER-I

1 Given PDE

$$\int \frac{\partial V}{\partial t} + \frac{1}{2} \frac{\partial \tilde{h}}{\partial s^2} \frac{\partial^2 V}{\partial s^2} + (v-s) \frac{\partial V}{\partial s} - vV = 0, \quad 0 \le s,$$

$$V(s,T) = V_T(s), \quad 0 \le s$$

$$V(s,T) = V_T(s), \quad 0 \le s$$

Using following transformation

$$\begin{cases} \zeta = \frac{S}{S+lm} \\ T = T-t \end{cases}$$

$$V(S_{1}+) = (S+lm) V(\zeta_{1},T)$$

Then we have
$$S = \frac{lm q}{1-q}$$
 and $S + lm = \frac{lm}{1-q}$ 3

Now $\frac{dq}{dq} = 1 \cdot \frac{1}{S+P_m} + S \cdot \left(\frac{-1}{(S+P_m)^2}\right) \cdot 1$

$$\frac{dQ}{dS} = \frac{P_m}{(S+P_m)^2}$$

$$\Rightarrow \left[\frac{dq}{ds} = \frac{(1-q)^2}{r_m}\right] \rightarrow 4$$
 \(\text{Vsmf (3) }\)

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Now we have

$$\frac{\partial V}{\partial t} = \frac{\partial}{\partial t} \left[\left(S + Pm \right) V \left(q_i \tau \right) \right] = \left(S + Pm \right) \frac{\partial V \left(q_i \tau \right)}{\partial t}$$

$$= (S+Pm) \frac{\partial V}{\partial t} \cdot \frac{\partial T}{\partial t} = -(S+Pm) \frac{\partial V}{\partial t} \Big|$$

Using 3, we have

$$\frac{\partial V}{\partial t} = -\frac{\rho_m}{1-Q} \frac{\partial \overline{V}}{\partial C}$$

Now
$$\frac{\partial V}{\partial S} = \frac{\partial}{\partial S} \left[\left(S + Rm \right) \overline{V(S_1, \overline{c})} \right] = 1 \cdot \overline{V(S_1, \overline{c})} + \left(S + Rm \right) \frac{\partial \overline{V(S_1, \overline{c})}}{\partial S}$$

$$= V + \frac{\rho_m}{1-q} \cdot \frac{\partial V}{\partial q} \cdot \frac{(1-q)^2}{\rho_m} \int u \sin \frac{3}{q} \frac{1}{q} \frac{1}{q} \frac{\partial u}{\partial q} \frac{\partial u}{\partial$$

$$=) \sqrt{\frac{\partial V}{\partial S}} = V + (1-8) \frac{\partial V}{\partial 9}$$

Now
$$\frac{\partial^2 V}{\partial S^2} = \frac{\partial}{\partial S} \left[\frac{\partial V}{\partial S} \right] = \frac{\partial}{\partial S} \left[\overline{V} + (1 - 9) \frac{\partial \overline{V}}{\partial 9} \right]$$

$$=\frac{\partial \overline{V}}{\partial q},\frac{\partial q}{\partial s}+\frac{\partial}{\partial s}\left[(1-q),\frac{\partial \overline{V}}{\partial q}\right]$$

$$(\overline{3})$$

$$= \frac{\partial^{2} v}{\partial s^{2}} = \frac{\partial v}{\partial s} \cdot \frac{\partial s}{\partial s} + \frac{\partial}{\partial s} \left[(i + \frac{\partial}{\partial s}) \frac{\partial v}{\partial s} \right] \cdot \frac{\partial s}{\partial s}$$

$$= \frac{\partial v}{\partial s} \cdot \frac{\partial s}{\partial s} - \frac{\partial v}{\partial s} \cdot \frac{\partial s}{\partial s} + (i + \frac{\partial}{s}) \cdot \frac{\partial v}{\partial s^{2}} \cdot \frac{\partial s}{\partial s}$$

$$= \frac{\partial v}{\partial s} \cdot \frac{\partial s}{\partial s} - \frac{\partial v}{\partial s} \cdot \frac{\partial s}{\partial s} + (i + \frac{\partial}{s}) \cdot \frac{\partial v}{\partial s^{2}} \cdot \frac{\partial s}{\partial s}$$

$$=) \left[\frac{\partial^2 V}{\partial S^2} - \frac{(1-4)^3}{P_m} \frac{\partial^2 V}{\partial 4^2} \right]$$

Now let
$$\overline{\mathcal{T}}(G) = \overline{\mathcal{T}}(S(G)) = \overline{\mathcal{T}}(\frac{\rho_m G}{1-G})$$

Now putting all values in PDECD, De get

$$\frac{\partial \overline{V}}{\partial \overline{L}} = \frac{\mathcal{F}^{2}(q) \cdot q^{2}(1-q)^{2}}{2} \cdot \frac{\partial^{2} \overline{V}}{\partial q^{2}} + (\gamma-5) q(1-q) \frac{\partial \overline{V}}{\partial q}
- [\gamma(1-q) + Sq] \overline{V}$$

$$0 \le q < 1, \ T \ge 0$$

Now assuming V is smooth function of E_q .

then E_q^n also holds at $E_q^n = 1$.

Now
$$V(S,T) = (S+P_m) \overline{V(S_10)} = \overline{V(S_10)} \frac{P_m}{1-S_1}$$

Hence the condition $V(S,T) = V_T(S)$ can be re-written as:

$$V(\varsigma_{10}) = V_{+}\left(\frac{\rho_{m}\varsigma_{1}}{1-\varsigma_{1}}\right) \cdot \frac{1-\varsigma_{1}}{\rho_{m}}$$

This is extent Printial Condition for PDE (5).

Boundary Conditions

At 9=0 and 9=1, PDE 5 degenerates to O.D. E at boundaries

$$=) \frac{dV(o_i\tau)}{V(o_i\tau)} = -v\cdot d\tau$$

Similarly [at
$$G_1=1$$
]
$$\frac{\partial V(1,T)}{\partial T} = -SV(1,T)$$

$$= \frac{dV(1,T)}{V(1,T)} = -8dT$$

$$= \int \overline{V(1,\tau)} = \overline{V(1,0)} e^{-S\tau}$$

Hence for PDE (5), the two Solution of ODE provide boundary values and no. B.C. are needed for the PDE to have unique solution.

Mence the final PDE becomes!

$$\frac{\partial \overline{V}}{\partial t} = \frac{1}{2} \overline{\tau_{(q)}^2} \frac{3^2 (1-q)^2}{3q^2} + (\gamma-5) \frac{3(1-q)}{3q} \frac{3\overline{V}}{3q} - [\gamma(1-q)+8q] \overline{V}$$

$$0 \le q \le 1, 0 \le \overline{T}$$

$$\overline{V}(\varsigma_1,0) = \frac{1-\varsigma_1}{\rho_m} V_{\overline{T}}(\frac{\rho_m \varsigma_1}{1-\varsigma_1}), 0 \leq \varsigma_1 \leq 1 \longrightarrow \overline{I} \cdot C$$

$$V(0,T) = V(0,0) e^{-\gamma T}$$

$$V(1,T) = V(1,0) e^{-8T}$$

Hence B.S. PDE is transformed from infinite Domain $S \gtrsim 0$ to finite domain $D \leq G \leq 1$.

(2)

American option pricing problem can be formulated as variational problem with class of comparison functions defined as?

 $K = \int_{0}^{\infty} v e^{-\frac{1}{2}t} \frac{\partial v}{\partial x} dx \text{ piecewise } e^{-\frac{1}{2}t}$ $V(\pi_{1}\tau) \geq g(\pi_{1}\tau) + \pi_{1}\tau$ $V(\pi_{1}0) = g(\pi_{1}\tau)$ $V(\pi_{max_{1}}\tau) = g(\pi_{max_{1}}\tau)$ $V(\pi_{min_{1}}\tau) = g(\pi_{min_{1}}\tau) \qquad y$ where $g(\pi_{1}\tau) = \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} e^{-\frac{1}{2}(\eta_{8}-1)} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} \int_{0}^{\infty} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} m_{4}x \int_{0}^{\infty} e^{-\frac{1}{2}(\eta_{8}-1)} \int_{0}^{\infty} v dx$ $e^{-\frac{1}{2}(\eta_{8}-1)} \frac{1}{2} + u_{4} \int_{0}^{\infty} v dx \int_{0}^{\infty}$

Now, let y be the exact solution of American ofton Pricing problem formulated as Linear Complimentarity Problem. As solution of Partial Differential Inequality, y is C^2 -smooth on Continuation region and $J \in K$.

Also for OER

Also from can complimentarity condition, we have

$$\int_{M_{min}}^{M_{max}} \left(\frac{\partial y}{\partial T} - \frac{\partial^2 y}{\partial x^2} \right) (y-y) dx = 0$$
 - (2)

Now Subtracting (2) from (1), we get

$$\int_{M_{min}}^{M_{max}} \left(\frac{\partial y}{\partial T} - \frac{\partial^2 y}{\partial x^2} \right) (U - y) dx > 0$$

=)
$$\int_{M_{min}}^{R_{max}} \frac{\partial y}{\partial t} (v-y) dx - \int_{M_{min}}^{M_{max}} (v-y) \cdot \frac{\partial^2 y}{\partial n^2} dx > 0$$

Now at boundary points ofmin and ofmax, we have

and V(Hmax It) = g(Hmax It)

three, we have.

Putting this in Egn (3), we get

$$I(y;v) = \int_{x_{min}}^{x_{max}} \left(\frac{\partial y}{\partial T} \cdot (v \cdot y) + \frac{\partial y}{\partial x} \left(\frac{\partial v}{\partial x} - \frac{\partial y}{\partial x} \right) \right) dx > 0 + v \in K$$

Now, this inequality holds $\forall v \in K$. Hence choosing v = y we have I(y;y) = 0 and the integral I(y;v) takes its minimum value i.e.

min
$$I(y;v) = I(y;y) = 0$$
 $v \in K$

(a) We now consider $\hat{y} \in K$. That is not \mathbb{Z}^2 -smooth on continuation region. We want to faid $\hat{y} \in K$ 8.+ $I(\hat{y};v) > 0$ + $v \in K$ | This is called weak formulation and inf $I(\hat{y};v) = 0$ | Weak formulation

Now, consider the following approximations for \hat{y} and v: $\hat{y} = \sum_{i} w_{i}(t) \, \hat{q}_{i}(x) \, \int_{y}^{y} the basis functions.$ $v = \sum_{i} v_{i}(t) \, \hat{q}_{i}(x) \, \int_{y}^{y} the basis functions.$

This formulation separates the inclipendent variables Z and X. there the same X-grid is applied for all Z, which there the same X-grid is applied for all Z, which results in sectangular grid in (X,Z)-plane. The time dependence is incorporated in coeff. functions wi and V_i . Since V_i superescut V_i -grid, we consider Semi-discretization of V_i as follows: $\int_{Z} \left(\sum_{i=1}^{N} d_i \right) \left(\sum_{i=1}^{N} (V_i - w_i) P_i \right) + \left(\sum_{i=1}^{N} w_i P_i \right) \left(\sum_{i=1}^{N} (V_i - w_i) P_i \right) P_i dX \neq 0$

= [= = dwi (v;-wi)] QiQ; dx + = = = wi(v;-w;)] QiQ; dx >0

Now writing the inequality into vector-mateix form, we have

$$\left(\frac{dw}{dt}\right)^T B(v-\omega) + \omega^T A(v-\omega) > 0$$

$$=) \left[(V-\omega)^{T} \left(B \frac{d\omega}{dt} + A\omega \right) > O \right] \rightarrow (\#)$$

where materies A and B are constructed using assembling algorithm 8.4.

$$a_{ij} = \int q_i q_j dx$$
 and $b_{ij} = \int q_i' q_j' dx$

Where D is domain.

For equidistant steps, we have for gias hat functions

$$A = \frac{1}{h} \begin{pmatrix} 1 & -1 \\ -1 & 2 & -1 \\ -1 & 2 \end{pmatrix}$$

$$B = \frac{h}{6} \begin{pmatrix} 2 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ 1 & 1 & 2 \end{pmatrix}$$

Now time is discretized.

Define
$$\omega^{(n)} = \omega(T_n)$$
 and $v^{(n)} = v(T_n)$

Now considering.

$$\frac{W_{i,n+1} - W_{in}}{\Delta \pi^{2}} = \frac{\partial \left(W_{i+1,n+1} - 2W_{i,n+1} + W_{i-1,n+1}\right) + \Delta \pi^{2}}{\Delta \pi^{2}}$$

$$\left(1-\partial\right) \left(W_{i+1,n} - 2W_{in} + W_{i-1,n}\right)$$

$$\Delta \pi^{2}$$

Now using this approximation in (#), we get

$$\left(v^{(n+1)} - \omega^{(n+1)} \right)^{T} \left(B \frac{1}{\Delta T} \left(\omega^{(n+1)} - \omega^{(n)} \right) + \partial A \omega^{(n+1)} + (1-\partial) A \omega^{(n)} \right) \geqslant 0$$

$$+ u$$

Now for 0= yz, we have crank-Nicolson method.

$$=) \left[\left(v^{(n+1)} - \omega^{(n+1)} \right)^{\mathsf{T}} \left(\left(B + \Delta \mathsf{T} \, \partial A \right) \omega^{(n+1)} + \left(\Delta \mathsf{T} \left(l - \theta \right) A - B \right) \omega^{(n)} \right) \geqslant 0 \right]$$



then we have.

$$\left(v^{(n+1)}-w^{(n+1)}\right)^{T}\left(cw^{(n+1)}-r\right)\geqslant0$$

which is fully discretized version of $\mathbb{Z}(\hat{y}; v)$ 70

Side Conditions

We have
$$\hat{y}(n,\tau) \geqslant g(n,\tau)$$

For hat functions
$$di$$
 ($di(\pi_i)=1$)
and $\pi=\pi_j$, we and $di(\pi_j)=0$ for $j\neq i$)

have $W_j(\tau) > g(\gamma_j, \tau)$.

then with T= In, we have

$$\left[\mathcal{W}^{(n)} \geqslant g^{(n)} \right]$$

(3)

For Showing equivalence of FEM and FDM, consider FEM and FDM ression American option pricing problem as follows:

FEM

Construct
$$\omega$$
 8.+.

 $\forall v \gg g$.

 $(v-\omega)^T(c\omega-r) \gg o$, $\omega \gg g$.

FDM

Find
$$W$$
 g, t :
$$Cw-r \geqslant 0, W \geqslant g$$

$$(cw-r)^{T}(w-g) = 0$$

the Proof of equivalence of FDM and FEM is as follows:

(FDM) => (FEM.)

ut w solve FDM, so w > g and.

$$(v-w)^{\mathsf{T}}(cw-r) = (v-g)^{\mathsf{T}}(cw-r) - (w-g)^{\mathsf{T}}(cw-r)$$

NOW, CW-770 and (W-8) (CW-r) =0

=)
$$(v-\omega)^{T}(c\omega-r)=(v-g)^{T}(c\omega-r)>0$$
 + $v>g$.

Lit w solve FEM, so wing sit.

Now suppose Kth component of CW-r is regative and make vx arbitrarily large. Then LHS becomes make vx arbitrarily large. Then LHS becomes arbitrarily small, which is contractiction

CW-770.

Now in (fem), we have.

$$(v-w)^{T}(cw-r) \neq 0$$

Put $v=g$, then $[(w-g)^{T}(cw-r) \leq 0] \rightarrow (\#)$

Using (*) and (#), we conclude that

$$\left[\left(\omega - g \right)' \left(c \omega - r \right) = 0 \right]$$

Hence N solves FDM.

combining both parts, we have FEM and FDM formulation are equivalent i.e. solution of one toroblem Batisfies another as well.