A Variant of the Kochen-Specker Theorem Locating Value Indefiniteness

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Between Kochen-Specker and Quantum Indeterminism

Kochen-Specker Theorem

- ▶ Impossibility of a (complete) noncontextual hidden variable theory for states in $d \ge 3$ Hilbert space
- "Quantum contextuality"
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There is a gap between this and the common interpretation of quantum measurements

Eigenvalue-Eigenstate link

A system in a state $|\psi\rangle$ has a definite property of an observable A if and only if $|\psi\rangle$ is an eigenstate of A.

Can we formally close this gap and show the *extent* of value indefiniteness?

The Kochen-Specker Theorem

A *context* in \mathbb{C}^n is a set of *n* compatible (commuting) observables.

In $n \ge 3$ Hilbert space there is a finite set of (one-dimensional projection) observables $\mathcal O$ such that

the following three are in contradiction:

- Every observable is assigned a definite value of 0 or 1;
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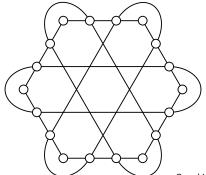
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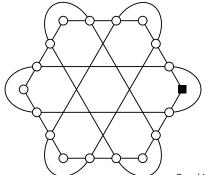
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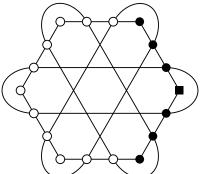
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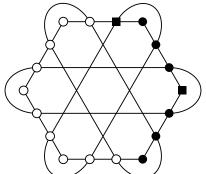
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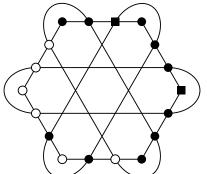
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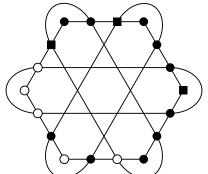
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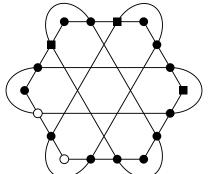
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The Extent of Value Indefiniteness

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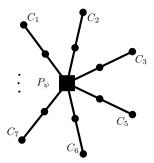
Either, we reject:

- QM: But then we depart from quantum theory;
- NC: Definite values depends on measurement context;
- VD: Some observables are value indefinite.

If we insist that *value definite* observables behave noncontextually, then only guaranteed that *some* observables are value indefinite

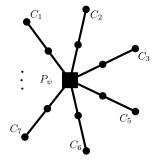
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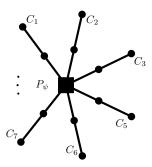


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 - VD: Every observable is assigned a defined value.
 - VD': One observable is assigned a defined value.
 - VD": An observable *P* assigned 1, and a *non-compatible* observable *P'* value definite.
- If system is in state $|\psi\rangle$, reasonable to expect $\nu(P_{\psi})=1$.
 - One direction of eigenvalue-eigenstate link.
- ► Intuitively, expect everything outside this 'star' to be value indefinite.
- ► Need to localise all the assumptions if we wish to go further formally.



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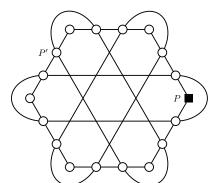
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Admissibility of v

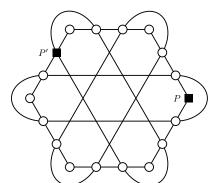
A value assignment function v is admissible if for every context $C \subset \mathcal{O}$:

- (a) if there exists a $P \in C$ with v(P) = 1, then v(P') = 0 for all $P' \in C \setminus \{P\}$;
- (b) if there exists a $P \in C$ with v(P') = 0 for all $P' \in C \setminus \{P\}$, then v(P) = 1.

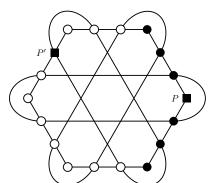
Admissibility provides a way of deducing the value definiteness of observables.



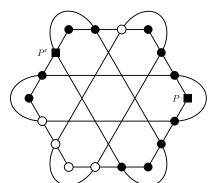
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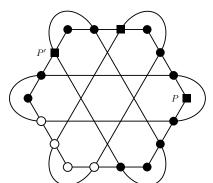
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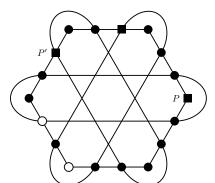
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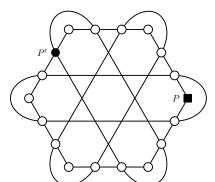
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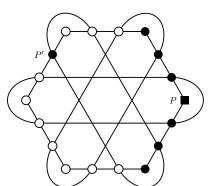


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- Need more careful interlinking of observables to obtain a contradiction.
- ► To prove for all P' not compatible with P we either
 - (a) Need to consider \mathcal{O} as the set of all projectors on \mathbb{C}^n ;
 - (b) Give a procedure to find $\mathcal{O} = \mathcal{O}(P')$ for a given P'.



Localised Value Indefiniteness: A Theorem

Theorem

Let $n \geq 3$. If an observable P on \mathbb{C}^n is assigned the value 1, then no other incompatible observable can be consistently assigned a definite value at all - i.e., is value indefinite.

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Let $n \geq 3$ and $|\psi\rangle, |\phi\rangle \in \mathbb{C}^n$ be states such that $0 < |\langle \psi | \phi \rangle| < 1$. Then there is a finite set of observables $\mathcal O$ containing P_ψ and P_ϕ for which there is no admissible value assignment function on $\mathcal O$ such that $v(P_\psi) = 1$ and P_ϕ is value definite.

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We prove in 3 steps:

- 1. We first prove the explicit case that $|\langle \psi | \phi \rangle| = \frac{1}{\sqrt{2}}$
- 2. We prove a reduction for $0 < |\langle \psi | \phi \rangle| < \frac{1}{\sqrt{2}}$ to the first case.
- 3. We prove a reduction for the last case of $\frac{1}{\sqrt{2}} < |\langle \psi | \phi \rangle| < 1$ case.



Interpretation and Conclusion

To interpret physically, need to connect P_{ψ} to the state of a system

Eigenstate value definiteness

If a system S is in a state $|\psi\rangle$, then $v_S(P_\psi)=1$ for the value assignment function v_S modelling the system.

Can then conclude:

- For any state $|\psi\rangle$ can locate precisely the value indefinite observables
- Almost all observables are value indefinite (measure-one set)

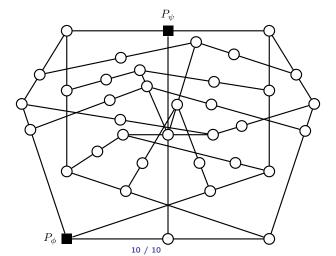
General Interpretation

If a system is in a state $|\psi\rangle$, then the result of measuring an observable A is indeterministic unless $|\psi\rangle$ is an eigenstate of A.

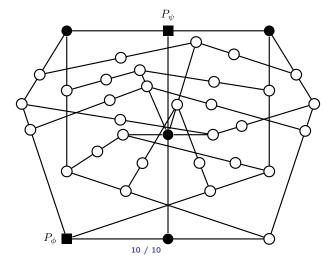
Thank You!

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- A. A. Abbott, C. S. Calude, J. Conder & K. Svozil. Strong Kochen-Specker theorem and incomputability of quantum randomness. PRA, 86(062109), 2012

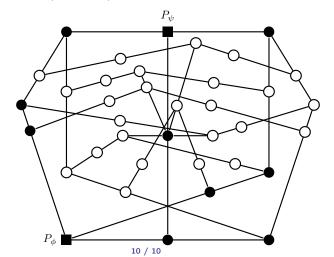
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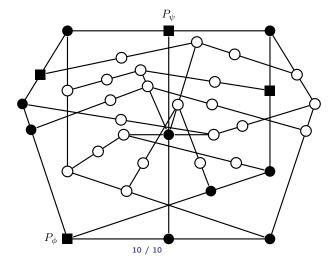
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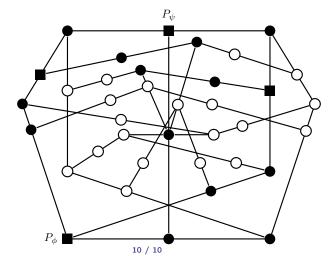
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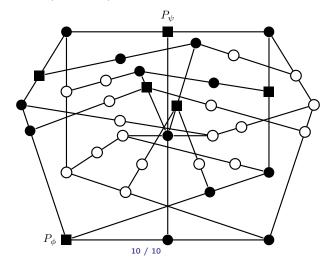
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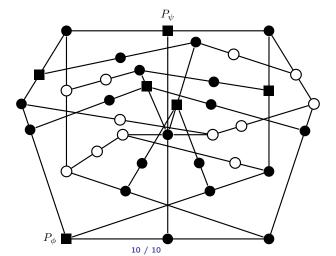
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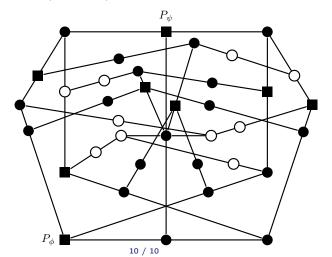
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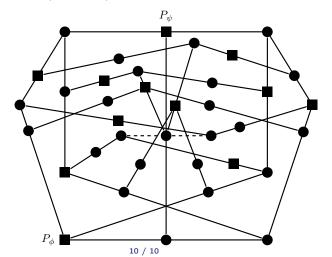
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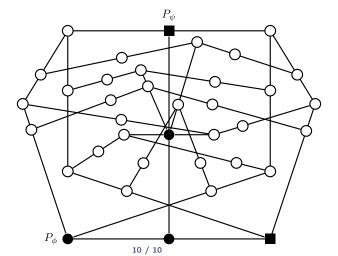


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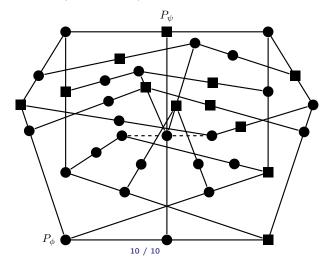
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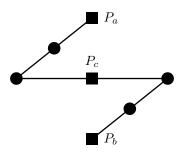
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First Reduction: Contraction (Back)

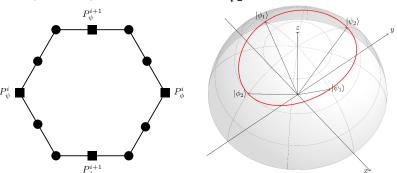
If $v(P_{\psi}) = v(P_{\phi}) = 1$ and $0 < |\langle \psi | \phi \rangle| < \frac{1}{\sqrt{2}}$, then we can find a $|\phi'\rangle$ with $\langle \psi | \phi' \rangle = \frac{1}{\sqrt{2}}$ and $v(P_{\phi'}) = 1$ under any admissible v.



The above diagram is realisable for $|\langle a|b\rangle|<|\langle a|c\rangle|<1.$

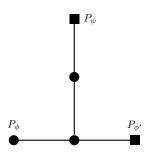
Second Reduction: Expansion (Back)

If $v(P_{\psi}) = v(P_{\phi}) = 1$ and $\frac{1}{\sqrt{2}} < |\langle \psi | \phi \rangle| < 1$, then we can find a finite sequence of states $|\psi_1\rangle, |\phi_1\rangle; \cdots, |\psi_n\rangle, |\phi_n\rangle$ such that for all i $v(P_{\psi}^i) = v(P_{\phi}^i) = 1$ and $\langle \psi_n | \phi_n \rangle = \frac{1}{\sqrt{2}}$ under any admissible v.



Completing the Proof (Back)

For the reductions we assumed $v(P_{\phi})=1$. If $v(P_{\phi})=0$, we can easily find $|\phi'\rangle$ with $v(P_{\phi'})=1$ and apply the reasoning above.



As a consequence, the set of value indefinite observables has measure 1: almost all observables are value indefinite.