

1 Coexistence of alleles: insights of Modern
2 Coexistence Theory into the maintenance of
3 genetic diversity

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1 Introduction

The question of how genetic variation is maintained, despite the effects of selection and drift, continues to be central to the study of evolutionary biology (Walsh & Lynch, 2018). Classical explanations include overdominance (heterozygote advantage) or frequency-dependent selection, but in the modern era of genomic data, all patterns of variation that exceed the expected variation under neutrality tend to be categorized broadly as balancing selection, regardless of the evolutionary mechanism (Mitchell-Olds *et al.*, 2007). One of the evolutionary mechanisms coined under balancing selection is sexually antagonistic selection, which occurs when the direction of natural selection on traits or loci differs between the sexes (Connallon & Hall, 2018).

Sexually antagonistic selection has been identified as a powerful engine of speciation that in some cases can maintain polymorphisms of otherwise dis-advantageous alleles in a population (Gavrilets, 2014). The effect of sexually antagonistic selection, however, has been generally studied under strong simplifying assumptions such as constant population sizes and homogeneous environments (e.g., Kidwell *et al.* (1977); Pamilo (1979); Immler *et al.* (2012)). Few studies have explored the effect of sexually antagonistic selection on the maintenance of polymorphism with more realistic assumptions. Exceptions include Connallon *et al.* (2018) who found that classical predictions break down when fluctuations in the environment combined with life-history traits allow local adaptations and promote the maintenance of genetic diversity. The effect of environmental fluctuations without local adaptation, however, has not been studied in the context of sexually

antagonistic selection to the best of our knowledge.

The contribution of environmental fluctuations to genetic variability remains a debated issue in evolutionary biology. Classic theoretical models predict that temporal fluctuations in environmental conditions are unlikely to maintain a genetic polymorphism (Hedrick, 1974; 1986). However, other studies have found that fluctuating selection can maintain genetic variance on sex-linked traits (Reinhold, 2000), or in populations where generations overlap (Ellner & Hairston Jr, 1994; Ellner & Sasaki, 1996). Similarly, temporal changes in population sizes have been shown to mitigate the effect of genetic drift in small populations (Pemberton *et al.*, 1996), and in annual plant systems (Nunney, 2002). Thus, both fluctuations in selection and population sizes could dramatically change the effect of sexually antagonistic selection in the maintenance of genetic diversity.

Importantly, progress requires more than just identifying if fluctuations can maintain genetic diversity in a population, but to quantify how exactly they contribute to its maintenance (Ellner *et al.*, 2016). Modern coexistence theory (MCT) provides a powerful conceptual framework to do so (Chesson, 2000b; 1994; Barabás *et al.*, 2018). Although its core ideas were formalized in an ecological context (Chesson, 1994; 2000a), this framework provides the necessary tools to examine the relative contributions of fluctuations to diversity maintenance, which can also be applied to evolutionary contexts (Ellner & Sasaki, 1996; Reinhold, 2000). From an ecological perspective, polymorphism of sexually antagonistic alleles is equivalent to the coexistence of species, and the fixation of either one of the alleles in a population is equivalent to competitive exclusion. The coexistence of alleles, thus, can be examined through the same lens as the coexistence of competing species.

Here, we seek to explicitly apply recent advances in MCT to the question of how polymorphism is maintained under sexually antagonistic selection. We examined how fluctuations in selection values, fluctuations in population sizes, and their interactions can stabilize or hinder the coexistence of alleles. In particular, we examined i) Can fluctuations in population sizes and selection values allow sexually antagonistic alleles to coexist when differences in their fitness would typically not allow them to? and ii) What is the relative contribution of different types of fluctuations that allow two sexually antagonistic alleles to be maintained in a population? Our study provides the tools to analyze evolutionary dynamics from a novel perspective and contributes to answering long-lasting questions regarding the effect of non-constant environments on genetic diversity.

2 Methods

We first present a model that describes the evolutionary dynamics of sexually antagonistic alleles and show how changes in allele frequencies can be expressed in terms of growth rates, a necessary condition for analyses done using MCT. We continue by simulating different scenarios of alleles invading a population, where we allowed population sizes, selection values, both, or neither to vary. Finally, we examine the results of our simulations through a MCT lens by calculating the contribution of each of these fluctuations in the coexistence of alleles across the parameter space of sexually antagonistic selection.

Population dynamics of sexually antagonistic alleles

As most population genetic models of sex-dependent selection, our model considered evolution at single, biallelic locus with frequency and density independent effects on the relative fitness of females and males (Wright, 1942; Kidwell *et al.*, 1977; Immler *et al.*, 2012).[†] We examined the dynamics of two sexually antagonistic alleles, j and k , that affect fitness in the haploid state. We assumed allele j always has a high fitness in females ($w_{jf} = 1$), but variable fitness in males ($w_{jm} < 1$); and allele k has a high fitness in males ($w_{km} = 1$), but variable fitness in females ($w_{kf} < 1$). The selection against allele j in males is therefore $S_m = 1 - w_{jm}$, and the selection against allele k in females is $S_f = 1 - w_{kf}$.

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The frequency of each allele in each sex at the beginning of a life-cycle at time t is given by:

$$p_{jm,t} = \frac{n_{jm,t}}{N_{m,t}} \quad (1)$$

$$p_{jf,t} = \frac{n_{jf,t}}{N_{f,t}} \quad (2)$$

$$p_{km,t} = \frac{N_{m,t} - n_{jm,t}}{N_{m,t}} \quad (3)$$

$$p_{kf,t} = \frac{N_{f,t} - n_{jf,t}}{N_{f,t}} \quad (4)$$

where $N_{m,t}$ and $N_{f,t}$ are the numbers of males and females in a population at time t , $n_{jf,t}$ is the number of females f with allele j , and $n_{jm,t}$ is the number of males m with allele j at time t , respectively.

The individuals in the population mate at random before selection occurs, and there-

fore the frequency of offspring with allele j after mating, $p'_{j,t}$ can be expressed as:

$$p'_{j,t} = \frac{(N_{m,t}n_{jf,t} + N_{f,t}n_{jm,t})}{2N_fN_m}. \quad (5)$$

Selection acts upon these offspring in order to determine the allelic frequencies in females and males in the next generation, $t + 1$. As an example the frequency of females with allele j after selection is given by:

$$p'_{jf,t+1} = \frac{n_{jf,t+1}}{N'_{f,t+1}} = \frac{p'_j w_{jf}}{p'_j w_{jf} + (1 - p'_j) w_{kf}} \quad (6)$$

The logarithmic growth rate of j in females, is therefore given by the number of females with allele j after selection, divided by the original number of females carrying allele j :

$$r_{jf,t} = \ln \left(\frac{n'_{jf,t+1}}{n_{jf,t}} \right) \quad (7)$$

An equivalent expression for the per capita growth rate of allele j in males m can be obtained by exchanging f for m across the various subscripts in this expression.

Allelic coexistence in a sexual population, however, is ultimately influenced by growth and establishment of an allele across both sexes. Therefore, the full growth rate of allele j across the entire population of females *and* males is given by:

$$r_j = \ln \left(\frac{n'_{jf,t+1} + n'_{jm,t+1}}{n_{jf,t} + n_{jm,t}} \right). \quad (8)$$

100 An equivalent expression describes r_k , the growth rate of allele k .

101 Selection maintains both alleles in the population under the condition that:

$$\frac{S_m}{1 + S_m} < S_f < \frac{S_m}{1 - S_m} \quad (9)$$

102 Thus, the maintenance of polymorphism of sexually antagonistic alleles is solely deter-
103 mined by the values of S_m and S_f . Note that in our model, the values S_m and S_f can take
104 are bounded from 0 to 1. Therefore the parameter space of sexually antagonistic selection
105 is within the range $0 < S_m, S_f < 1$. Classic theoretical models predict that in constant
106 environments, only in ≈ 0.38 of the selection parameter space alleles can coexist (Kidwell
107 *et al.*, 1977; Pamilo, 1979; Connallon *et al.*, 2018). If fluctuations in population sizes or se-
108 lection values have an effect on the coexistence of sexually antagonistic alleles, it would
109 be reflected in increases or decreases of the proportion of the parameter space of selection
110 where polymorphism is maintained.

111 Simulations

112 Typically, MCT would require decomposing alleles' growth rates (e.g., Eqn. 8) analytically
113 to examine the relative contributions of different types of fluctuations to their coexistence
114 (Barabás *et al.*, 2018). Although we present an analytical approach in the Supporting Infor-
115 mation, our general solution is not easily interpretable and soon becomes mathematically
116 intractable (S1 Supporting Information). Thus, we opted for an extension of MCT that
117 provides the flexibility to examine the contributions of different processes to coexistence
118 using simulations (Ellner *et al.*, 2019; Shoemaker *et al.*, 2020).

For each simulation, we examined coexistence outcomes across the selection parameter space of sexually antagonistic selection ($0 < S_m, S_f < 1$). To do so, we partitioned the parameter space into a grid of 50×50 , which yielded 2500 pairwise combinations of different w_{jm} and w_{kf} values. For each pairwise combination of w_{jm} and w_{kf} , as we detail in the next sections, we iterated our model while controlling the effect size of fluctuations in fitness values (σ_w), fluctuations in population sizes (σ_g) and their correlations (ρ_w and ρ_g respectively). Then, we performed simulations of each allele invading a population, determined coexistence outcomes, and the relative contribution of each type of fluctuation. Finally, we calculated for each simulation, the proportion of the parameter space that allowed alleles to coexist.

We explored all of the combinations of low, intermediate and high fluctuations in fitness values and population sizes, with different extents of correlations between fluctuations (Table 1). As a control simulation, we set $\sigma_w = 0.001$ and $\sigma_g = 0.001$, with no correlation between fluctuations. For each one of the factorial combinations of σ_g , σ_w , ρ_g and ρ_w (Table 1), we performed invasion simulations across the parameter space of selection. We ran ten replicates per parameter combination, which resulted in 3780 simulations.

Timeseries

To incorporate the effects of fluctuations into our population dynamics model we generated independent timeseries of fluctuations in fitness values and population sizes. In the case of fluctuations in selection values, for a given value of w_{jm} and w_{kf} (i.e., a fixed point in the selection parameter space), we generated a timeseries of 500 timesteps made

up of correlated fluctuations of w_{jm} and w_{kf} . We controlled the effect size of fluctuations in fitness values (σ_w) and its correlation (ρ_w) by using the Cholesky factorization of the variance-covariance matrix:

$$C_w = \begin{bmatrix} \sigma_w^2 & \rho_w \sigma_w^2 \\ \rho_w \sigma_w^2 & \sigma_w^2 \end{bmatrix} \quad (10)$$

We multiplied Eqn. 10 by a (2×500) matrix of random numbers from a normal distribution with mean 0 and unit variance, which yielded γ_j and γ_k . Then, we calculated the value of w_{jm} at time $t + 1$ as $w_{jm,t+1} = w_{jm}^{\gamma_{j,t}}$. We calculated the value of $w_{kf,t+1}$ analogously.

Similarly, we generated a timeseries of 500 timesteps made up of correlated fluctuations in population sizes. We chose values of N_m and N_f of 200 individuals each as the initial value of population sizes throughout our simulations. We performed a Cholesky factorization of the variance-covariance matrix, controlling the effect size of fluctuations in population sizes with σ_g and their correlation with ρ_g . Similar to our previous approach, we multiplied this factorization by a random matrix of uncorrelated random variables, which yielded γ_m and γ_f . Finally, we calculated the number of males in the population at time $t + 1$ as $N_{m,t+1} = N_m + \gamma_{m,t}$. We calculated the value of $N_{f,t+1}$ analogously. We bounded the values population sizes could take so there were no negative population sizes, since that would not be biologically plausible. We did not impose an upper bound to the values population sizes could take.†

Finally, we performed simulations where our population dynamics model (Eqns. 1

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to 8) iterated over 500 timesteps while allowing selection values and population sizes to fluctuate in each timestep. We started each simulation with the initial values of N_m and N_f described before and equal frequencies of allele j and allele k in each sex. For each timestep t in our simulations, the values of w_{jm} , w_{kf} , N_m and N_f used to calculate allele's frequencies in timestep t (e.g., Eqn. 6), corresponded to the t values calculated in each timeseries, as described previously. This approach yielded a final timeseries that captured the dynamics of sexually antagonistic alleles, with fluctuating values of selection and population sizes.

Invasion simulations

Modern coexistence theory has shown that coexistence is promoted by mechanisms that give species a population growth rate advantage over other species when they become rare (Chesson, 1982; 2003; Barabás *et al.*, 2018). Typically, one species is held at its *resident* state, as given by its steady-state abundances while the rare species is called the *invader*. In the context of alleles in a population, an allele is an *invader* when a mutation occurs that introduces that allele into a population in which it is absent (e.g., if in a population with only k alleles, a random mutation made one individual carry the j allele). Within sexually antagonistic selection, each allele has two pathways of invasion, depending on whether the mutation arises in a female or in a male. If an allele's *invasion growth rate* (or the average instantaneous population growth rate when rare) is positive, it buffers it against extinction, maintaining its persistence in the population. Coexistence, and hence polymorphism, occurs when both alleles have positive invasion growth rates.

To study the dynamics of sexually antagonistic alleles through this framework, we used the timeseries that captured the dynamics of our population model as a template to perform invasion simulations of both alleles. We allowed each allele to invade via two different pathways: males and females. We explored all potential combinations of each allele invading through a different pathway (e.g., allele j invading through males, and allele k invading through females, and so on). This yielded four types of invasion.

For each timestep in the timeseries, we performed simulations of the two alleles invading separately via their respective pathway. To simulate invasion, we set the density of the invading allele to one individual. For example, if allele j was invading via males, then we would set $n_{jm,i} = 1$ and $n_{jf,i} = 0$. Note that each invasion simulation was independent of the iteration that we used to generate the timeseries, therefore we denoted the initial timestep in an invasion simulation with the subscript i . We also set the resident allele, in this case k , to the corresponding value of the timeseries minus one individual, $n_{km,i} = N_{m,t} - 1$ and $n_{kf,i} = N_{f,t}$. Then, we iterated our model one timestep, $i + 1$, and calculated the logarithmic growth rate of j allele invading as:

$$r_j = \ln \left(\frac{n_{jm,i+1} + n_{jf,i+1}}{1} \right) \quad (11)$$

Correspondingly, the logarithmic growth rate of the k allele as a resident would be given by:

$$r_k = \ln \left(\frac{n_{km,i+1} + n_{kf,i+1}}{n_{km,i} + n_{kf,i}} \right) \quad (12)$$

We treated each timestep of the timeseries independently, and hence we performed

500 invasion simulations. We then calculated, for each allele invading via a different pathway, its mean invasion growth rate as the average of the 500 invasion growth rates. We also calculated the mean growth rate of the resident allele as the average of the 500 resident growth rates. We determined alleles to be coexisting if both of alleles had positive mean invasion growth rates, which is often referred to as the mutual invasibility criterion (Barabás *et al.*, 2018).

Functional decomposition

Our invasion simulations tell us whether or not sexually antagonistic alleles can coexist in a determined point of the selection parameter space. However, we also quantified the relative contributions of fluctuating selection and population sizes into the predicted coexistence outcome. Therefore, we used an extension of MCT that provides the flexibility to analyze the contributions of different processes to coexistence using *functional decomposition* (Ellner *et al.*, 2016; 2019; Shoemaker *et al.*, 2020).

We applied the functional decomposition approach by breaking up the average growth rate of each allele into a null growth rate in the absences of fluctuations in all selected variables, a set of main effect terms that represent the effect of only one variable fluctuating, and a set of two-way interaction terms representing the effect of variables fluctuating simultaneously (Ellner *et al.*, 2019). In our simulations, this is a function of four variables: the number of males in the population (N_m), the number of females in the population (N_f), the fitness of allele j in males (w_{jm}), and the fitness of allele k in females (w_{kf}). As an example, if only N_m and N_f were fluctuating, the growth rate of allele j when it is the

invader at timestep t could be decomposed into:

$$r_{j,t}(N_m, N_f) = \mathcal{E}_j^0 + \mathcal{E}_j^{N_m} + \mathcal{E}_j^{N_f} + \mathcal{E}_j^{N_m N_f} \quad (13)$$

Where \mathcal{E}^0 is the null growth rate when N_m and N_f are set to their averages. Terms with superscripts represent the marginal effects of letting all superscripted variables vary while fixing all the other variables at their average values. For example, the term \mathcal{E}^{N_m} expresses the contribution of fluctuations in N_m when N_f is at its average, without the contribution when both variables are set to their averages :

$$\mathcal{E}_j^{N_m} = r_{j,t}(N_m, \overline{N_f}) - \mathcal{E}_j^0 \quad (14)$$

If we average Eqn. 13 across the timesteps in our simulation, we get a partition of the average population growth rate into the variance-free growth rate, the main effects of variability in N_m , the main effects of variability in N_f , and the interaction between variability in N_m and N_f

$$\bar{r}_j = \mathcal{E}_j^0 + \overline{\mathcal{E}_j^{N_m}} + \overline{\mathcal{E}_j^{N_f}} + \overline{\mathcal{E}_j^{N_m N_f}} \quad (15)$$

However, in our simulations w_{jm} and w_{kf} also fluctuated, therefore the full functional decomposition of the growth rate of allele j as an invader is found in Table 2, as well as a brief description of the meaning of each term. The implementation and interpretation of the functional decomposition of the invasion growth rates of each allele are identical to each other. We calculated the value of each of the terms in Table 2 by performing another

234 set of invasion simulations as described previously, but instead of allowing all variables
 235 to fluctuate, systematically setting the required variables to their means and subtracting
 236 the corresponding \mathcal{E} values.

237 The functional decomposition approach further requires the *comparison* of each term,
 238 to understand if how it affects invaders and residents. This is because fluctuations can
 239 promote coexistence by helping whichever allele is rare, or they can hurt whichever allele
 240 is common. Therefore, to understand the role of each type of fluctuation, it is necessary
 241 to compare how it affects invader *and* resident growth rates. In the example presented
 242 in Eqn. 15, if allele j is invading, then allele k is at it's resident state and there exists an
 243 analogue decomposition of \bar{r}_k with the exact same terms. Therefore we can express the
 244 difference between contributions of fluctuations in N_m as:

$$\Delta_j^{N_m} = \bar{\mathcal{E}}_j^{N_m} - \bar{\mathcal{E}}_k^{N_m} \quad (16)$$

245 If $\Delta_j^{N_m}$ is positive, then fluctuations in the male population benefit allele j when it is
 246 rare more than what they benefit k as a resident. If $\Delta_j^{N_m}$ is negative, then fluctuations
 247 benefit k as a resident more than j as an invader, and if it is minimal, then fluctuations
 248 have an equal effect in j and k . Therefore, for each allele invading via a different pathway,
 249 we calculated 16 Δ values, one for each one of the \mathcal{E} terms in Table 2. However, since the
 250 magnitude of each one of these values could vary considerably across the parameter space
 251 of selection, to make them comparable, we normalized each Δ value by dividing it by the
 252 square root of the sum of the squares of the 16 Δ values. For example, the normalized

253 value of Eqn. 16 would be given by:

$$\delta_j^{N_m} = \frac{\Delta_j^{N_m}}{\sqrt{\sum_{d=1}^{16} (\Delta_d)^2}} \quad (17)$$

254 This normalization bounded δ values from -1 to 1 .

255 3 Results

256 Our simulations showed that fluctuations in selection and population sizes could both
257 increase the proportion of allelic coexistence in the parameter space compared to classic
258 theoretical expectations (Fig. 1A and B). As a baseline, we show in Fig. 1C the outcome
259 of the control simulation, which matches previous findings that without fluctuations, al-
260 leles can coexist in only ≈ 0.38 of the selection parameter space (Connallon & Hall, 2018).
261 Importantly, we also found that the extent and relative contribution of each type of fluc-
262 tuation, differed from each other.

263 When *only* population sizes fluctuated,

Figures and tables

Table 1: Parameters used in our simulations to control the effect size of fluctuations in population sizes (σ_g) and selection values σ_w , as well as their respective correlations (ρ_g and ρ_w).

Parameter	Values	Description
σ_w	0.001, 0.1, 0.3, 0.5, 0.7, 0.9	Effect size of fluctuations in fitness values
σ_g	0.001, 1, 10, 20, 30, 50	Effect size of fluctuations in population sizes
ρ_w	-0.75, 0, 0.75	Correlation between fluctuations in fitness values
ρ_g	-0.75, 0, 0.75	Correlation between fluctuation in population sizes

Table 2: Functional decomposition of the growth rate of allele j .

Term	Formula	Meaning
\mathcal{E}_j^0	$\bar{r}_j(\bar{N}_m, \bar{N}_f, \bar{w}_{jm}, \bar{w}_{kf})$	Growth rate at mean population size and fitness values.
$\bar{\mathcal{E}}_j^{N_m}$	$\bar{r}_j(N_m, \bar{N}_f, \bar{w}_{jm}, \bar{w}_{kf}) - \mathcal{E}_j^0$	Main effect of fluctuations in N_m
$\bar{\mathcal{E}}_j^{N_f}$	$\bar{r}_j(\bar{N}_m, N_f, \bar{w}_{jm}, \bar{w}_{kf}) - \mathcal{E}_j^0$	Main effect of fluctuations in N_f
$\bar{\mathcal{E}}_j^{w_{jm}}$	$\bar{r}_j(\bar{N}_m, \bar{N}_f, w_{jm}, \bar{w}_{kf}) - \mathcal{E}_j^0$	Main effect of fluctuations in w_{jm}
$\bar{\mathcal{E}}_j^{w_{kf}}$	$\bar{r}_j(\bar{N}_m, \bar{N}_f, \bar{w}_{jm}, w_{kf}) - \mathcal{E}_j^0$	Main effect of fluctuations in w_{kf}
$\bar{\mathcal{E}}_j^{N_m, N_f}$	$\bar{r}_j(N_m, N_f, \bar{w}_{jm}, \bar{w}_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{N_f}]$	Interaction of fluctuations in N_m and N_f
$\bar{\mathcal{E}}_j^{w_{jm}, w_{kf}}$	$\bar{r}_j(\bar{N}_m, \bar{N}_f, w_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{w_{jm}} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations in w_{jm} and w_{kf}
$\bar{\mathcal{E}}_j^{N_m, w_{jm}}$	$\bar{r}_j(N_m, \bar{N}_f, w_{jm}, \bar{w}_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{w_{jm}}]$	Interaction of fluctuations in N_m and w_{jm}
$\bar{\mathcal{E}}_j^{N_m, w_{kf}}$	$\bar{r}_j(N_m, \bar{N}_f, \bar{w}_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations in N_m and w_{kf}
$\bar{\mathcal{E}}_j^{N_f, w_{jm}}$	$\bar{r}_j(\bar{N}_m, N_f, w_{jm}, \bar{w}_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{jm}}]$	Interaction of variation in N_f and w_{jm}
$\bar{\mathcal{E}}_j^{N_f, w_{kf}}$	$\bar{r}_j(\bar{N}_m, N_f, \bar{w}_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations N_f and w_{kf}
$\bar{\mathcal{E}}_j^{N_m, w_{jm}, w_{kf}}$	$\bar{r}_j(N_m, \bar{N}_f, w_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{w_{jm}} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations in N_m , w_{jm} , and w_{kf}
$\bar{\mathcal{E}}_j^{N_f, w_{jm}, w_{kf}}$	$\bar{r}_j(\bar{N}_m, N_f, w_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{jm}} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations in N_f , w_{jm} , and w_{kf}
$\bar{\mathcal{E}}_j^{N_m, N_f, w_{jm}}$	$\bar{r}_j(N_m, N_f, w_{jm}, \bar{w}_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{jm}}]$	Interaction of variation in N_m , N_f , and w_{jm}
$\bar{\mathcal{E}}_j^{N_m, N_f, w_{kf}}$	$\bar{r}_j(N_m, N_f, \bar{w}_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of fluctuations in N_m , N_f , and w_{kf}
$\bar{\mathcal{E}}_j^{N_m, N_f, w_{jm}, w_{kf}}$	$\bar{r}_j(N_m, N_f, w_{jm}, w_{kf}) - [\mathcal{E}_j^0 + \bar{\mathcal{E}}_j^{N_m} + \bar{\mathcal{E}}_j^{N_f} + \bar{\mathcal{E}}_j^{w_{jm}} + \bar{\mathcal{E}}_j^{w_{kf}}]$	Interaction of variation in N_f , N_m , w_{jm} , and w_{kf}

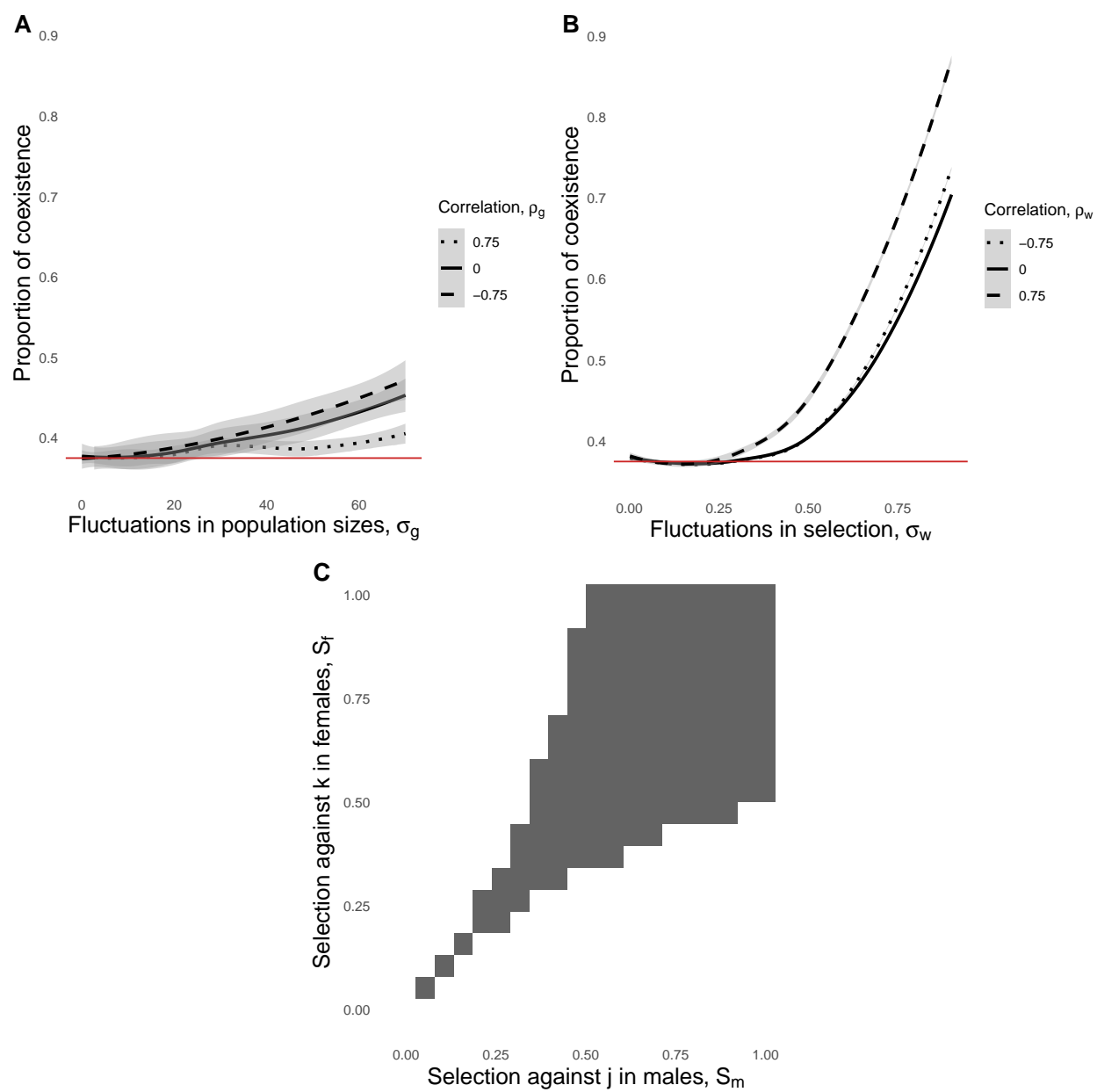


Figure 1

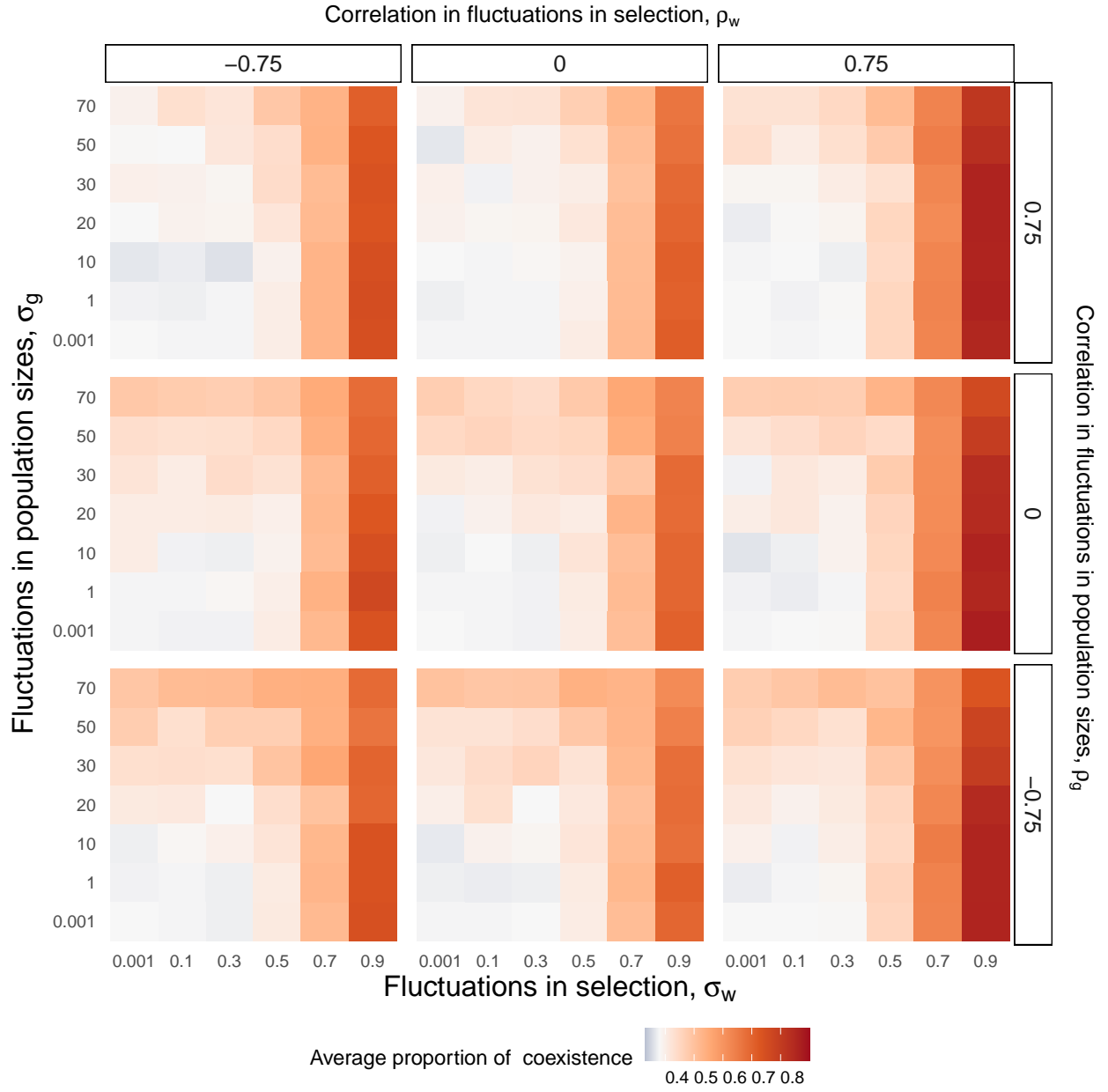
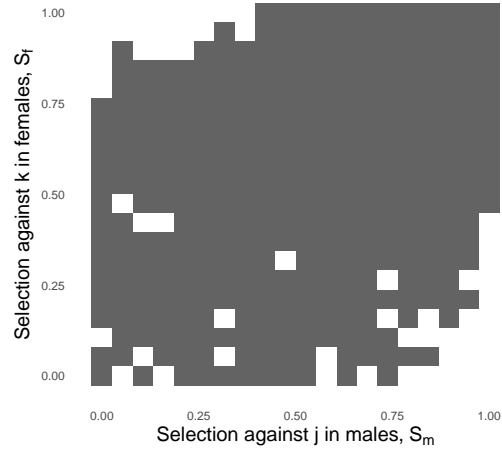


Figure 2

A



B

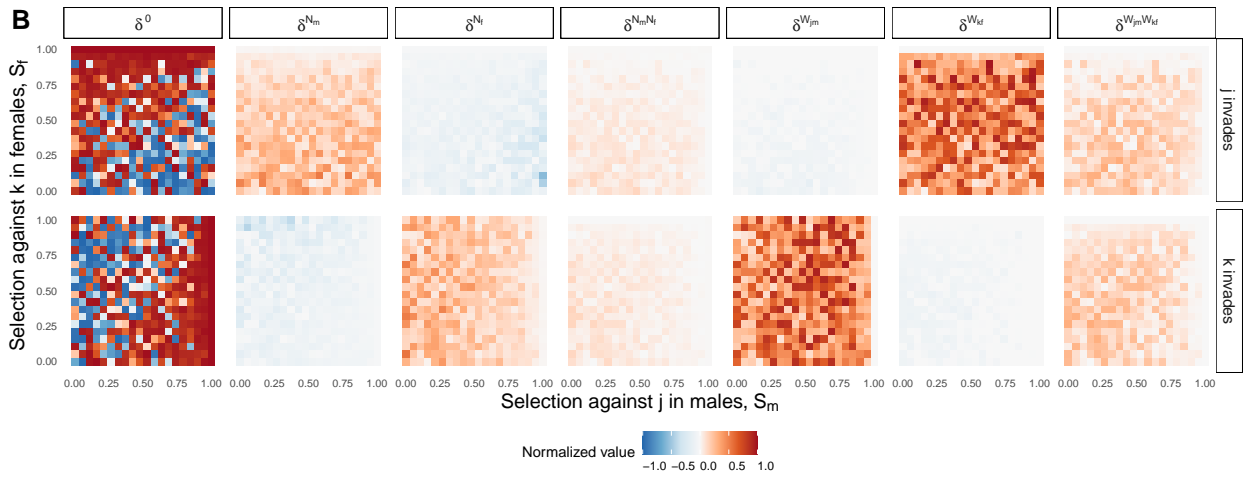


Figure 3

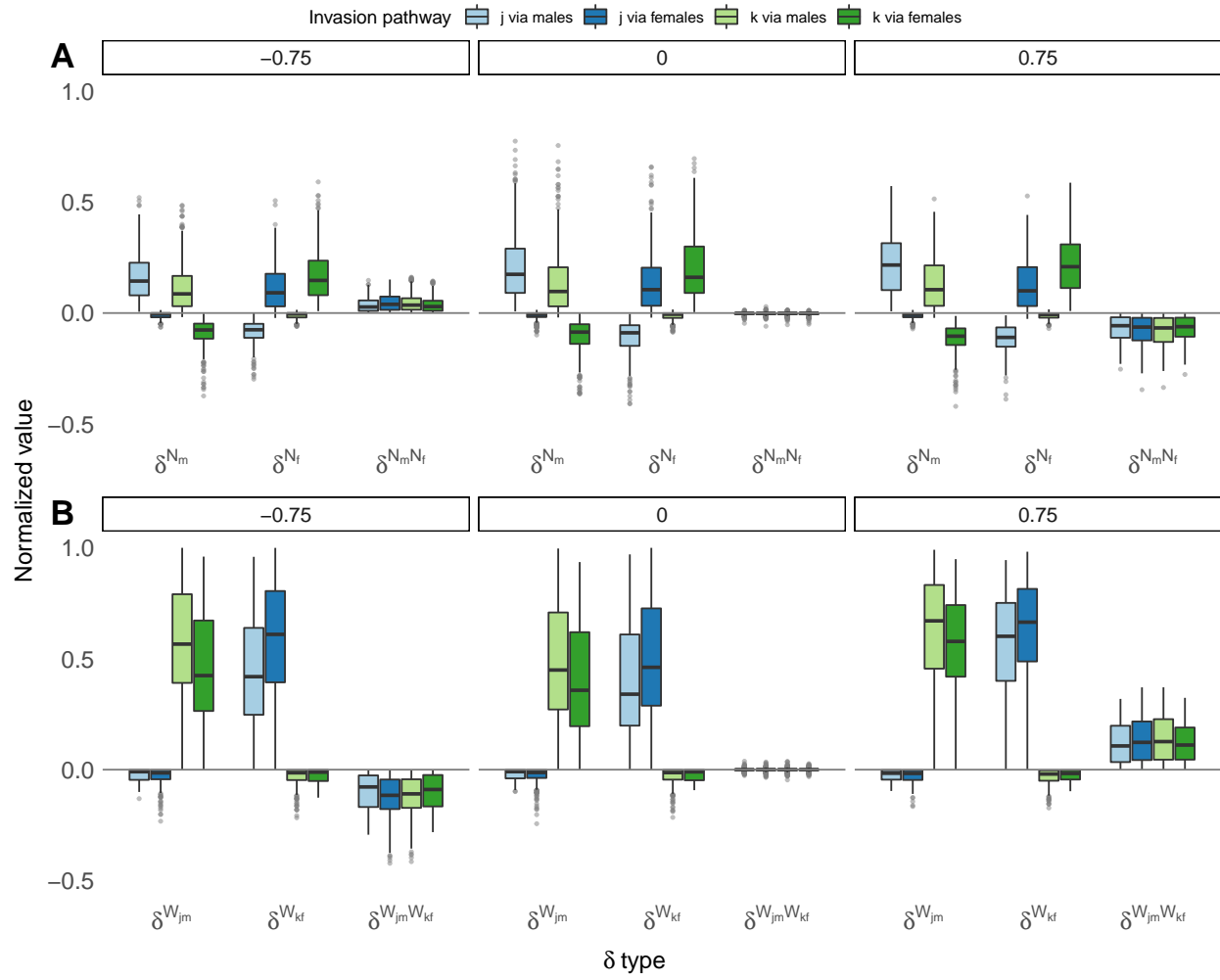


Figure 4

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