

Theory

In this section we will examine the necessary in understanding the michelson-morley interferometer. At its most basic level, we are interested in understanding what happens when light waves collide. In this experiment we will be assuming that the colliding are identical in all aspects except phase. Their wavelengths and frequencies are identical. Let us assume that our light wave moves along the optical axis, it may then be described as,

$$\mathbf{E}_i = E_0 \cos(\omega t - kx).$$

Where ω is the frequency, k the wave number and E_0 the amplitude of the wave. When the wave is measured, it has been transmitted and reflected once. We therefore multiply the wave amplitude by the coefficients of transmission and reflection, given by the Fresnel Relations.^[grif]

$$|\mathbf{E}_i| = \sqrt{RT} \cdot E_0 \cdot \cos(\omega t + \rho_i).$$

Where ρ_i is the phase of our wave, at the point where our detector lies. This phase is clearly related to the path length in the following way,

$$\rho_i = \frac{2\pi}{\lambda} S_i.$$

For our two waves we obtain,

$$|\mathbf{E}_1| = \sqrt{RT} \cdot E_0 \cdot \cos(\omega t + \rho_1)$$

$$|\mathbf{E}_2| = \sqrt{RT} \cdot E_0 \cdot \cos(\omega t + \rho_2)$$

Where we have used the fact that transmission and reflection does not impact the frequency of light. If the optics are aligned correctly, we will be able to measure the overlapping wave on our detector. This wave is given as the sum of \mathbf{E}_1 and \mathbf{E}_2 . Its intensity is,

$$I = c\epsilon_0 |\mathbf{E}_1 + \mathbf{E}_2|^2 = c\epsilon_0 RT (\cos(\omega t + \rho_1) + \cos(\omega t + \rho_2))^2.$$

In practice, we are only able to measure the temporal averaging of this, as the frequency is a small quantity. The average of a periodic function with period τ is,

$$\langle f \rangle = \frac{1}{\tau} \int_0^\tau f(t) dt.$$

This gives us,

$$\begin{aligned} \langle I \rangle &= \frac{1}{2\pi} \int_0^{2\pi} [\cos(\omega t + \rho_1) + \cos(\omega t + \rho_2)]^2 d(\omega t) \\ &= 1 + \cos(\rho_1 - \rho_2). \end{aligned}$$

Now let $\Delta\rho = \rho_1 - \rho_2$ and also assume that the coefficients of transmissions and reflection equal 0.5. We then obtain,

$$I = \frac{1}{4} c\epsilon_0 E_0^2 (1 + \cos \Delta\rho).$$

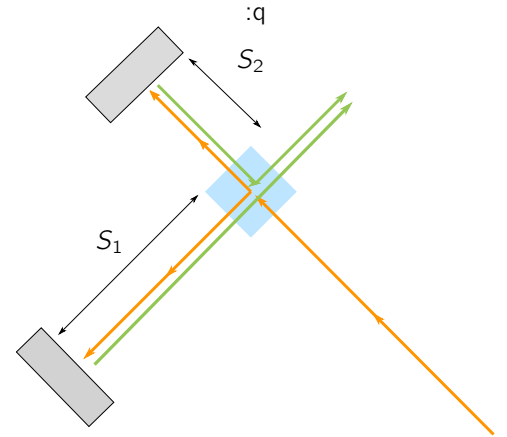


Figure 1: When the light is the incident light hits the beamsplitter, part of it is reflected and the remainder transmitted. Each lightbeam then travels a distance before hitting a mirror. The difference between these distances affects their relative phasedifference. We call it Δs .

Piezoelectric

We make use of a piezoelectric ring chip, which has the property that it expands when the potential over it increases. This allows us to slightly increase and decrease the path difference Δs in a controlled fashion. The relationship between the potential applied over the piezoelectric and the expansion is nearly linear, as seen in **fig. 2**. As a result, the path difference should be proportional to the voltage

$$V \propto \Delta s$$

It is particularly smart to change the potential over the piezoelement linearly. By doing this we are able to get displacement from the piezoelement spec-sheet **fig 2**. We fit these datapoints to determine the relationships between contraction and expansion, S_C and S_E , these functions will give half of the path change. By extensions the phase change will be,

$$\Delta\rho_{E/C}(V) = \frac{4\pi}{\lambda} S_{E/C}(V).$$

A measurement of the intensity, should then change sinusoidally in time.

Course of action

For this experiment, we took great care in setting up the interferometer correctly, as the quality of our data solely on this. The laser initially strikes two mirrors, which are used to direct it and correct it horizontally. It then strikes the beam splitter, which is aligned perpendicularly to the ray. The light is then split in two beams of equal intensity. One is transmitted while the other is reflected. These beams then strike mirrors, which send them back towards the beam splitter, where they are gathered and directed towards a detector. Before striking the detector, they pass through a concave lens. This spreads the light more evenly over the sensors in the detector. On one of the mirrors, is mounted a piezoelectric element. Applying a voltage over this element allows us to vary the path distance as the piezoelectric expands and contracts.

Results

In our experiment we measure the induced voltage in the detector as a function of the changing potential over a piezoelectric element. We are interested in relating this changing voltage to the phase change, and thereby ultimately to the path change. In order to figure out how. As the piezoelectric, behaves differently for increasing and decreasing voltage we have split our data accordingly. For both datasets, we have fit our data to functions of the form,

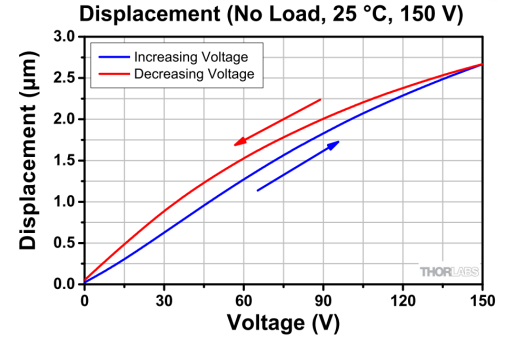


Figure 2