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SEVERO OCHOA SUMMER SCHOOL IN
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ENERGY**

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SUPERCONDUCTIVITY

(Tutorial)



UNIVERSITÀ
DEGLI STUDI
DI GENOVA

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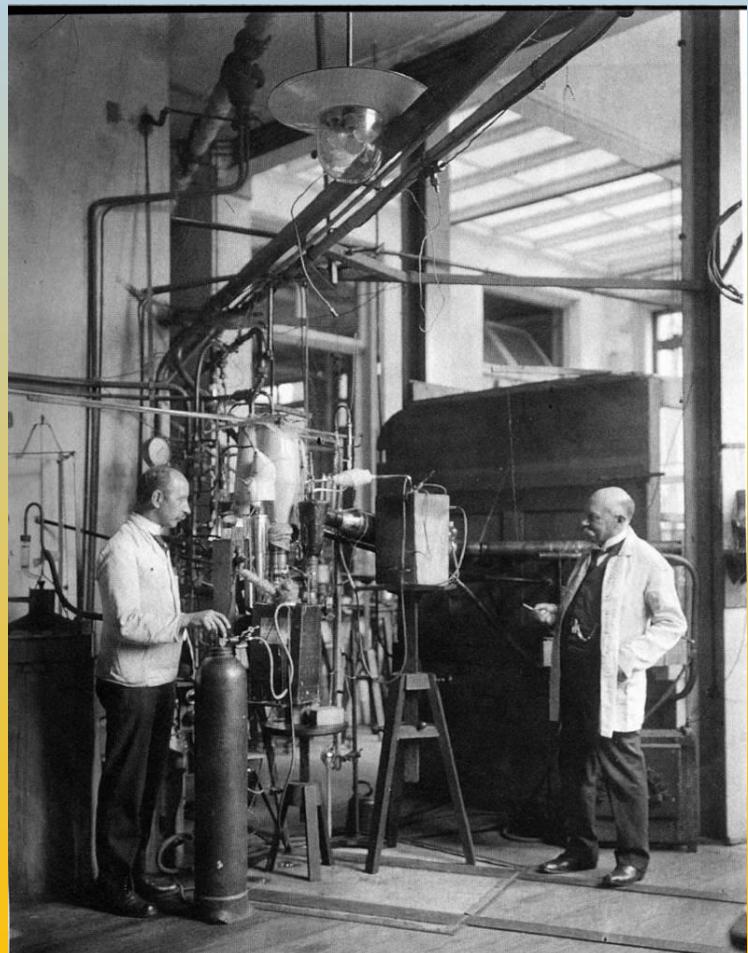


Outline

- ⇒ **Story of superconductivity phenomenon** going through the discovery of its main properties.
- ⇒ **Microscopic theory of superconductivity** and main parameters which characterize the superconducting state.
- ⇒ **Magnetic properties** and main mechanisms which allow a superconductors to carry superconducting current.
- ⇒ **Main superconducting materials.**
- ⇒ Peculiarities and issues of **novel superconducting materials.**

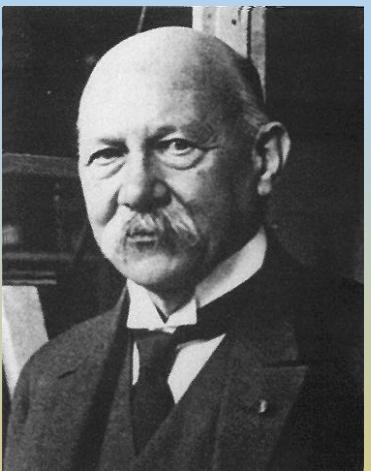
Kammerlingh Onnes, Leiden

1908: Liquefaction of Helium



1911: discovery of superconductivity

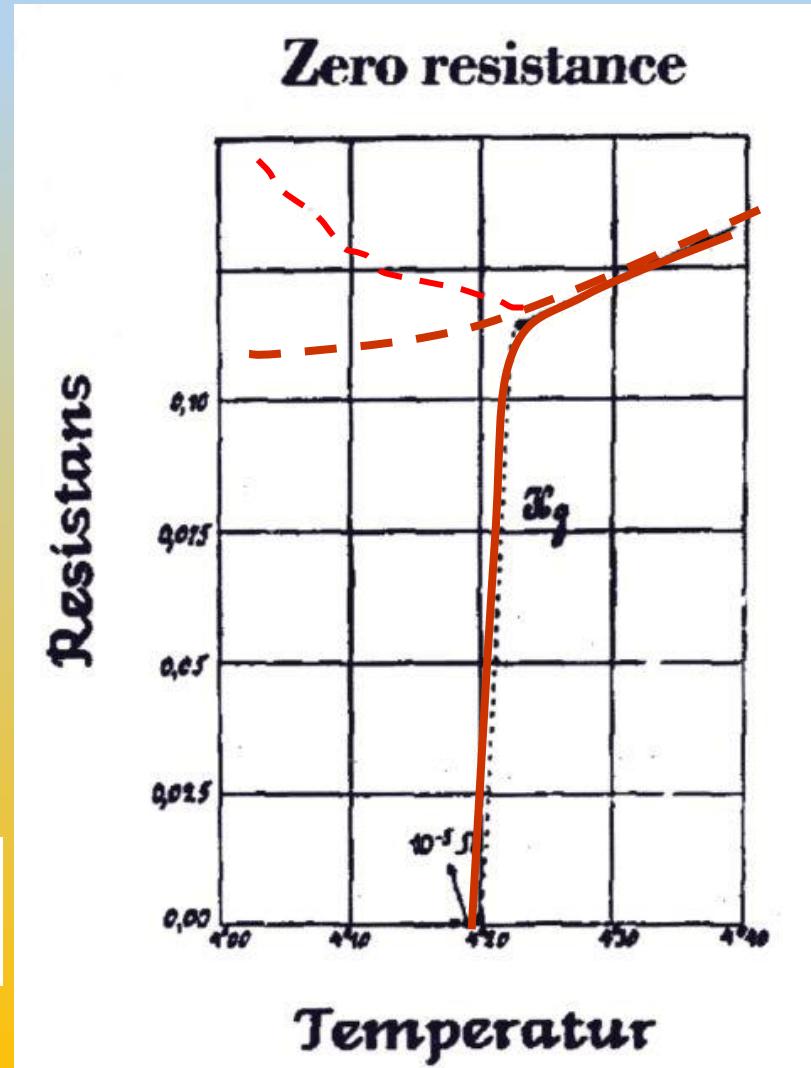
1913: Nobel price



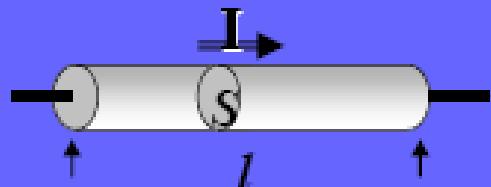
1913

⇒ He was measuring the resistivity of Hg to investigate the behavior of resistivity at low temperature

⇒ He noticed that the electrical resistance dropped to zero at 4.2K



The consequences of the Zero Resistance State ($R = 0$)



$$V = RI$$

$$R = \rho \frac{l}{S}$$

$R = 0$ implies no Joule dissipation in cable conducting current

$R = 0$ implies that current can flow in a loop without decay

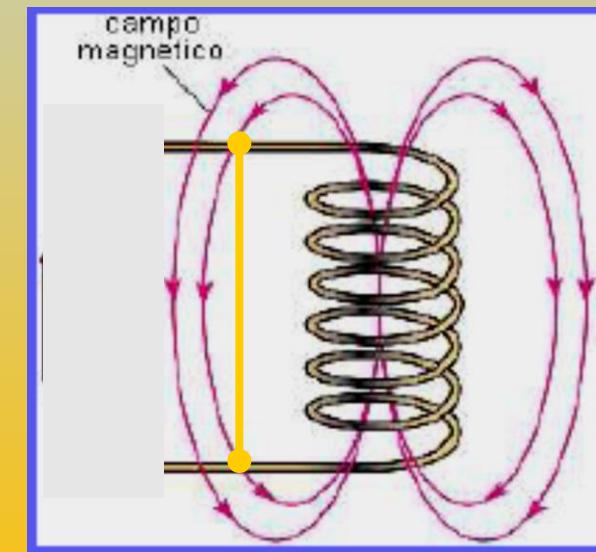
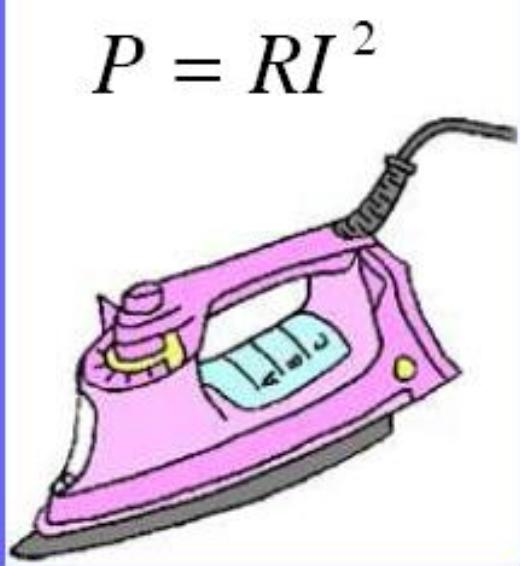
$R = 0$ implies that large magnetic field can be realized through solenoids with large current flowing in persistent mode



$$I = I_0 e^{-\frac{t}{\tau}}$$

$$\tau = \frac{L}{R}$$

$$P = RI^2$$



Kammerlingh Onnes's Visions

Construction of a 10 T Magnet with Hg und Pb Wires



Presented at 3° International Congress on
Refrigeration , Chikago 1913

Superconducting Magnet Coil

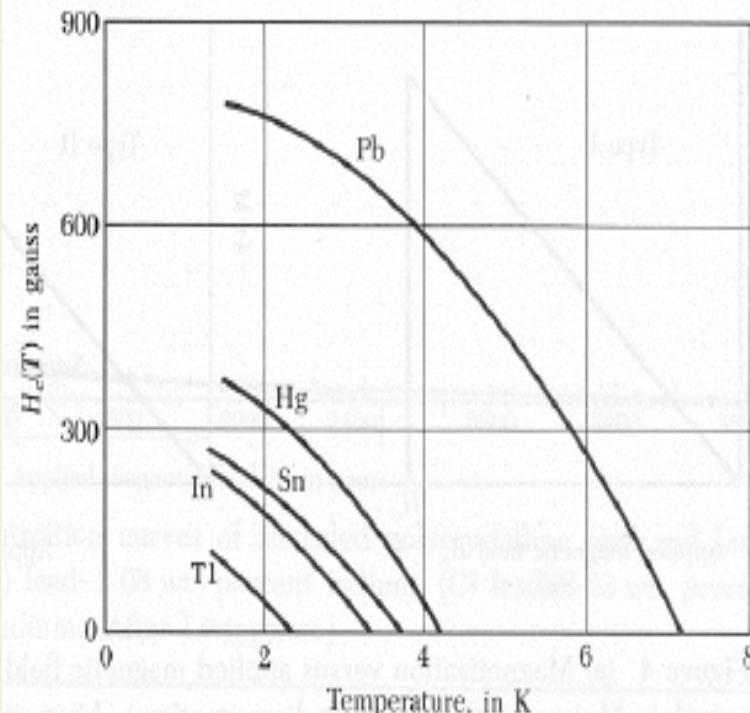
Experiments with Hg-and Pb wires failed

The coil lost superconducting properties already at small Current densities and at Magnetic Fields of several 100 Gauss.

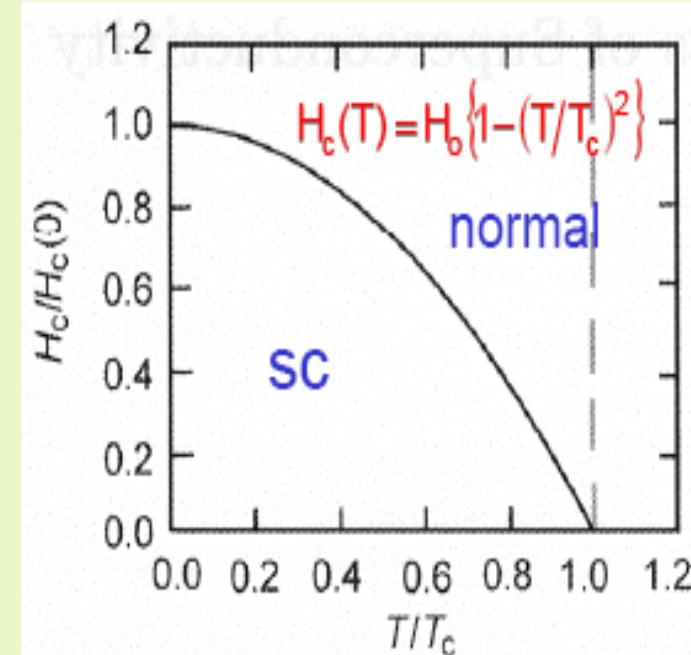
Superconductivity can be destroyed also by an external magnetic field , which is called **critical field, H_c**
(Kamerlingh Onnes 1914)

Critical Field

Temperature dependence of $H_c(T)$



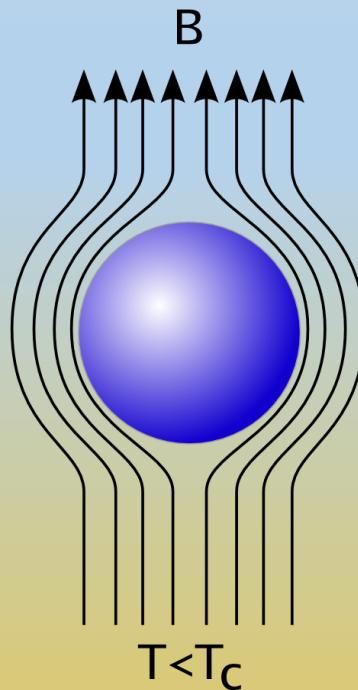
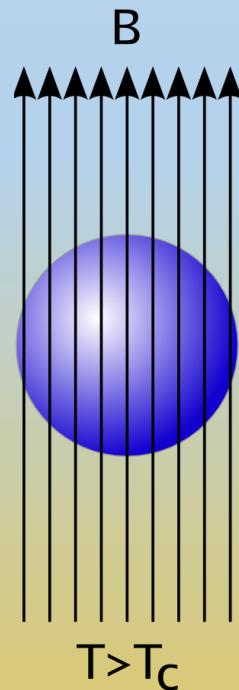
All curves can be collapsed onto a similar curve after re-scaling



Empirically :

$$H_c(T) = H_c(0) \left[1 - \left(\frac{T}{T_c} \right)^2 \right].$$

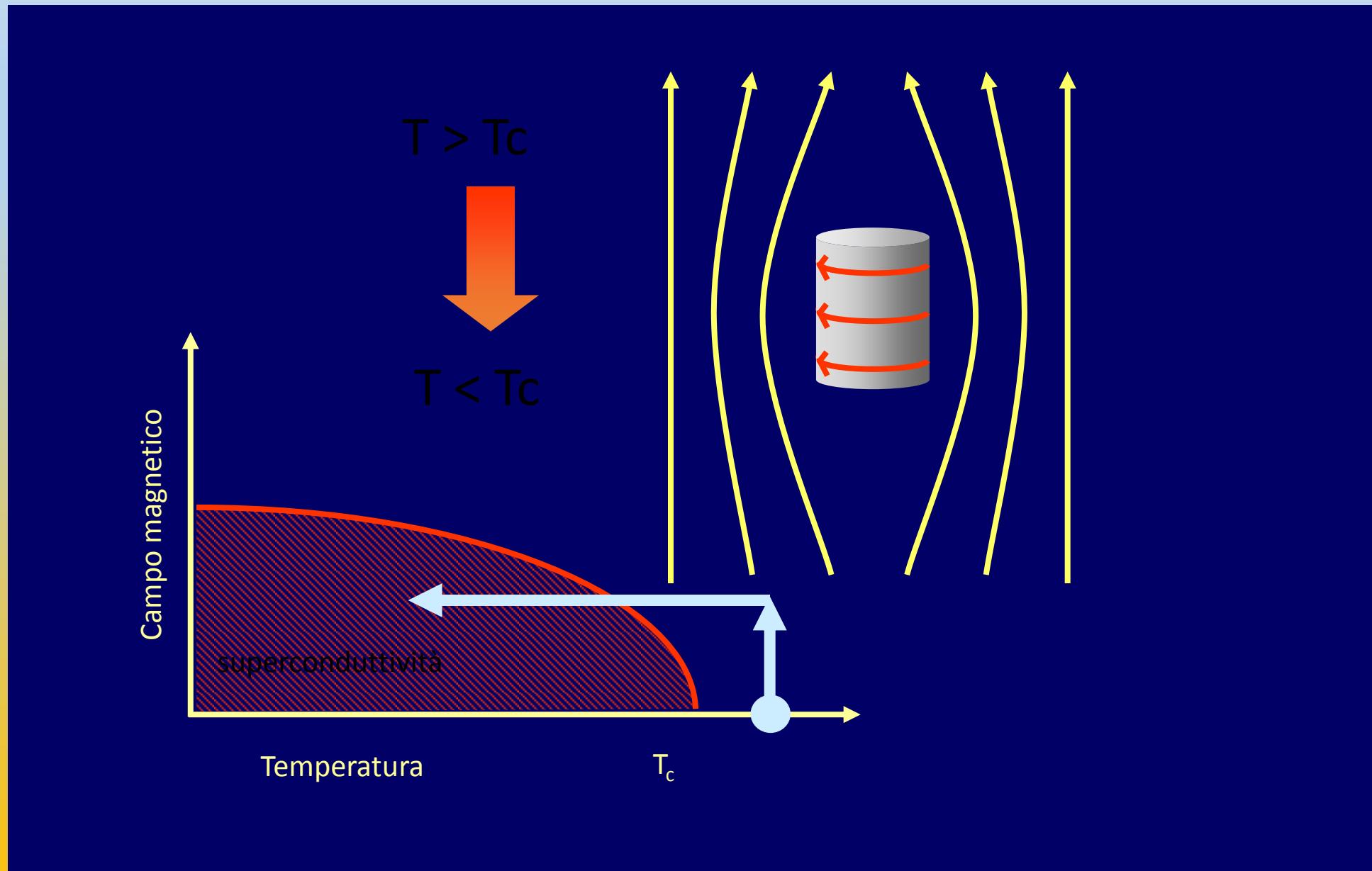
1933: Meissner-Ochsenfeld effect



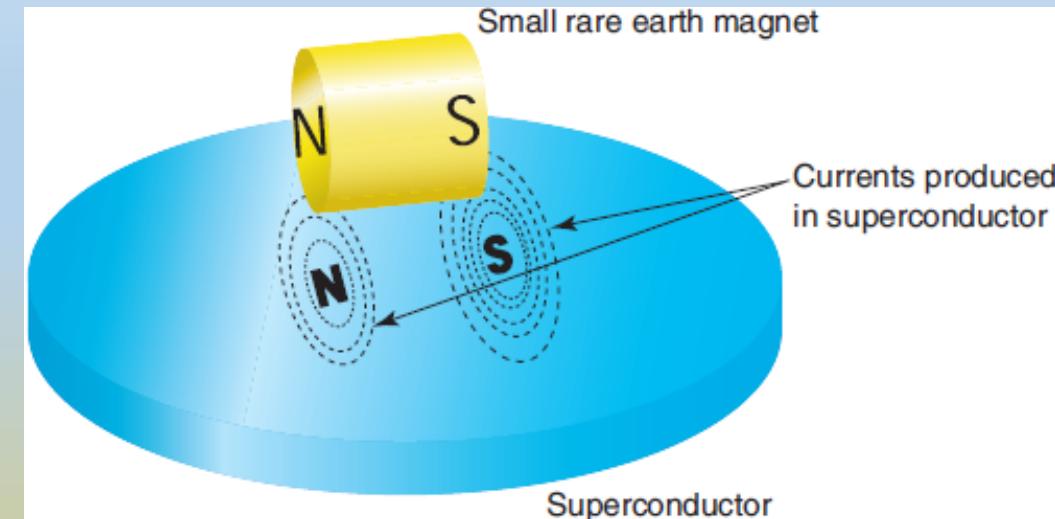
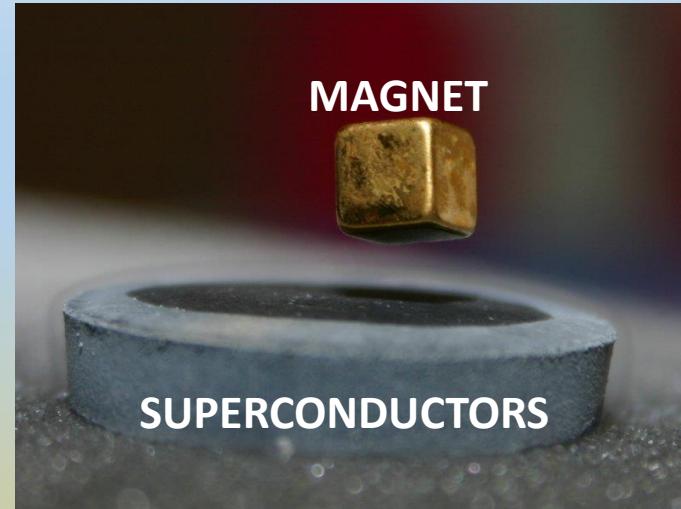
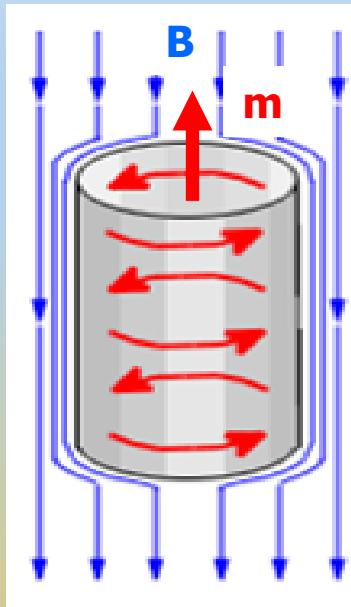
Magnetic field does not penetrate the superconductor

Ideal conductor!

Ideal diamagnetic!



The consequences of the perfect diamagnetism $\chi=-1$



The expulsion of the magnetic field, generates **great magnetic moments**, able to levitate extremely large masses like the coaches of a train



ELECTRODINAMIC LEVITATION
The future (2025) maglev train between **Tokyo and Osaka** will be the **fastest train** in the world with a speed approximately **1000 km/h.**

The quantum coherence

1950: Ginzburg-Landau Phenomenology Ψ -Theory of
Superconductivity

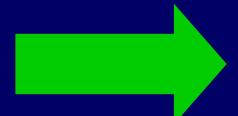
$T > T_c$



$T < T_c$



A macroscopic wave function describes the system as a whole



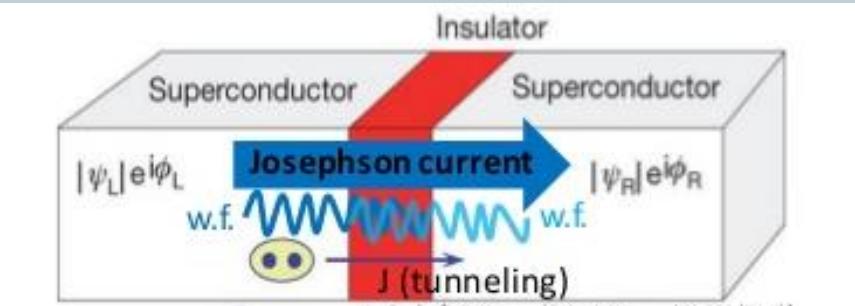
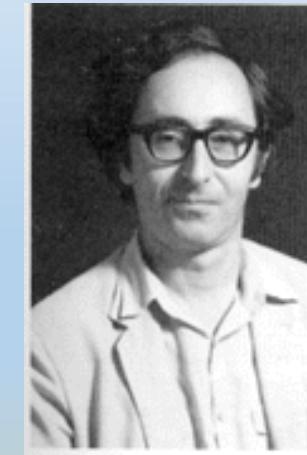
Macrscopic quantum effects

1962: Josephson effect

Super electrons cross the insulating junction between the two superconductors as a consequence of the phase difference between the two superconductors

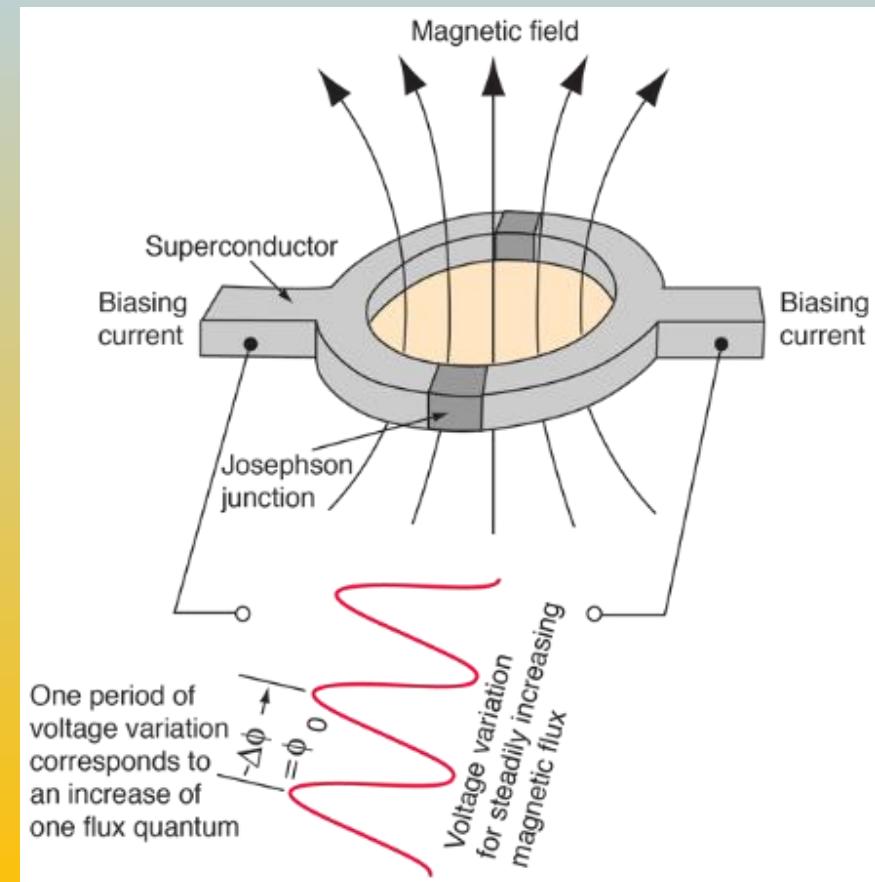


1973



SQUID superconducting quantum interference device

it is based on the presence of two superconductive junctions and on the interference due to the phase difference across the two junctions



SQUID for biomagnetisms



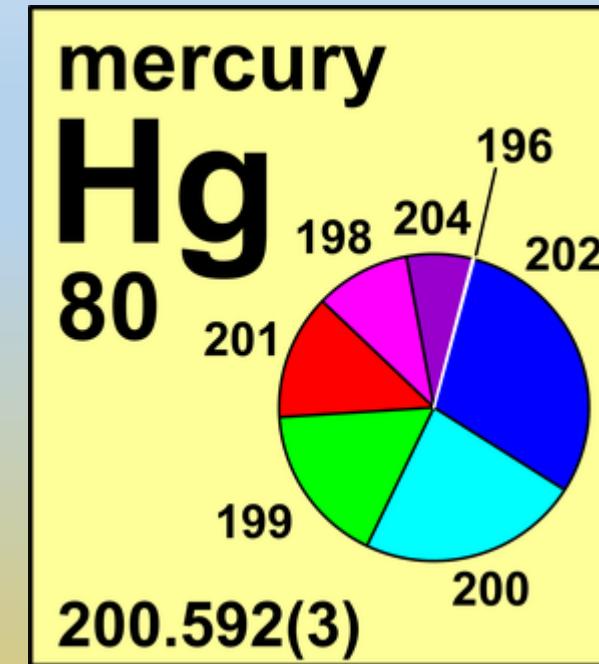
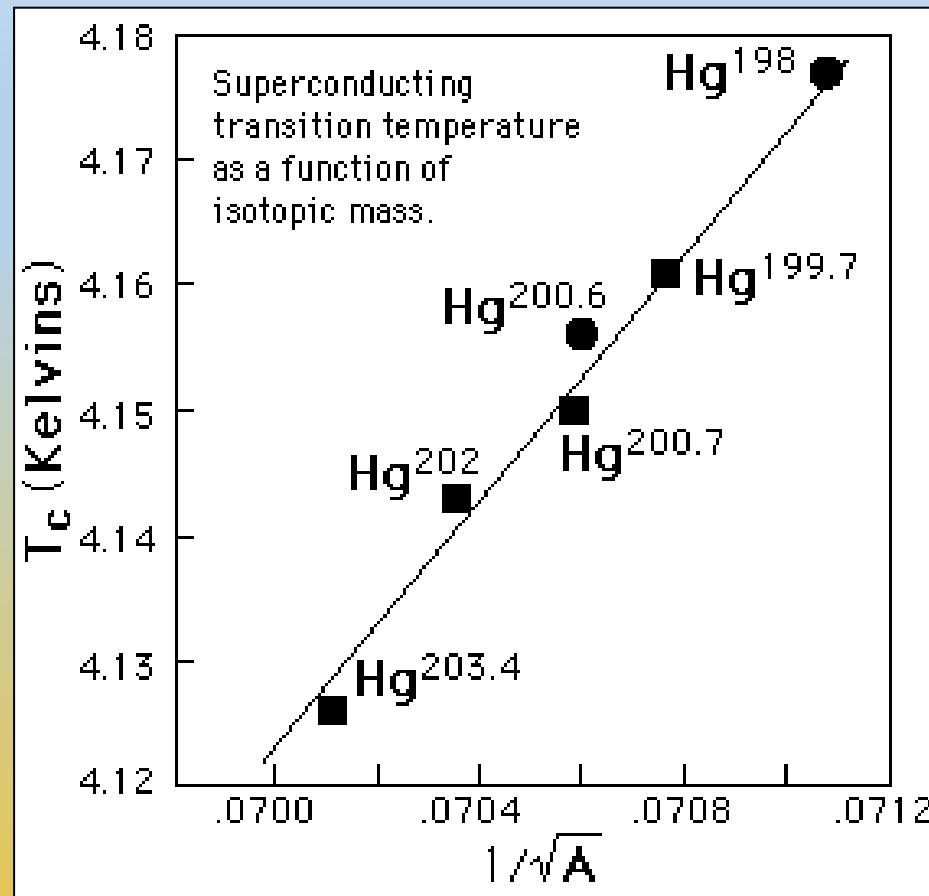
SQUID sensitivity	10^{-14} T
Heart magn. field	10^{-10} T
Brain magn. field	10^{-13} T
Earth magn. field	10^{-5} T
Fridge magnet	10^{-2} T

1957: Microscopic theory of superconductivity (Bardeen, Cooper, Schrieffer)



1972

1950: Isotopic effect

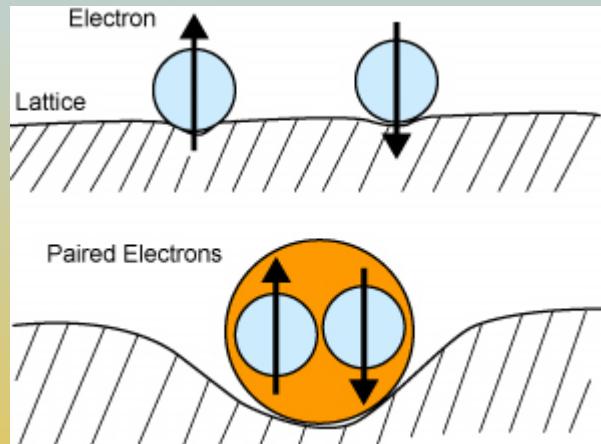


Isotopic effect suggests that the ion lattice plays a crucial role

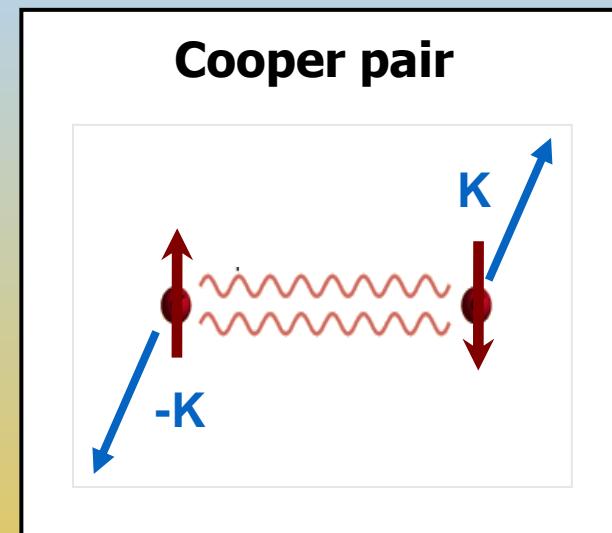
$$T_c \propto M^{-1/2}, \quad H_c \propto M^{-1/2}.$$

BCS- Theory of superconductivity

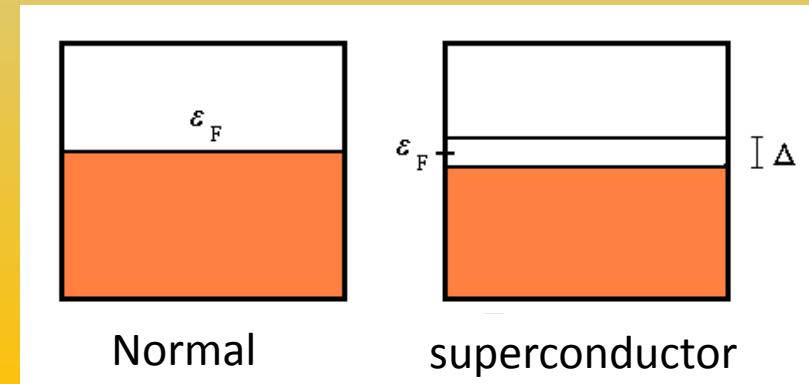
■ in the presence of an **attractive interaction** provided by the lattice



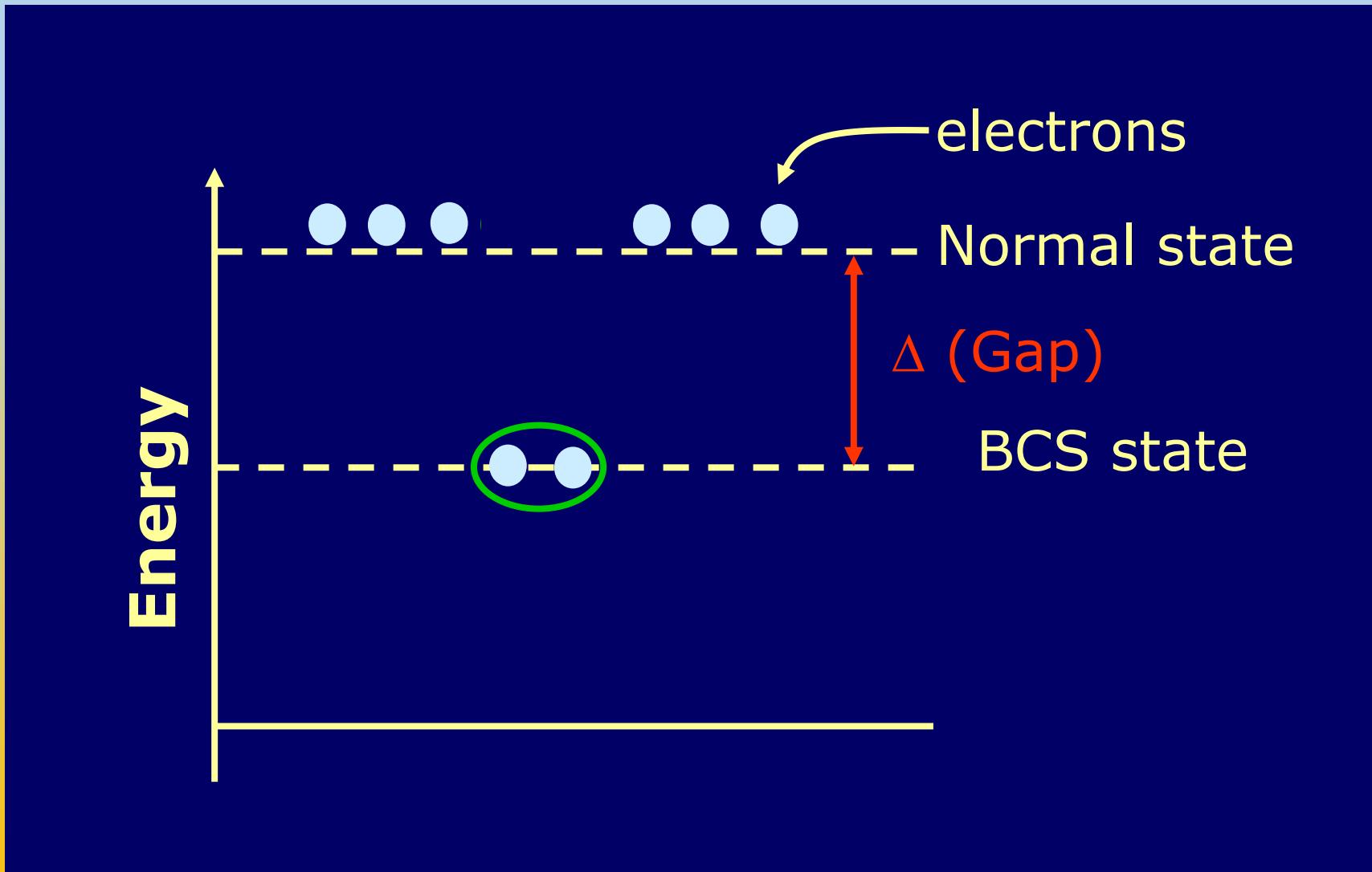
■ electrons form pairs of opposite impulse and spin: the **Cooper pairs**

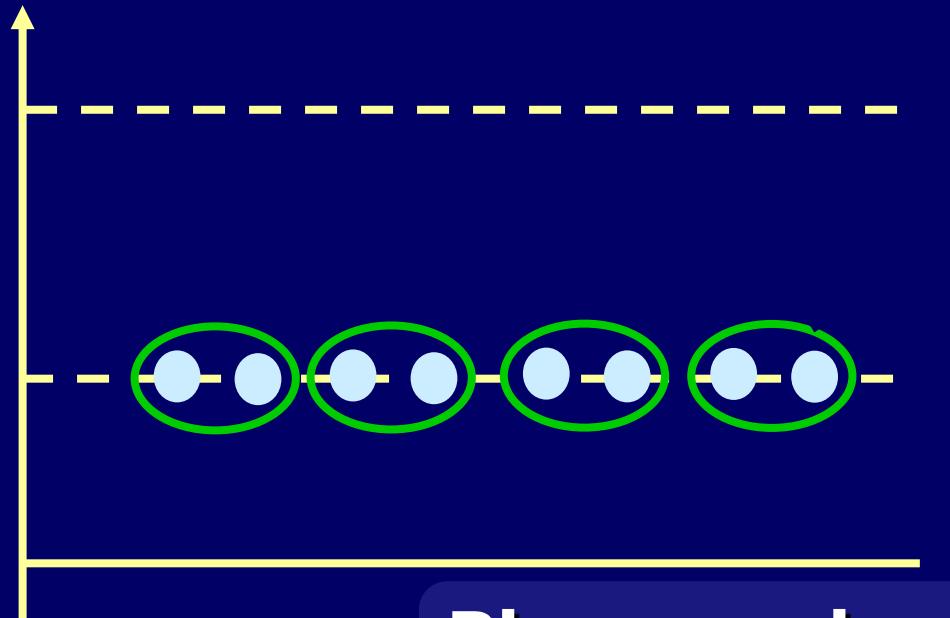


■ An **energy gap opens** in the energy spectrum which stabilizes the superconducting state



Energy gap



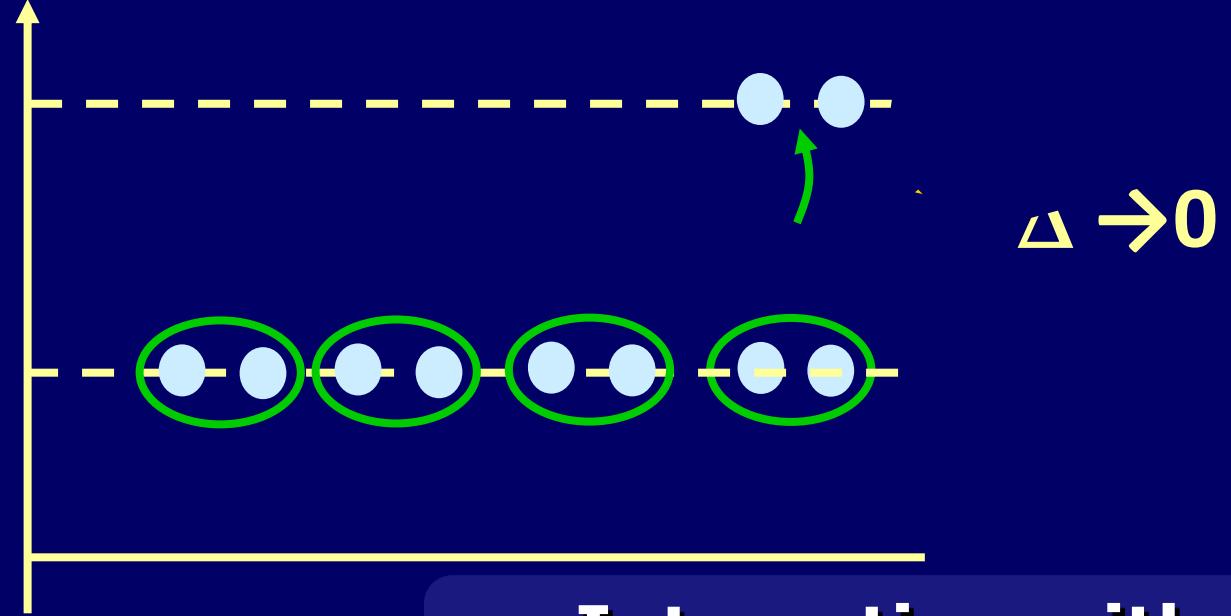


$T < T_c$

**Phonons don't have
energy enough to
break the Cooper pairs**



$\rho = 0$

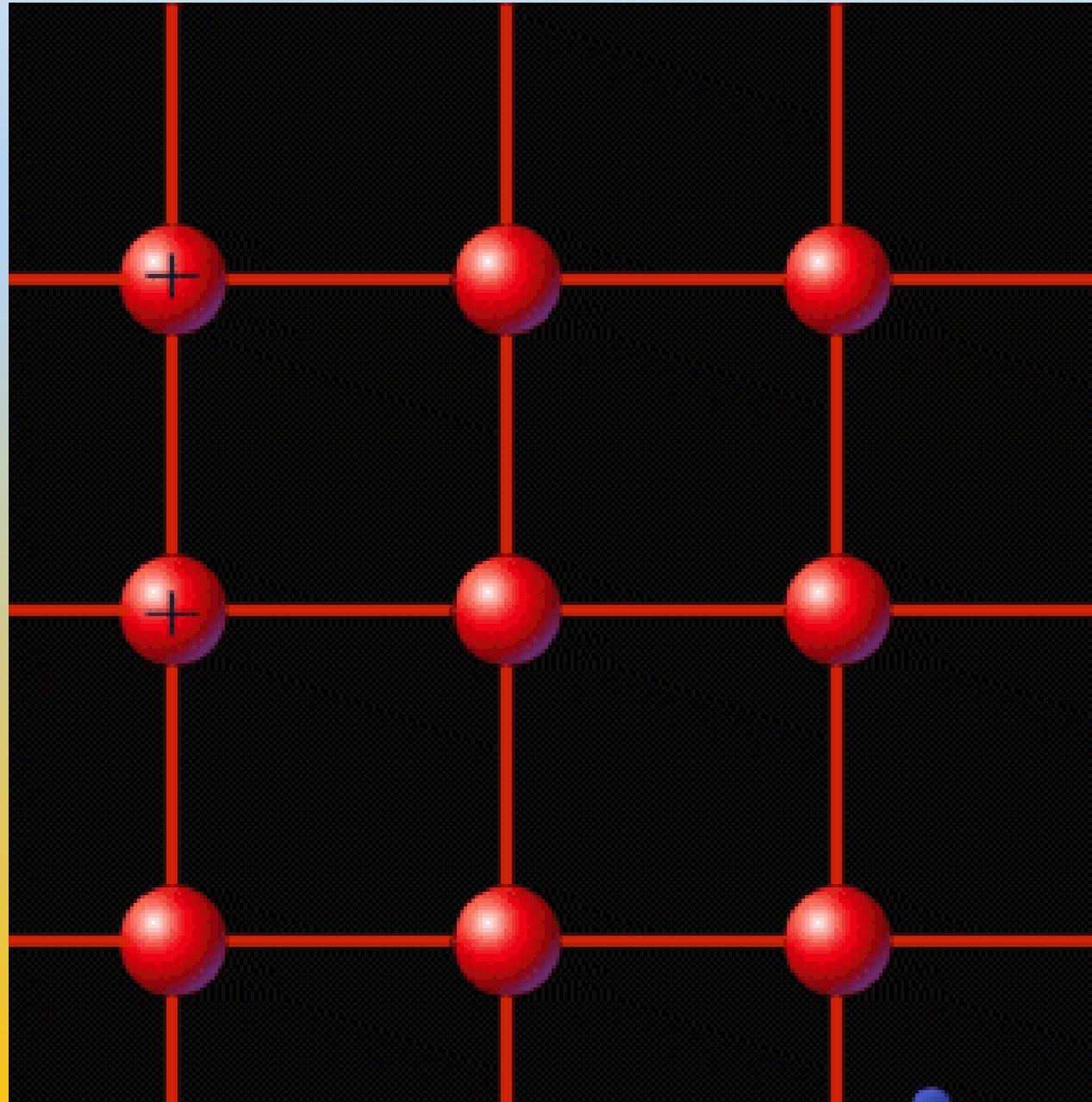


A $T \rightarrow T_c$

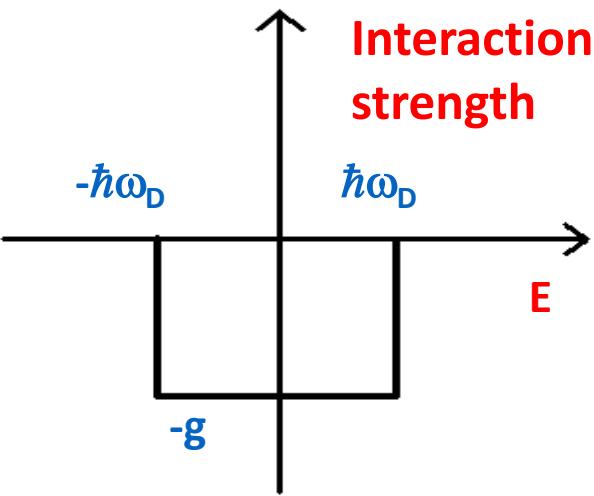
**Interaction with
phonons break the
Cooper pairs**

\downarrow
 $\rho \neq 0$

Electron pairing mechanisms

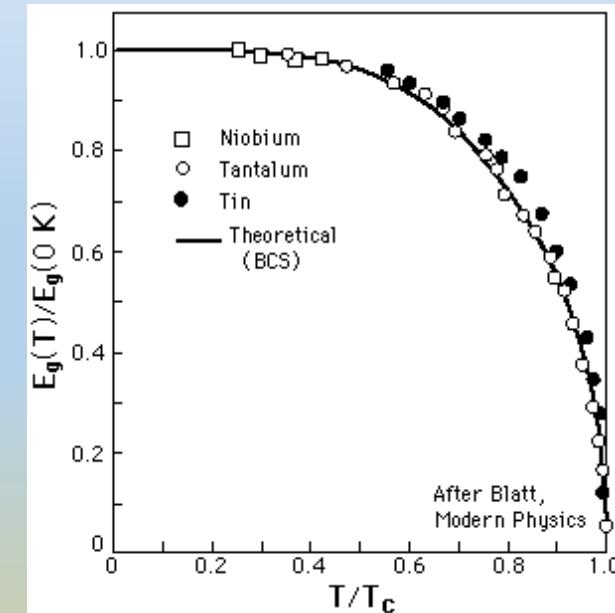


BCS theory predictions



Energy Gap

$$\Delta = \hbar\omega_D \exp\left(-\frac{1}{g}\right)$$



Critical temperature $T_c \sim 1-20 \text{ K}$

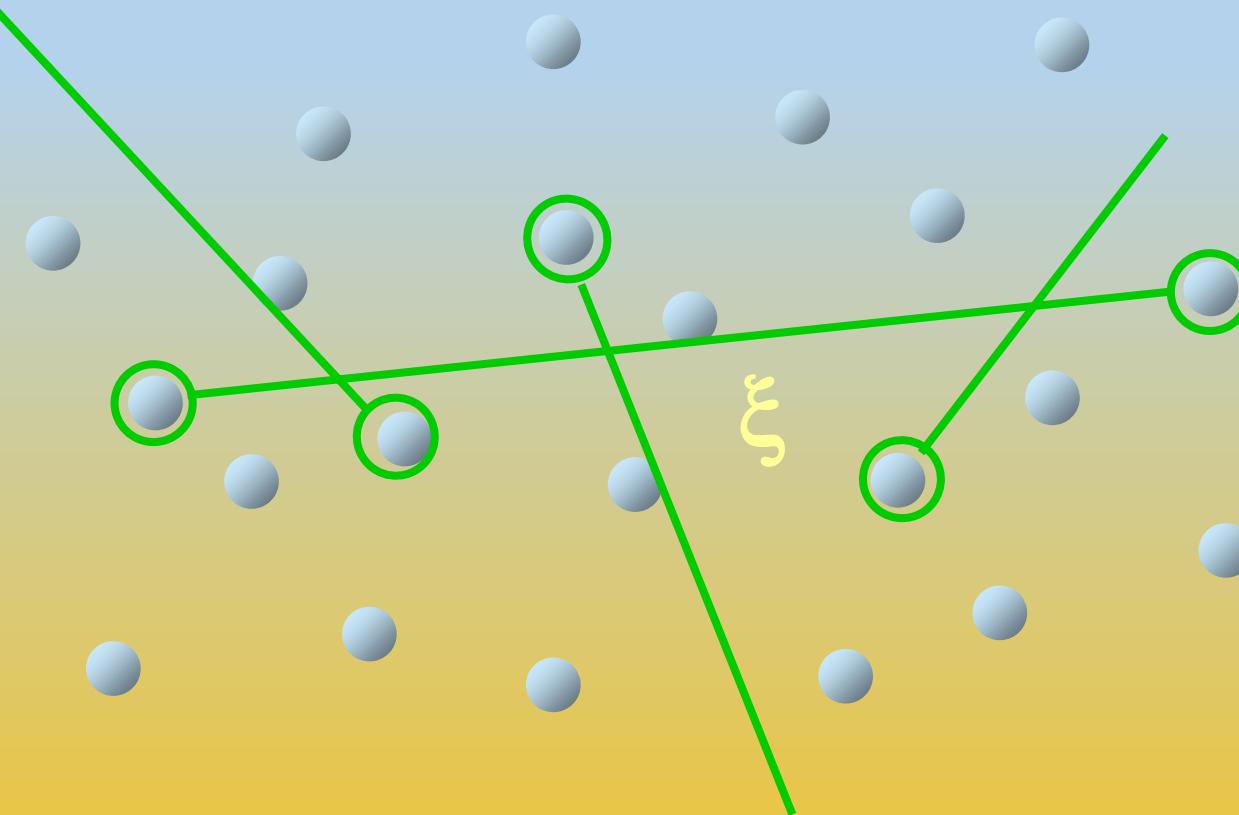
$T_{c\max} \sim 30-40 \text{ K}$

$$T_c = \Theta_D \exp\left(-\frac{1}{g}\right)$$

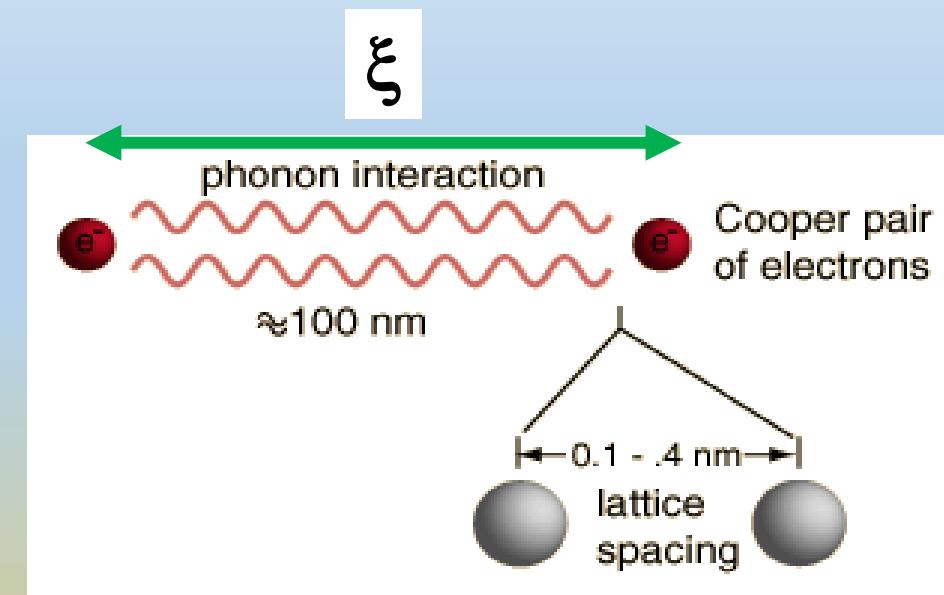
Θ_D Debye Temperature

$\Theta_D \sim 100-400 \text{ K}$

Coherence length



$\Rightarrow \xi \sim \text{size of the Cooper pair}$

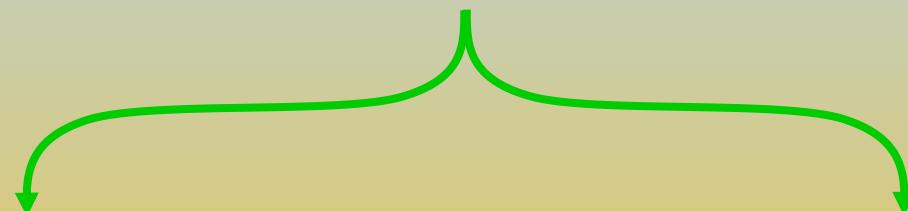


For the superconducting elements the coherence length can be as large as hundreds nm, much larger than the lattice spacing.

⇒ between the two electrons of a pair, there may be thousands of electrons belonging to other pairs.

MAGNETIC PROPERTIES

**2 classes of superconductors
with different properties**



TYPE I

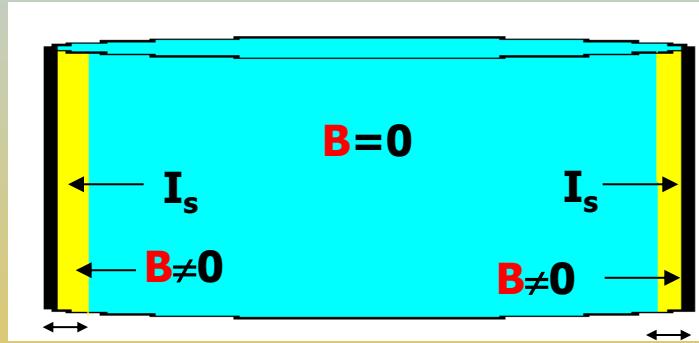
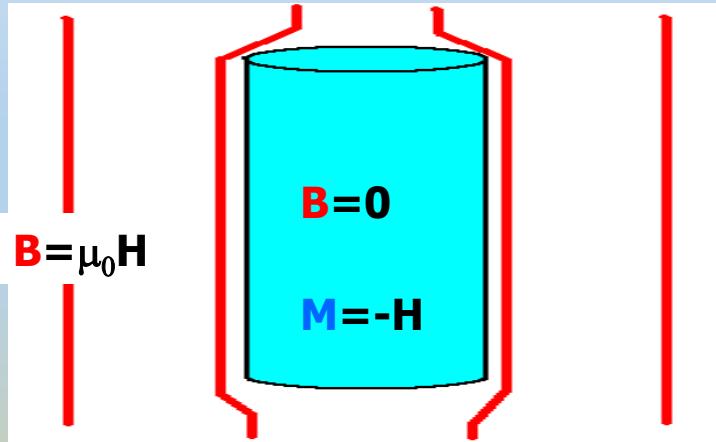
(elements)

TYPE II

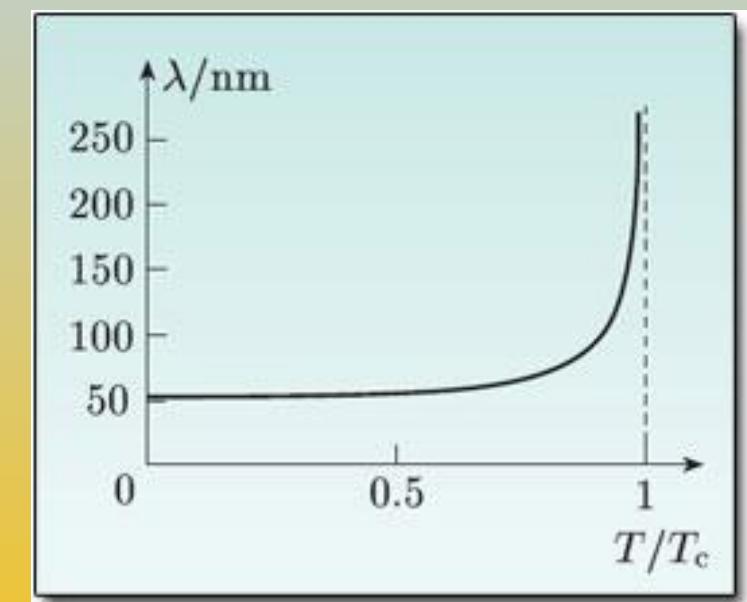
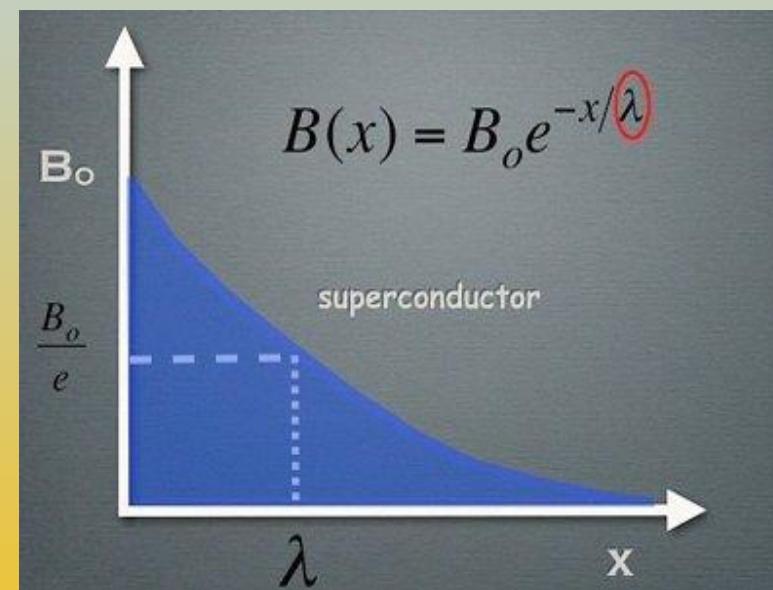
(alloys, composites, ...)

London penetration depth

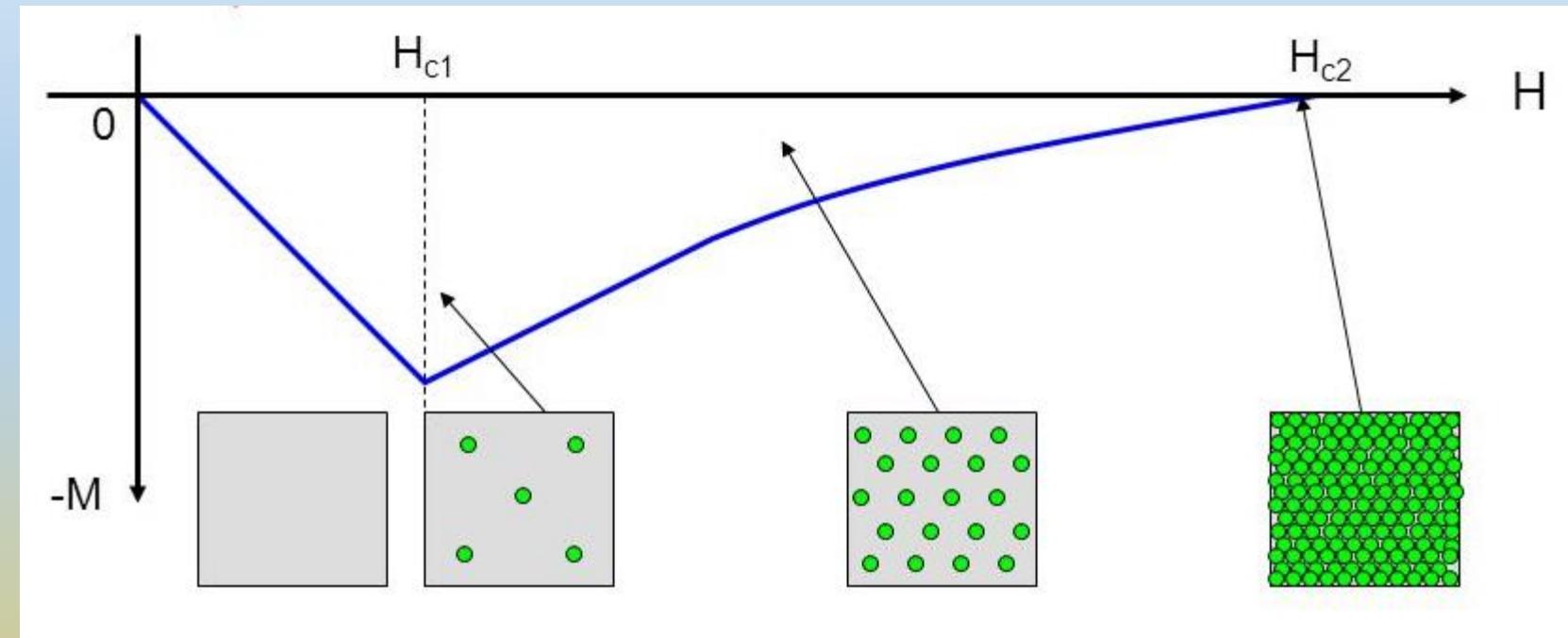
1935: Brothers London theory



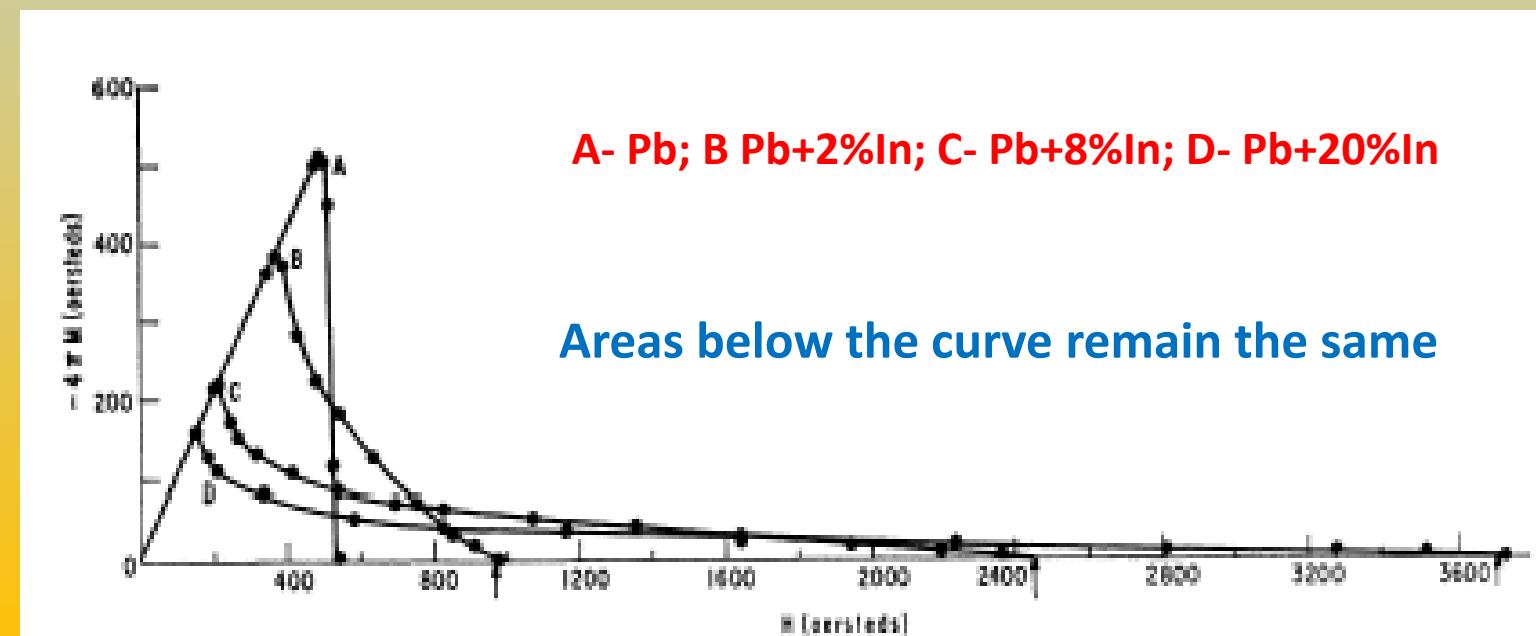
The magnetic field does not change abruptly to zero within the superconducting material. In the superficial thickness where the superconducting currents flow, the magnetic field goes progressively to zero

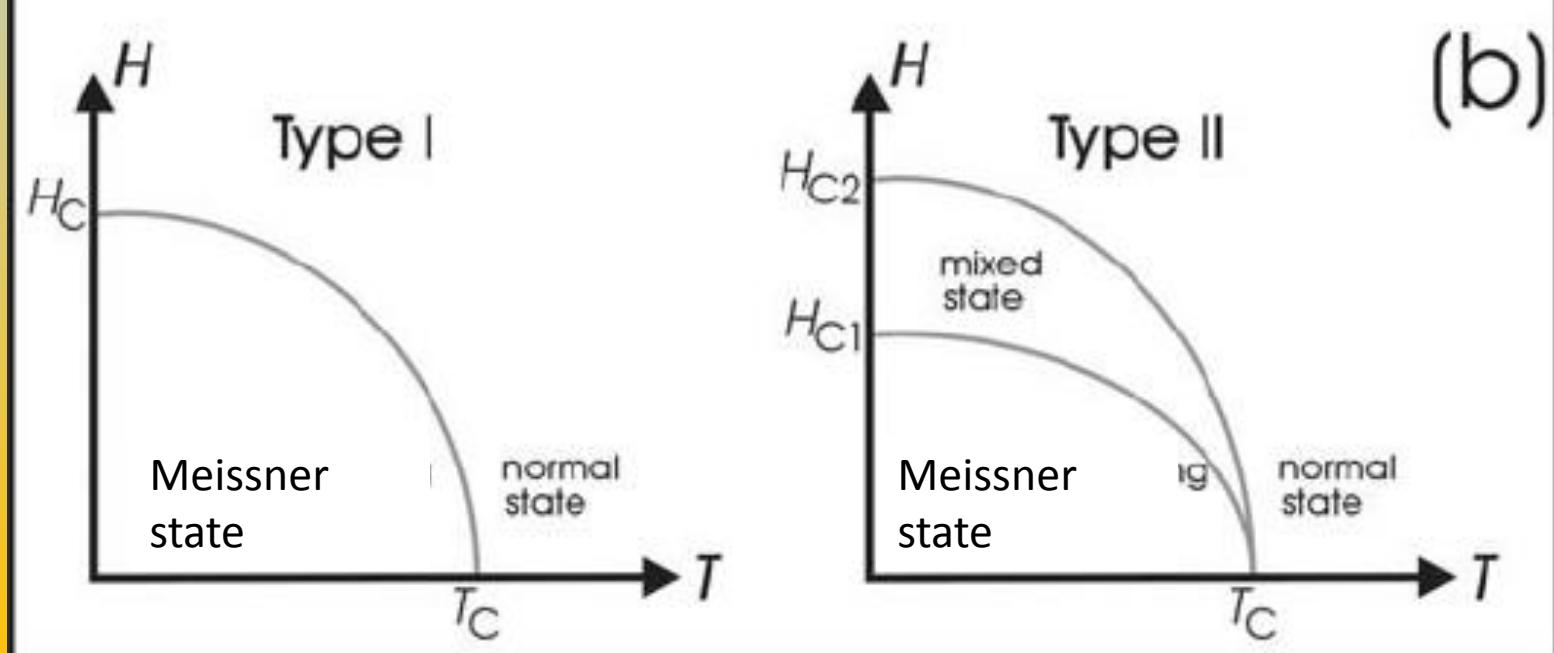
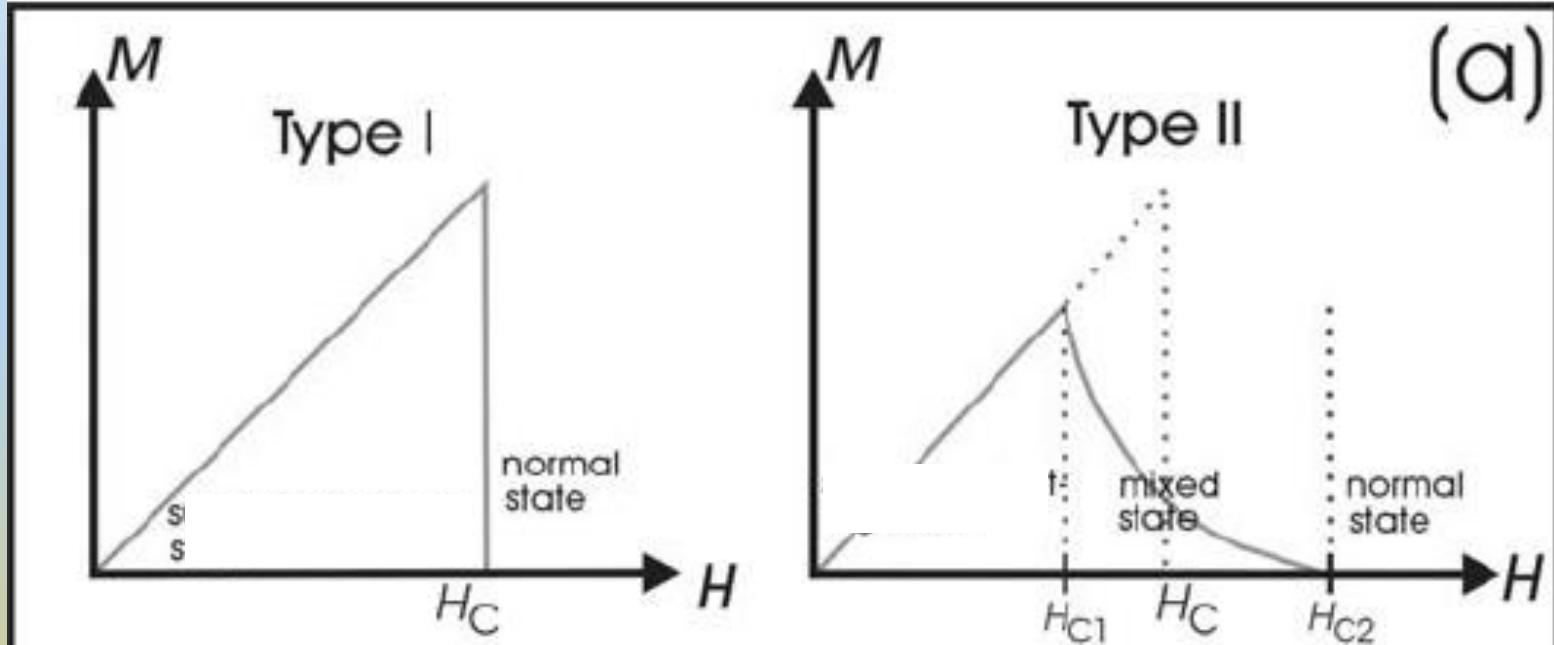


Magnetization curves



$H_{c2} > H_c$
Type II superconductors
can be high field
superconductors





Who is of I or II type?

Type I

Type II

Material	T_c (K)	ξ (nm)	λ (nm)	κ (λ/ξ)
Cd	0.56	760	110	0.14
Al ^a	1.18	550	40	0.03
In ^a	3.41	360	40	0.11
Sn ^a	3.72	180	42	0.23
Ta	4.4	93	35	0.38
Pb ^a	7.20	82	39	0.48
Nb ^a	9.25	39	50	1.28
Pb-In	7.0	30	150	5.0
Pb-Bi	8.3	20	200	10
Nb-Ti	9.5	4	300	75
Nb-N	16	5	200	40
PbMo ₆ S ₈ (Chevrel)	15	2	200	100
V ₃ Ga (A15)	15	= 2.5	90	= 35
V ₃ Si (A15)	16	3	60	20
Nb ₃ Sn (A15)	18	3	65	22
Nb ₃ Ge (A15)	23.2	3	90	30
K ₃ C ₆₀	19	2.6	240	92
Rb ₃ C ₆₀	29.6	2.0	247	124
(La _{0.925} Sr _{0.075}) ₂ CuO ₄ ^b	37	2.0	200	100
YBa ₂ Cu ₃ O ₇	89	1.8	170	95
HgBaCaCuO	126	2.3		
HgBa ₂ Ca ₂ Cu ₃ O _{8+δ}	131			100
MgB ₂	39	10	80	10
SmFeAsO _{0.7} F _{0.3}	58	2.5	200	80
KBaFe ₂ As ₂	38	3	200	67
Fe(Se,Te)	18	1.5	400	300

$$\kappa = \frac{\lambda}{\xi} \quad \text{Ginzburg - Landau parameter}$$

$$\kappa < \frac{1}{\sqrt{2}} \quad \text{Type I} ; \quad \kappa > \frac{1}{\sqrt{2}} \quad \text{Type II}$$

With disorder

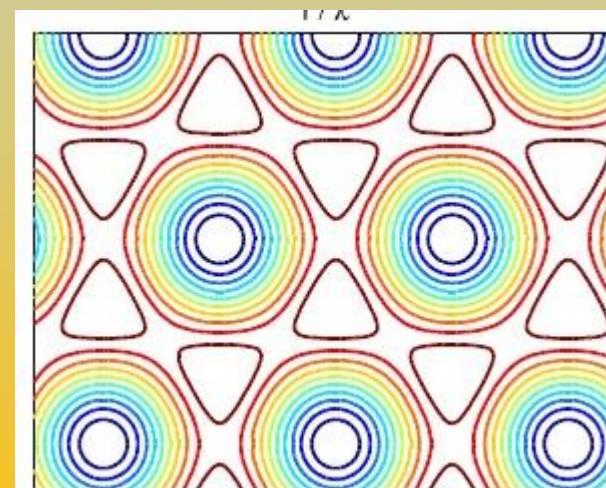
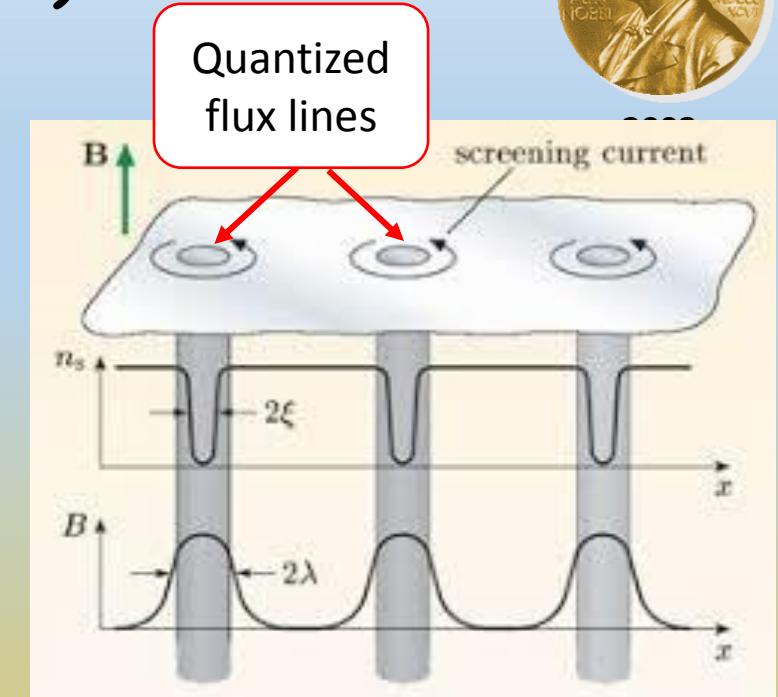
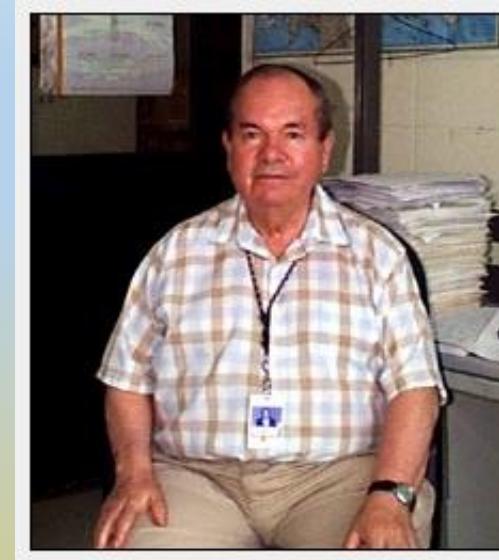
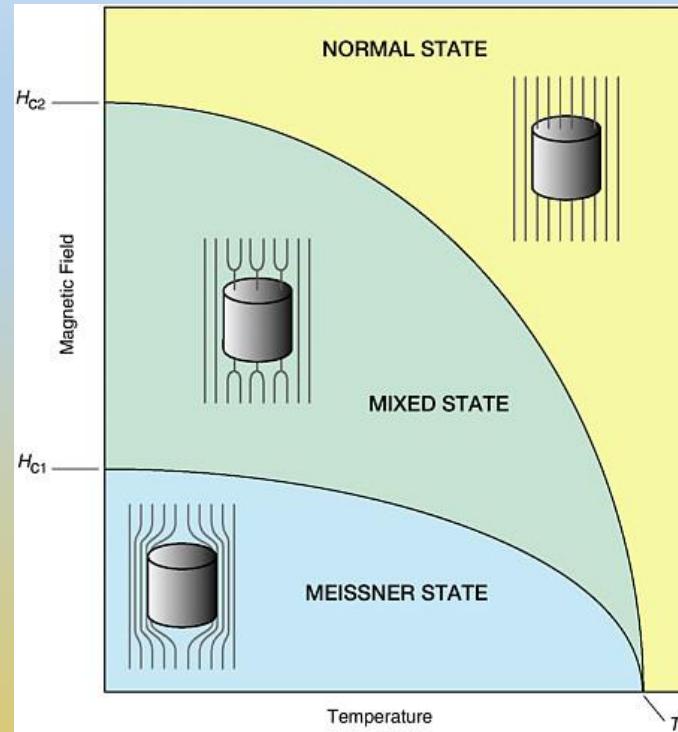
→ Type II

$$\lambda = \left(\frac{m}{2e^2 \mu_0 n_s} \right)^{1/2}$$

$$\xi \propto \frac{n_s^{1/3}}{T_c}$$

$$\kappa = \frac{\lambda}{\xi}$$

1957: What the mixed state is (Abrikosov)



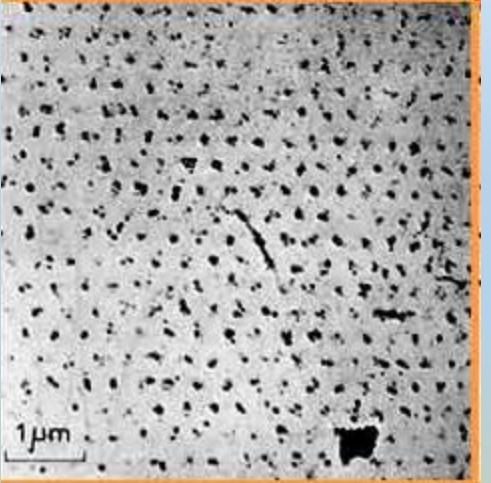
⇒ Quantized flux lines
(fluxons/vortices) enter the superconductor

⇒ Fluxons form an exagonal lattice

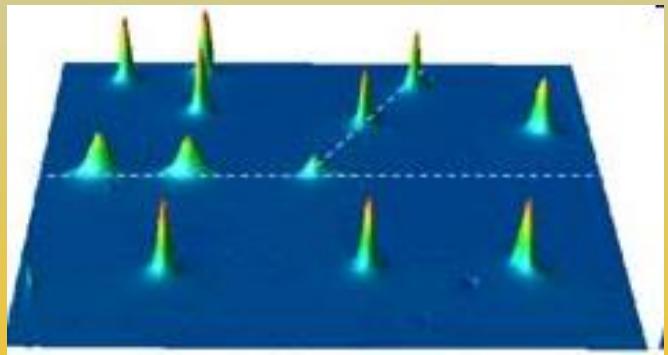
Abrikosov solution near H_{c2} –
Contours of B

$$\Phi_0 = \frac{h}{2e} = 2.067 \times 10^{-15} Wb$$

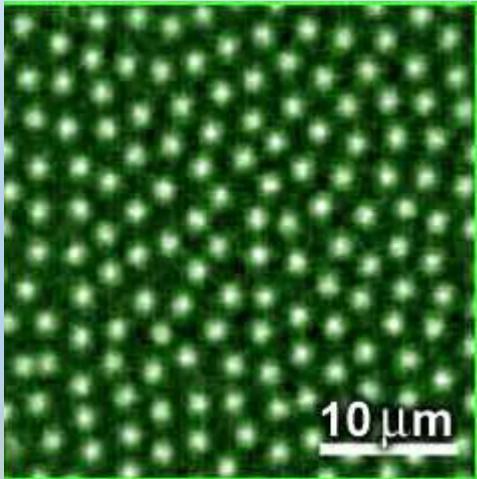
A graph showing the radial profile of the magnetic field B (red curve) and the density of fluxons n_s (blue curve) as a function of radial distance r . The field B has sharp peaks at $r = 0$ and $r = \lambda$, and a minimum at $r = \xi$. The fluxon density n_s has peaks at $r = \xi$ and $r = 2\xi$, and a minimum at $r = \lambda$. A current j is indicated by a circular arrow below the graph.



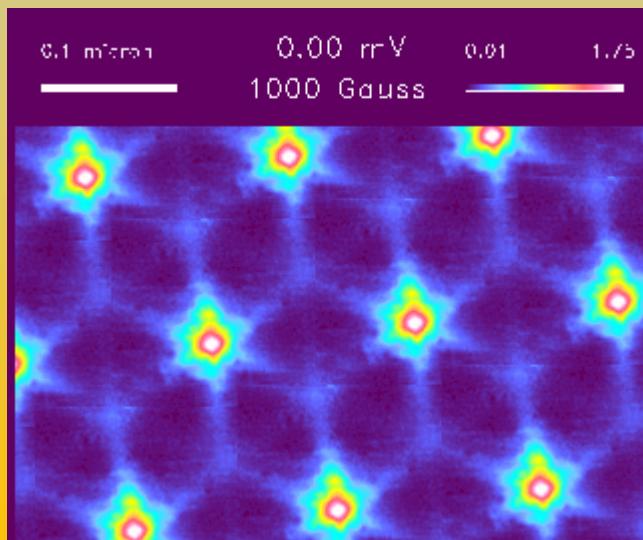
U. Essmann and H. Trauble
Max-Planck Institute, Stuttgart
[Physics Letters 24A, 526 \(1967\)](#)

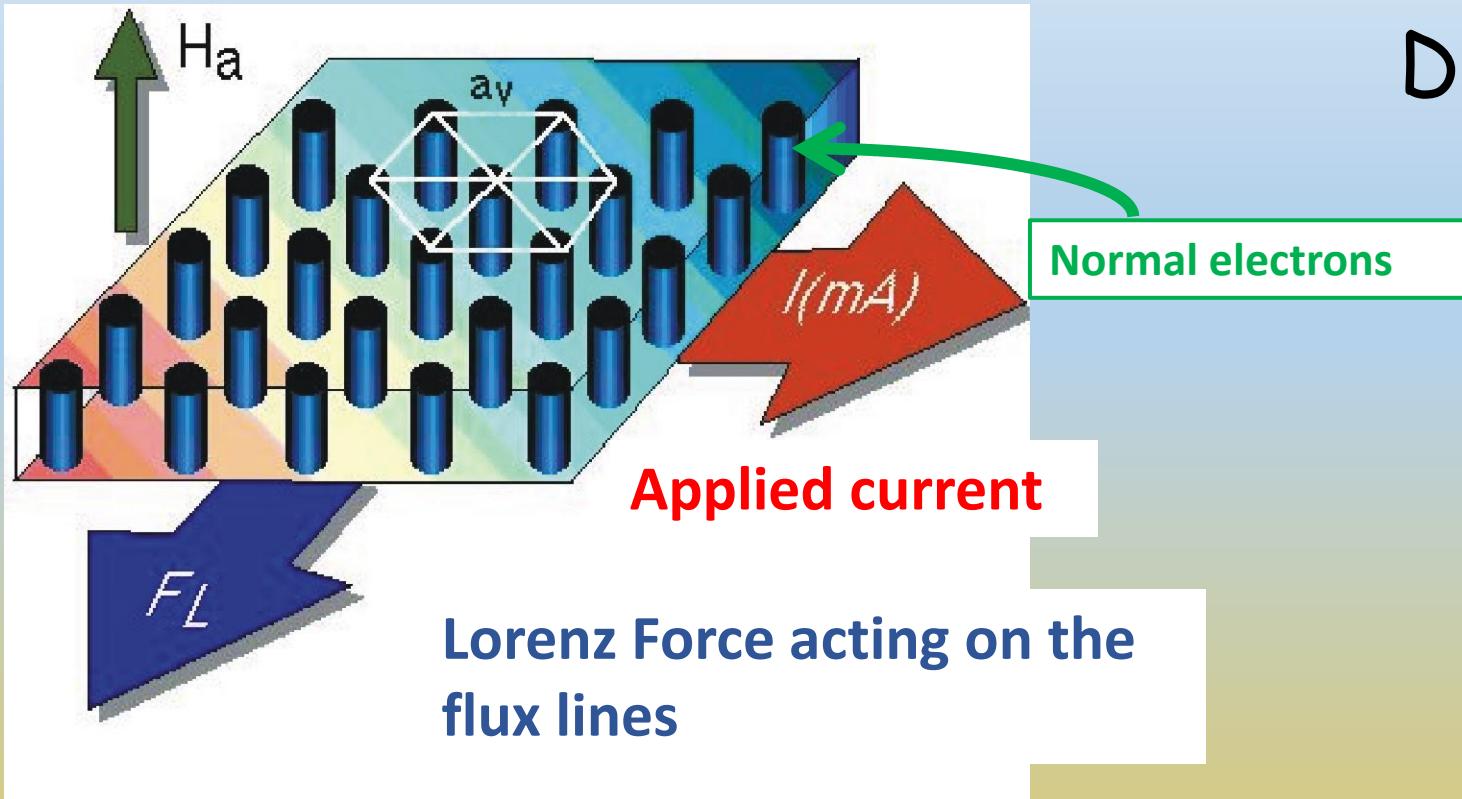


Scanning SQUID Microscopy of half-integer vortex, 1996
J. R. Kirtley et al. [IBM Thomas J. Watson Research Center](#)
[Phys. Rev. Lett. 76, 1336 \(1996\)](#)



Magneto-optical image
of Vortex lattice, 2001
P.E. Goa et al.
University of Oslo
[Supercond. Sci. Technol. 14, 729 \(2001\)](#)

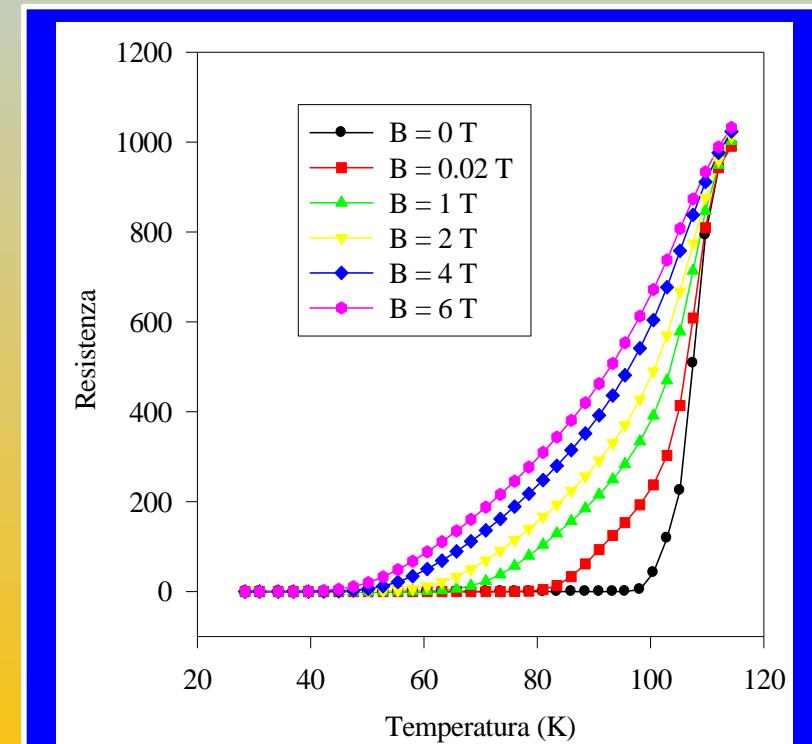




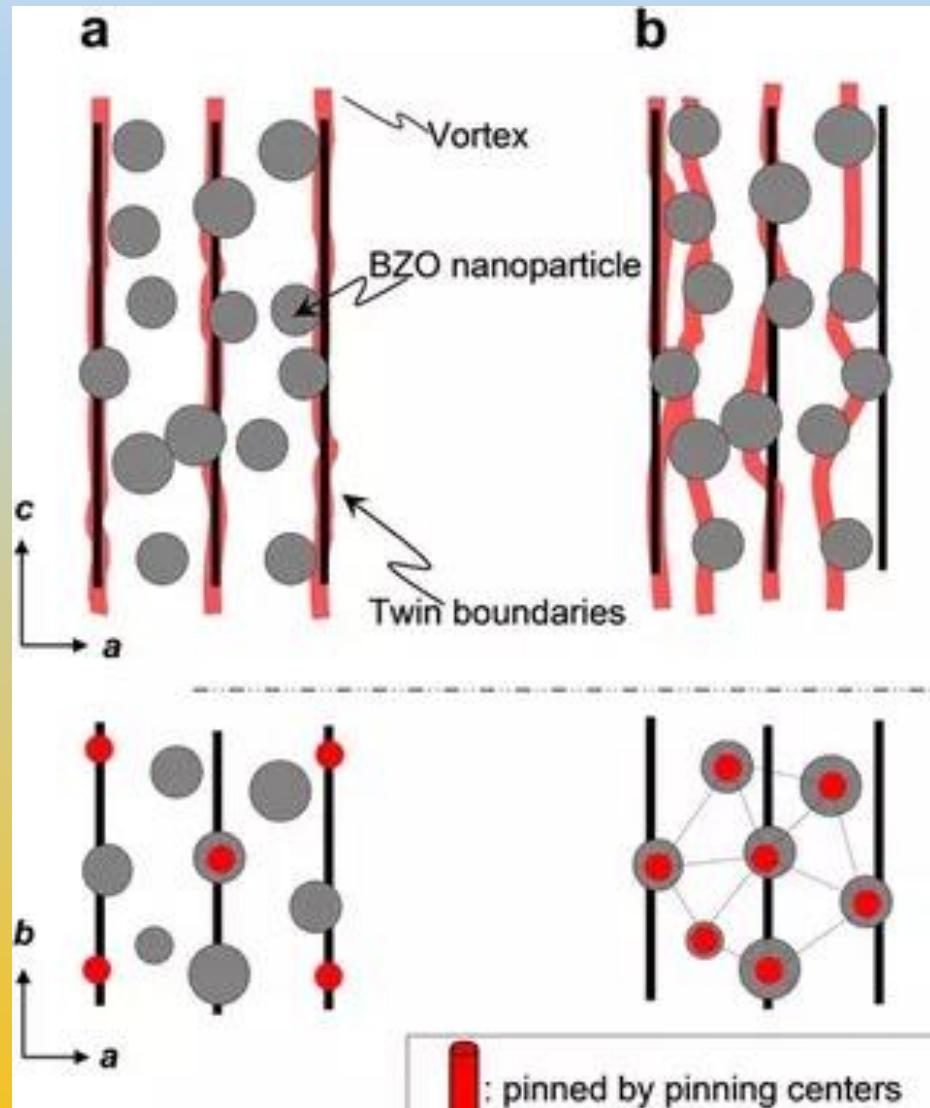
Dissipation in the mixed state

- ⇒ an **applied current** exerts a force on the vortices
- ⇒ the motion of the fluxons whose core is made up of **normal electrons** causes dissipation

- ⇒ The system is superconducting but it **is resistive and dissipates**



Vortex pinning



The flux motion reduces the current that the superconductor can carry without dissipation.

SOLUTION: to introduce defects which pin the vortices

Critical current I_c :

the maximum current that the superconductors can carry without dissipation

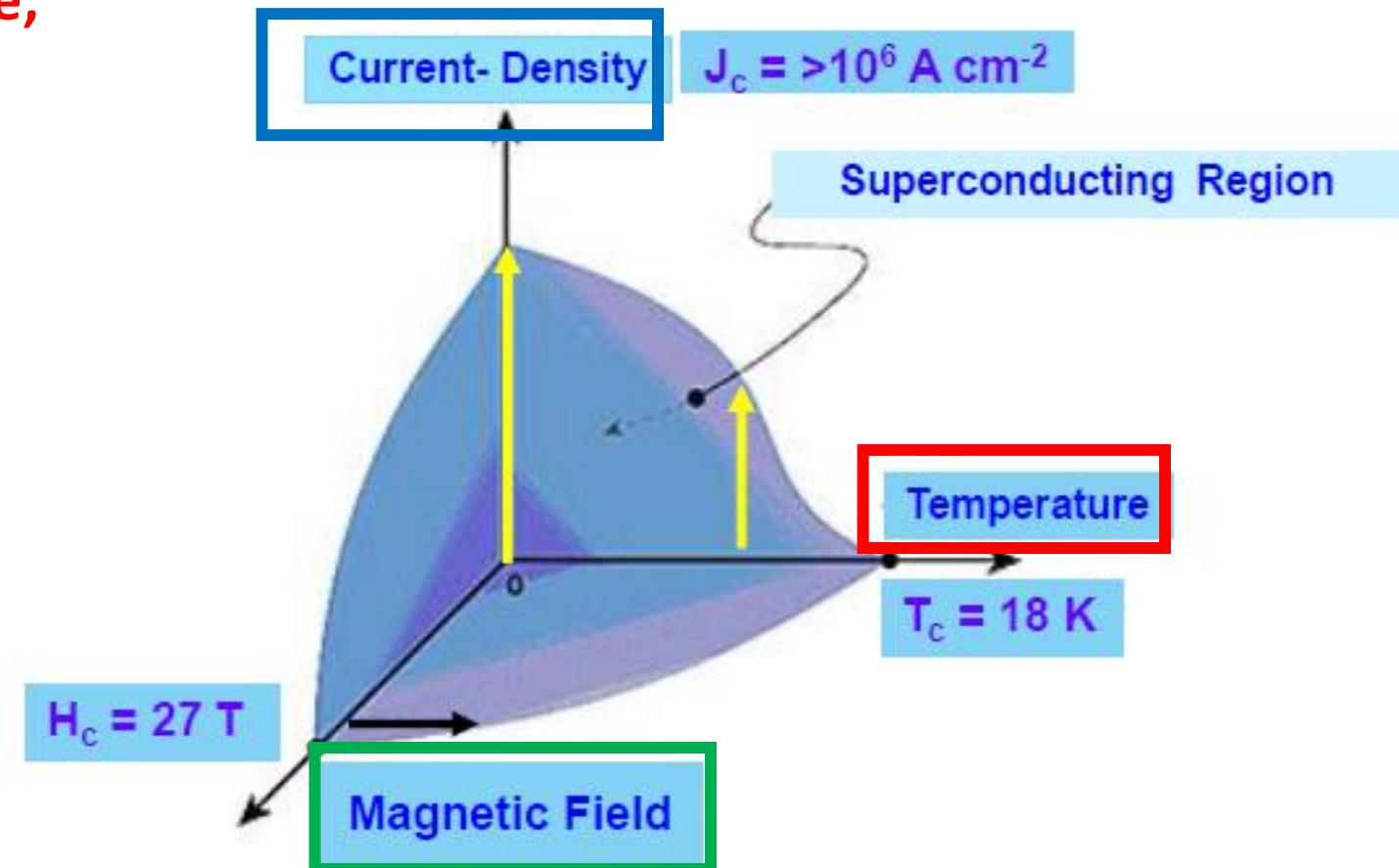
The region in which a superconductors can work without dissipation
is limited by the three parameters:

Nb₃Sn (NbTi)

The critical temperature;

The critical field;

The critical current.

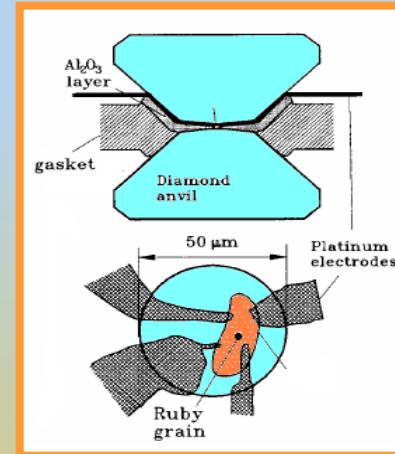
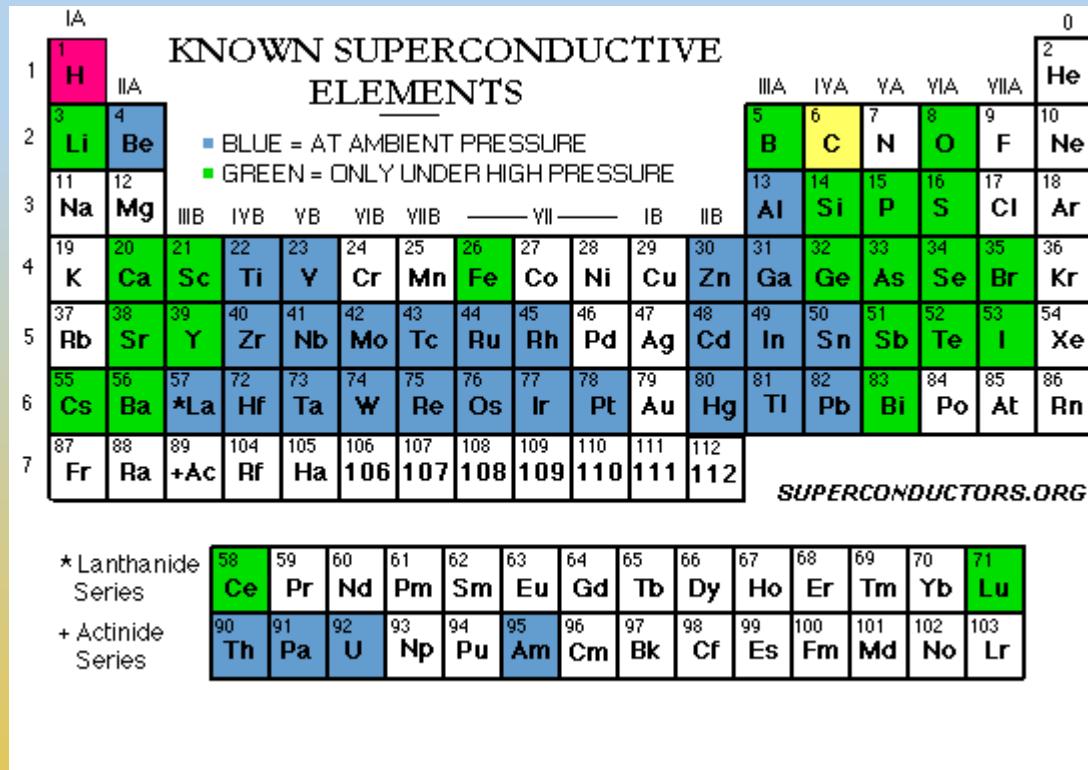


High Current and high Field Applications since 1961
J.E.Kunzler et al., Phys.Rev.Lett. 6:3 (1961) 89

The superconducting materials

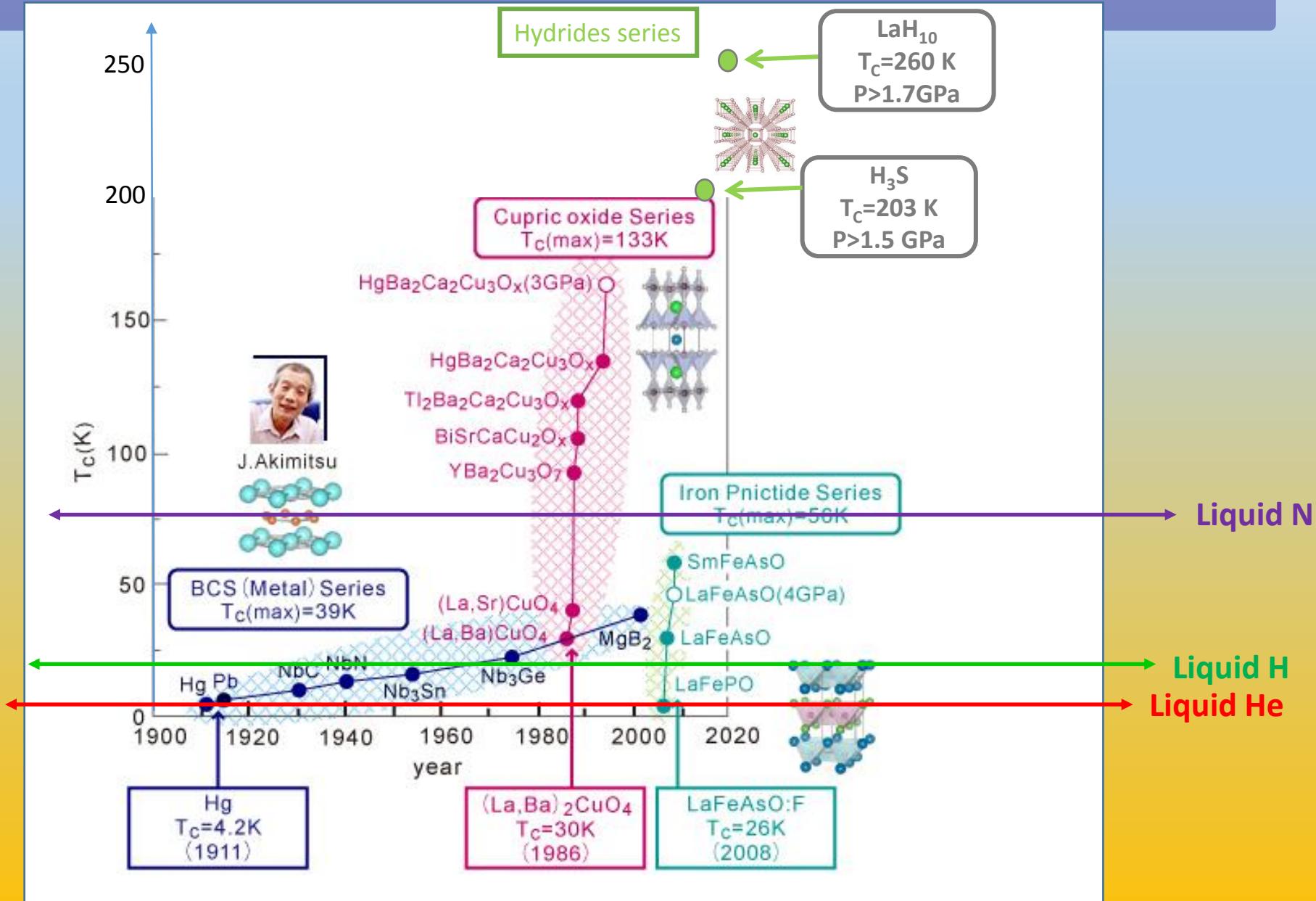
Superconductive elements

Future?



Record Pressure:
250 GPa
2,500,000 atmospheres

Story of the superconducting materials

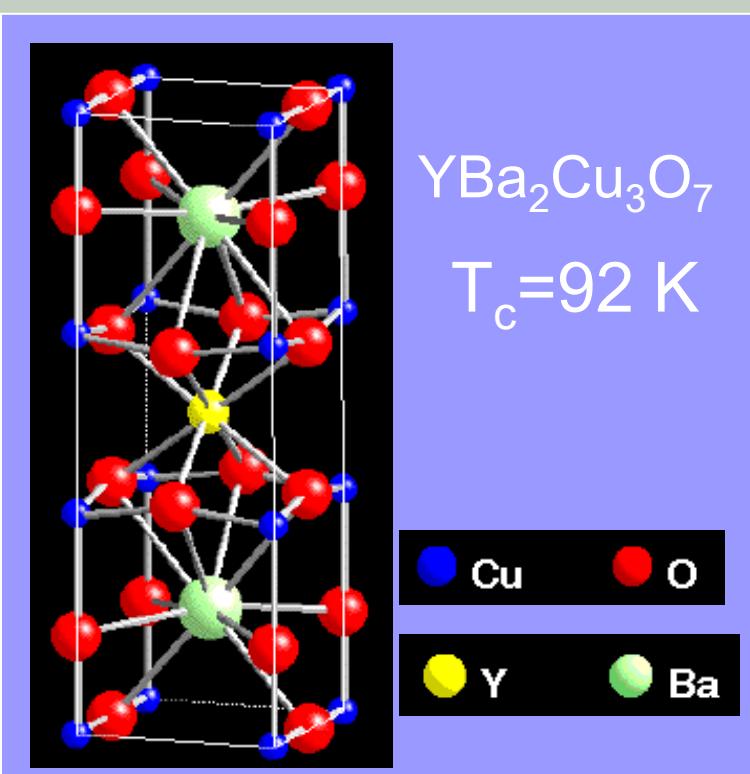


High Temperature superconductors (HTS)

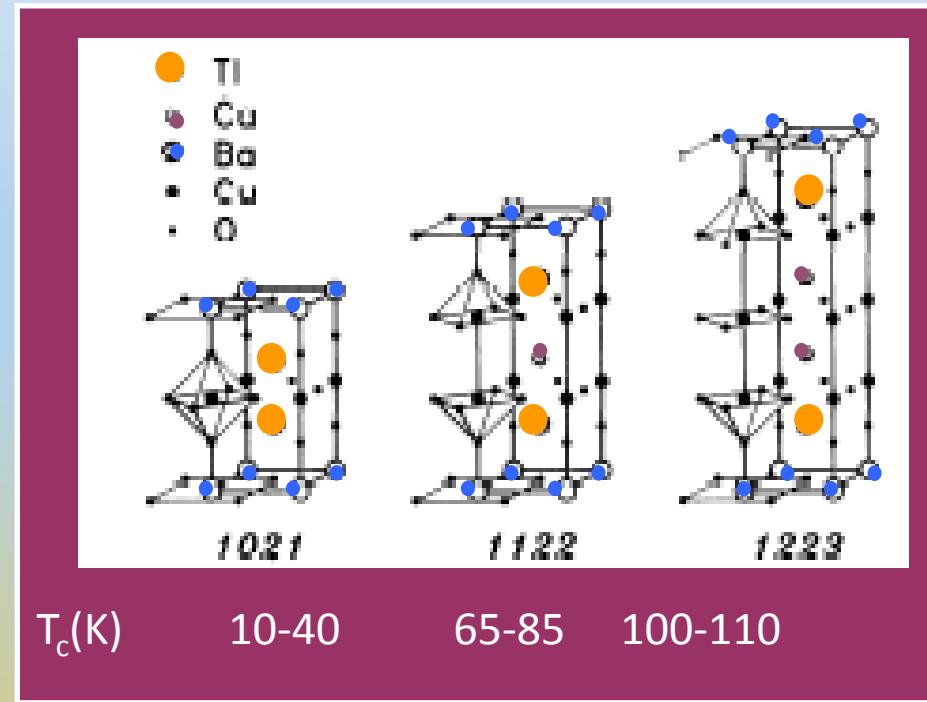
1986 Bednorz and Müller 1986



1987

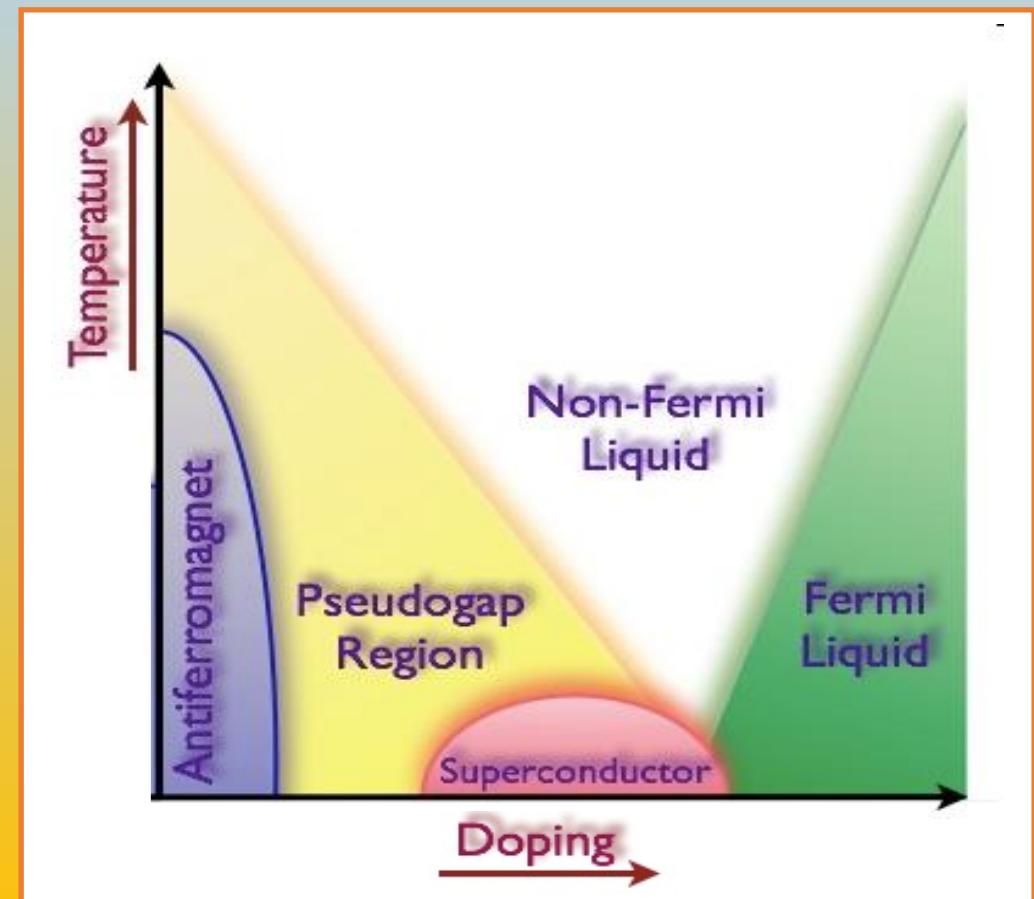


- Ceramic materials characterized by CuO_2 layers spaced by a reservoir block
- Conduction (and superconductivity) occurring in the CuO_2 layers
- Maximum $T_c = 165 \text{ K}$ in $\text{HgBa}_2\text{Ca}_2\text{Cu}_3\text{O}_{8+\delta}$ under pressure



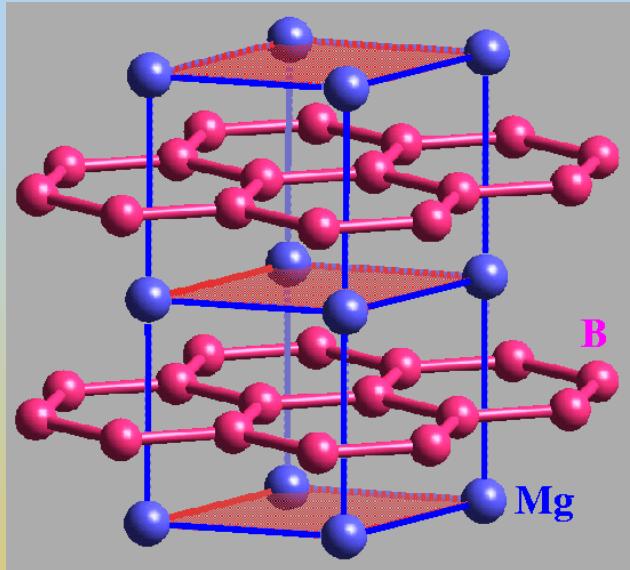
■ The increasing of CuO_2 layer increases T_c

- Complex phase diagram as a function of the doping in the CuO_2 planes
- The coupling between electron is determined by unconventional mechanism (not mediated by the phonons) not yet completely clarified



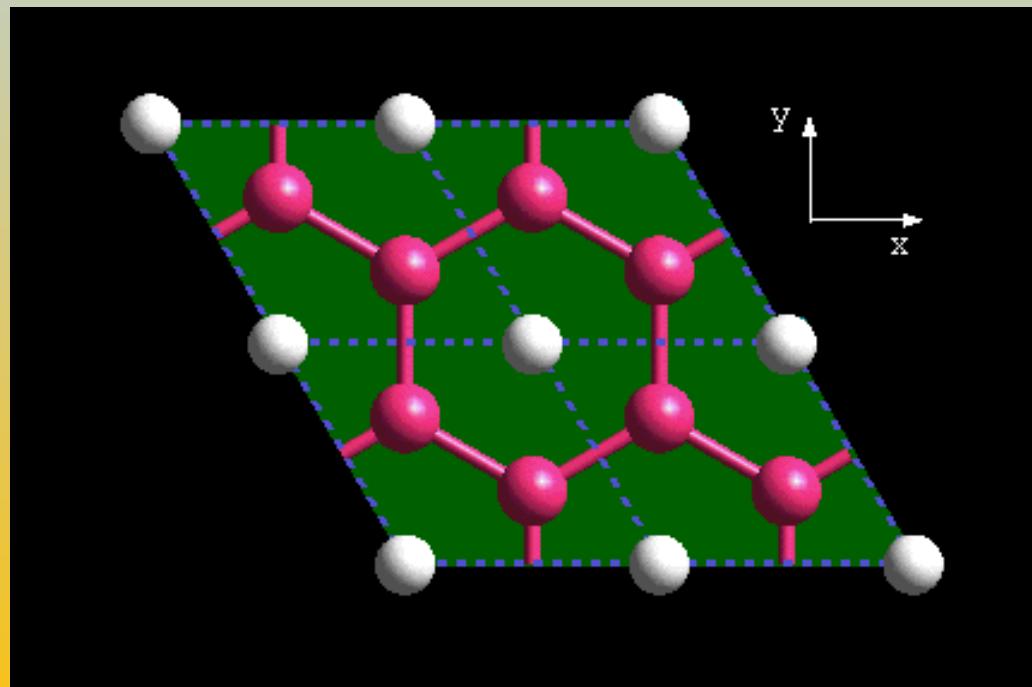
MgB₂

Akimitsu (2001)



- $T_c = 39$ K
- Layered structure
- Metallic system
- Electron pairing mediated by the optical mode of B ions
- Conventional superconductivity

$T_c = 39$ K is too near the maximum T_c predicted by the BCS theory!!



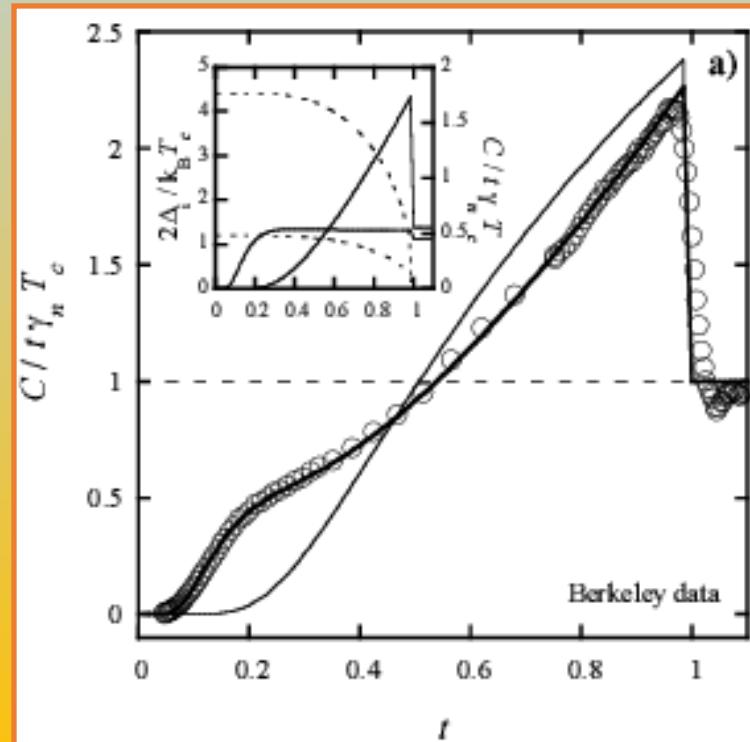
MgB₂ first example of a two-gap superconductor

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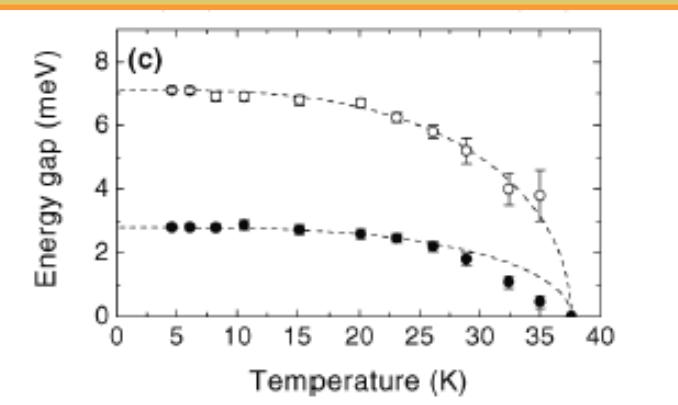
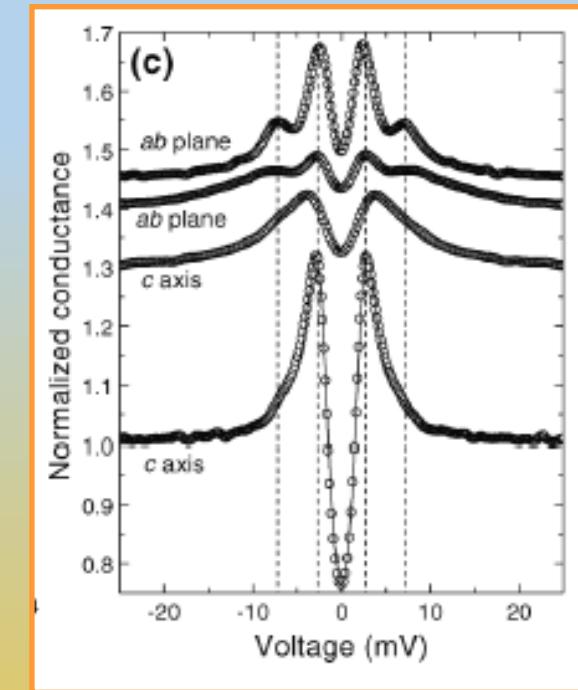
23 JULY 2001

Specific Heat of Mg¹¹B₂: Evidence for a Second Energy Gap

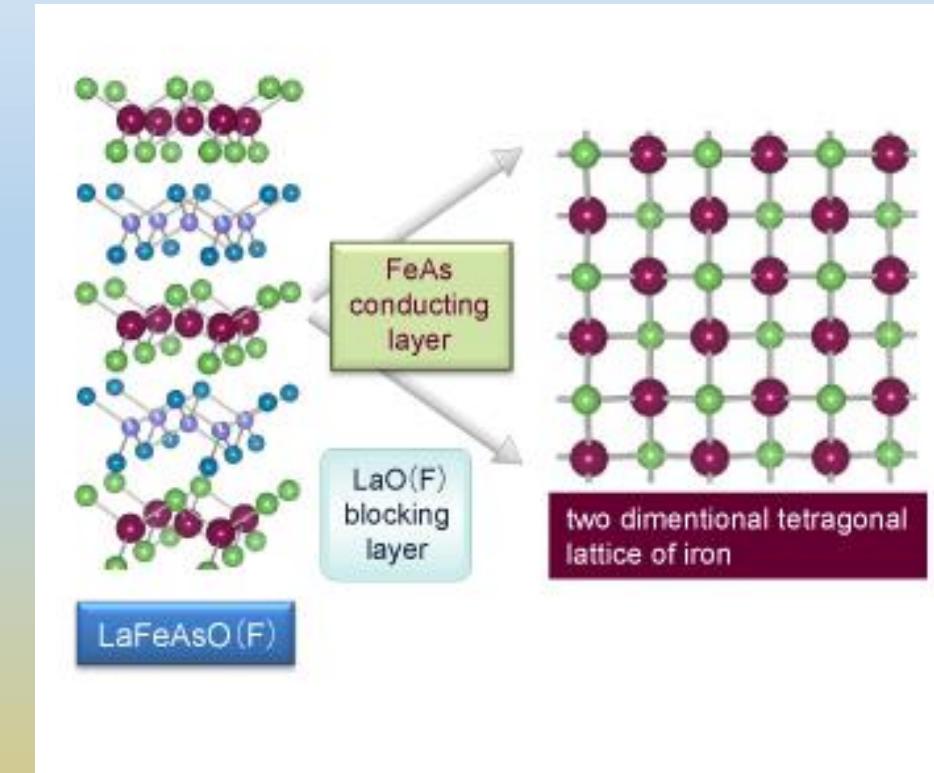
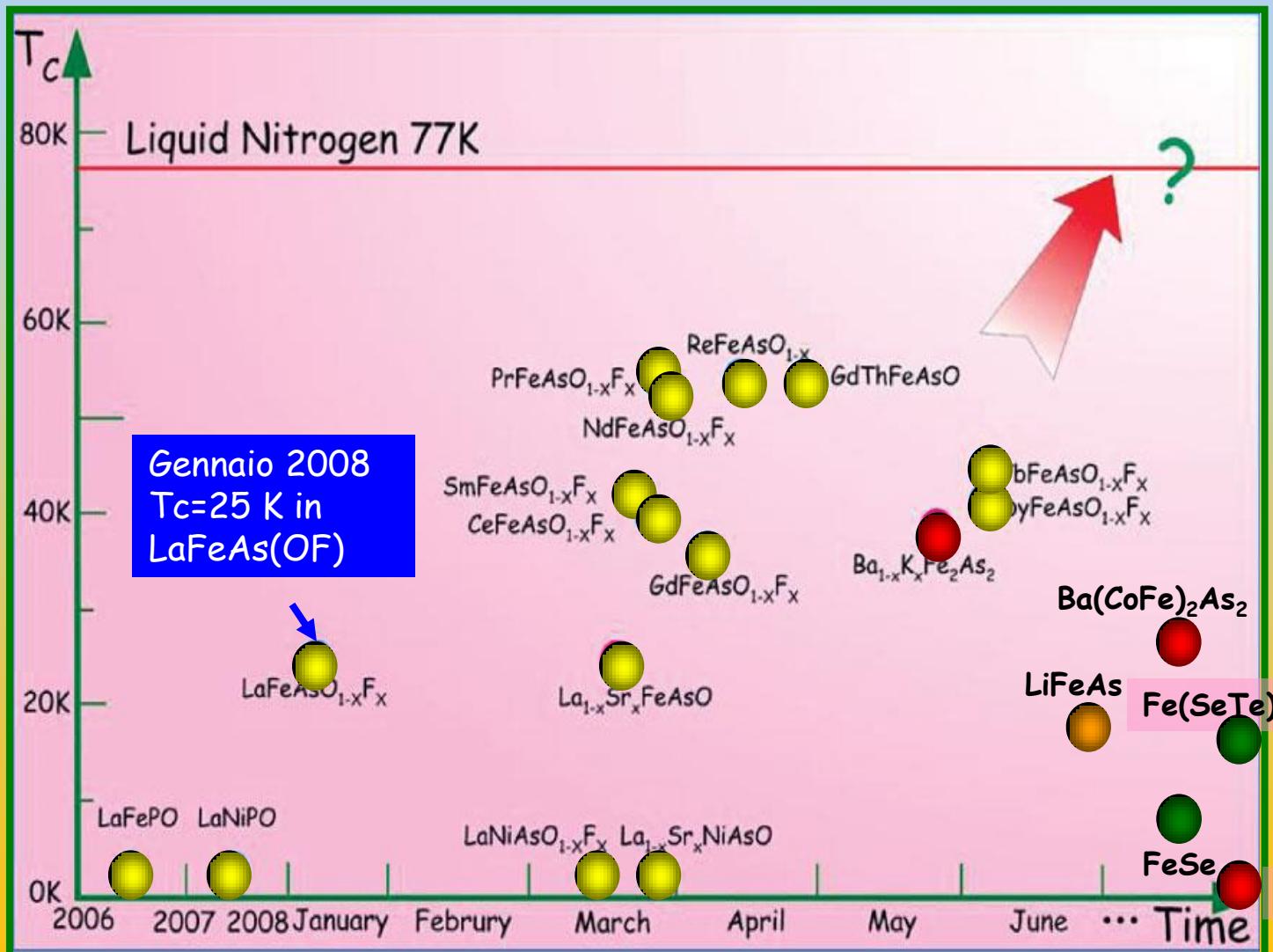


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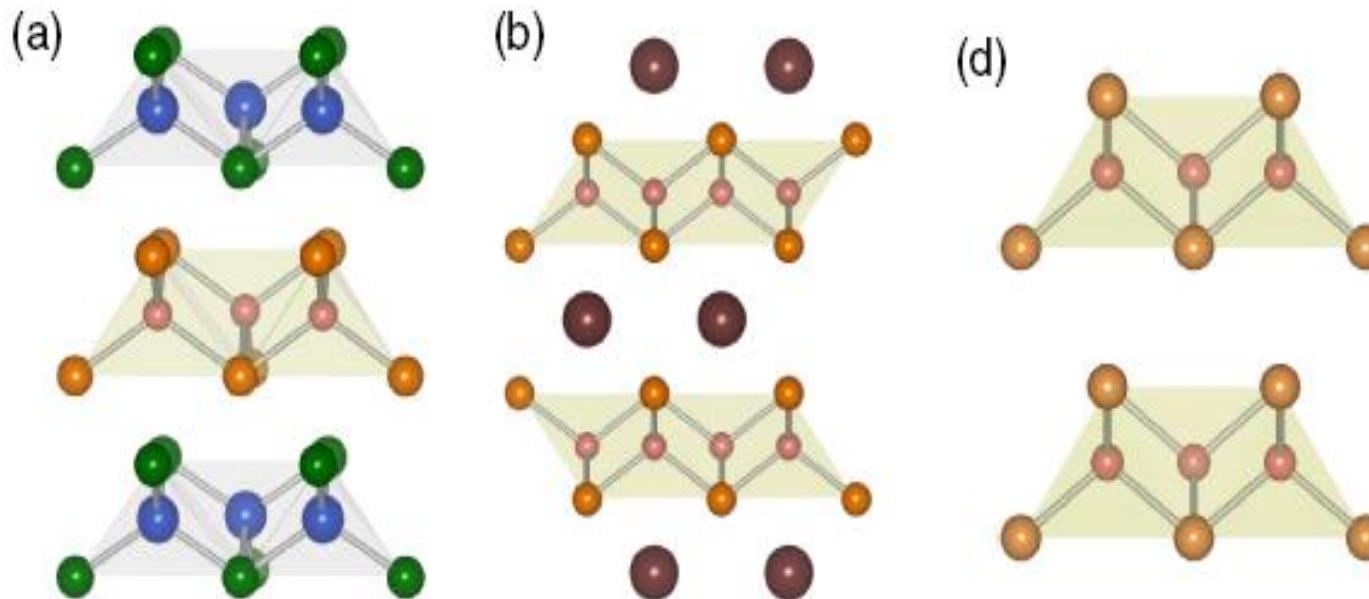
Direct Evidence for Two-Band Superconductivity in MgB₂ Single Crystals from Directional Point-Contact Spectroscopy in Magnetic Fields



Iron based superconductors



The most famous families



$REFeAsO$
1111

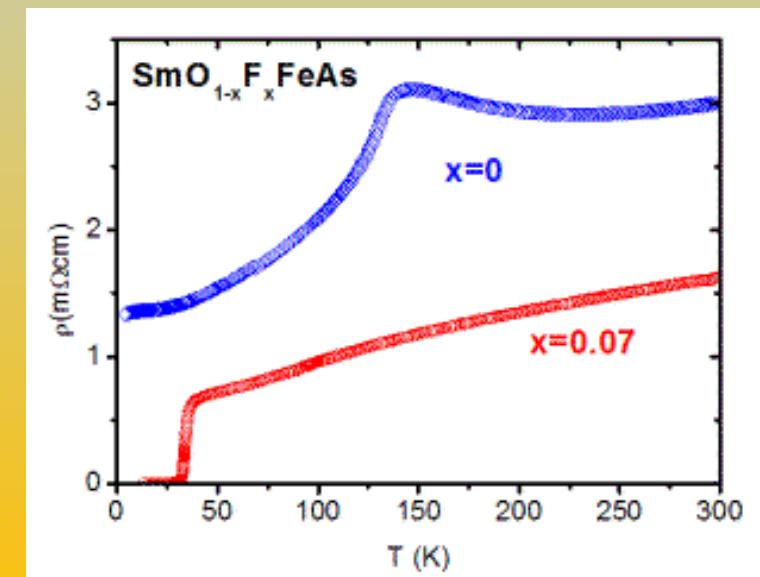
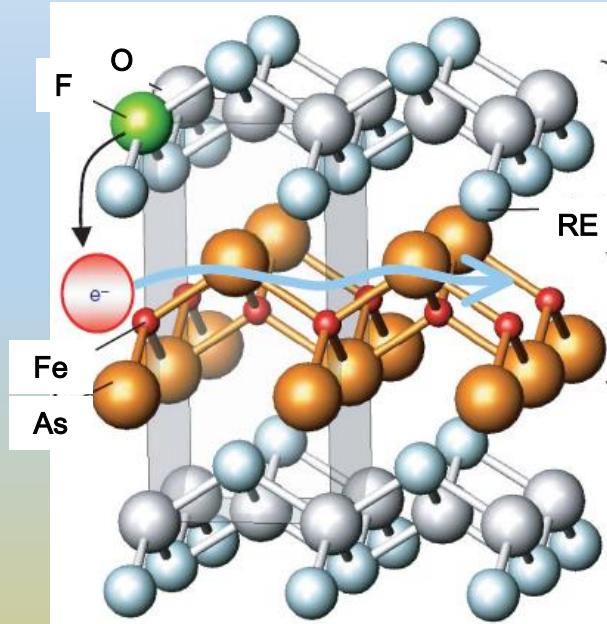
$T_{c,\max} = 58\text{ K}$

AFe_2As_2
122

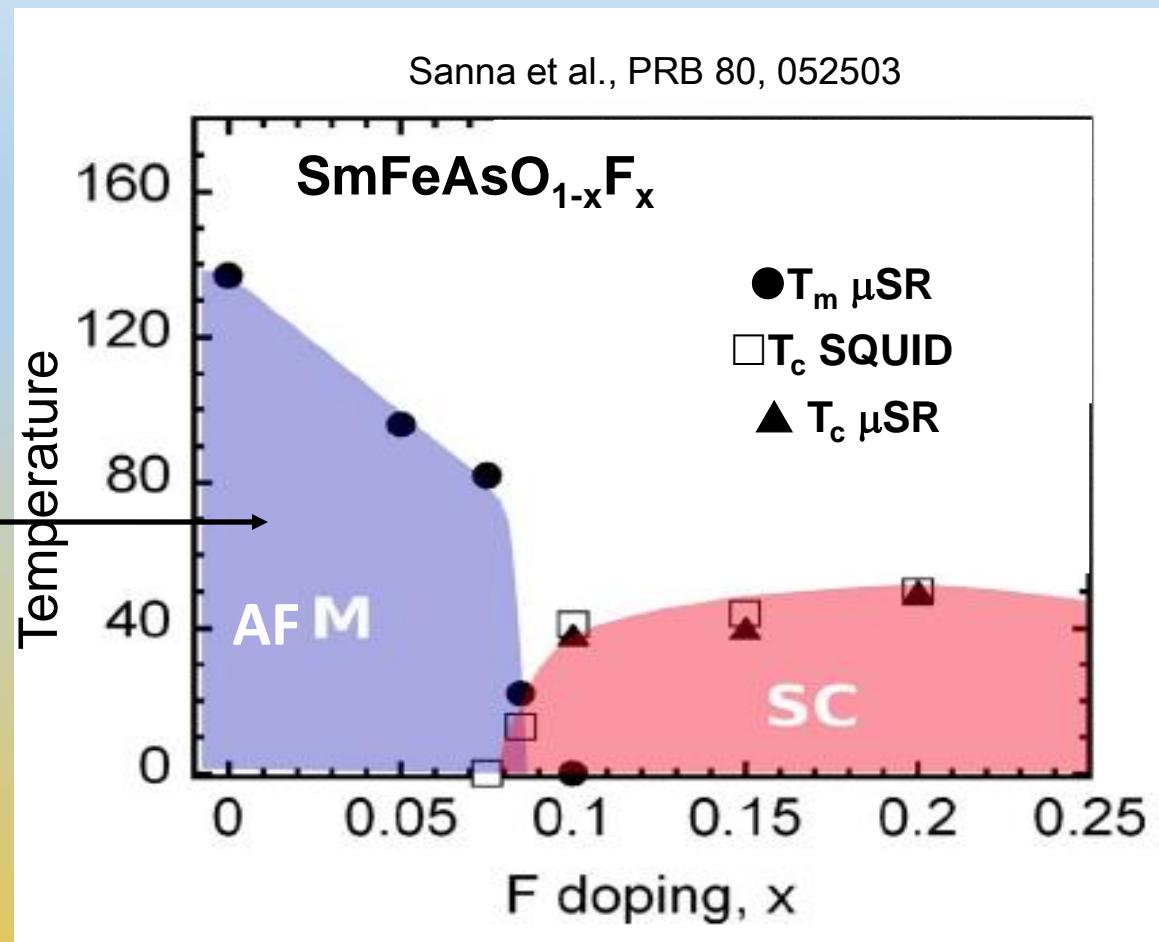
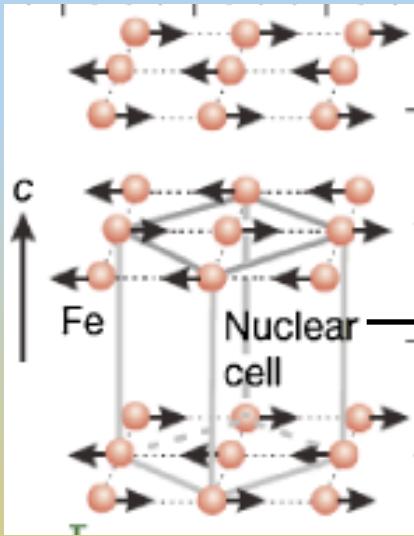
$T_{c,\max} = 40\text{ K}$

$Fe(Se,Te)$
11

$T_{c,\max} = 21\text{ K}$



Phase diagram as a function of doping

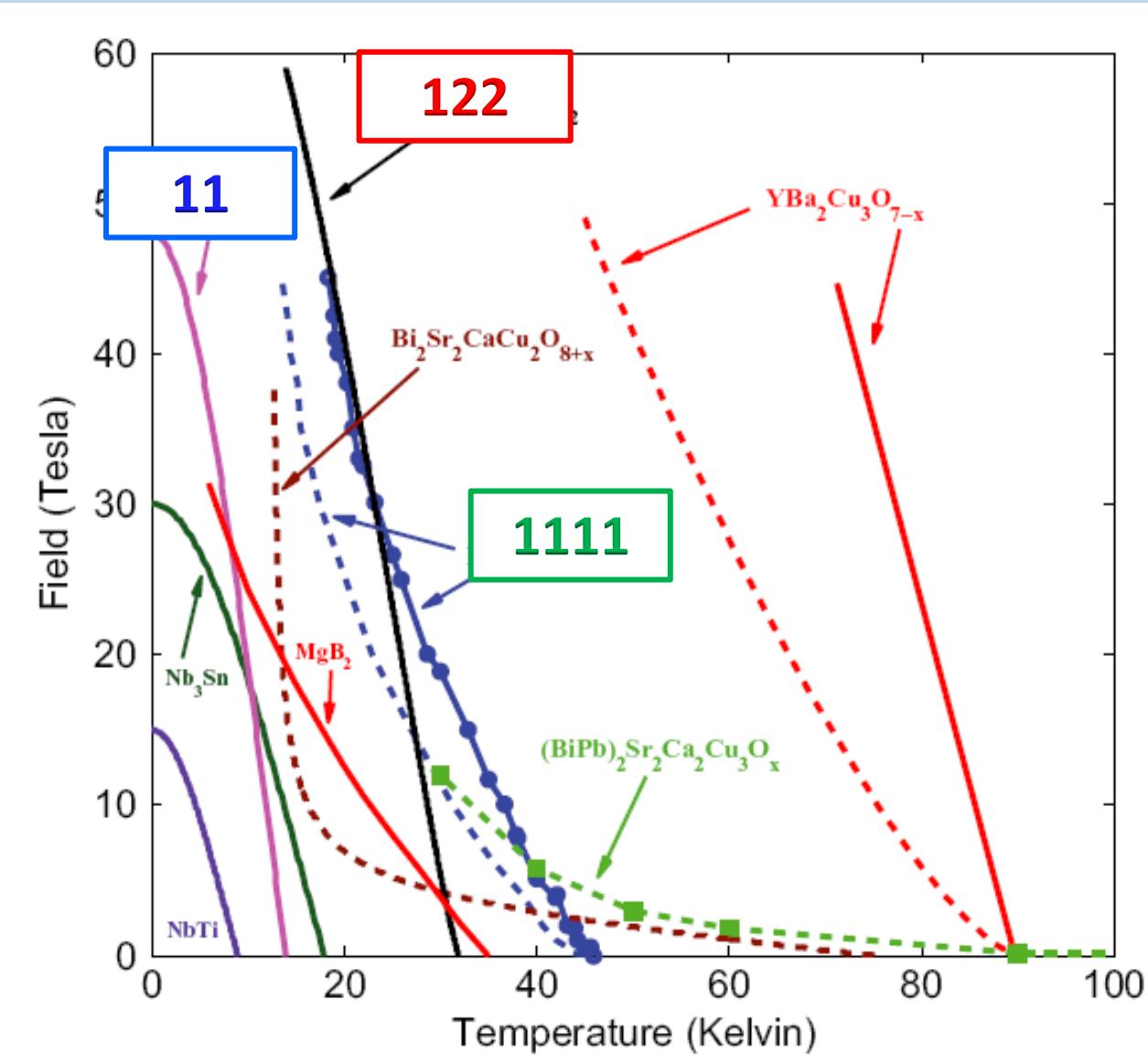


Analogies with
HTS
and with MgB_2

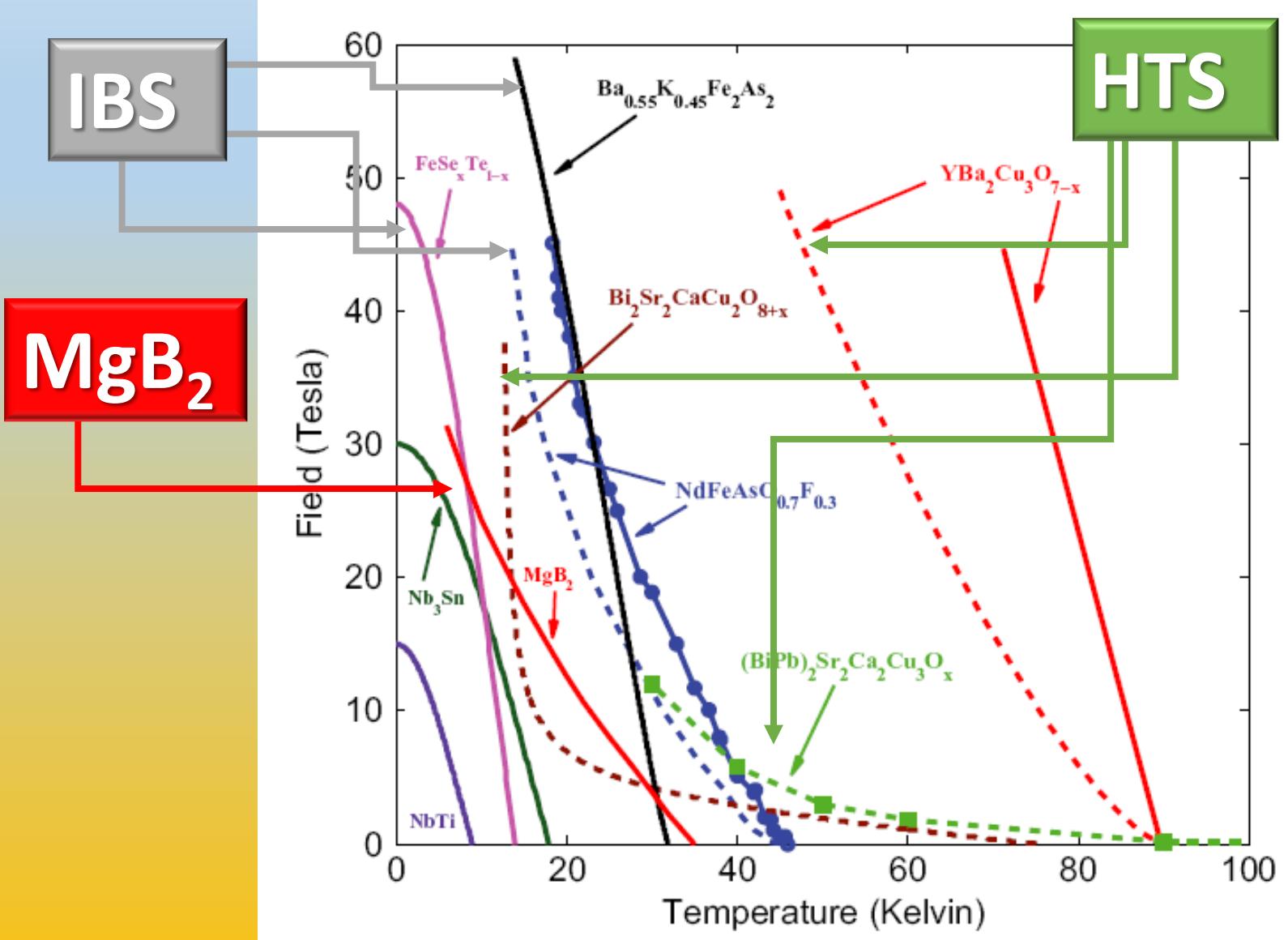
- superconductivity occurs upon doping
- proximity with magnetic ordering
- Anisotropic structure

- two-gap superconductivity

Upper critical Fields

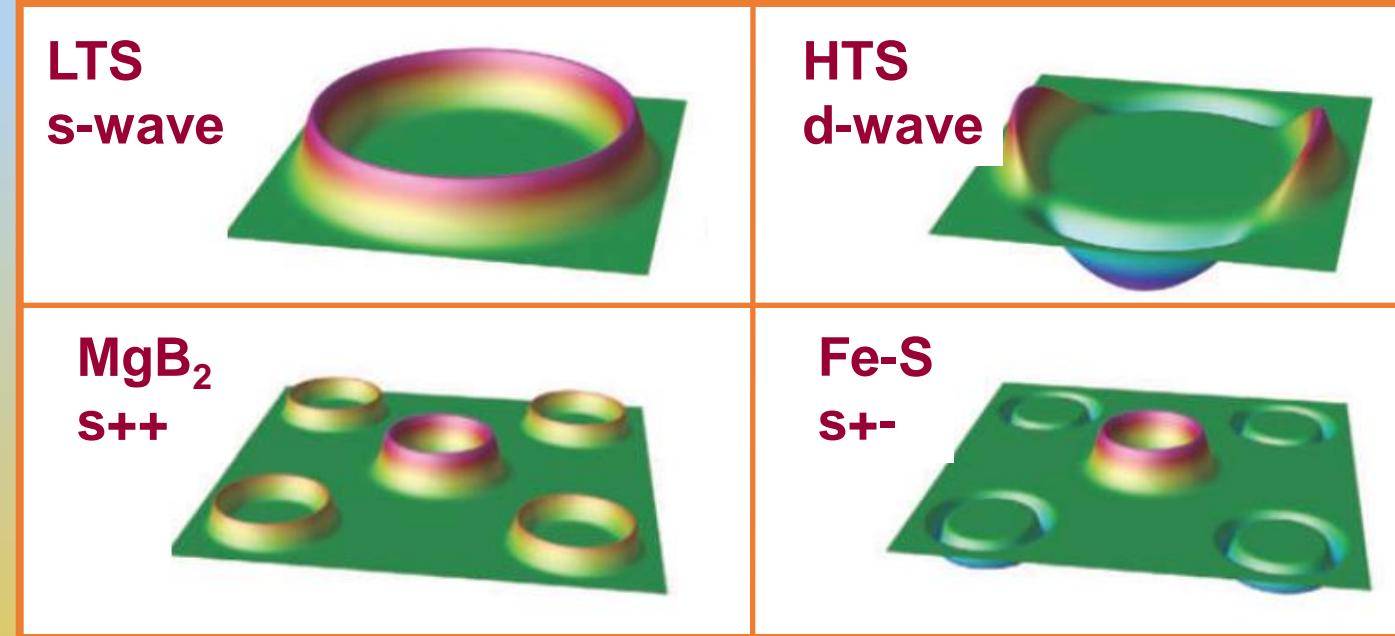


Upper critical Fields



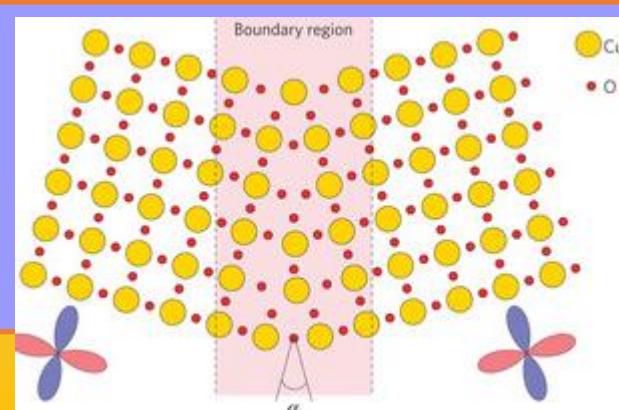
Peculiarities and issues of novel superconducting materials

1: Not conventional order parameter (Cooper pair wave function)



What does it imply
a not conventional order parameter ?

Weak coupling between grains



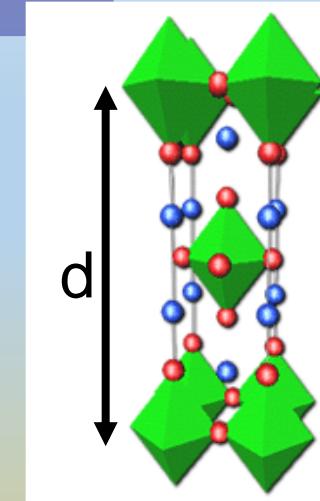
Peculiarities and issues of novel superconducting materials

2: Small coherence length

$$\xi \propto \frac{n_s^{1/2}}{T_c}$$

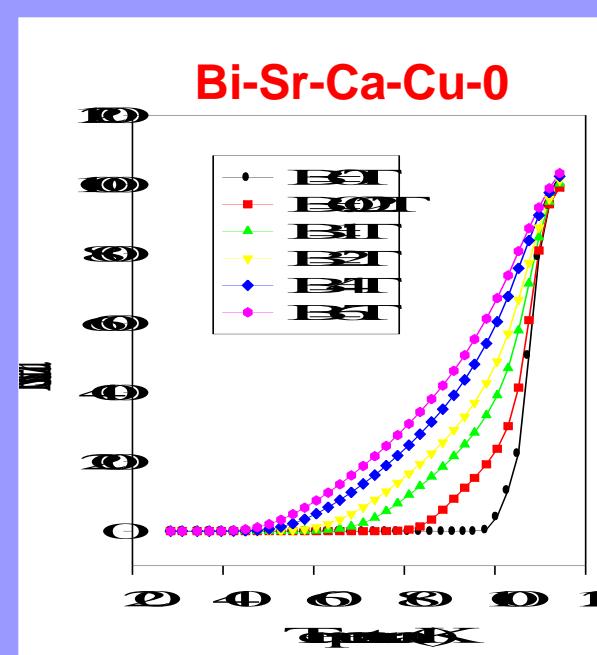
3: Large structural anisotropy

$$\xi_c < \frac{\xi_{ab}}{\gamma} < d$$



**Small ξ
and large anisotropy:
What does it imply ?**

Giant dissipation in presence
of an applied magnetic field





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THANK YOU FOR YOUR ATTENTION