

Scientific Computation

Spring, 2019

Lecture 15

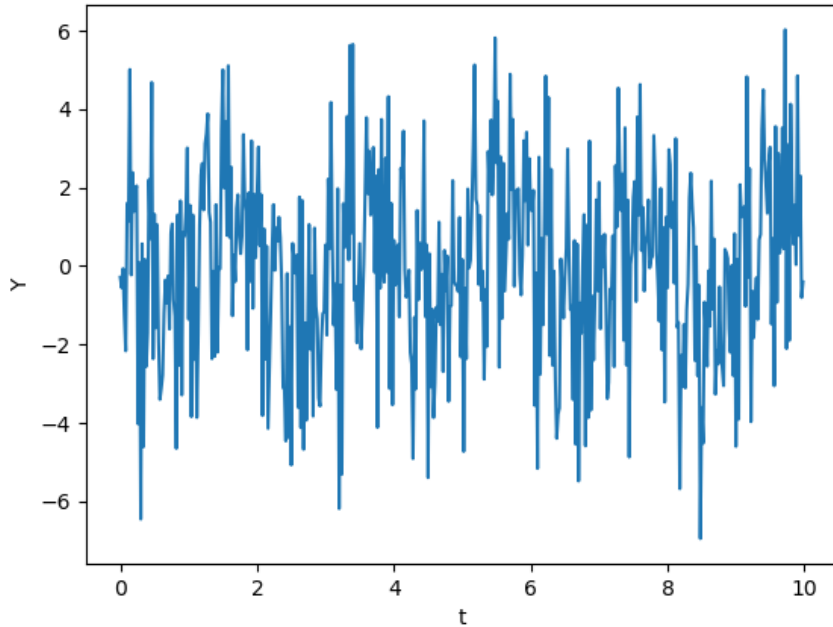
Mastery material

- The mastery material has a reading component (and coding+discussion components)
- I would like to release the reading component this Thursday along with HW3
- Please let me know as soon as possible if you object to this

Data analysis

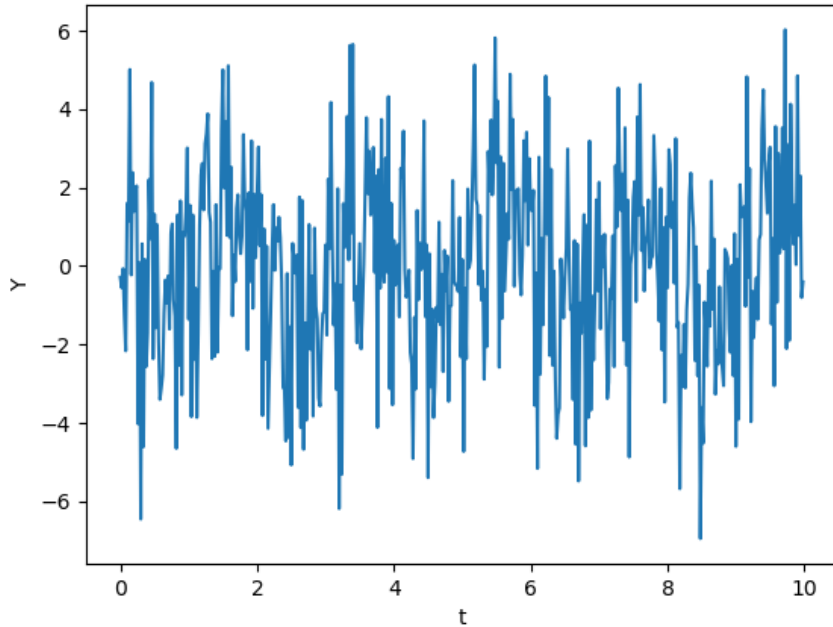
- Last few lectures: analyze datasets arranged as tables or matrices
- Today (and tomorrow): Something simpler! We'll look at data arranged in 1D arrays
- Typical example, a time series: $f(t_i)$, $t_i = i \Delta t$, $i = 0, 1, \dots, N_t - 1$
- Could also have data in space (along a line): $f(x_i)$, $x_i = i \Delta x$, $i = 0, 1, \dots, N_x - 1$
- How do we extract trends or, more generally, “information”, from such data?

An example



- What should we do with a signal that looks like this?
- First step is to build a description
 - E.g. compute mean and rms (variance)
- But what next?

An example



- What should we do with a signal that looks like this?
- First step is to build a description
 - E.g. compute mean and rms (variance)
- But what next?
 - We can think about the *spectrum*
 - Underlying idea: decompose signal into waves with different frequencies
 - And check which frequencies hold the most “energy”

Fourier series

- **Periodic functions can be expanded as Fourier series:**

$$f(x) = \sum_{n=-\infty}^{\infty} c_n \exp(inx)$$

$$c_n = (1/2\pi) \int_{-\pi}^{\pi} f(x) \exp(-inx) dx$$

- **This expression considers a period = 2π**
- **This can be modified to L through a simple change of variables ($y = xL/(2\pi)$)**
- **Convergence theory considers general functions integrable on the interval $[-\pi, \pi)$**
- **In practice, Fourier series are only used for periodic functions**
- **For *infinitely differentiable functions* (e.g. a Gaussian), we have exponential convergence:** $c_n \sim \exp(-\mu|n|)$, μ is a positive constant

Fourier series

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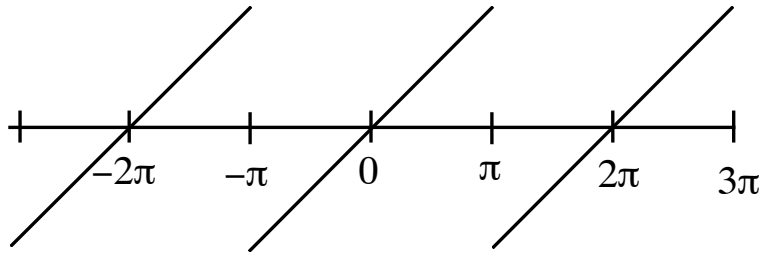
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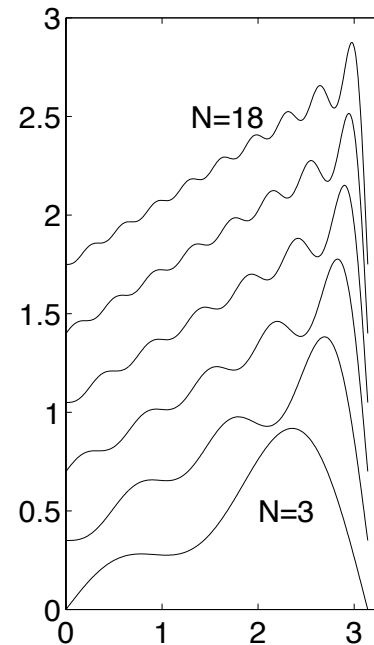
- **We will largely ignore the formal convergence theory and try to build intuition**
- **The rate-of-convergence of the series depends on the smoothness of the function**
 - **In the interior of the domain, we will simply assume the function and all of its derivatives are continuous**
 - **Let's look at a few examples with non-smooth behavior at the boundaries...**

Fourier series

- **Example 1: sawtooth function:**

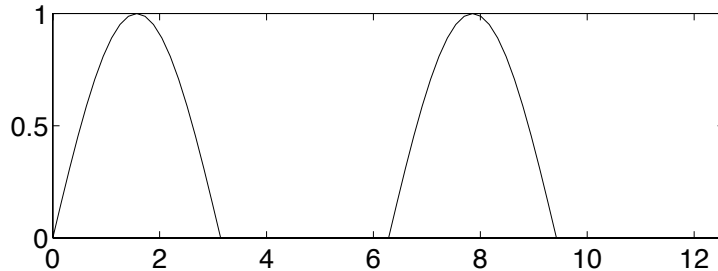


- **Function is discontinuous at “boundaries”:** $f(-\pi) \neq f(\pi)$
- **“Convergence” is slow,** $|c_n| = \frac{1}{|n|}$
- **Figure shows approximation retaining 1st 2N terms in series ($|n| < N$)**



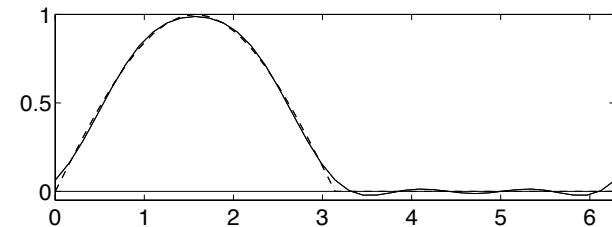
Fourier series

- **Example 2: half-wave rectifier:**



$$f(t) \equiv \begin{cases} \sin(t), & 0 < t < \pi \\ 0, & \pi < t < 2\pi \end{cases}$$

- **Function's derivative is discontinuous at “boundaries”:** $f'(-\pi) \neq f'(\pi)$
- **Convergence is slow,** $|c_n| \sim \frac{1}{|n|^2}$
- **Figure shows approximation retaining 1st 4 terms in series (dashed curve is the approximation):**



Fourier series

- **A general result:**

- *If:*

- 1.

$$f(\pi) = f(-\pi), f^{(1)}(\pi) = f^{(1)}(-\pi), \dots, f^{(k-2)}(\pi) = f^{(k-2)}(-\pi)$$

2. $f^{(k)}(x)$ is integrable

then the coefficients of the Fourier series

$$f(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

have the upper bounds

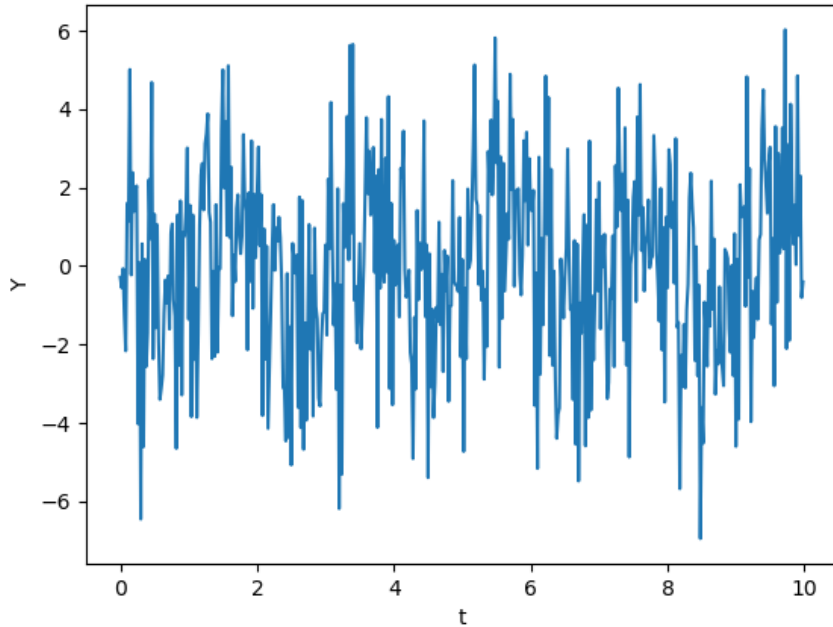
$$|a_n| \leq F/n^k; \quad |b_n| \leq F/n^k$$

for some sufficiently large constant F , which is independent of n .

- **Proof: repeated application of integration by parts**

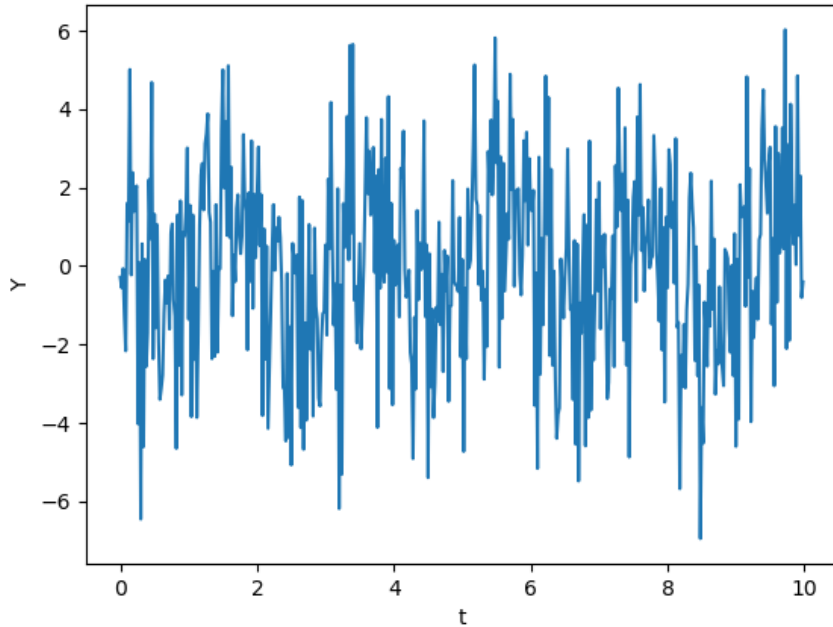
- **Note: this definition of a Fourier series is equivalent to previous with complex exponential**

An example



- This signal is periodic ($T=10$)
- How do we compute its Fourier coefficients?

An example



- This signal is periodic
- How do we compute its Fourier coefficients?
- We will use the Discrete Fourier Transform (DFT)
- Available in `numpy.fft`, and `scipy.fftpack`

Discrete Fourier transform

Now, our function is represented on a N-point discrete, equispaced grid:

$$t_j = j\Delta t, j = 0, 1, \dots, N - 1$$

We have a truncated Fourier series at the jth point:

$$f(t_j) = \sum_{n=-N/2}^{N/2-1} c_n \exp(i2\pi n t_j / T) = \sum_{n=-N/2}^{N/2-1} c_n \exp(i2\pi j n / N)$$

with $N\Delta t = T$, and the inverse transform is now a discrete sum:

$$c_n = \frac{1}{N} \sum_{j=0}^{N-1} f_j \exp(-i2\pi n t_j / T) = \frac{1}{N} \sum_{j=0}^{N-1} f_j \exp(-i2\pi j n / N)$$

- **np.fft.fft computes c, with the (1/N) factor omitted, but returns it in a ... strange order:**

$$\text{np.fft.fft}(f) = c_{\text{np}} = N \left(c_0, c_1, \dots, c_{N/2-1}, c_{-N/2}, c_{-N/2+1}, \dots, c_{-1} \right)$$

Discrete Fourier transform

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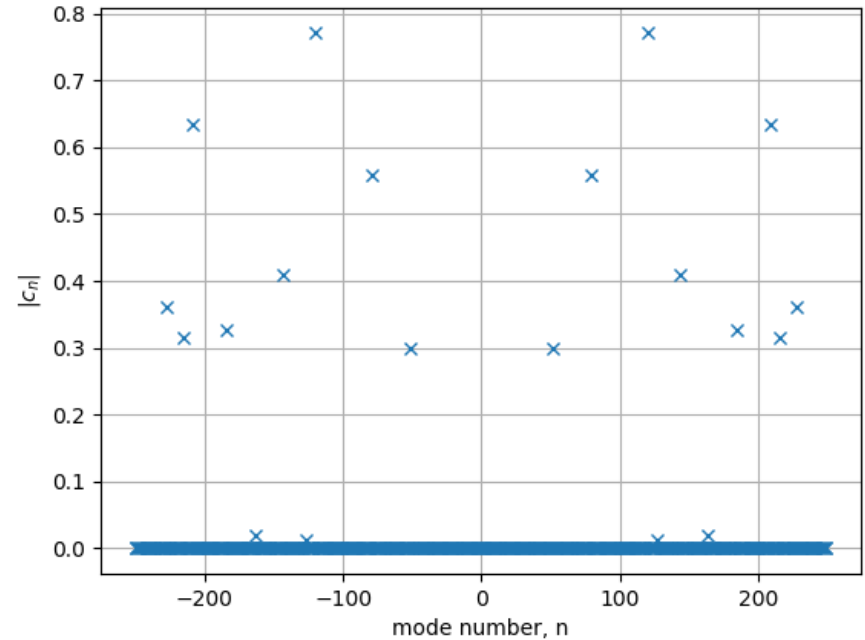
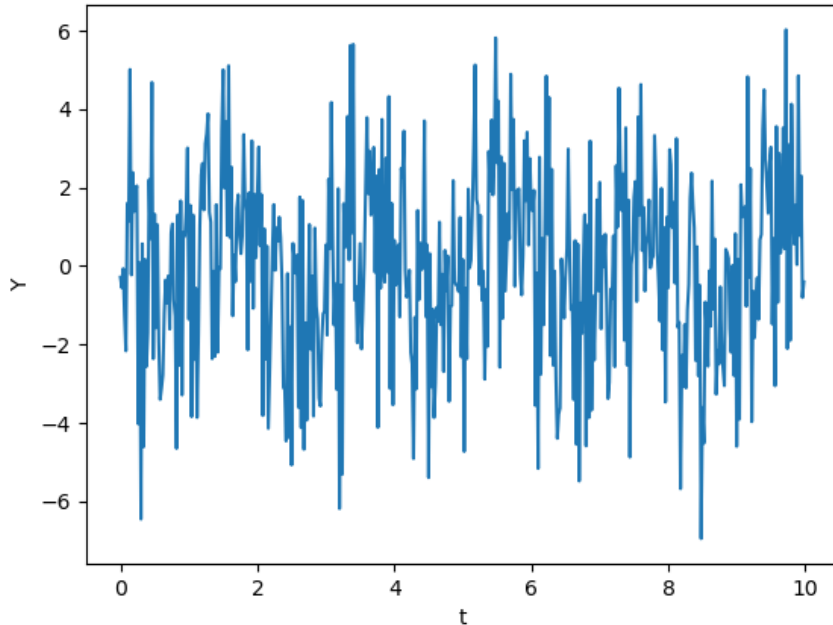
$$\text{np.fft.fft}(f) = c_{\text{np}} = N (c_0, c_1, \dots, c_{N/2-1}, c_{-N/2}, c_{-N/2+1}, \dots, c_{-1})$$

- **But we have a tool to help:**

$$\text{np.fft.fftshift}(c_{\text{np}})/N = c_{-N/2}, c_{-N/2+1}, \dots, c_{N/2-1}$$

- **Let's go back to our example!**

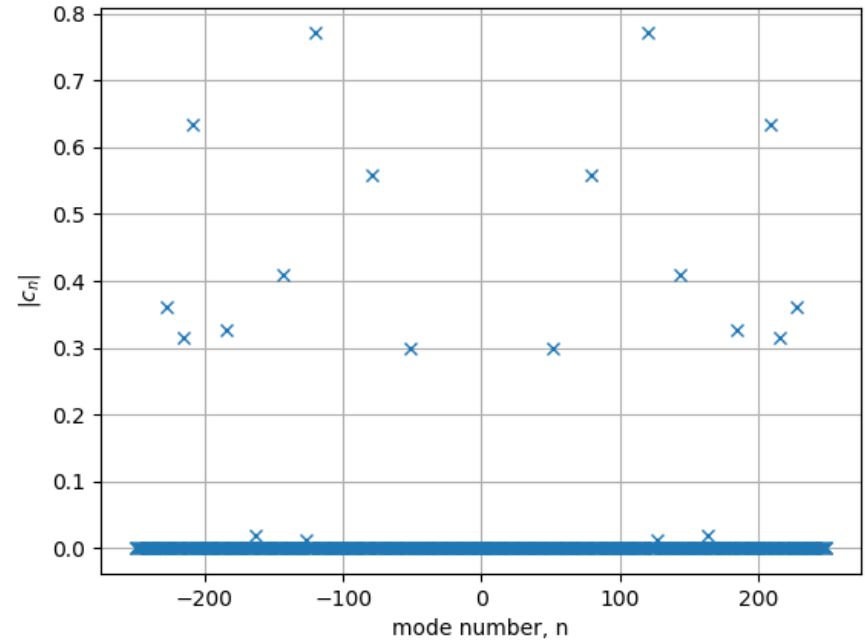
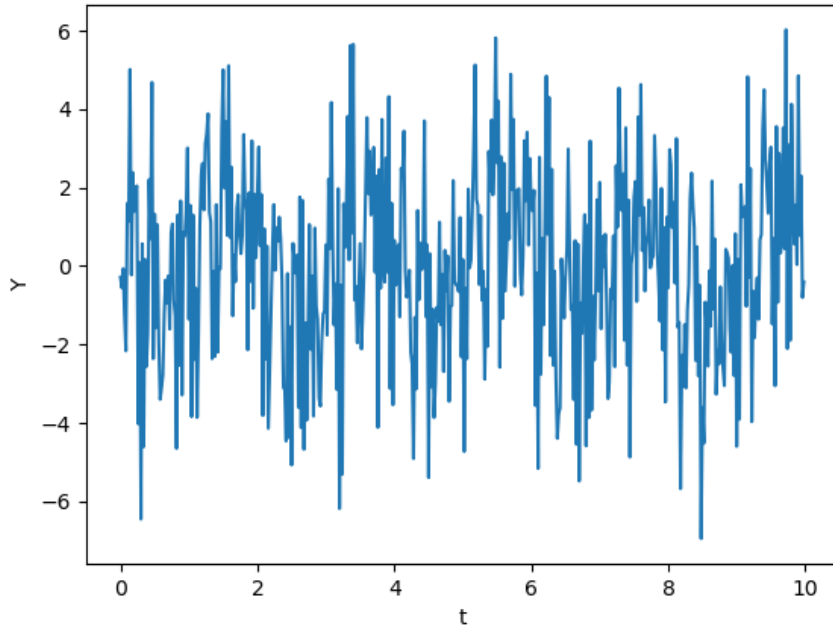
An example



```
c = np.fft.fft(Y)
c = np.fft.fftshift(c)/Nt
n = np.arange(-Nt/2, Nt/2)
```

```
plt.figure()
plt.plot(n, np.abs(c), 'x')
plt.xlabel('mode number, n')
plt.ylabel('$|c_n|$')
plt.grid()
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$$Y = \sum_{m=1}^{10} a_m \sin(2\pi f_m t + \phi_m)$$

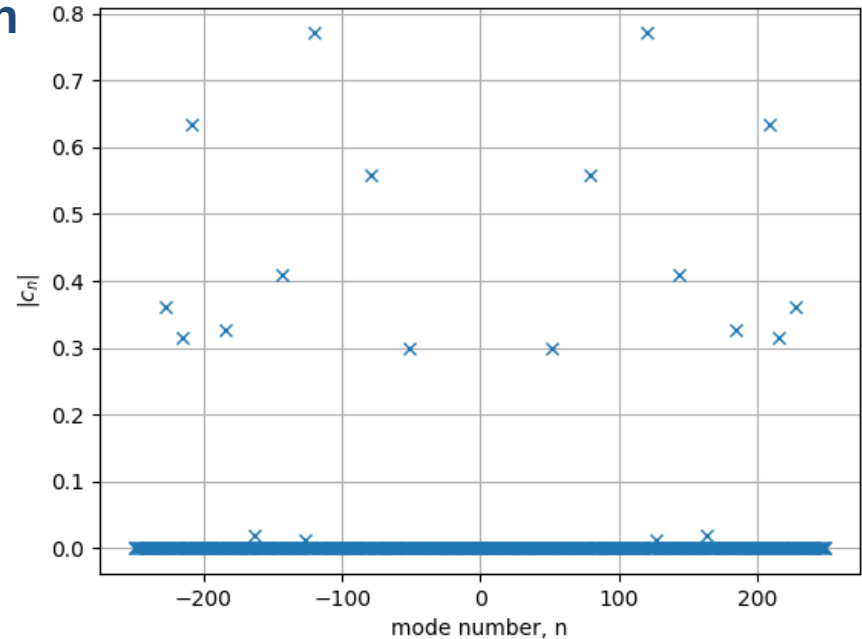
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- Each f_m is a *frequency*
- $|c_n|^2$ is the “energy” in a mode with frequency = n/T
- For real-valued data: $c_n = c_{-n}^*$, so only $n \geq 0$ needed to be shown in figure (should use rfft instead of fft)
- It’s essential that the timespan of the signal is an integer multiple of each frequency

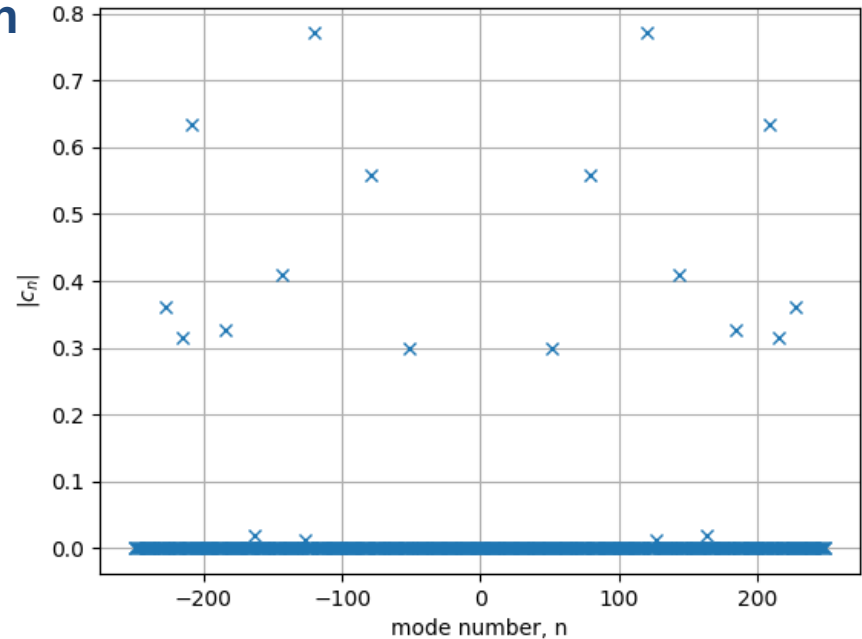


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- $(Nt/2-1)/T = \Delta t/2-1/T$ is the highest frequency that can be “resolved”
- Rule-of-thumb: >2 points/period of highest-frequency component are needed (here, period = $1/f$)

Another example

- Let's take two sine waves, with frequency $f=2/T$ and $2.5/T$. Here, T is the timespan of the signal and $1/f$ is the period of the wave:
- Problem setup:

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In [3]: T = 5
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```
In [4]: Nt = 100
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```
In [5]: t = np.linspace(0,T,Nt+1)
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In [6]: t = t[:-1]
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In [7]: f1 = 2/T
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In [8]: f2 = 2.5/T
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In [10]: Y1 = np.sin(2*np.pi*f1*t)
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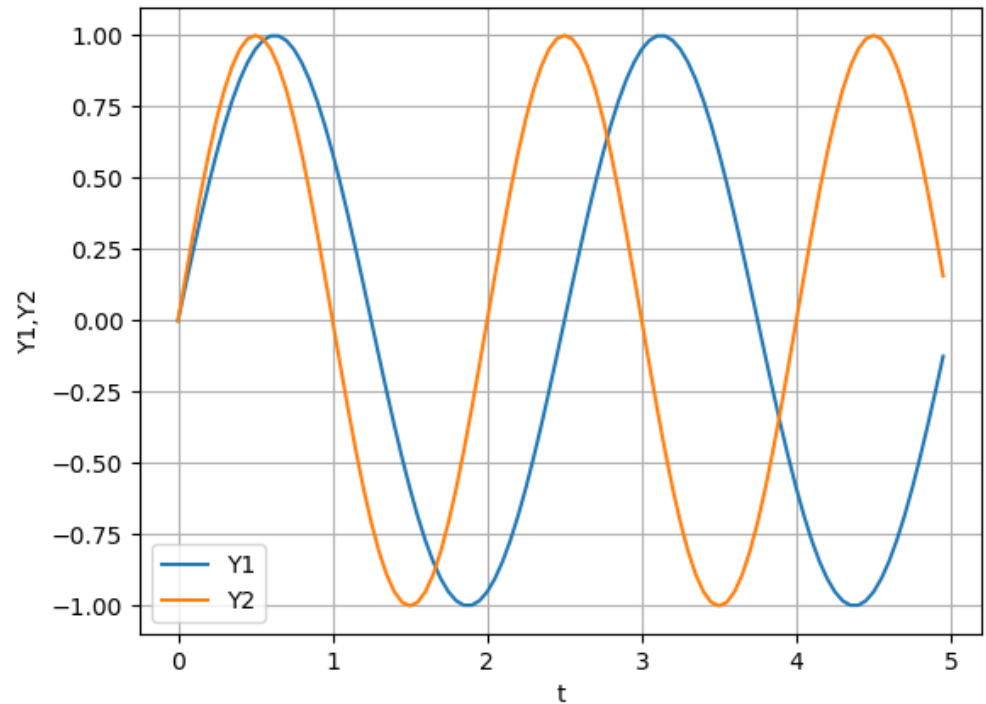
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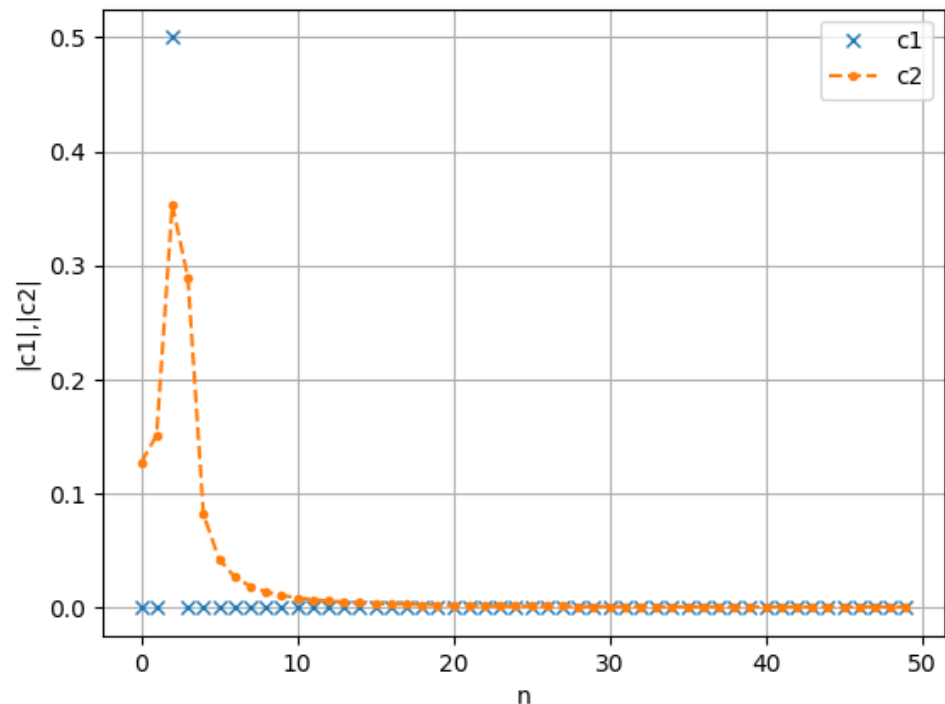
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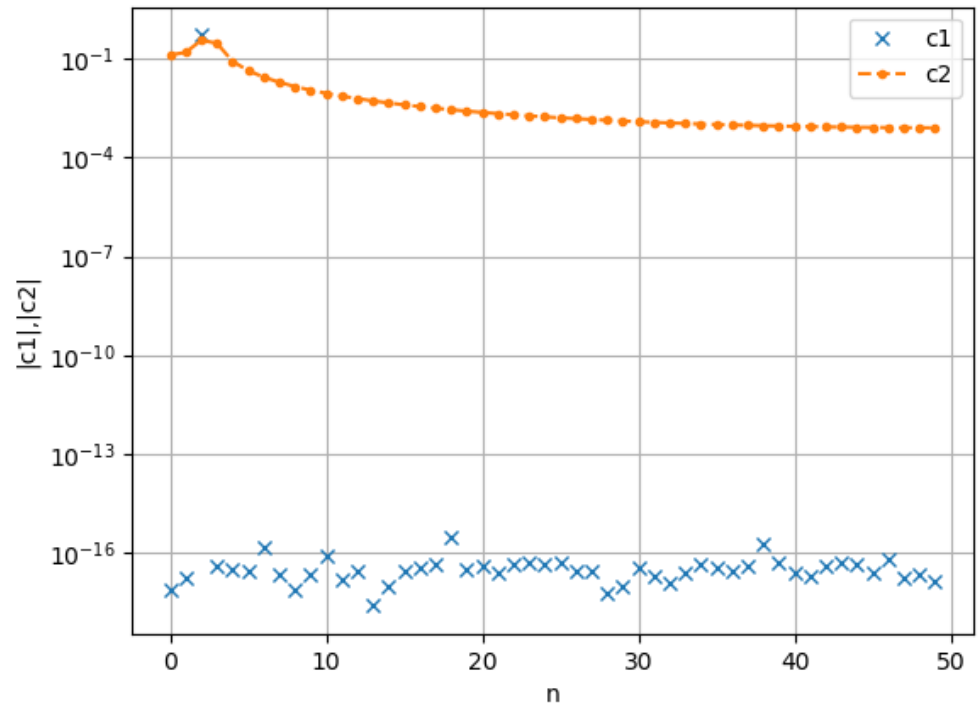
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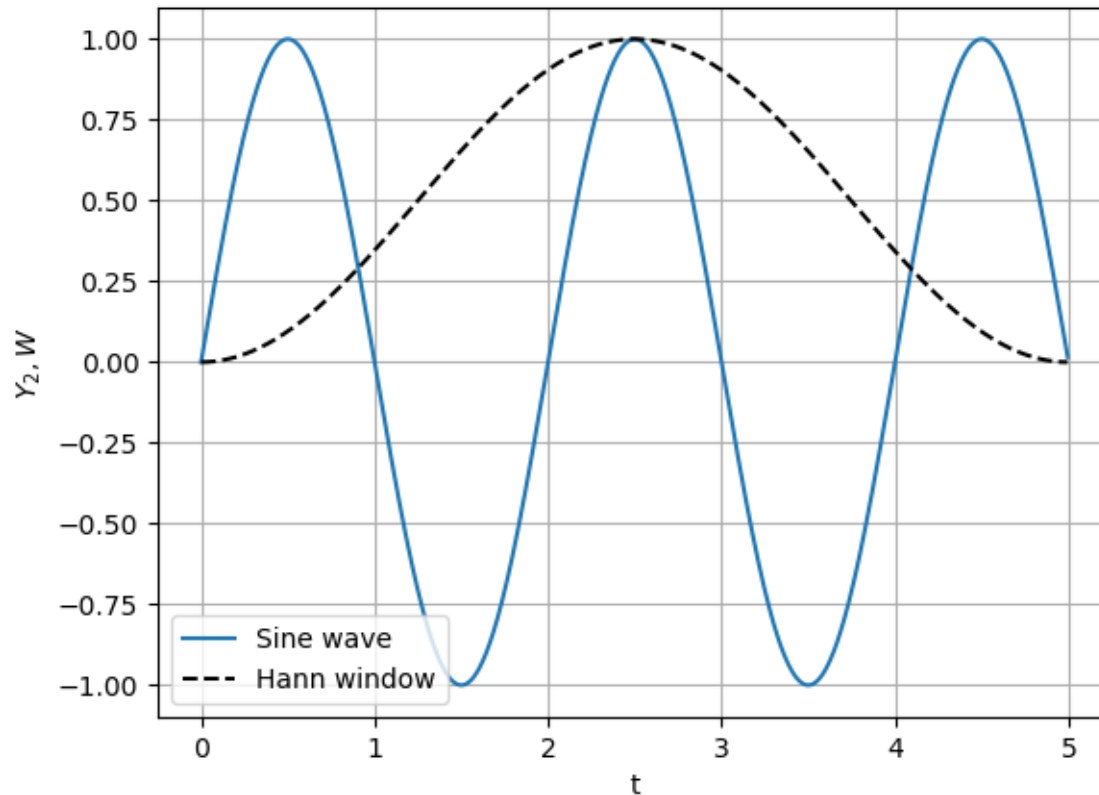
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The 2 waves, note the “spread” of the energy across several frequencies for the 2nd wave and the (very) slow decay due to the “discontinuity” at $t=5$

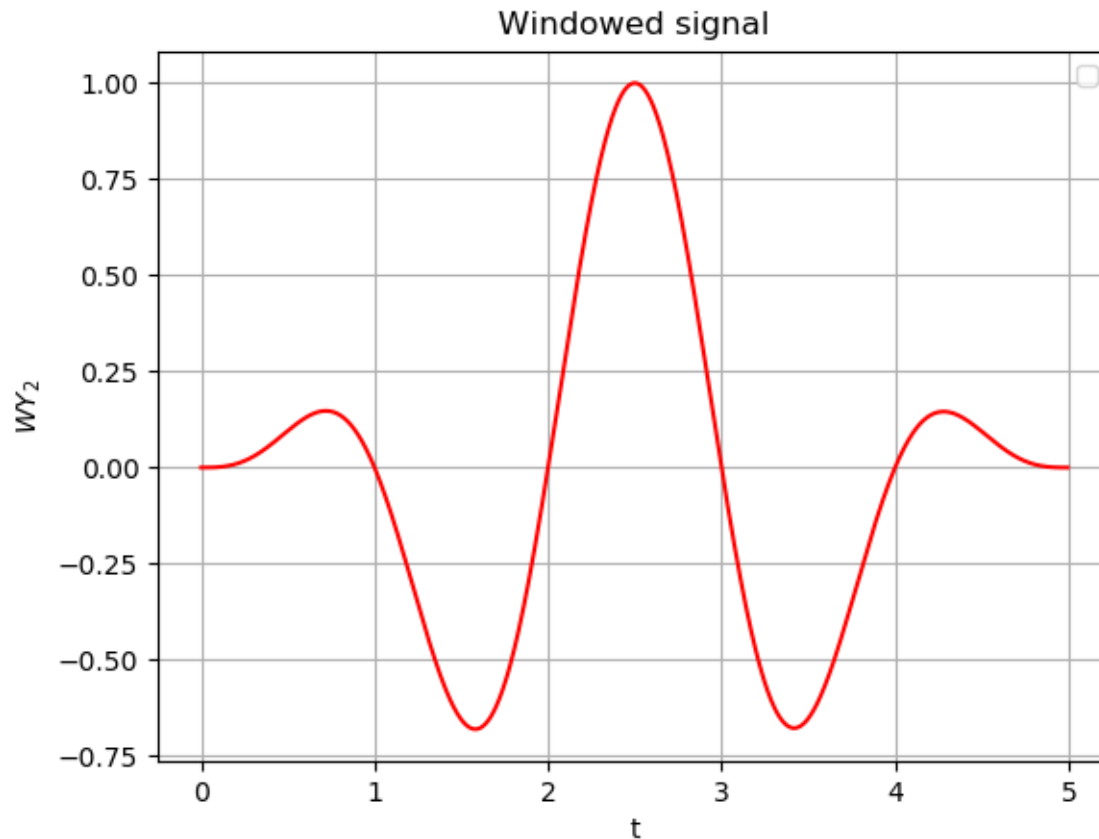
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- There is no reason to generally expect signals to be periodic so this is an important issue
- The standard “fix” is to use windowing:
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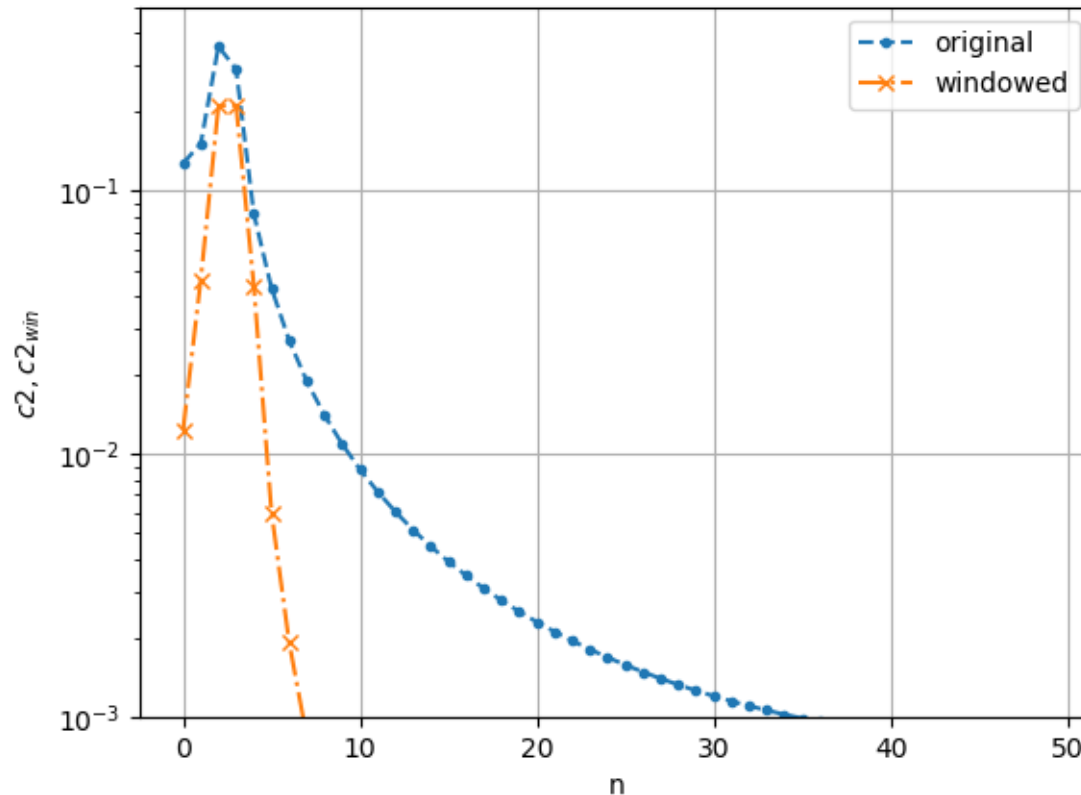
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- The standard “fix” is to use windowing:
- The spectrum of the windowed signal “looks” more like a simple wave
- We have lost energy – there is no perfect solution



Data analysis

- In practice, we don't have to go through the windowing process ourselves
- Signal processing tools exist which:
 - Break the signal up into overlapping segments
 - Window the signal within each segment and compute the spectrum
 - Average the spectra from each segment (typically $|c|^2$ rather than $|c|$)
- This produces an estimate of the *autospectral density*
- And can be computed using *Welch's method*, `scipy.signal.welch`

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In [201]: w2, Pxx2 = sig.welch(Y2)
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In [203]: w2 = w2*Nt/T
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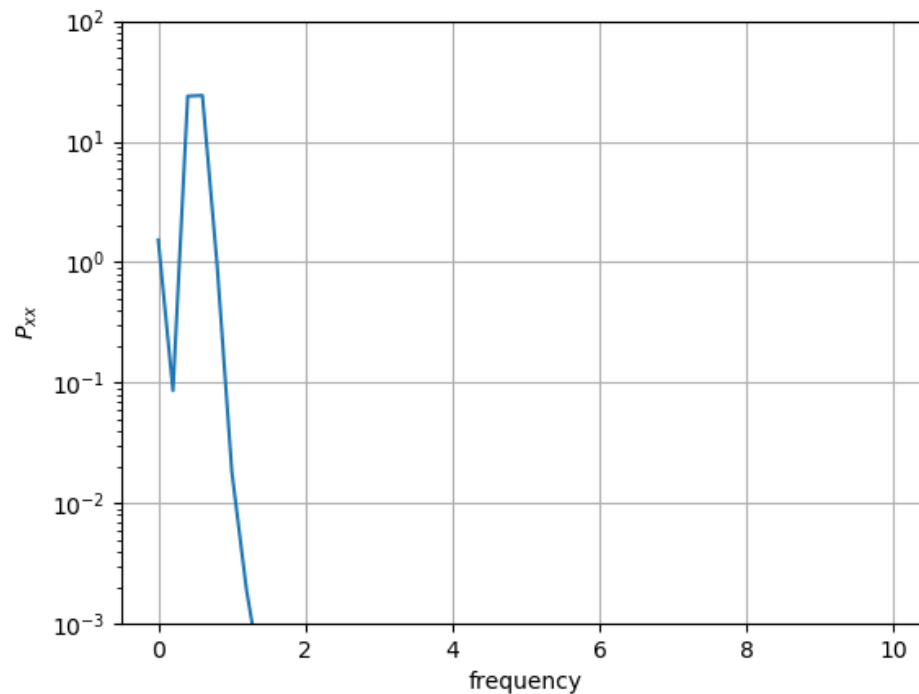
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Notes:

- Consider a signal of length T with step Δt as a distribution of “energy” across a range of frequencies
- We need $\Delta t < 2/f_{\min}$ to *resolve* the highest frequency components
- Also need $T \gg 1/f_{\max}$ to ensure “slow” components are contained within the signal
 - Not the case in our previous example
 - Often, slow components have the largest amplitude

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 - Not the case in our previous example
 - Often, slow components have the largest amplitudes
- What about the running time?
 - Direct evaluation of the sum in the DFT requires $O(N^2)$ operations
 - But the FFT uses a divide and conquer approach and $O(N \log_2 N)$ operations are needed