# 1 QUANTUM MECHANICS FORMULAE

### 1.1 Postulates

$$p_{a_i} = p(a_i, |b_j\rangle) = ||P_i|b_j\rangle||^2$$
$$(\Delta A)_{\psi}^2 = \langle A^2 \rangle_{\psi} - \langle A \rangle_{\psi}^2$$

## 1.2 Temporal evolution

$$U(t,t_0) = e^{-iH(t-t_0)/\hbar} = \sum_{i} e^{-iE_i(t-t_0)/\hbar} |i\rangle\langle i|$$

$$S: |\psi(t)\rangle = U(t,t_0) |\psi(t_0)\rangle, \quad H: A(t) = U^{\dagger}(t,t_0)AU(t,t_0)$$

$$\frac{\mathrm{d}}{\mathrm{d}t} \langle A_H(t)\rangle_{\psi} = \frac{1}{i\hbar} \langle [A_H,H]\rangle_{\psi} + \left\langle \frac{\partial A}{\partial t} \right\rangle_{\psi}$$

## 1.3 Infinite well

$$[0,L] \Rightarrow \begin{cases} \phi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) & x \in [0,L] \\ \phi_n(x) = 0 & x \notin [0,L] \end{cases}$$
$$\left[-\frac{L}{2}, \frac{L}{2}\right] \Rightarrow \begin{cases} \phi_n(x) = \sqrt{\frac{2}{L}} \cos\left(\frac{n\pi x}{L}\right) & n \text{ odd} \\ \phi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) & n \text{ even} \end{cases}$$
$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

## 1.4 HARMONIC OSCILLATOR (1D)

Gaussian: 
$$\psi(x) = \frac{1}{(2\pi\sigma_x^2)^{1/4}} \exp\left[-\frac{x^2}{4\sigma_x^2}\right]$$
  
 $\phi_n(x) = C_n H_n(\tilde{x}) \exp\left[-\frac{\tilde{x}^2}{2}\right], \quad E_n = \hbar\omega\left(n + \frac{1}{2}\right)$ 

$$\tilde{x} = \frac{x}{a_0}, \quad a_0 = \sqrt{\frac{\hbar}{m\omega}}, \quad C_n = \left(\frac{1}{\pi a_0^2}\right)^{1/4} \left(\frac{1}{2^n n!}\right)^{1/2}$$

$$H_0(\tilde{x}) = 1 \qquad H_2(\tilde{x}) = 4\tilde{x}^2 - 2$$

$$H_1(\tilde{x}) = 2\tilde{x} \qquad H_3(\tilde{x}) = 8\tilde{x}^3 - 12\tilde{x}$$

## 1.5 Symmetries

$$D_{\hat{n}}(\phi) = e^{-i\phi \, \hat{n} \cdot \vec{\sigma}/2} = \cos \frac{\phi}{2} \mathbb{1} - i \sin \frac{\phi}{2} \, \hat{n} \cdot \vec{\sigma}$$

$$\hat{n} = \left(\cos \phi \sin \theta, \sin \phi \sin \theta, \cos \theta\right)$$

$$T_a = e^{-ia\hat{k}} = e^{-i\vec{a} \cdot \vec{p}/\hbar} = e^{\alpha(a^{\dagger} - a)}$$

$$\pi |l, m\rangle = (-1)^l |l, m\rangle, \quad P_a = P_b P_c (-1)^l$$

$$\Theta |l, m\rangle = (-1)^m |l, -m\rangle$$

$$[T_a,T_{a'}]$$
  $\checkmark$ ,  $[T_a,\pi]$   $X$ ,  $[D_{\hat{n}}(\phi),D_{\hat{n}'}(\phi)']$   $X$ ,  $[D_{\hat{n}}(\phi),\pi]$   $\checkmark$ 

### 1.6 Angular Momentum

$$J_{\pm} = J_x \pm J_{-}, \quad J_x = \frac{1}{2}(J_{+} + J_{-}), \quad J_y = \frac{1}{2i}(J_{+} - J_{-})$$
$$[J_i, J_j] = i\hbar\epsilon_{ijk}J_k, \quad [J_z, J_{\pm}] = \pm\hbar J_{\pm}$$
$$[J_{+}, J_{-}] = 2\hbar J_z, \qquad [J^2, J_{z,\pm}] = 0$$
$$J_{\pm}|j, m\rangle = \hbar\sqrt{j(j+1) - m(m\pm 1)}|j, m\pm 1\rangle$$
$$J^2|j, m\rangle = \hbar^2 j(j+1)|j, m\rangle$$
$$J_z|j, m\rangle = \hbar m|j, m\rangle$$

## 1.7 Addition of Angular Momentum

$$\mathcal{H}_{J_1} \otimes \mathcal{H}_{J_1} = \mathcal{H}_{J^{(1)}} \oplus \cdots \oplus \mathcal{H}_{J^{(n)}}, \quad \dim(\mathcal{H}_J) = 2J + 1$$
  
 $J = J_1 + J_2, J_1 + J_2 - 1, \dots, |J_1 - J_2| \ge 0$ 

# 1.8 Pauli matrices (spin-1/2)

$$\sigma_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_{y} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_{z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\sigma_{i}\sigma_{j} = \delta_{ij}\mathbb{I} + i\epsilon_{ijk}\sigma_{k}$$

$$\sigma_{\hat{n}} = \hat{n} \cdot \begin{pmatrix} \sigma_{x} \\ \sigma_{y} \\ \sigma_{z} \end{pmatrix} = \begin{pmatrix} \cos \theta & e^{-i\phi} \sin \theta \\ e^{i\phi} \sin \theta & -\cos \theta \end{pmatrix}$$

$$|+\rangle_{\hat{n}} = \begin{pmatrix} \cos \frac{\theta}{2} e^{-i\frac{\varphi}{2}} \\ \sin \frac{\theta}{2} e^{i\frac{\varphi}{2}} \end{pmatrix}, \quad |-\rangle_{\hat{n}} = \begin{pmatrix} -\sin \frac{\theta}{2} e^{-i\frac{\varphi}{2}} \\ \cos \frac{\theta}{2} e^{i\frac{\varphi}{2}} \end{pmatrix}$$

$$|+\rangle_{x} = (1,1)/\sqrt{2}, \quad |-\rangle_{x} = (1,-1)/\sqrt{2}, \quad \sigma_{x} \mid \pm \rangle = |\mp\rangle$$

$$|+\rangle_{y} = (1,i)/\sqrt{2}, \quad |-\rangle_{y} = (i,1)/\sqrt{2}, \quad \sigma_{y} \mid \pm \rangle = \pm i \mid \mp\rangle$$

$$|+\rangle_{z} = (1,0), \quad |-\rangle_{z} = (0,1), \quad \sigma_{z} \mid \pm \rangle = \pm i \mid \pm\rangle$$

# 1.9 Spin-1 matrices

$$J_{x} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \begin{vmatrix} +\hbar \rangle_{x} = (1, \sqrt{2}, 1)/2 \\ |-\hbar \rangle_{x} = (1, -\sqrt{2}, 1)/2 \\ |0 \rangle_{x} = (1, 0, -1)/\sqrt{2} \end{vmatrix}$$

$$J_{y} = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad \begin{vmatrix} +\hbar \rangle_{y} = (-1, -\sqrt{2}i, 1)/2 \\ |-\hbar \rangle_{y} = (-1, \sqrt{2}i, 1)/2 \\ |0 \rangle_{y} = (1, 0, 1)/\sqrt{2} \end{vmatrix}$$

$$J_{z} = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad \begin{vmatrix} +\hbar \rangle_{z} = (1, 0, 0) \\ |-\hbar \rangle_{z} = (0, 0, 1) \\ |0 \rangle_{z} = (0, 1, 0) \end{vmatrix}$$

## 1.10 Spin-1 squared matrices

$$\begin{split} J_x^2 &= \frac{\hbar^2}{2} \begin{pmatrix} 1 & 0 & 1 \\ 0 & 2 & 0 \\ 1 & 0 & 1 \end{pmatrix}, & \left| \frac{\hbar^2}{2} \right| = (1,0,1)/\sqrt{2} \\ \left| \frac{\hbar^2}{2} \right| = (0,1,0) \\ \left| 0 \right\rangle &= (1,0,-1)/\sqrt{2} \\ J_y^2 &= \frac{\hbar^2}{2} \begin{pmatrix} 1 & 0 & -1 \\ 0 & 2 & 0 \\ -1 & 0 & 1 \end{pmatrix}, & \left| \frac{\hbar^2}{2} \right\rangle &= (1,0,-1)/\sqrt{2} \\ \left| \frac{\hbar^2}{2} \right\rangle &= (0,1,0) \\ \left| 0 \right\rangle &= (1,0,1)/\sqrt{2} \\ J_z^2 &= \hbar^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}, & \left| \frac{\hbar^2}{2} \right\rangle &= (0,0,1) \\ \left| \frac{\hbar^2}{2} \right\rangle &= (0,0,1) \\ \left| 0 \right\rangle &= (0,1,0) \end{split}$$

## 1.11 Trigonometry

$$\sin^2\theta = \frac{1}{2}(1 - \cos 2\theta) \qquad \cos\frac{\pi}{3} = \sin\frac{\pi}{6} = \frac{1}{2}$$
$$\cos^2\theta = \frac{1}{2}(1 + \cos 2\theta) \qquad \cos\frac{\pi}{4} = \sin\frac{\pi}{4} = \frac{1}{\sqrt{2}}$$
$$\sin(\theta \pm \phi) = \sin\theta\cos\phi \pm \cos\theta\sin\phi$$
$$\cos(\theta \pm \phi) = \cos\theta\cos\phi \mp \sin\theta\sin\phi \qquad \cos\frac{\pi}{6} = \sin\frac{\pi}{3} = \frac{\sqrt{3}}{2}$$

## 1.12 Eigenstuff in SU(2) (RMT)

$$\begin{pmatrix} a & c - id \\ c + id & b \end{pmatrix} = \frac{a + b}{2} \mathbb{1} + \frac{a - b}{2} \sigma_z + c\sigma_x + d\sigma_y$$
$$= \varepsilon \mathbb{1} + \vec{v} \cdot \vec{\sigma}$$
$$\Rightarrow \text{vaps} : \lambda_{\pm} = \varepsilon \pm ||\vec{v}||$$

## 1.13 The good stuff

$$\int_{-\infty}^{\infty} e^{-ax^2} x^2 dx = \frac{1}{2a} \sqrt{\frac{\pi}{a}}, \quad \int_{-\infty}^{\infty} e^{-ax^2} x^4 dx = \frac{3}{4a^2} \sqrt{\frac{\pi}{a}}$$
$$\int_{0}^{\infty} e^{-ar} r^n dr = \frac{n!}{(a)^{n+1}}, \quad 0 \le n$$

### 1.14 Time-dependent perturbation theory

$$H = H_0 + V \quad \Rightarrow |\psi_S(t)\rangle \equiv \sum_n e^{-iE_n t/\hbar} c_n^{(1)}(t) |n\rangle$$

$$c_n^{(1)}(t) = -\frac{i}{\hbar} \int_0^t e^{i\omega_{n0}t} \langle n|V(x,t)|0\rangle \,\mathrm{d}t \,, \quad p_{0\to n} = \left\|c_n^{(1)}\right\|^2$$

$$c_n^{(2)}(t) = -\frac{i}{\hbar} \iint_0^{t,t'} e^{i\omega_{n0}t} \sum_m \langle n|V|m\rangle \, \langle m|V|0\rangle \,\mathrm{d}t \,\mathrm{d}t'$$

#### 1.15 FOCK BASIS

$$x = \sqrt{\frac{\hbar}{2m\omega}} \left( a^{\dagger} + a \right), \quad p = i\sqrt{\frac{\hbar m\omega}{2}} \left( a^{\dagger} - a \right)$$
$$a \left| \alpha \right\rangle = \alpha \left| \alpha \right\rangle, \quad \left| \alpha \right\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \left| n \right\rangle = e^{-\frac{|\alpha|^2}{2}} e^{\alpha a^{\dagger}} \left| 0 \right\rangle$$

### Position

$$\begin{array}{lll} x \left| 0 \right\rangle \rightarrow \left| 1 \right\rangle, & x^2 \left| 0 \right\rangle \rightarrow \left| 0 \right\rangle + \sqrt{2} \left| 2 \right\rangle \\ x \left| 1 \right\rangle \rightarrow \left| 0 \right\rangle + \sqrt{2} \left| 3 \right\rangle, & x^2 \left| 1 \right\rangle \rightarrow 3 \left| 1 \right\rangle + \sqrt{6} \left| 3 \right\rangle \\ x \left| 2 \right\rangle \rightarrow \sqrt{2} \left| 1 \right\rangle + \sqrt{3} \left| 3 \right\rangle, & x^2 \left| 2 \right\rangle \rightarrow \sqrt{2} \left| 0 \right\rangle + 5 \left| 2 \right\rangle + 2\sqrt{3} \left| 4 \right\rangle \\ x \left| 3 \right\rangle \rightarrow \sqrt{3} \left| 2 \right\rangle + \sqrt{4} \left| 4 \right\rangle, & x^2 \left| 3 \right\rangle \rightarrow \sqrt{6} \left| 1 \right\rangle + 7 \left| 3 \right\rangle + 2\sqrt{5} \left| 5 \right\rangle \\ x \left| 4 \right\rangle \rightarrow \sqrt{4} \left| 3 \right\rangle + \sqrt{5} \left| 5 \right\rangle, & x^2 \left| 4 \right\rangle \rightarrow 2\sqrt{3} \left| 2 \right\rangle + 9 \left| 4 \right\rangle + \sqrt{30} \left| 6 \right\rangle \\ x \left| 5 \right\rangle \rightarrow \sqrt{5} \left| 4 \right\rangle + \sqrt{6} \left| 6 \right\rangle, & x^2 \left| 5 \right\rangle \rightarrow 2\sqrt{5} \left| 3 \right\rangle + 11 \left| 5 \right\rangle + \sqrt{42} \left| 7 \right\rangle \end{array}$$

### Momentum

$$\begin{split} p & | 0 \rangle \rightarrow | 1 \rangle \;, & p^2 & | 0 \rangle \rightarrow - | 0 \rangle + \sqrt{2} \, | 2 \rangle \\ p & | 1 \rangle \rightarrow - | 0 \rangle + \sqrt{2} \, | 3 \rangle \;, & p^2 & | 1 \rangle \rightarrow -3 \, | 1 \rangle + \sqrt{6} \, | 3 \rangle \\ p & | 2 \rangle \rightarrow -\sqrt{2} \, | 1 \rangle + \sqrt{3} \, | 3 \rangle \;, & p^2 & | 2 \rangle \rightarrow \sqrt{2} \, | 0 \rangle - 5 \, | 2 \rangle + 2 \sqrt{3} \, | 4 \rangle \\ p & | 3 \rangle \rightarrow -\sqrt{3} \, | 2 \rangle + \sqrt{4} \, | 4 \rangle \;, & p^2 & | 3 \rangle \rightarrow \sqrt{6} \, | 1 \rangle - 7 \, | 3 \rangle + 2 \sqrt{5} \, | 5 \rangle \\ p & | 4 \rangle \rightarrow -\sqrt{4} \, | 3 \rangle + \sqrt{5} \, | 5 \rangle \;, & p^2 & | 4 \rangle \rightarrow 2 \sqrt{3} \, | 2 \rangle - 9 \, | 4 \rangle + \sqrt{30} \, | 6 \rangle \\ p & | 5 \rangle \rightarrow -\sqrt{5} \, | 4 \rangle + \sqrt{6} \, | 6 \rangle \;, & p^2 & | 5 \rangle \rightarrow 2 \sqrt{5} \, | 3 \rangle - 11 \, | 5 \rangle + \sqrt{42} \, | 7 \rangle \end{split}$$