

Arbitrary-Lagrangian-Eulerian Finite-Element Hydrodynamics

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April 1, 2024

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Chapter 1

The Lagrangian Step

1.1 Lagrangian governing equations

We consider Lagrangian fluid particles, for which we define their position $\mathbf{x}^+ = \mathbf{x}^+(t, \mathbf{y})$, density $\rho^+ = \rho^+(t, \mathbf{y})$, velocity $\mathbf{u}^+ = \mathbf{u}^+(t, \mathbf{y})$, and internal energy $e^+ = e^+(t, \mathbf{y})$. The vector \mathbf{y} is the location of each fluid particle at time zero and thus serves to differentiate between the different particles. The Eulerian counterparts for the density, velocity, and internal energy are, respectively, $\rho = \rho(t, \mathbf{x})$, $\mathbf{u} = \mathbf{u}(t, \mathbf{x})$, and $e = e(t, \mathbf{x})$. The vector \mathbf{x} is a location in Eulerian space. Also consider the volume Ω_0 as the set of all \mathbf{y} vectors that make up the initial domain. The control volume $\Omega^+ = \Omega^+(t, \Omega_0)$ is then defined by

$$\Omega^+ = \{\mathbf{x}^+ : \mathbf{y} \in \Omega_0\}. \quad (1.1)$$

Note that $\Omega^+(0, \Omega_0) = \Omega_0$.

The governing equations for the Lagrangian fluid particles are derived in my fluid-mechanics notes (see section on kinematics, Lagrangian governing equations, etc.). These are shown below

$$\frac{\partial \mathbf{x}^+}{\partial t} = \mathbf{u}^+, \quad (1.2)$$

$$\frac{\partial \rho^+}{\partial t} = -\rho^+ (\nabla \cdot \mathbf{u})_{\mathbf{x}=\mathbf{x}^+}, \quad (1.3)$$

$$\rho^+ \frac{\partial \mathbf{u}^+}{\partial t} = (\nabla \cdot \boldsymbol{\sigma})_{\mathbf{x}=\mathbf{x}^+}, \quad (1.4)$$

$$\rho^+ \frac{\partial e^+}{\partial t} = (\boldsymbol{\sigma} : \nabla \mathbf{u})_{\mathbf{x}=\mathbf{x}^+}. \quad (1.5)$$

In the above, $\boldsymbol{\sigma} = \boldsymbol{\sigma}(t, \mathbf{x})$ is the stress tensor.

A note on notation. The products that involve a tensor $\boldsymbol{\tau}$ can be expressed in Einstein notation as

$$\nabla \cdot \boldsymbol{\tau} = \frac{\partial \tau_{ij}}{\partial x_j}, \quad (1.6)$$

$$\boldsymbol{\tau} \cdot \nabla f = \tau_{ij} \frac{\partial f}{\partial x_j}, \quad (1.7)$$

$$\mathbf{g} \cdot \boldsymbol{\tau} \cdot \nabla f = g_i \tau_{ij} \frac{\partial f}{\partial x_j}, \quad (1.8)$$

$$\boldsymbol{\tau} : \nabla \mathbf{g} = \tau_{ij} \frac{\partial g_i}{\partial x_j}. \quad (1.9)$$

where f is a scalar and \mathbf{g} a vector. In these notes we'll mostly be using indices i and j for FE expansions, rather than for Einstein notation.

1.2 Lagrangian finite elements

We introduce a Lagrangian basis function $\Phi_i^+ = \Phi_i^+(t, \mathbf{y})$ and an Eulerian basis function $\Phi_i = \Phi_i(t, \mathbf{x})$. These are related to each other as any other Lagrangian-Eulerian pair, namely

$$\Phi_i^+(t, \mathbf{y}) = \Phi_i(t, \mathbf{x}^+(t, \mathbf{y})). \quad (1.10)$$

We now introduce the Lagrangian variable $f^+ = f^+(t, \mathbf{y})$ and the Eulerian counterpart $f = f(t, \mathbf{x})$, and they also satisfy

$$f^+(t, \mathbf{y}) = f(t, \mathbf{x}^+(t, \mathbf{y})). \quad (1.11)$$

The expansion of an Eulerian variable in terms of basis functions is as follows

$$f = \sum_i^n F_i \Phi_i, \quad (1.12)$$

where $F_i = F_i(t)$. Plugging in \mathbf{x}^+ for \mathbf{x} in the above, and using eqs. (1.10) and (1.11) gives

$$f^+ = \sum_i^n F_i \Phi_i^+. \quad (1.13)$$

Thus, both the Lagrangian and Eulerian variables share the same finite-element coefficients F_i .

As shown in my fluid mechanics notes, we also have

$$\frac{\partial \Phi_i^+}{\partial t} = \left(\frac{\partial \Phi_i}{\partial t} + \mathbf{u} \cdot \nabla \Phi_i \right)_{\mathbf{x}=\mathbf{x}^+}, \quad (1.14)$$

where $\mathbf{u} = \mathbf{u}(t, \mathbf{x})$ is the Eulerian counterpart to \mathbf{u}^+ . We'll introduce the restriction that Φ_i^+ is constant in time, that is $\partial \Phi_i^+ / \partial t = 0$, which gives

$$\frac{\partial \Phi_i}{\partial t} + \mathbf{u} \cdot \nabla \Phi_i = 0. \quad (1.15)$$

Thus, F_i in eq. (1.13) accounts for the time dependence of F^+ , whereas Φ_i^+ accounts for the dependence on \mathbf{y} .

1.3 Finite element expansion

We introduce the coefficients $\hat{\mathbf{x}}_i = \hat{\mathbf{x}}_i(t)$, $\hat{\mathbf{u}}_i = \hat{\mathbf{u}}_i(t)$ and $\hat{e}_i = \hat{e}_i(t)$, as well as the Lagrangian basis functions $\phi_i^+ = \phi_i^+(\mathbf{y}) \in L^2$, and $w_i^+ = w_i^+(\mathbf{y}) \in H^1$. We note that $\hat{\mathbf{x}}_i$ and $\hat{\mathbf{u}}_i$ are each vectors, e.g., the components of $\hat{\mathbf{u}}_i$ are $\hat{u}_{i,\alpha} = \hat{u}_{i,\alpha}(t)$ for $\alpha = x, y, z$. We also note that ϕ_i^+ and w_i^+ have Eulerian

counterparts $\phi_i = \phi_i(t, \mathbf{x})$ and $w_i = w_i(t, \mathbf{x})$, respectively. The coefficients are used in the following expansions

$$\mathbf{x}^+ = \sum_j^{N_w} \hat{\mathbf{x}}_j w_j^+, \quad (1.16)$$

$$\mathbf{u}^+ = \sum_j^{N_w} \hat{\mathbf{u}}_j w_j^+, \quad (1.17)$$

$$e^+ = \sum_j^{N_\phi} \hat{e}_j \phi_j^+. \quad (1.18)$$

We note that the expansion coefficients are the same for the Lagrangian and Eulerian variables, as shown in section 1.2. For example, for the Eulerian velocity, we have

$$\mathbf{u} = \sum_j^{N_w} \hat{\mathbf{u}}_j w_j. \quad (1.19)$$

1.4 Semi-discrete Lagrangian governing equations

1.4.1 Position and Jacobian

Plugging in eqs. (1.16) and (1.17) in eq. (1.2) gives

$$\sum_j^{N_w} \frac{d\hat{\mathbf{x}}_j}{dt} w_j^+ = \sum_j^{N_w} \hat{\mathbf{u}}_j w_j^+. \quad (1.20)$$

To satisfy the equation above, we'll require

$$\frac{d\hat{\mathbf{x}}_j^+}{dt} = \hat{\mathbf{u}}_j. \quad (1.21)$$

We now introduce the vectors \mathbf{X} and \mathbf{U} , whose components are $\hat{\mathbf{x}}_i$ and $\hat{\mathbf{u}}_i$, respectively. Thus, the above is written as

$$\frac{d\mathbf{X}}{dt} = \mathbf{U}. \quad (1.22)$$

1.4.2 Density

We introduce the Jacobian matrix $\mathbf{J}^+ = \mathbf{J}^+(t, \mathbf{y})$, which is defined as

$$\mathbf{J}^+ = \frac{\partial \mathbf{x}^+}{\partial \mathbf{y}}. \quad (1.23)$$

It's determinant is denoted by $J^+ = J^+(t, \mathbf{y})$, and it satisfies the following equation

$$\frac{\partial J^+}{\partial t} = J^+ (\nabla \cdot \mathbf{u})_{\mathbf{x}=\mathbf{x}^+}. \quad (1.24)$$

Thus, Equation (1.3) can be re-written as

$$\frac{1}{\rho^+} \frac{\partial \rho^+}{\partial t} = -\frac{1}{J^+} \frac{\partial J^+}{\partial t}, \quad (1.25)$$

or

$$\frac{\partial J^+ \rho^+}{\partial t} = 0. \quad (1.26)$$

Since $J^+ = 1$ at the initial time, we obtain the density according to

$$\rho^+ = \frac{\rho_0^+}{J^+}, \quad (1.27)$$

where $\rho_0^+ = \rho^+(0, \mathbf{y})$.

To compute J^+ , we first plug in eq. (1.16) in the definition of the Jacobian matrix, which gives

$$\mathbf{J}^+ = \frac{\partial}{\partial \mathbf{y}} \sum_j^{N_w} \hat{\mathbf{x}}_j w_j^+ = \sum_j^{N_w} \hat{\mathbf{x}}_j \nabla_{\mathbf{y}} w_j^+. \quad (1.28)$$

We then simply compute the determinant of the above to obtain J^+ . Note that for any function \mathbf{x}^+ and its corresponding \mathbf{u}^+ , whether it be an exact analytical expression or a finite-element expansion as given by eq. (1.16), one obtains eq. (1.24). Thus, the derivation of eq. (1.26) from eq. (1.3) still holds whether one plans to represent J^+ using an analytical expression or a finite-element expansion.

1.4.3 Velocity

Plugging in eq. (1.27) in eq. (1.4) we get

$$\rho_0^+ \frac{\partial \mathbf{u}^+}{\partial t} = (\nabla \cdot \boldsymbol{\sigma})_{\mathbf{x}=\mathbf{x}^+} J^+. \quad (1.29)$$

We then multiply both sides of the above by the basis functions for velocity and integrate over all space to obtain

$$\int_{\Omega_0} \rho_0^+ \frac{\partial \mathbf{u}^+}{\partial t} w_i^+ dV_y = \int_{\Omega_0} (\nabla \cdot \boldsymbol{\sigma})_{\mathbf{x}=\mathbf{x}^+} w_i^+ J^+ dV_y. \quad (1.30)$$

For the left-hand side we have

$$\begin{aligned} \int_{\Omega_0} \rho_0^+ \frac{\partial \mathbf{u}^+}{\partial t} w_i^+ dV_y &= \int_{\Omega_0} \rho_0^+ \sum_j^{N_w} \frac{d\hat{\mathbf{u}}_j}{dt} w_j^+ w_i^+ dV_y, \\ &= \sum_j^{N_w} \frac{d\hat{\mathbf{u}}_j}{dt} \int_{\Omega_0} \rho_0^+ w_i^+ w_j^+ dV_y, \\ &= \sum_j^{N_w} \frac{d\hat{\mathbf{u}}_j}{dt} m_{\mathcal{V},ij}, \end{aligned} \quad (1.31)$$

where

$$m_{\mathcal{V},ij} = \int_{\Omega_0} \rho_0^+ w_i^+ w_j^+ dV_y \quad (1.32)$$

is a mass bilinear form (which is independent of time). For the right-hand side we have

$$\begin{aligned}
\int_{\Omega_0} (\nabla \cdot \boldsymbol{\sigma})_{\mathbf{x}=\mathbf{x}^+} w_i^+ J^+ dV_y &= \int_{\Omega_0} (\nabla \cdot \boldsymbol{\sigma} w_i)_{\mathbf{x}=\mathbf{x}^+} J^+ dV_y \\
&= \int_{\Omega^+} \nabla \cdot \boldsymbol{\sigma} w_i dV_x \\
&= - \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i dV_x.
\end{aligned} \tag{1.33}$$

The second equality above follows from integration by substitution. Combining results we have

$$\sum_j^{N_w} \frac{d\hat{\mathbf{u}}_j}{dt} m_{\mathcal{V},ij} = - \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i dV_x. \tag{1.34}$$

We introduce the matrix $\mathbf{M}_{\mathcal{V}}$, whose components are $m_{\mathcal{V},ij}$. Thus, the left-hand side of eq. (1.34) can be written as $\mathbf{M}_{\mathcal{V}} d\mathbf{U}/dt$. We also introduce the vector bilinear form

$$\mathbf{f}_{ij} = \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i \phi_j dV_x. \tag{1.35}$$

This is a *vector* bilinear form since \mathbf{f}_{ij} has components $f_{ij,\alpha} = f_{ij,\alpha}(t)$, for $\alpha = x, y, z$, where α denotes the first index of $\boldsymbol{\sigma}$. We introduce the force matrix \mathbf{F} , whose components are \mathbf{f}_{ij} . We also expand the field with constant value of one as follows

$$1 = \sum_i^{N_\phi} \hat{\mathbf{1}}_i \phi_i. \tag{1.36}$$

If we define the vector $\hat{\mathbf{1}}$ as that with components $\hat{\mathbf{1}}_i$, we can show that

$$\begin{aligned}
\mathbf{F} \hat{\mathbf{1}} &= \sum_j^{N_\phi} \mathbf{f}_{ij} \hat{\mathbf{1}}_j \\
&= \sum_j^{N_\phi} \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i \phi_j dV_x \hat{\mathbf{1}}_j \\
&= \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i \left(\sum_j^{N_\phi} \hat{\mathbf{1}}_j \phi_j \right) dV_x \\
&= \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i dV_x.
\end{aligned} \tag{1.37}$$

The above is the negative of the right-hand side of eq. (1.34). Thus, combining all together we get

$$\mathbf{M}_{\mathcal{V}} \frac{d\mathbf{U}}{dt} = -\mathbf{F} \hat{\mathbf{1}}. \tag{1.38}$$

We note that since both the Lagrangian and Eulerian velocities share the same coefficients \mathbf{U} , we now have a solution for both.

1.4.4 Energy

Plugging in eq. (1.27) in eq. (1.5) we get

$$\rho_0^+ \frac{\partial e^+}{\partial t} = (\boldsymbol{\sigma} : \nabla \mathbf{u})_{\mathbf{x}=\mathbf{x}^+} J^+. \quad (1.39)$$

We then multiply both sides of the above by the basis functions for energy and integrate over all space to obtain

$$\int_{\Omega_0} \rho_0^+ \frac{\partial e^+}{\partial t} \phi_i^+ dV_y = \int_{\Omega_0} (\boldsymbol{\sigma} : \nabla \mathbf{u})_{\mathbf{x}=\mathbf{x}^+} \phi_i^+ J^+ dV_y. \quad (1.40)$$

For the left-hand side we have

$$\begin{aligned} \int_{\Omega_0} \rho_0^+ \frac{\partial e^+}{\partial t} \phi_i^+ dV_y &= \int_{\Omega_0} \rho_0^+ \sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} \phi_j^+ \phi_i^+ dV_y, \\ &= \sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} \int_{\Omega_0} \rho_0^+ \phi_j^+ \phi_i^+ dV_y, \\ &= \sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} m_{\mathcal{E},ij} \end{aligned} \quad (1.41)$$

where

$$m_{\mathcal{E},ij} = \int_{\Omega_0} \rho_0^+ \phi_j^+ \phi_i^+ dV_y \quad (1.42)$$

is a mass bilinear form (which is independent of time). For the right-hand side we have

$$\begin{aligned} \int_{\Omega_0} (\boldsymbol{\sigma} : \nabla \mathbf{u})_{\mathbf{x}=\mathbf{x}^+} \phi_i^+ J^+ dV_y &= \int_{\Omega_0} (\boldsymbol{\sigma} : \nabla \mathbf{u} \phi_i)_{\mathbf{x}=\mathbf{x}^+} J^+ dV_y \\ &= \int_{\Omega^+} \boldsymbol{\sigma} : \nabla \mathbf{u} \phi_i dV_x. \end{aligned} \quad (1.43)$$

Combining results we have

$$\sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} m_{\mathcal{E},ij} = \int_{\Omega^+} \boldsymbol{\sigma} : \nabla \mathbf{u} \phi_i dV_x. \quad (1.44)$$

We now show that

$$\boldsymbol{\sigma} : \nabla \mathbf{u} = \boldsymbol{\sigma} : \nabla \left(\sum_k^{N_w} \hat{\mathbf{u}}_k w_k \right) = \sum_k^{N_w} \hat{\mathbf{u}}_k \cdot \boldsymbol{\sigma} \cdot \nabla w_k, \quad (1.45)$$

and hence the previous result is written as

$$\sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} m_{\mathcal{E},ij} = \sum_k^{N_w} \hat{\mathbf{u}}_k \cdot \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_k \phi_i dV_x. \quad (1.46)$$

The above is finally re-written as

$$\sum_j^{N_\phi} \frac{d\hat{e}_j}{dt} m_{\mathcal{E},ij} = \sum_k^{N_w} \hat{\mathbf{u}}_k \cdot \mathbf{f}_{k,i}. \quad (1.47)$$

Note that in the above there is a dot product in the right-hand side, that is, the right-hand side expanded out is

$$\sum_k^{N_w} \hat{\mathbf{u}}_k \cdot \mathbf{f}_{ki} = \sum_k^{N_w} \sum_{\alpha=x,y,z} \hat{u}_{k,\alpha} f_{ki,\alpha}. \quad (1.48)$$

We now introduce the vector \mathbf{E} whose components are \hat{e}_i . We also introduce the matrix $\mathbf{M}_{\mathcal{E}}$ whose components are $m_{\mathcal{E},ij}$. Thus, eq. (1.47) can be succinctly written as

$$\mathbf{M}_{\mathcal{E}} \frac{d\mathbf{E}}{dt} = \mathbf{F}^T \cdot \mathbf{U}. \quad (1.49)$$

Note again that on the right-hand side above there is a matrix-vector product *and* a dot product. We also note that since both the Lagrangian and Eulerian internal energies share the same coefficients \mathbf{E} , we now have a solution for both.

1.5 Momentum and energy conservation

We'll now define the internal energy $IE = IE(t)$, the kinetic energy $KE = KE(t)$, and the momentum $P_{\mathbf{n}} = P_{\mathbf{n}}(t)$ along a constant \mathbf{n} direction.

$$\begin{aligned} IE &= \int_{\Omega^+} \rho e dV_x \\ &= \int_{\Omega_0} \rho^+ e^+ J^+ dV_y \\ &= \int_{\Omega_0} \rho_0^+ e^+ dV_y \\ &= \int_{\Omega_0} \rho_0^+ \sum_j^{N_\phi} \hat{e}_j \phi_j^+ dV_y \\ &= \int_{\Omega_0} \rho_0^+ \sum_j^{N_\phi} \hat{e}_j \phi_j^+ \left(\sum_i^{N_\phi} \hat{\mathbf{i}}_i \phi_i^+ \right) dV_y \\ &= \sum_i^{N_\phi} \sum_j^{N_\phi} \hat{\mathbf{i}}_i \int_{\Omega_0} \rho_0^+ \phi_i^+ \phi_j^+ dV_y \hat{e}_j \\ &= \sum_i^{N_\phi} \sum_j^{N_\phi} \hat{\mathbf{i}}_i m_{\mathcal{E},ij} \hat{e}_j \\ &= \hat{\mathbf{i}}^T \mathbf{M}_{\mathcal{E}} \mathbf{E} \end{aligned} \quad (1.50)$$

$$\begin{aligned}
KE &= \int_{\Omega^+} \frac{1}{2} \rho \mathbf{u} \cdot \mathbf{u} dV_x \\
&= \int_{\Omega_0} \frac{1}{2} \rho^+ \mathbf{u}^+ \cdot \mathbf{u}^+ J^+ dV_y \\
&= \int_{\Omega_0} \frac{1}{2} \rho_0^+ \mathbf{u}^+ \cdot \mathbf{u}^+ dV_y \\
&= \int_{\Omega_0} \frac{1}{2} \rho_0^+ \left(\sum_i^{N_w} \hat{\mathbf{u}}_i w_i^+ \right) \cdot \left(\sum_j^{N_w} \hat{\mathbf{u}}_j w_j^+ \right) dV_y \\
&= \sum_i^{N_w} \sum_j^{N_w} \frac{1}{2} \hat{\mathbf{u}}_i \cdot \int_{\Omega_0} \rho_0^+ w_i^+ w_j^+ dV_y \hat{\mathbf{u}}_j \\
&= \sum_i^{N_w} \sum_j^{N_w} \frac{1}{2} \hat{\mathbf{u}}_i \cdot m_{\mathcal{V},ij} \hat{\mathbf{u}}_j \\
&= \frac{1}{2} \mathbf{U}^T \cdot \mathbf{M}_{\mathcal{V}} \mathbf{U}.
\end{aligned} \tag{1.51}$$

$$\begin{aligned}
P_{\mathbf{n}} &= \int_{\Omega^+} \rho \mathbf{u} \cdot \mathbf{n} dV_x \\
&= \int_{\Omega_0} \rho^+ \mathbf{u}^+ \cdot \mathbf{n} J^+ dV_y \\
&= \int_{\Omega_0} \rho_0^+ \mathbf{u}^+ \cdot \mathbf{n} dV_y \\
&= \int_{\Omega_0} \rho_0^+ \left(\sum_j^{N_w} \hat{\mathbf{u}}_j w_j^+ \right) \cdot \left(\sum_i^{N_w} \hat{\mathbf{n}}_i w_i^+ \right) dV_y \\
&= \sum_i^{N_w} \sum_j^{N_w} \hat{\mathbf{n}}_i \cdot \int_{\Omega_0} \rho_0^+ w_i^+ w_j^+ dV_y \hat{\mathbf{u}}_j \\
&= \sum_i^{N_w} \sum_j^{N_w} \hat{\mathbf{n}}_i \cdot m_{\mathcal{V},ij} \hat{\mathbf{u}}_j \\
&= \mathbf{N}^T \cdot \mathbf{M}_{\mathcal{V}} \mathbf{U}.
\end{aligned} \tag{1.52}$$

The total energy is conserved, as shown below

$$\begin{aligned}
\frac{d}{dt}(IE + KE) &= \hat{\mathbf{1}}^T \mathbf{M}_{\mathcal{E}} \frac{d\mathbf{E}}{dt} + \mathbf{U}^T \cdot \mathbf{M}_{\mathcal{V}} \frac{d\mathbf{U}}{dt} \\
&= \hat{\mathbf{1}}^T \mathbf{F}^T \cdot \mathbf{U} - \mathbf{U}^T \cdot \mathbf{F} \hat{\mathbf{1}} \\
&= 0.
\end{aligned} \tag{1.53}$$

The momentum along a constant direction is conserved, as shown below

$$\begin{aligned}
\frac{dP_{\mathbf{n}}}{dt} &= \mathbf{N}^T \cdot \mathbf{M}_{\mathcal{V}} \frac{d\mathbf{U}}{dt} \\
&= -\mathbf{N}^T \cdot \mathbf{F} \hat{\mathbf{l}} \\
&= -\sum_i^{N_w} \sum_j^{N_\phi} \hat{\mathbf{n}}_i \cdot \mathbf{f}_{ij} \hat{\mathbf{l}}_j \\
&= -\sum_i^{N_w} \sum_j^{N_\phi} \hat{\mathbf{n}}_i \cdot \int_{\Omega^+} \boldsymbol{\sigma} \cdot \nabla w_i \phi_j dV_x \hat{\mathbf{l}}_j \\
&= -\int_{\Omega^+} \boldsymbol{\sigma} : \nabla \mathbf{n} dV_x \\
&= 0.
\end{aligned} \tag{1.54}$$

1.6 The reference element

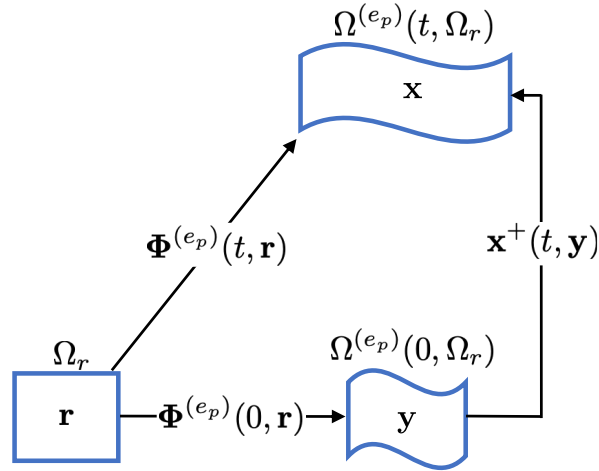


Figure 1.1: Schematic of the three domains Ω_r , $\Omega^{(e_p)}(t, \Omega_r)$, $\Omega^{(e_p)}(0, \Omega_r)$.

We introduce the reference element as the unit square in 2D or the unit cube in 3D. The domain of this reference element is labelled as Ω_r and it doesn't change with time. We introduce the function $\Phi^{(e_p)} = \Phi^{(e_p)}(t, \mathbf{r})$, which maps from points \mathbf{r} in Ω_r to points in the finite element e_p of the physical space. The evolving domain of the finite element e_p is given by the function $\Omega^{(e_p)} = \Omega^{(e_p)}(t, \Omega_r)$. A depiction of these domains and their mappings is shown in fig. 1.1. Whereas for Ω^+ we had $\Omega^+(0, \Omega_0) = \Omega_0$, for $\Omega^{(e_p)}$ the analogue does not hold, that is, $\Omega^{(e_p)}(0, \Omega_r) \neq \Omega_r$.

The mapping functions $\Phi^{(e_p)}$ and \mathbf{x}^+ are related to each other as follows

$$\Phi^{(e_p)}(t, \mathbf{r}) = \mathbf{x}^+(t, \Phi^{(e_p)}(0, \mathbf{r})). \tag{1.55}$$

The Jacobian $\mathbf{J}^{(e_p)} = \mathbf{J}^{(e_p)}(t, \mathbf{r})$ is defined as

$$\mathbf{J}^{(e_p)} = \frac{\partial \Phi^{(e_p)}}{\partial \mathbf{r}}, \tag{1.56}$$

with its determinant labeled as $J^{(e_p)} = J^{(e_p)}(t, \mathbf{r})$. Using eq. (1.55) in the definition of $\mathbf{J}^{(e_p)}$ we get

$$\begin{aligned}\mathbf{J}^{(e_p)} &= \left(\frac{\partial \mathbf{x}^+}{\partial \mathbf{y}} \right)_{\mathbf{y}=\Phi^{(e_p)}(0, \mathbf{r})} \frac{\partial \Phi^{(e_p)}(0, \mathbf{r})}{\partial \mathbf{r}} \\ &= (\mathbf{J}^+)_{\mathbf{y}=\Phi^{(e_p)}(0, \mathbf{r})} \mathbf{J}_0^{(e_p)},\end{aligned}\tag{1.57}$$

where $\mathbf{J}_0^{(e_p)} = \mathbf{J}^{(e_p)}(0, \mathbf{r})$. Taking the determinant of the above gives

$$J^{(e_p)} = (J^+)_{\mathbf{y}=\Phi^{(e_p)}(0, \mathbf{r})} J_0^{(e_p)},\tag{1.58}$$

where $J_0^{(e_p)} = J^{(e_p)}(0, \mathbf{r})$.

A Lagrangian variable $f^+ = f^+(t, \mathbf{y})$ is related to $f = f(t, \mathbf{x})$ according to the following

$$f^+(t, \mathbf{y}) = f(t, \mathbf{x}^+(t, \mathbf{y})).\tag{1.59}$$

In an analogous manner, $f^{(e_p)} = f^{(e_p)}(t, \mathbf{r})$ is related to $f = f(t, \mathbf{x})$ according to

$$f^{(e_p)}(t, \mathbf{r}) = f(t, \Phi^{(e_p)}(t, \mathbf{r})).\tag{1.60}$$

Examples of these reference-element functions include those for density $\rho^{(e_p)} = \rho^{(e_p)}(t, \mathbf{r})$, velocity $\mathbf{u}^{(e_p)} = \mathbf{u}^{(e_p)}(t, \mathbf{r})$, and internal energy $e^{(e_p)} = e^{(e_p)}(t, \mathbf{r})$. Using integration by substitution and then eq. (1.60) we show

$$\begin{aligned}\int_{\Omega^{(e_p)}} f dV_x &= \int_{\Omega_r} f(t, \Phi^{(e_p)}(t, \mathbf{r})) J^{(e_p)} dV_r \\ &= \int_{\Omega_r} f^{(e_p)} J^{(e_p)} dV_r.\end{aligned}\tag{1.61}$$

In other words, integrals over elements at any time can be computed as integrals over the reference space.

If the integrand contains a derivative, a bit of extra care is required. To show this, we'll use index notation for the sake of clarity. Consider as an example a term of the form

$$(\boldsymbol{\sigma} \cdot \nabla f)_{\mathbf{x}=\Phi^{(e_p)}} = \left(\sigma_{ij} \frac{\partial f}{\partial x_j} \right)_{\mathbf{x}=\Phi^{(e_p)}} = \sigma_{ij}^{(e_p)} \left(\frac{\partial f}{\partial x_j} \right)_{\mathbf{x}=\Phi^{(e_p)}}.\tag{1.62}$$

We first note that

$$\frac{\partial f^{(e_p)}}{\partial r_k} = \left(\frac{\partial f}{\partial x_i} \right)_{\mathbf{x}=\Phi^{(e_p)}} \frac{\partial x_i^{(e_p)}}{\partial r_k} = \left(\frac{\partial f}{\partial x_i} \right)_{\mathbf{x}=\Phi^{(e_p)}} J_{ik}^{(e_p)}.\tag{1.63}$$

Upon multiplying both sides by the inverse of $\mathbf{J}^{(e_p)}$, we get

$$\left(\frac{\partial f}{\partial x_j} \right)_{\mathbf{x}=\Phi^{(e_p)}} = \frac{\partial f^{(e_p)}}{\partial r_k} \left(J^{(e_p)} \right)_{kj}^{-1}.\tag{1.64}$$

Thus, we now have

$$(\boldsymbol{\sigma} \cdot \nabla f)_{\mathbf{x}=\Phi^{(e_p)}} = \sigma_{ij}^{(e_p)} \frac{\partial f^{(e_p)}}{\partial r_k} \left(J^{(e_p)} \right)_{kj}^{-1} = \sigma_{ij}^{(e_p)} \left[\left(J^{(e_p)} \right)^{-1} \right]_{jk}^T \frac{\partial f^{(e_p)}}{\partial r_k}.\tag{1.65}$$

In tensor notation, the above is written as

$$(\boldsymbol{\sigma} \cdot \nabla f)_{\mathbf{x}=\boldsymbol{\Phi}^{(e_p)}} = \boldsymbol{\sigma}^{(e_p)} \cdot \left[\left(\mathbf{J}^{(e_p)} \right)^{-1} \right]^T \cdot \nabla_{\mathbf{r}} f^{(e_p)}. \quad (1.66)$$

Thus, for the force matrix \mathbf{f}_{ij} we can now write

$$\begin{aligned} \int_{\Omega^{(e_p)}} \boldsymbol{\sigma} \cdot \nabla w_i \phi_j dV_x &= \int_{\Omega_{\mathbf{r}}} (\boldsymbol{\sigma} \cdot \nabla w_i \phi_j)_{\mathbf{x}=\boldsymbol{\Phi}^{(e_p)}} J^{(e_p)} dV_r \\ &= \int_{\Omega_{\mathbf{r}}} \boldsymbol{\sigma}^{(e_p)} \cdot \left[\left(\mathbf{J}^{(e_p)} \right)^{-1} \right]^T \cdot \nabla_{\mathbf{r}} w_i^{(e_p)} \phi_j^{(e_p)} J^{(e_p)} dV_r. \end{aligned} \quad (1.67)$$

We also note that we can evaluate eq. (1.27) at $\mathbf{y} = \boldsymbol{\Phi}^{(e_p)}(0, \mathbf{r})$ to obtain

$$\rho^{(e_p)} = \frac{\rho_0^{(e_p)} J_0^{(e_p)}}{J^{(e_p)}}. \quad (1.68)$$

As with the other variables, we can define a reference basis function $w^{(e_p)}$ so that it satisfies

$$w_j^{(e_p)}(t, \mathbf{r}) = w_j^+(t, \boldsymbol{\Phi}^{(e_p)}(0, \mathbf{r})). \quad (1.69)$$

Now, as mentioned earlier, the Lagrangian basis functions are independent of time, and as a result the reference basis functions are so as well. That is, $w^{(e_p)} = w^{(e_p)}(\mathbf{r})$. Consider the expansion in eq. (1.16). Plugging in $\boldsymbol{\Phi}^{(e_p)}(0, \mathbf{r})$ for \mathbf{y} gives

$$\boldsymbol{\Phi}^{(e_p)} = \sum_j^{N_w} \hat{\boldsymbol{\Phi}}_j w_j^{(e_p)}. \quad (1.70)$$

Thus, both the Lagrangian and reference variables share the same finite-element coefficients.

1.7 Temporal integration

We now integrate forward in time the semi-discrete eqs. (1.22), (1.38) and (1.49), which we repeat below for convenience

$$\mathbf{M}_{\mathcal{V}} \frac{d\mathbf{U}}{dt} = -\mathbf{F}\hat{\mathbf{1}}. \quad (1.38)$$

$$\mathbf{M}_{\mathcal{E}} \frac{d\mathbf{E}}{dt} = \mathbf{F}^T \cdot \mathbf{U}. \quad (1.49)$$

$$\frac{d\mathbf{X}}{dt} = \mathbf{U}. \quad (1.22)$$

The equations are integrated using the RK2-average scheme of Dobrev et al. [2012], which consists of the following for the first stage

$$\begin{aligned} \mathbf{U}^{n+1/2} &= \mathbf{U}^n - \frac{\Delta t}{2} (\mathbf{M}_{\mathcal{V}})^{-1} \mathbf{F}^n \hat{\mathbf{1}}, \\ \mathbf{E}^{n+1/2} &= \mathbf{E}^n + \frac{\Delta t}{2} (\mathbf{M}_{\mathcal{E}})^{-1} (\mathbf{F}^n)^T \cdot \mathbf{U}^{n+1/2}, \\ \mathbf{X}^{n+1/2} &= \mathbf{X}^n + \frac{\Delta t}{2} \mathbf{U}^{n+1/2}, \end{aligned} \quad (1.71)$$

and the following for the second stage

$$\begin{aligned}
\mathbf{U}^{n+1} &= \mathbf{U}^n - \Delta t (\mathbf{M}_{\mathcal{V}})^{-1} \mathbf{F}^{n+1/2} \hat{\mathbf{1}}, \\
\mathbf{E}^{n+1} &= \mathbf{E}^n + \Delta t (\mathbf{M}_{\mathcal{E}})^{-1} \left(\mathbf{F}^{n+1/2} \right)^T \cdot \bar{\mathbf{U}}^{n+1/2}, \\
\mathbf{X}^{n+1} &= \mathbf{X}^n + \Delta t \bar{\mathbf{U}}^{n+1/2}.
\end{aligned} \tag{1.72}$$

In the above, $\bar{\mathbf{U}}^{n+1/2} = (\mathbf{U}^n + \mathbf{U}^{n+1})/2$. In particular, this scheme is used since it conserves total energy, that is, $(IE + KE)^{n+1} - (IE + KE)^n = 0$. To prove this we first show that for the internal energy we have

$$\begin{aligned}
IE^{n+1} - IE^n &= \hat{\mathbf{1}}^T \mathbf{M}_{\mathcal{E}} (\mathbf{E}^{n+1} - \mathbf{E}^n) \\
&= \Delta t \hat{\mathbf{1}}^T \left(\mathbf{F}^{n+1/2} \right)^T \cdot \bar{\mathbf{U}}^{n+1/2}
\end{aligned} \tag{1.73}$$

For the kinetic energy we have

$$\begin{aligned}
KE^{n+1} - KE^n &= \frac{1}{2} \left[(\mathbf{U}^{n+1})^T \cdot \mathbf{M}_{\mathcal{V}} \mathbf{U}^{n+1} - (\mathbf{U}^n)^T \cdot \mathbf{M}_{\mathcal{V}} \mathbf{U}^n \right] \\
&= \frac{1}{2} \left[(\mathbf{U}^{n+1})^T \mathbf{M}_{\mathcal{V}} \cdot \mathbf{U}^{n+1} - (\mathbf{U}^n)^T \mathbf{M}_{\mathcal{V}} \cdot \mathbf{U}^n \right] \\
&= \frac{1}{2} (\mathbf{U}^{n+1} - \mathbf{U}^n)^T \mathbf{M}_{\mathcal{V}} \cdot (\mathbf{U}^{n+1} + \mathbf{U}^n) \\
&= (\mathbf{U}^{n+1} - \mathbf{U}^n)^T \mathbf{M}_{\mathcal{V}} \cdot \bar{\mathbf{U}}^{n+1/2} \\
&= \left[-\Delta t (\mathbf{M}_{\mathcal{V}})^{-1} \mathbf{F}^{n+1/2} \hat{\mathbf{1}} \right]^T \mathbf{M}_{\mathcal{V}} \cdot \bar{\mathbf{U}}^{n+1/2} \\
&= -\Delta t \hat{\mathbf{1}}^T \left(\mathbf{F}^{n+1/2} \right)^T \left[(\mathbf{M}_{\mathcal{V}})^{-1} \right]^T \mathbf{M}_{\mathcal{V}} \cdot \bar{\mathbf{U}}^{n+1/2} \\
&= -\Delta t \hat{\mathbf{1}}^T \left(\mathbf{F}^{n+1/2} \right)^T \left[(\mathbf{M}_{\mathcal{V}})^T \right]^{-1} \mathbf{M}_{\mathcal{V}} \cdot \bar{\mathbf{U}}^{n+1/2} \\
&= -\Delta t \hat{\mathbf{1}}^T \left(\mathbf{F}^{n+1/2} \right)^T (\mathbf{M}_{\mathcal{V}})^{-1} \mathbf{M}_{\mathcal{V}} \cdot \bar{\mathbf{U}}^{n+1/2} \\
&= -\Delta t \hat{\mathbf{1}}^T \left(\mathbf{F}^{n+1/2} \right)^T \cdot \bar{\mathbf{U}}^{n+1/2}.
\end{aligned} \tag{1.74}$$

Thus, adding eq. (1.73) and eq. (1.74) leads to total energy conservation.

Chapter 2

The Re-mesh Step

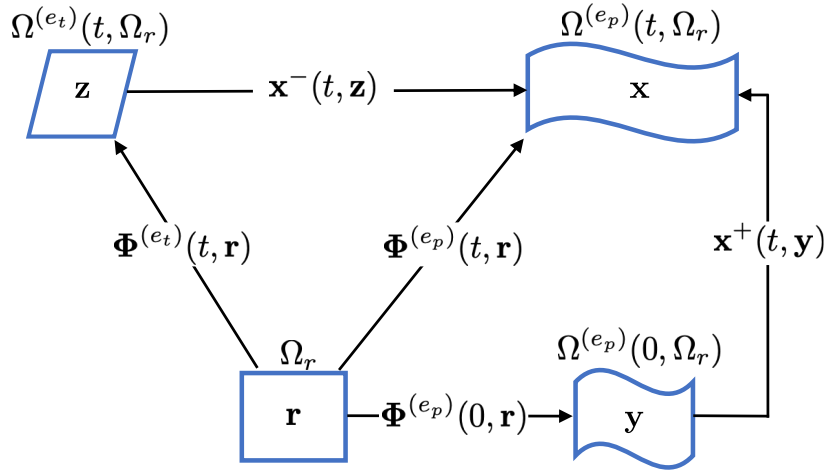


Figure 2.1: Schematic of the four domains Ω_r , $\Omega^{(e_p)}(t, \Omega_r)$, $\Omega^{(e_p)}(0, \Omega_r)$, $\Omega^{(e_t)}(t, \Omega_r)$.

We introduce a new space, the target space, which is divided into target elements, where each corresponds to a physical element e_p . Consider a mapping $\Phi^{(e_t)} = \Phi^{(e_t)}(t, \mathbf{r})$ from a point \mathbf{r} in the reference element to a point in the target element. Also consider the mapping $\mathbf{x}^- = \mathbf{x}^-(t, \mathbf{z})$ from a point \mathbf{z} in the target space to a point in the physical space. Note that $\Phi^{(e_p)}$, $\Phi^{(e_t)}$, and \mathbf{x}^- are related to each other according to

$$\Phi^{(e_p)}(t, \mathbf{r}) = \mathbf{x}^-(t, \Phi^{(e_t)}(t, \mathbf{r})). \quad (2.1)$$

We define the Jacobians as follows

$$\mathbf{J}^{(e_t)} = \frac{\partial \Phi^{(e_t)}}{\partial \mathbf{r}}, \quad (2.2)$$

$$\mathbf{J}^- = \frac{\partial \mathbf{x}^-}{\partial \mathbf{z}}. \quad (2.3)$$

where $\mathbf{J}^{(e_t)} = \mathbf{J}^{(e_t)}(t, \mathbf{r})$ and $\mathbf{J}^- = \mathbf{J}^-(t, \mathbf{z})$. Taking the derivative of eq. (2.1) we get

$$\frac{\partial \Phi^{(e_p)}}{\partial \mathbf{r}} = \left(\frac{\partial \mathbf{x}^-}{\partial \mathbf{z}} \right)_{\mathbf{z}=\Phi^{(e_t)}} \frac{\partial \Phi^{(e_t)}}{\partial \mathbf{r}}, \quad (2.4)$$

which we write as

$$\mathbf{J}^{(e_p)} = (\mathbf{J}^-)_{\mathbf{z}=\Phi^{(e_t)}} \mathbf{J}^{(e_t)}. \quad (2.5)$$

Multiplying both sides by the inverse of $\mathbf{J}^{(e_t)}$ we finally get

$$(\mathbf{J}^-)_{\mathbf{z}=\Phi^{(e_t)}} = \mathbf{J}^{(e_p)} \left(\mathbf{J}^{(e_t)} \right)^{-1}. \quad (2.6)$$

Combining eq. (1.55) and eq. (2.1) we get

$$\mathbf{x}^-(t, \Phi^{(e_t)}(t, \mathbf{r})) = \mathbf{x}^+(t, \Phi^{(e_p)}(0, \mathbf{r})). \quad (2.7)$$

We also define a target basis function $w^{(e_t)} = w^{(e_t)}(t, \mathbf{z})$ so that it satisfies

$$w^-(t, \Phi^{(e_t)}(t, \mathbf{r})) = w^+(t, \Phi^{(e_p)}(0, \mathbf{r})). \quad (2.8)$$

Consider the expansion in eq. (1.16). Plugging in $\Phi^{(e_p)}(0, \mathbf{r})$ for \mathbf{y} gives

$$\mathbf{x}^-(t, \Phi^{(e_t)}(t, \mathbf{r})) = \sum_j^{N_w} \hat{\mathbf{x}} w^-(t, \Phi^{(e_t)}(t, \mathbf{r})). \quad (2.9)$$

Assuming this holds for any $\Phi^{(e_t)}$ we get

$$\mathbf{x}^- = \sum_j^{N_w} \hat{\mathbf{x}} w^-. \quad (2.10)$$

Thus, both the Lagrangian and the target variables share the same finite-element coefficients.

To obtain the relaxed mesh, one minimizes the following function

$$F(\mathbf{X}) = \sum_{e_t \in \mathcal{M}_t} \int_{\Omega^{(e_t)}} \mu(\mathbf{J}^-) dV_z \quad (2.11)$$

Chapter 3

The Re-map Step

Bibliography

V. A. Dobrev, T. V. Kolev, and R. N. Rieben. High-order curvilinear finite element methods for Lagrangian hydrodynamics. *SIAM Journal on Scientific Computing*, 34(5), 2012. doi: <https://doi.org/10.1137/120864672>.