Campos electromagnéticos dependientes del tiempo



$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_o}$$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{\nabla} \mathbf{x} \vec{E} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \vec{xB} = \mu_0 \vec{J}$$

Ecuaciones de Maxwell para campos estáticos

Electricidad

dos fenómenos independientes

Magnetismo

Ley de Faraday-Henry



Michael Faraday (1791-1867)



Joseph Henry (1797-1878)



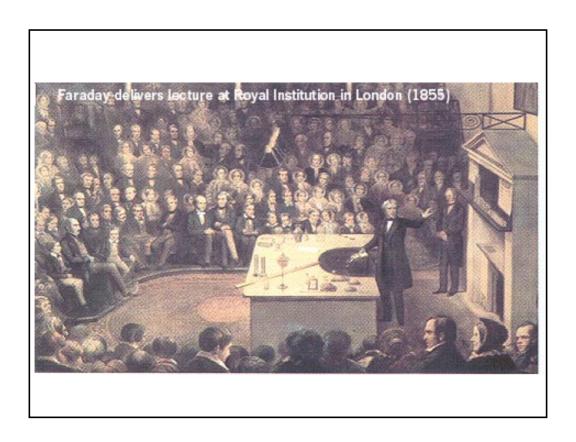


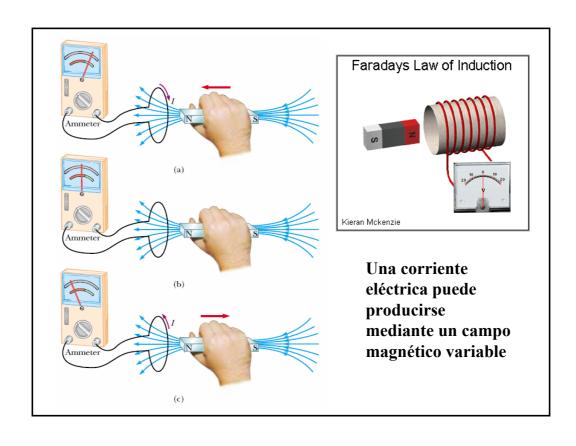




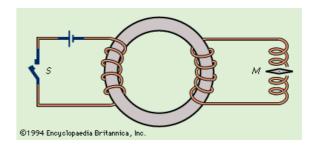




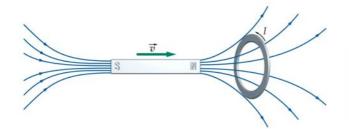




Experimento de Faraday

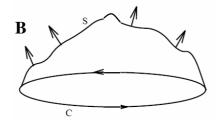


Un campo magnético variable induce una fem en el circuito secundario



$$\varepsilon = -\frac{d\Phi_B}{dt}$$

Enunciado general de la ley de Faraday-Henry



$$\mathcal{E} = -\frac{d\phi_{\rm m}}{dt}$$

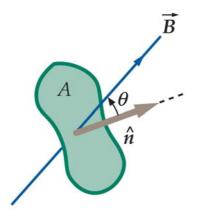
$$\phi_{\rm m} = \int_{S} \vec{B} \cdot \hat{n} \, dA = \int_{S} B_n \, dA$$

$$\mathcal{E} = \oint_C \overrightarrow{E}_{\text{nc}} \cdot d\overrightarrow{\ell} = -\frac{d}{dt} \int_S \overrightarrow{B} \cdot \hat{n} \, dA = -\frac{d\phi_{\text{m}}}{dt}$$

En una bobina con N vueltas, todas de la misma área

$$\mathbf{\varepsilon} = -N \frac{d\Phi_B}{dt}$$

 $\Phi_{\rm B}(t)$?

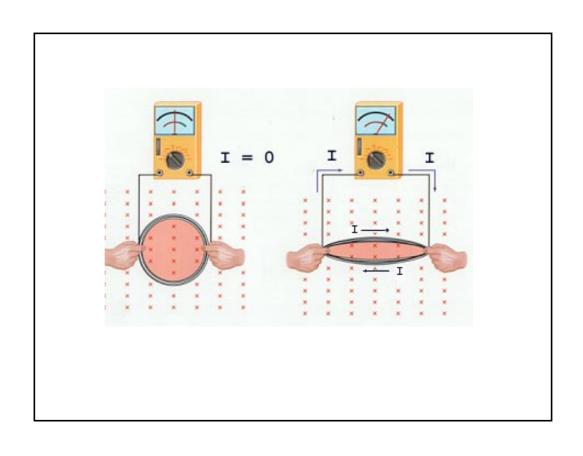


$$\boldsymbol{\mathcal{E}} = -\frac{d}{dt} \left(BA \cos \theta \right)$$

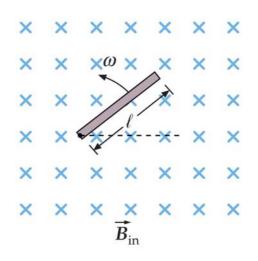
$$-B(t)$$

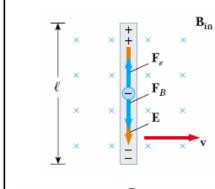
$$-A(t)$$

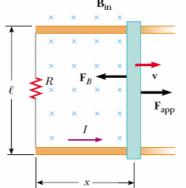
$$- A(t)$$
$$- \theta(t)$$



fem de movimiento







$$qE = qvB$$

$$E = vB$$

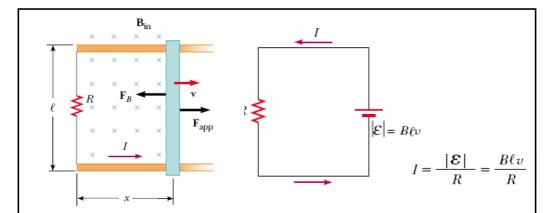
$$\Delta V = E\ell = B\ell v$$

Si se invierte v, se invierte la polaridad de V

$$\Phi_B = B\ell x$$

$$\boldsymbol{\varepsilon} = -\frac{d\Phi_B}{dt} = -\frac{d}{dt} \left(B\ell \boldsymbol{x} \right) = -B\ell \frac{d\boldsymbol{x}}{dt}$$

$$\varepsilon = -B\ell v$$



Consideración energética: si v = cte —

$$\longrightarrow$$
 $F_{\text{app}} = I\ell I$

La potencia entregada por la fuerza aplicada es:

$$\mathcal{P} = F_{\rm app} v = \left(I\ell B\right) v = \frac{B^2 \, \ell^2 \, v^2}{R} = \frac{\mathcal{E}^2}{R}$$

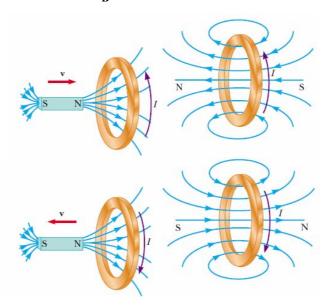
Energía Mecánica Energía Energía Térmica

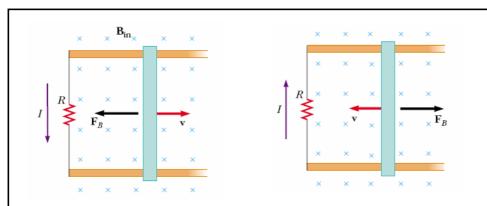
Ley de Lenz



Heinrich Lenz (1804-1865)

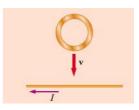
La polaridad de una fem inducida es tal que tiende a producir una corriente eléctrica que creara un flujo magnético que se opone al cambio de $\Phi_{\rm B}$ a través del lazo.

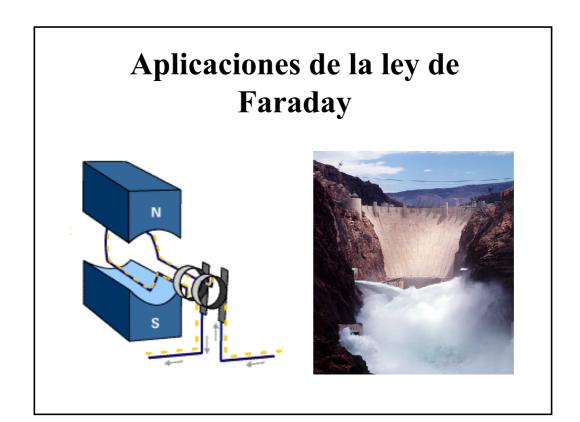


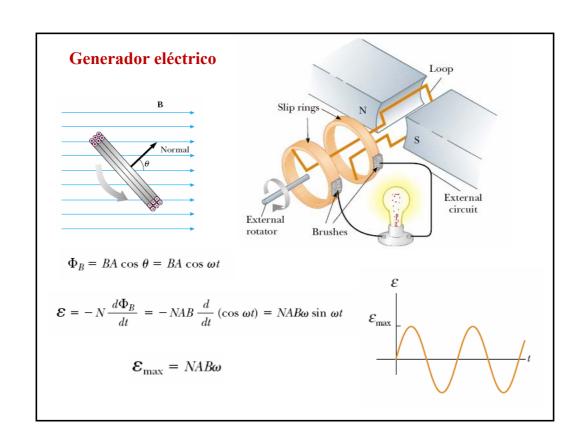


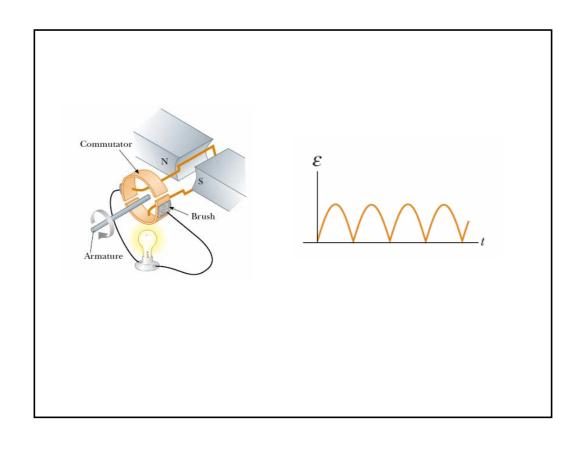
Qué sucede si el sentido de I es contrario al obtenido utilizando la ley de Lenz en los dos casos anteriores ?

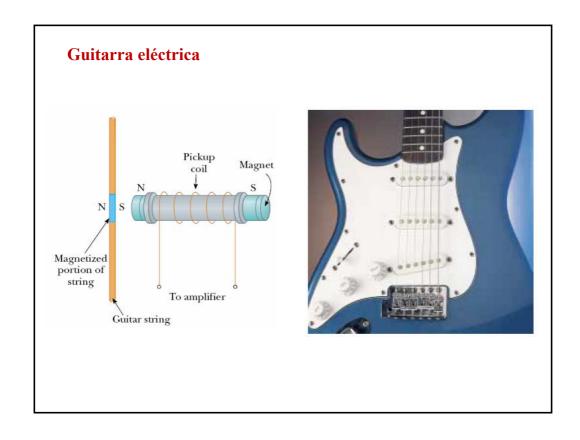
Determinar el sentido de la corriente inducida en la espira circular.

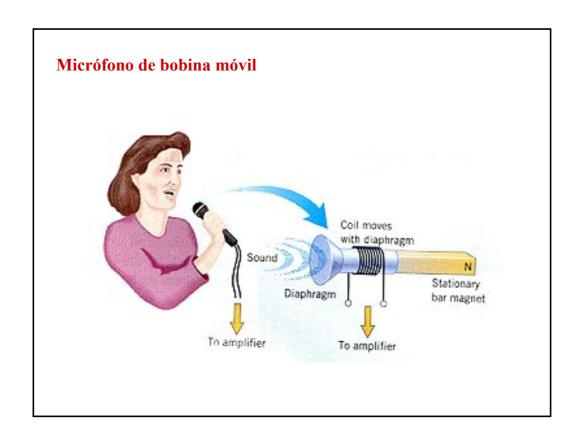










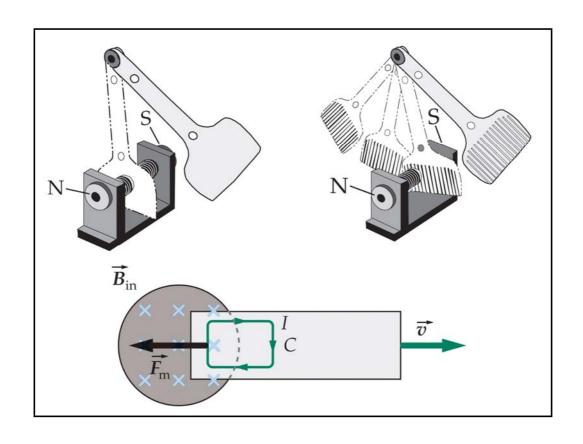


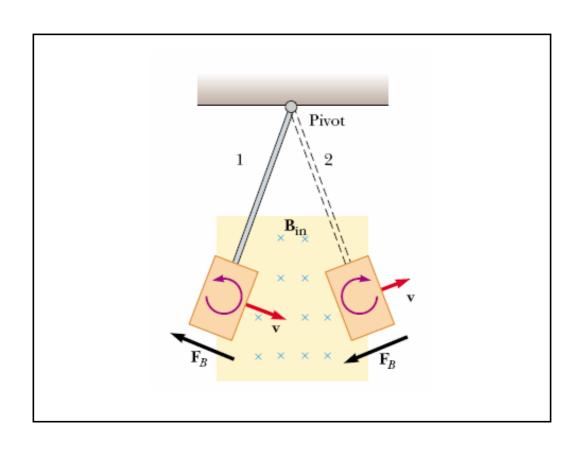
Corrientes de Foucault

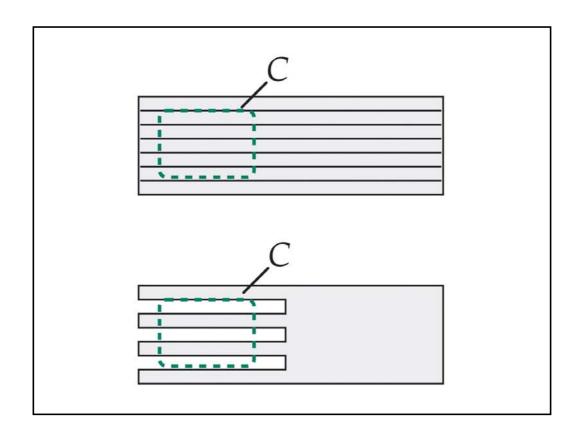




Jean Leon Foucault (1819-1868)

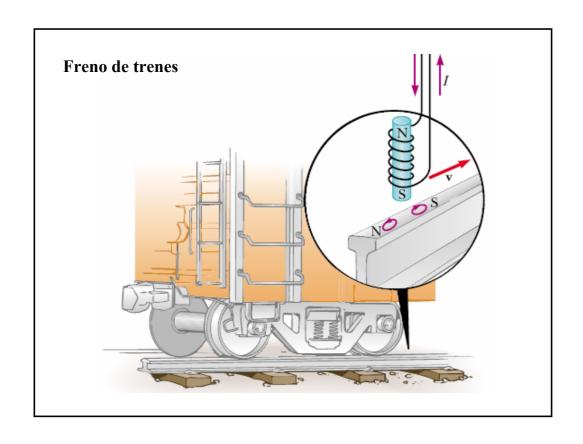


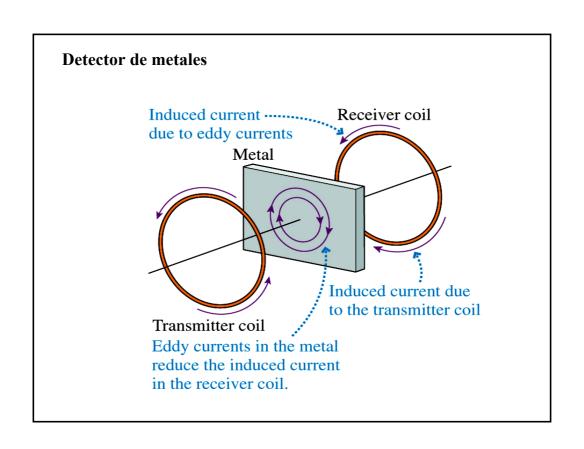


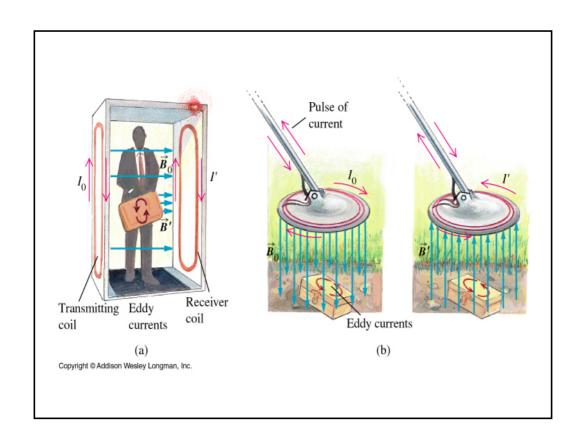


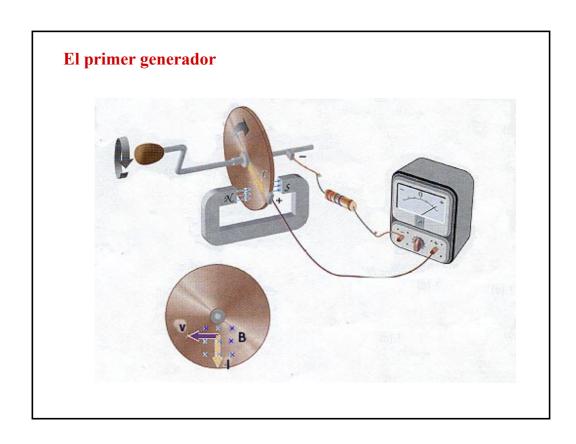
Son útiles para amortiguar oscilaciones mecánicas





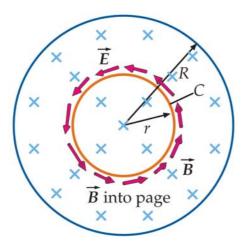


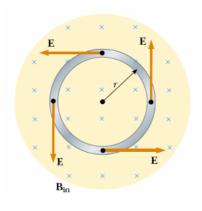






fem inducidas y campos eléctricos





$$\mathcal{E} = -\frac{d\phi_{\rm m}}{dt}$$

$$\phi_{\rm m} = \int_{S} \vec{B} \cdot \hat{n} \, dA = \int_{S} B_{n} \, dA$$

$$\mathcal{E} = \oint_{C} \overrightarrow{E}_{nc} \cdot d\overrightarrow{\ell} = -\frac{d}{dt} \int_{S} \overrightarrow{B} \cdot \hat{n} \, dA = -\frac{d\phi_{m}}{dt}$$

El campo eléctrico inducido no es conservativo

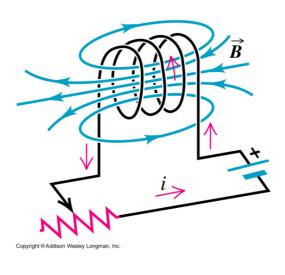
$$\oint_C \vec{E} \cdot \vec{dl} = -\frac{d}{dt} \int_S \vec{B} \cdot \vec{dA}$$

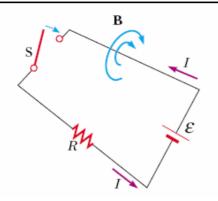
$$\begin{cases}
\int_{\mathcal{L}} \vec{E} \cdot d\vec{l} &= \int_{\mathcal{S}} \nabla \wedge \vec{E} \cdot d\vec{s} \\
\frac{d}{dt} \int_{\mathcal{S}} \vec{B} \cdot d\vec{s} &= \int_{\mathcal{S}} \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}
\end{cases} \rightarrow \nabla \wedge \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

La ley de Faraday-Henry en forma diferencial







S se cierra \Longrightarrow I (t)

$$\bigcup_{\mathbf{B}(\mathbf{t})}$$

Variación temporal de Φ_B a través del área del circuito fem autoinducida

$$\varepsilon = -\frac{d\Phi_B}{dt}$$

Definimos:

$$\Phi_{\rm B} = L I$$

L: coeficiente de autoinductancia
o inductancia

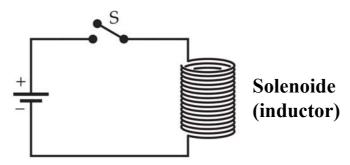
$$\mathbf{\varepsilon}_L = -L \frac{dI}{dt}$$

$$L = \Phi_B / I = -\mathbf{E}_L / (dI/dt)$$

L

$$[L] = T m^2 / A = Wb / A = V s / A = H (Henry)$$

La inductancia de un circuito depende de su geometría



Lenz's law emf
$$--\frac{1}{2} \begin{vmatrix} \frac{1}{2} - \frac{1}{2} - \frac{1}{2} \end{vmatrix} = -\frac{1}{2} \begin{vmatrix} \frac{1}{2} - \frac{1}{2} - \frac{1}{2} \end{vmatrix}$$

$$L = \frac{N\Phi_B}{I}$$

$$B = \mu_0 nI = \mu_0 \frac{N}{\ell} I$$

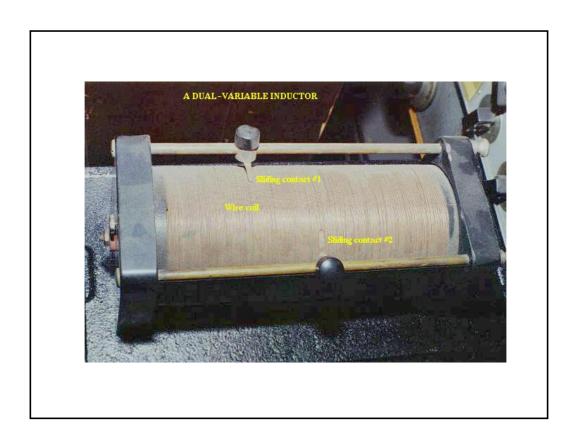
$$n = N/\ell$$

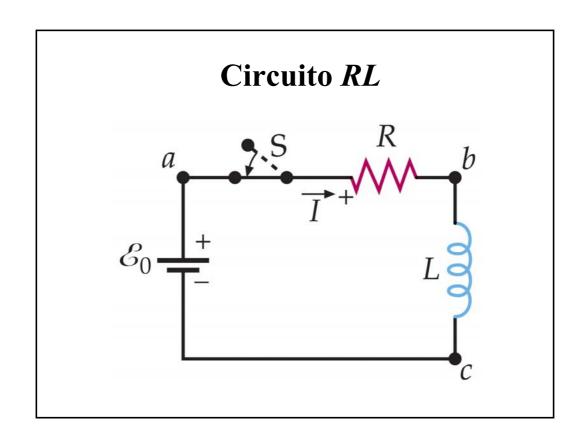
$$\Phi_B = BA = \mu_0 \frac{NA}{\ell} I$$

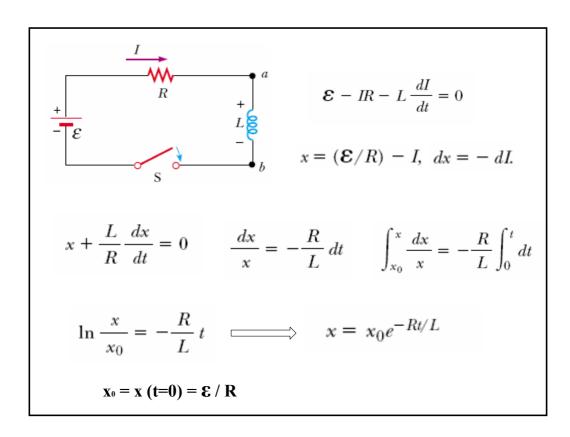
$$L = \frac{N\Phi_B}{I} = \frac{\mu_0 N^2 A}{\ell}$$

$$L = \mu_0 \frac{(n\ell)^2}{\ell} A = \mu_0 n^2 A \ell = \mu_0 n^2 V$$

$$V = A\ell$$





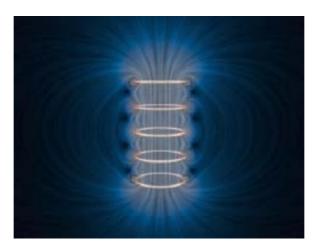


$$\frac{\mathcal{E}}{R} - I = \frac{\mathcal{E}}{R} e^{-Rt/L} \qquad I = \frac{\mathcal{E}}{R} (1 - e^{-Rt/L})$$

$$I = \frac{\mathcal{E}}{R} (1 - e^{-t/\tau})$$

$$\tau = \frac{L}{R}$$
Constante de tiempo del circuito RL
$$\tau = \frac{L}{R}$$

Energía del Campo Magnético



La fem inducida evita que la batería establezca una corriente: la batería efectúa trabajo contra el inductor

Energía suministrada por la batería

Energía almacenada en el inductor

$$\mathcal{E} - IR - L \frac{dI}{dt} = 0$$
 $I\mathcal{E} = I^2R + LI \frac{dI}{dt}$

$$I\mathbf{\varepsilon} = I^2 R + L I \frac{dI}{dt}$$

U: energía almacenada en $\frac{dU}{dt} = LI \frac{dI}{dt}$ el inductor

$$\frac{dU}{dt} = LI \frac{dI}{dt}$$

$$U = \int dU = \int_0^I LI \, dI = L \int_0^I I \, dI \qquad U = \frac{1}{2} LI^2 \qquad U = \frac{1}{2} C(\Delta V)^2.$$

$$U = \frac{1}{2} L I^2$$

$$U = \frac{1}{2}C(\Delta V)^2.$$

Análogo eléctrico

Para calcular la densidad de energía almacenada en el campo magnético, consideramos un solenoide:

$$L = \mu_0 n^2 A \ell$$

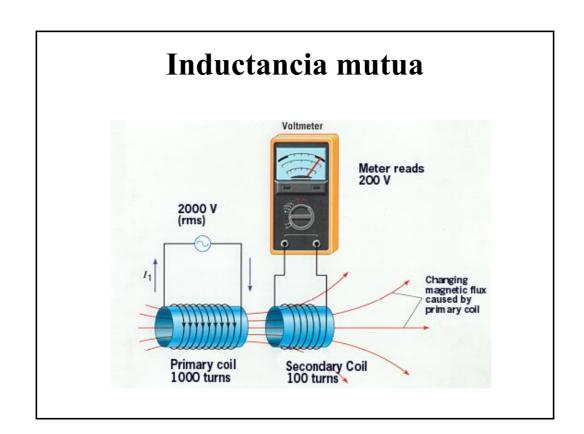
El campo magnético dentro $B = \mu_0 nI$ del solenoide es

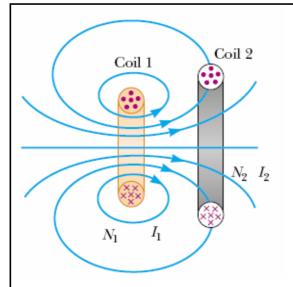
$$U = \frac{1}{2} L I^2 = \frac{1}{2} \, \mu_0 n^2 A \, \ell \left(\frac{B}{\mu_0 n} \right)^2 = \frac{B^2}{2 \mu_0} \, A \ell$$

 $A \ell$ volumen del solenoide donde se encuentra confinado B

$$u_B = \frac{U}{A\ell} = \frac{B^2}{2\mu_0}$$

Es válida para cualquier región del espacio donde haya B





 Φ_{12}

Flujo magnético a través de la bobina 2 producido por la bobina 1.

Definimos inductancia mutua de la bobina 2 respecto a la 1

$$M_{12} \equiv \frac{N_2 \Phi_{12}}{I_1}$$

M12 depende de la geometría de ambos circuitos y de la orientación de uno respecto del otro.

La fem inducida en la bobina 2 es:

$$\mathbf{\mathcal{E}}_{2} = -N_{2} \frac{d\Phi_{12}}{dt} = -N_{2} \frac{d}{dt} \left(\frac{M_{12}I_{1}}{N_{2}} \right) = -M_{12} \frac{dI_{1}}{dt}$$

De la misma manera:

$$\mathbf{\varepsilon}_1 = -M_{21} \frac{dI_2}{dt}$$

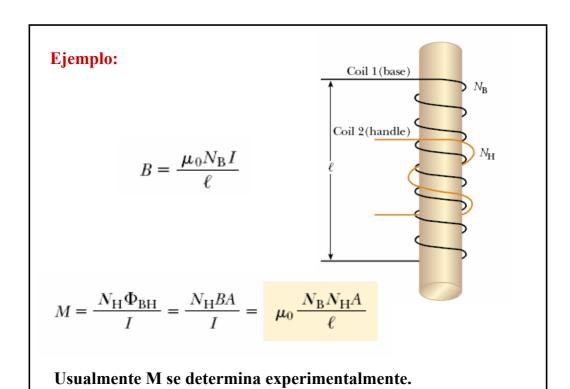
Se puede demostrar que $M_{12} = M_{21} = M_{21}$

$$M_{12} = M_{21} = M$$

$$\mathbf{\varepsilon}_2 = -M \frac{dI_1}{dt}$$
 $\mathbf{\varepsilon}_1 = -M \frac{dI_2}{dt}$

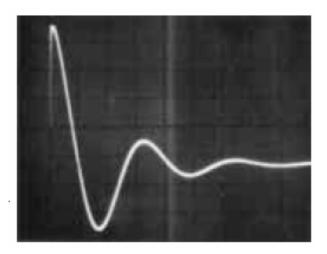
$$\varepsilon_1 = -M \frac{dI_2}{dt}$$

$$[M] = Hy$$





Oscilaciones eléctricas Circuitos LC y RLC



$$U = U_C + U_L = \frac{Q^2}{2C} + \frac{1}{2}LI^2$$

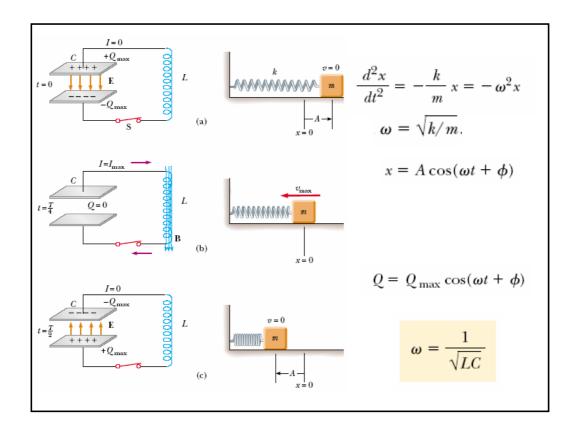
$$U = U_C + U_L = \frac{Q^2}{2C} + \frac{1}{2}LI^2$$

$$\frac{dU}{dt} = \frac{d}{dt} \left(\frac{Q^2}{2C} + \frac{1}{2}LI^2 \right) = \frac{Q}{C} \frac{dQ}{dt} + LI \frac{dI}{dt} = 0$$

$$I = dQ/dt. \qquad \longrightarrow \qquad \frac{Q}{C} + L \frac{d^2Q}{dt^2} = 0$$

$$\frac{d^2Q}{dt^2} = -\frac{1}{LC}Q$$

 $\frac{d^2Q}{dt^2} = -\frac{1}{LC}Q$ Ecuación diferencial del oscilador armónico



$$Q = Q_{\text{max}} \cos(\omega t + \phi) \qquad I = \frac{dQ}{dt} = -\omega Q_{\text{max}} \sin(\omega t + \phi)$$

$$I = 0 \text{ at } t = 0 \qquad 0 = -\omega Q_{\text{max}} \sin \phi$$

$$Q = Q_{\text{max}} \cos \omega t$$

$$I = -\omega Q_{\text{max}} \sin \omega t = -I_{\text{max}} \sin \omega t$$

$$I = -\omega Q_{\text{max}} \sin \omega t = -I_{\text{max}} \sin \omega t$$

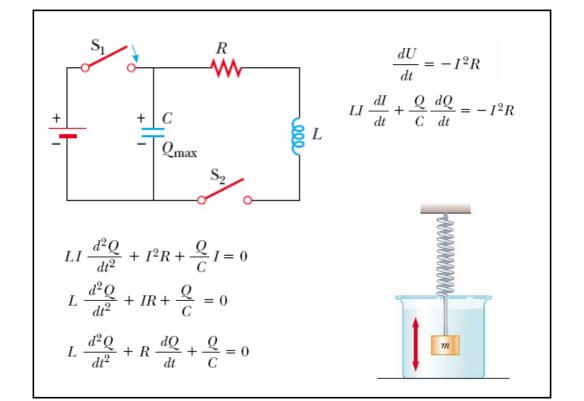
$$U = U_C + U_L = \frac{Q_{\text{max}}^2}{2C} \cos^2 \omega t + \frac{1}{2} L I_{\text{max}}^2 \sin^2 \omega t$$

$$U_C$$

$$\frac{Q_{\text{max}}^2}{2C} \qquad \frac{Q_{\text{max}}^2}{2C} = \frac{L I_{\text{max}}^2}{2}$$

$$U_L \qquad U = \frac{Q_{\text{max}}^2}{2C} (\cos^2 \omega t + \sin^2 \omega t) = \frac{Q_{\text{max}}^2}{2C}$$

$$\frac{L I_{\text{max}}^2}{2} \qquad Que \ ocurre \ si \ agrego \ R??$$



$$L \frac{d^2Q}{dt^2} + R \frac{dQ}{dt} + \frac{Q}{C} = 0 \iff m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

$$m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

$$Q \leftrightarrow x$$

$$I \leftrightarrow v_x$$

$$\Delta V \leftrightarrow F_x$$

$$m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx = 0$$

$$C \leftrightarrow 1/k$$

$$Q = Q_{\max} e^{-Rt/2L} \cos \omega_d t$$

$$I = \frac{dQ}{dt} \iff v_x = \frac{dx}{dt}$$

$$\frac{dI}{dt} = \frac{d^2Q}{dt^2} \iff a_x = \frac{dv_x}{dt} = \frac{d^2x}{dt^2}$$

donde

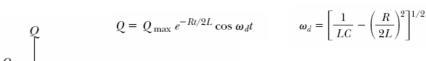
$$\omega_d = \left[\frac{1}{LC} - \left(\frac{R}{2L}\right)^2\right]^{1/2} \qquad U_L = \frac{1}{2}LI^2 \iff K = \frac{1}{2}mv^2$$

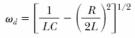
$$U_C = \frac{1}{2}\frac{Q^2}{C} \iff U = \frac{1}{2}kx^2$$

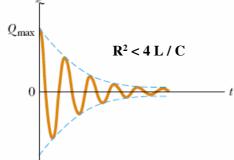
$$U_L = \frac{1}{2} L I^2 \iff K = \frac{1}{2} m v^2$$

$$U_C = \frac{1}{2} \frac{Q^2}{C} \iff U = \frac{1}{2} kx^2$$

$$I^2R \leftrightarrow bv^2$$

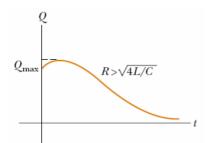






El sistema esta subamortiguado

 $R^2 = 4 L / C$ amortiguamiento crítico



 $R^2 > 4 L/C$ sobreamortiguado