

Our goal is to study the homogeneous magnetostatic problem on the exterior domain of a triangulated torus. That means that for the unbounded domain  $\Omega \subseteq \mathbb{R}^3$  we have  $\mathbb{R}^3 \setminus \Omega$  is a triangulated torus. We also need a piecewise straight (i.e. triangulated) closed curve around the torus. (TBD: Define the "triangulated torus" more rigorous)

Let  $B$  be a magnetic field on the domain  $\Omega$ . We then have the following boundary value problem:

$$\operatorname{curl} B = 0, \quad (1)$$

$$\operatorname{div} B = 0 \text{ in } \Omega \quad (2)$$

$$B \cdot n = 0 \text{ on } \partial\Omega \text{ and} \quad (3)$$

$$\int_{\gamma} B \cdot dl = C_0 \quad (4)$$

where  $n$  is the outward normal vector field on  $\partial\Omega$  and  $C_0 \in \mathbb{R}$ . We want to prove existence and uniqueness of solutions. In order to do so we will need to introduce Sobolev spaces of differential forms and basics from simplicial topology.

At first, let us introduce some basic notions about differential forms. We follow the brief introduction given by Arnold (cf. [1, Sec. 6.1]), but less details will be given. Let  $V$  be a real vector space with  $\dim V = n$  and  $\operatorname{Alt}^k V$  be the space of  $k$ -alternating maps from  $V^k$  to  $\mathbb{R}$ . For  $\omega \in \operatorname{Alt}^k V$ ,  $\mu \in \operatorname{Alt}^l V$  we define the wedge product  $\omega \wedge \mu \in \operatorname{Alt}^{k+l} V$

$$(\omega \wedge \mu)(v_1, \dots, v_k, v_{k+1}, \dots, v_{k+l}) = \sum_{\substack{i_1 < \dots < i_k \\ i_{k+1} < \dots < i_{k+l}}} \operatorname{sgn}(i_1, \dots, i_{k+l}) \omega(v_{i_1}, \dots, v_{i_k}) \mu(v_{i_{k+1}}, \dots, v_{i_{k+l}})$$

where  $\operatorname{sgn}(i_1, \dots, i_{k+l})$  is the sign of the permutation  $(1, \dots, k+l) \mapsto (i_1, \dots, i_{k+l})$ . This definition is not very intuitive. TBD: Examples in 3D.

Let  $\{u_i\}_{i=1}^n$  be any basis of  $V$  and  $\{u^i\}_{i=1}^n$  the corresponding dual basis. Then

$$\{u^{i_1} \wedge u^{i_2} \wedge \dots \wedge u^{i_k} \mid 1 \leq i_1 < \dots < i_k \leq n\}$$

is a basis of  $\operatorname{Alt}^k V$ . In particular,  $\dim \operatorname{Alt}^k V = \binom{n}{k}$ .

Given an inner product  $\langle \cdot, \cdot \rangle_V$  on  $V$  we obtain an inner product on  $\operatorname{Alt}^k V$  by defining

$$\langle u^{i_1} \wedge u^{i_2} \wedge \dots \wedge u^{i_k}, u^{j_1} \wedge \dots \wedge u^{j_k} \rangle_{\operatorname{Alt}^k V} := \det [(\langle u_{i_k}, u_{i_l} \rangle_V)_{1 \leq k, l \leq n}]$$

which is then extended to all of  $\text{Alt}^k V$  by linearity. We denote with  $\|\cdot\|_{\text{Alt}^k V}$  the induced norm. From this definition it follows directly that for a orthonormal basis  $b_1, \dots, b_n$  the corresponding basis  $b^{i_1} \wedge b^{i_2} \wedge \dots \wedge b^{i_k}$ ,  $1 \leq i_1 < \dots < i_k \leq n$  is an orthonormal basis of  $\text{Alt}^k V$ .

$\text{Alt}^n V$  is one-dimensional and so we have to choose a basis. We say that two orthonormal bases of  $V$  have the same orientation if the change of basis has positive determinant. That divides the orthonormal bases into two classes with different orientation. We choose one of these classes and call these orthonormal bases positively oriented. Take  $\omega \in \text{Alt}^n V$ . Then  $\omega(b_1, \dots, b_n)$  is the same for any positively oriented orthonormal basis. We now define the *volume form*  $\text{vol} \in \text{Alt}^n V$  by requiring it to be 1 on all positively oriented orthonormal bases. Using this volume form we can now define the *Hodge star operator*  $*$  :  $\text{Alt}^k V \rightarrow \text{Alt}^{n-k} V$  via the property

$$\omega \wedge \mu = \langle * \omega, \mu \rangle_{\text{Alt}^{n-k} V} \text{vol} \quad \forall \omega \in \text{Alt}^k V, \mu \in \text{Alt}^{n-k} V.$$

The Hodge star is an isometry, we have  $** = (-1)^{k(n-k)} \text{Id}$  and

$$\omega \wedge * \mu = \langle \omega, \mu \rangle_{\text{Alt}^k V} \text{vol} \quad \forall \omega, \mu \in \text{Alt}^k V.$$

In particular in  $\mathbb{R}^3$ , we have  $** = \text{Id}$  i.e.  $*$  is self-inverse.

Now we will move on to differential forms. We will mostly deal with the case  $V = \mathbb{R}^n$  and denote  $\text{Alt}^k \mathbb{R}^n$  just as  $\text{Alt}^k$ . Let  $\Omega \subseteq \mathbb{R}^n$  be a Lipschitz domain and denote the space of differential forms of degree  $k$  on  $\Omega$  as  $\Lambda^k(\Omega)$ . We extend the Hodge star operator to differential forms  $*$  :  $\Lambda^k(\Omega) \rightarrow \Lambda^{n-k}(\Omega)$  simply by applying it pointwise.

Then we define the  $L_p$ -norm of a  $k$ -form  $\omega$  for  $1 \leq p < \infty$  as (cf. [3])

$$\|\omega\|_{L_p^k(\Omega)} := \left( \int_{\Omega} \|\omega\|_{\text{Alt}^k}^p \right)^{1/p}$$

and for  $p = \infty$  as

$$\text{ess sup}_{x \in \Omega} \|\omega(x)\|_{\text{Alt}^k}.$$

$L_p^k(\Omega)$  are the spaces of  $k$ -forms s.t. the corresponding  $L_p$ -norm is finite. For  $p = 2$  we obtain a Hilbert space (cf. [1, Sec. 6.2.6]) with the  $L_2$  inner product

$$\langle \omega, \nu \rangle := \int_{\Omega} \langle \omega, \nu \rangle_{\text{Alt}^k}. \quad (5)$$

**Proposition 1.** *The Hodge star operator  $*$  :  $L_2^k(\Omega) \rightarrow L_2^{n-k}(\Omega)$  is a Hilbert space isometry.*

*Proof.* This follows directly from the definition of the inner product (5) and the fact that  $*$  is an isometry when applied to alternating forms  $\text{Alt}^k$ .  $\square$

Our next goal is to extend the exterior derivative  $d$  of smooth differential forms in the weak sense (cf. [3]). Let  $\overset{\circ}{d} : L_2^k(\Omega) \rightarrow L_2^{k+1}(\Omega)$  be the exterior derivative as an unbounded operator with domain  $D(\overset{\circ}{d}) = C_0^\infty \Lambda^k(\Omega)$  which are the smooth compactly supported differential forms of degree  $k$ . Analogous, let  $\overset{\circ}{\delta} : L_2^k(\Omega) \rightarrow L_2^{k-1}(\Omega)$  be the codifferential operator  $\overset{\circ}{\delta} := (-1)^{n(k-1)+1} * \overset{\circ}{d} *$  also with domain  $C_0^\infty \Lambda^k(\Omega)$ .

Then the exterior derivative  $d\omega \in L_{k+1}^p(\Omega)$  is defined as the unique  $(k+1)$ -form in  $L_{k+1}^p(\Omega)$  s.t.

$$\int_{\Omega} d\omega \wedge * \phi = \int_{\Omega} \omega \wedge * \overset{\circ}{\delta} \phi \quad \forall \phi \in C_0^\infty \Lambda^k(\Omega).$$

Just as in the usual Sobolev setting we define the following spaces:

$$\begin{aligned} W_p^k(\Omega) &= \{ \omega \in L_p^k(\Omega) \mid d\omega \in L_p^{k+1}(\Omega) \}, \\ W_{p,loc}^k(\Omega) &= \{ \omega \text{ } k\text{-form} \mid \omega|_A \in W_p^k(A) \text{ for every open } A \subseteq \Omega \text{ s.t. } \overline{A} \subseteq \Omega \text{ is compact} \}. \end{aligned}$$

For  $\omega \in W_p^k(\Omega)$  for  $p < \infty$  we define the norm

$$\|\omega\|_{W_p^k} := \left( \|\omega\|_{L_p^k}^p + \|d\omega\|_{L_p^{k+1}}^p \right)^{1/p}$$

and for  $p = \infty$

$$\|\omega\|_{W_\infty^k} := \max \{ \|\omega\|_{L_\infty^k}, \|d\omega\|_{L_\infty^{k+1}} \}.$$

**Definition 1** ( $L^p$ -cohomology). We define the following subspaces of  $W_p^k(\Omega)$ ,  $1 \leq p \leq \infty$ :

$$\begin{aligned} \mathfrak{B}_k &:= dW_p^{k-1}(\Omega) \text{ and} \\ \mathfrak{Z}_k &:= \{ \omega \in W_p^k(\Omega) \mid d\omega = 0 \}. \end{aligned}$$

We call the  $k$ -forms in  $\mathfrak{B}_k$  exact and the forms in  $\mathfrak{Z}_k$  closed. Because  $d \circ d = 0$  we always have  $\mathfrak{B}_k \subseteq \mathfrak{Z}_k$ . Then we define the de Rham- or  $L^p$ -cohomology space  $H_{p,dR}^k(\Omega)$  as the quotient space

$$H_{p,dR}^k(\Omega) := \mathfrak{Z}_k / \mathfrak{B}_k.$$

We want to examine the Hilbert space  $L_2^k(\Omega)$  more closely (see [1, Sec. 6.2.6] for more details). We denote  $H^k(d; \Omega) := W_2^k(\Omega)$ . If the domain is clear

we will leave it out. Note that the above definition of the exterior derivative is in the Hilbert space setting equivalent to defining  $d$  as the adjoint of  $\mathring{d}$ .

Now we just define  $\delta := (-1)^{n(k-1)+1} * d *$  as in the smooth setting. We will show that this is the adjoint of  $\mathring{d}$ . Denote with  $D(\mathring{d}^*) \subseteq L_2^{k+1}(\Omega)$  the domain of the adjoint. Define

$$H^k(\delta; \Omega) := \{\omega \in L_2^k(\Omega) \mid * \omega \in H^{n-k}(d)\}.$$

Now take  $\omega \in H^{k+1}(\delta)$  and  $\phi \in C_0^\infty \Lambda^k(\Omega)$ . Then

$$\begin{aligned} \langle \delta \omega, \phi \rangle &= (-1)^{nk+1} \langle * d * \omega, \phi \rangle \\ &= (-1)^{nk+1} (-1)^{k(n-k)} \langle d * \omega, * \phi \rangle = (-1)^{nk+1} (-1)^{k(n-k)} \langle * \omega, \mathring{d} * \phi \rangle \\ &= (-1)^{nk+1} (-1)^{k(n-k)} (-1)^{n(n-k-1)+1} \langle * \omega, * \mathring{d} * \phi \rangle \\ &= (-1)^{n(n-1)+2} (-1)^{k(n-k)} \langle \omega, \mathring{d} * \phi \rangle \\ &= \langle \omega, \mathring{d} \phi \rangle \end{aligned}$$

where we used repeatedly that  $*$  is an isometry and  $** = (-1)^{k(n-k)} \text{Id}$ . This shows that  $H^{k+1}(\delta) \subseteq D(\mathring{d}^*)$  and that  $\mathring{d}^* \omega = \delta \omega$ . Now for the other inclusion assume that  $\omega \in D(\mathring{d}^*)$  and take  $\phi \in C_0^\infty \Lambda^{n-k}(\Omega)$  arbitrary.

$$\langle * \omega, \mathring{d} \phi \rangle = \pm \langle \omega, \mathring{d} * \phi \rangle = \pm \langle \mathring{d}^* \omega, * \phi \rangle = \pm \langle * \mathring{d}^* \omega, \phi \rangle.$$

Here we use  $\pm$  to mean that we choose the sign correctly, s.t. all the operations are correct. Then by choosing the sign appropriately we find that  $\pm * \mathring{d}^* \omega = d * \omega$  and therefore  $* \omega \in H^{n-k-1}(d)$  so we proved  $D(\mathring{d}^*) \subseteq H^{k+1}(\delta)$  and we are done.

We then define additionally the space  $\mathring{H}^k(d; \Omega)$  as the closure of  $C_0^\infty \Lambda^k(\Omega) \subseteq H^k(d; \Omega)$  w.r.t. the  $H^k(d)$ -norm i.e.  $\mathring{H}^k(d; \Omega)$  corresponds to  $k$ -forms in  $H^k(d; \Omega)$  being zero on the boundary. **TBD: There are several different ways to characterize zero boundary conditions in the  $L^2$  setting. We have to choose the one that works best.** Then we define the spaces

$$\begin{aligned} H_0^k(d; \Omega) &:= \{\omega \in H^k(d; \Omega) \mid d\omega = 0\} \\ \mathring{H}_0^k(d; \Omega) &:= \{\omega \in \mathring{H}^k(d; \Omega) \mid d\omega = 0\} \end{aligned}$$

i.e. the spaces of closed forms. We will use the analogous definition for  $H_0^k(\delta; \Omega)$  and  $\mathring{H}_0^k(\delta; \Omega)$  which we call coclosed forms. We then define the spaces of harmonic forms

$$\mathring{H}_0^k(d, \delta; \Omega) := \{\omega \in \mathring{H}^k(d; \Omega) \mid d\omega = 0, \delta\omega = 0\}.$$

With this one can prove the Hodge decomposition ([1, Lemma 1])

$$L_2^k(\Omega) = \overline{d\mathring{H}^{k-1}(d)} \oplus^\perp \mathring{H}_0^k(d, \delta) \oplus^\perp \overline{\delta H^{k+1}(\delta)} \quad (6)$$

and furthermore for the closed and coclosed forms respectively,

$$\mathring{H}_0^k(d) = \overline{d\mathring{H}^{k-1}(d)} \oplus^\perp \mathring{H}_0^k(d, \delta) \quad (7)$$

$$H_0^k(\delta) = \overline{\delta H^{k+1}(\delta)} \oplus^\perp \mathring{H}_0^k(d, \delta). \quad (8)$$

Before we can reformulate the boundary value problem in the language of differential forms we have to introduce some things from simplicial topology. A quick overview over the essential notions that we will use can for example be found in [1, Chapter 2]. A more thorough introduction can be found in [2, Chapter 4]. For the sake of brevity, we will follow the more intuitive approach taken by Arnold.

**Definition 2** (Affine simplex). Let  $x_0, x_1, \dots, x_k \in \mathbb{R}^n$  be affine independent. Then

$$[x_0, x_1, \dots, x_k] := \text{conv} \{x_0, \dots, x_k\}$$

is called an affine  $k$ -simplex.

We will assume all simplices to be affine.

**Definition 3** (Simplicial complex). A *simplicial complex*  $K$  is a collection of affine simplices s.t.

1.  $\sigma \in K \Rightarrow$  any face of  $\sigma$  is in  $K$ ,
2.  $\sigma, \tau \in K \Rightarrow \sigma \cap \tau$  is either empty or a face of both  $\sigma$  and  $\tau$ .

We call  $|K| := \bigcup \{\sigma \mid \sigma \in K\}$  the polyhedron of  $K$ .

For any topological space  $X$  a homeomorphism  $\tau : |K| \rightarrow X$  is called *triangulation* of  $X$ . Let  $\{x_1, x_2, \dots\}$  be the vertices in the simplicial complex  $K$ . We fix an ordering of the vertices for every simplex. That means for any  $k$  every  $k$  simplex  $\sigma$  has a designated representation in the form of

$$\sigma = [x_{i_0}, x_{i_1}, \dots, x_{i_k}].$$

**Definition 4** ( $k$ -chain). Let  $K$  be a simplicial complex. A formal linear combination of  $k$  simplices

$$c = \sum_{\sigma \in K \text{ } k \text{ simplex}} c_\sigma \sigma$$

with  $c_\sigma \in \mathbb{R}$  is called  $k$ -chain. The vector space of all  $k$ -chains is denoted by  $C_k^c$ .

These spaces of  $k$ -chains become now a chain complex by introducing the boundary operator  $\partial$ .

**Definition 5** (Boundary). For any simplex  $[x_{i_0}, x_{i_1}, \dots, x_{i_k}]$  we define the boundary

$$\partial[x_{i_0}, x_{i_1}, \dots, x_{i_k}] := \sum_{j=0}^k (-1)^j [x_{i_0}, x_{i_1}, \dots, \hat{x}_{i_j}, \dots, x_{i_k}]$$

where  $[x_{i_0}, x_{i_1}, \dots, \hat{x}_{i_j}, \dots, x_{i_k}]$  is the simplex without vertex  $x_{i_j}$ . We then extend the definition of the boundary operator linearly to  $k$ -chains  $c = \sum_{\sigma} c_{\sigma} \sigma$  by

$$\partial c := \sum_{\sigma} c_{\sigma} \partial \sigma.$$

.

A crucial property of the chain operator is the following.

**Proposition 2.**  $\partial \circ \partial = 0$ .

*Proof.* This can be proven by direct computation, analogous to [2, Chap.4, Lemma 1.6]  $\square$

We call a  $k$ -chain  $c$  a  $k$ -cycle if  $\partial c = 0$  and we call  $c$  a  $k$ -boundary if there exists a  $(k+1)$ -chain  $d$  s.t.  $c = \partial d$ . Let  $Z_c^k \subseteq C_c^k$  be the subspace of  $k$ -cycles and  $B_c^k \subseteq C_c^k$  the subspace of  $k$ -boundaries. We can now define the homology spaces of our simplicial complex.

**Definition 6** (Chain homology). The homology spaces  $H_c^k$  are the quotient spaces of cycles and boundaries i.e.

$$H_c^k := Z_c^k / B_c^k.$$

The homology spaces are independent of the chosen simplicial complex [**empty citation**].

Next, we define the space of  $k$ -cochains  $C^k(K)$  which is nothing else than the dual space of the space of  $k$ -chains i.e.

$$C^k(K) := \text{Hom}(C_c^k; \mathbb{R})$$

where  $\text{Hom}(X, Y)$  is the space of all vector space homomorphisms (i.e. linear mappings) from  $X$  to  $Y$ . We now introduce an operator between these spaces of cochains.

**Definition 7.** We define the operator  $\delta : C^k K \rightarrow C^{k+1}(K)$  via

$$(\delta f)(c) := f(\partial c).$$

We call a cochain  $f \in C^k(K)$  *closed* if  $\delta f = 0$  and we call  $f$  *exact* if there is a  $g \in C^{k+1}(K)$  s.t.  $f = \delta g$ .

We define the cohomology spaces analogous to homology spaces above.

**Definition 8** (Cohomology of cochains). Denote the space of closed  $k$ -cochains as  $Z^k(K)$  and  $B^k(K)$  the space of exact  $k$ -cochains. We then define the *cohomology spaces*  $H^k(K)$  as

$$H^k(K) := Z^k(K) / B^k(K).$$

We will later show that if we consider certain subspaces of cochains so called *p-summable* cochains that the  $L^p$ -cohomology defined above and the cohomology spaces of these *p-summable* cochains are isomorphic.

In order to show existence and uniqueness of solutions of the magneto-static problem we rely on a result about the isomorphism of a simplicial cohomology space  $H_p^k(K)$  and the  $L^p$ -cohomology space  $H_{p,dR}^k(\bar{\Omega})$ . This result was proven in [3]. In the diploma thesis of Nikolai Nowaczyk [4], which mostly is based on this paper, many additional details can be found. The result will be presented in the next section. It should be noted that even though the results in [3] are proved explicitly for smooth manifolds without boundary the results can be extended to Lipschitz manifolds with boundary (see the proof of Theorem 2 and the remark at the end in [3]). Therefore, we can apply the result to our case.

## 1 Isomorphism of Cohomology

Before we state the theorem about the isomorphism of we will first formulate a crucial assumption for this result to hold.

Because  $\bar{\Omega}$  from our problem is itself a polyhedron we can assume that  $\bar{\Omega}$  and  $|K|$  are equal as subsets of  $\mathbb{R}^n$  and we can simply use the identity as triangulation. However, we will use different metrics on  $|K|$  and  $\bar{\Omega}$ . We use the Euclidian metric on  $\bar{\Omega}$  and we use the standard simplicial metric on  $|K|$  (cf. [3, p.191]). This metric is defined as follows:

Choose some numbering of the vertices  $\{x_1, x_2, \dots\}$  and take  $f : |K| \rightarrow \ell^2$  where  $\ell^2$  is the Hilbert space of real-valued square-summable sequences s.t.  $f(x_i) = e_i$  with  $e_i \in \ell^2$  being the standard unit vectors and  $f$  is affine on every simplex. This mapping is unique.

Then we define the metric on  $|K|$  as the pullback  $g_S = f^*g$  where  $g$  is the standard metric in  $\ell^2$ . Let  $\langle \cdot, \cdot \rangle$  be the standard scalar product on  $\ell^2$ . Then for  $x \in |K|$  and  $\sum_{i=1}^n v_i \frac{\partial}{\partial x_i}, \sum_{j=1}^n w_j \frac{\partial}{\partial x_j} \in T_x|K|$  we have

$$\begin{aligned} g_S|_x \left( \sum_{i=1}^n v_i \frac{\partial}{\partial x_i}, \sum_{j=1}^n w_j \frac{\partial}{\partial x_j} \right) &= \left\langle \sum_{k=1}^{\infty} \sum_{i=1}^n v_i \frac{\partial f_k}{\partial x_i}(x) \frac{\partial}{\partial y_k}, \sum_{l=1}^{\infty} \sum_{j=1}^n w_j \frac{\partial f_l}{\partial x_j}(x) \frac{\partial}{\partial y_l} \right\rangle \\ &= \sum_{i,j=1}^n \sum_{k,l=1}^{\infty} v_i \frac{\partial f_k}{\partial x_i}(x) w_j \frac{\partial f_l}{\partial x_j}(x) \left\langle \frac{\partial}{\partial y_k}, \frac{\partial}{\partial y_l} \right\rangle \\ &= \sum_{i,j=1}^n \sum_{k=1}^{\infty} v_i \frac{\partial f_k}{\partial x_i}(x) w_j \frac{\partial f_k}{\partial x_j}(x) \\ &= \sum_{i,j=1}^n \sum_{k=1}^{\infty} v_i w_j (Df(x)^T Df(x))_{ij} \\ &= v^T Df(x)^T Df(x) w = \langle Df(x)v, Df(x)w \rangle, \end{aligned}$$

where  $D$  denotes the Jacobian. (TBD: This Jacobian as written here would technically be in  $\mathbb{R}^{\infty \times n}$ . Only finitely many lines are non-zero though, but this is not quite rigorous yet. )

We have two crucial assumptions on the triangulation for the result to hold (cf. [3, p.194]). We summarize them under *GKS-condition* named after the three authors of [3].

**Assumption 1** (GKS-condition). We will assume the following on the simplicial complex  $K$  and the triangulation  $\tau$ :

1. The star of every vertex in  $K$  contains at most  $N$  simplices.
2. For the differential of  $\tau$  we have constants  $C_1, C_2 > 0$  s.t.

$$\|d\tau|_x\| < C_1, \quad \|d\tau^{-1}|_{\tau(x)}\| < C_2,$$

where  $d$  denotes the differential in the sense of differential geometry and the norm is the operator norm w.r.t. the metrics on  $|K|$  and  $\overline{\Omega}$ .

The first assumption is equivalent to every vertex being contained in at most  $N$  simplices, which is fulfilled if we have a shape regular mesh.

Because  $\tau$  is just the identity in our case the second assumption says that for every  $x \in |K|$

$$\sup_{v \neq 0} \frac{\|v\|}{\sqrt{g_S|_x(v, v)}} = \sup_{v \neq 0} \frac{\|v\|}{\|Df(x)v\|} < C_1$$



and analogously

$$\sup_{v \neq 0} \frac{\|Df(x)v\|}{\|v\|} < C_1.$$

## 1.1 Statement of the Isomorphism

From now on we will assume that the GKS condition is fulfilled.

**Definition 9** (Induced map). Let  $V$  and  $W$  be real vector spaces,  $X \subseteq V$ ,  $Y \subseteq W$  be subspaces. For a linear map  $L : V \rightarrow W$  with  $L(X) \subseteq Y$  we define the induced map

$$[L] : V/X \rightarrow W/Y, [v] \mapsto [Lv].$$

It is easy to check that the induced map is well-defined using the definition of quotient space.

The first isomorphism is induced from a linear mapping from the so called *S-forms*  $S_p^k(K)$  to *p-summable k-cochains*  $C_p^k(K)$  which will both be defined next.

**Definition 10.** We define the following norm of a  $k$ -cochain  $f$

$$\|f\|_{C_p^k(K)} := \left( \sum_{c \text{ } k\text{-chain}} |f(c)|^p \right)^{1/p}.$$

and the space of *p-summable k-cochains*

$$C_p^k(K) := \{f \text{ } k\text{-cochain} \mid \|f\|_{C_p^k(K)} < \infty\}.$$

The idea is that for a  $k$ -form  $\theta$  we have the cochain  $I(\theta) \in C^k(K)$ ,  $I(\theta)(\sigma) = \int_\sigma \theta$  for any  $k$ -simplex  $\sigma$ . We want to show that this map induces an isomorphism on cohomology.

At first, we will discuss when and how this integration mapping is well-defined. In order to do so, we will use restriction operators that are introduced in [3, p.191]. Let  $\sigma, \tau \in K$  be simplices s.t.  $\tau < \sigma$  i.e.  $\tau$  is a face of  $\sigma$ . The pullback of the standard inclusion mapping  $j_{\tau, \sigma} : \tau \hookrightarrow \sigma$  gives us a restriction  $j_{\tau, \sigma}^* : W_\infty^k(\sigma) \rightarrow W_\infty^k(\tau)$ , but we have to check whether this is well-defined. In order to do so, take some  $\eta \in W_\infty^k(\sigma)$ . Before restriction we extend  $\eta$  to a neighborhood  $U \supseteq \sigma$  in the affine hull of  $\sigma$  to get  $\tilde{\eta} \in W_\infty^k(U)$ . Then we set  $j_{\tau, \sigma}^* \eta = j_{\tau, U}^* \tilde{\eta}$  which is now well-defined and independent of the extension  $\tilde{\eta}$  chosen. It should be stressed that this construction does not work for  $W_p^k$  with  $p < \infty$ .

Now take  $\eta \in W_{\infty,loc}^k(\overline{\Omega})$ . By applying this restriction operator repeatedly we can now find  $\eta_\sigma \in W_\infty^k(\sigma)$  for every  $\sigma \in K$ . Then we define

$$(I\eta)(\sigma) = \int_\sigma \eta_\sigma \text{ for all } k\text{-simplices } \sigma \in K.$$

So we have established that  $I : W_{\infty,loc}^k(\overline{\Omega}) \rightarrow C^k(K)$  is well-defined.

In the next step, we will use two operators  $\mathcal{R}$  and  $\mathcal{A}$  from [3]. The precise definition and details of their construction are not relevant for our purposes because we will only use the following properties (cf. [3, Thm.2]).

**Theorem 1.** *Assume that the triangulation  $\tau$  fulfills the GKS-condition. Then there exist linear mappings  $\mathcal{R} : L_{1,loc}^k(\overline{\Omega}) \rightarrow L_{1,loc}^k(\overline{\Omega})$ ,  $\mathcal{A} : L_{1,loc}^k(\overline{\Omega}) \rightarrow L_{1,loc}^{k-1}(\overline{\Omega})$  such that*

1.  $\mathcal{R}\omega - \omega = d\mathcal{A}\omega + \mathcal{A}d\omega$  for  $\omega \in W_{1,loc}^k(\overline{\Omega})$
2. for any  $1 \leq p \leq \infty$ ,  $\mathcal{R}(W_p^k(\overline{\Omega})) \subseteq S_p^k(\overline{\Omega})$ .

where

$$S_p^k(\overline{\Omega}) = \{\eta \in W_{\infty,loc}^k \mid I\eta \in C_p^k(K)\}.$$

Now define  $\bar{I} := I \circ \mathcal{R}$ . This gives a well-defined mapping from  $W_p^k(\overline{\Omega})$  to  $C_p^k(K)$ . Following the arguments of [3] and using the properties of  $\mathcal{R}$  and  $\mathcal{A}$  one can show that the induced homomorphism of cohomologies  $[\bar{I}] : H_{p,dR}^k(\overline{\Omega}) \rightarrow H_p^k(K)$  is an isomorphism. This isomorphism will play a crucial part in the proof of uniqueness and existence of solutions in the next section.

## 2 Existence and uniqueness of solutions

### 2.1 Reformulation of the problem

We will return now to the magnetostatic problem. In order to use the results above we will reformulate the problem in the notation of differential forms. There are two ways to identify a vector field with a differential form (cf. [1, Table 6.1 and p.70]) either as a 1-form or a 2-form. For a vector field  $B$  we define

$$\begin{aligned} F^1 B &:= B_1 dx_1 + B_2 dx_2 + B_3 dx_3 \text{ and} \\ F^2 B &:= B_2 dx_2 \wedge dx_3 - B_1 dx_1 \wedge dx_3 + B_3 dx_1 \wedge dx_2 \end{aligned}$$

as the corresponding 1-form and 2-form. Then the exterior derivative is  $dF^2\omega$  corresponds to the divergence, the codifferential  $\delta F^2\omega$  corresponds to the curl and the normal component being zero on the boundary corresponds to  $\omega \in \mathring{H}^2(d)$ .**[empty citation]**.

If we then use the association of 3-forms with scalars we have the corresponding boundary value problem without the integral condition for 2-forms: Find  $\omega \in \mathring{H}^2(d)$  s.t.

$$\delta\omega = 0, \tag{9}$$

$$d\omega = 0 \text{ in } \Omega. \tag{10}$$

Next, we have to add the integral condition. We use that we are in three dimensions so  $** = (-1)^{k(n-k)}\tilde{\nu} = \tilde{\nu}$  [1, p.66] for any  $k$ -form  $\tilde{\nu}$ . and observe

$$\begin{aligned} *F^2 B &= B_1 **dx_1 + B_2 **dx_2 + B_3 **dx_3 = B_1 dx_1 + B_2 dx_2 + B_3 dx_3 \\ &= F^1 B. \end{aligned}$$

Then we have

$$\int_{\gamma} *F^2 B = \int_{\gamma} F^1 B = \int_{\gamma} B \cdot dl.$$

In the last step we used the fact that the integration of a 1-form over a curve is equivalent to the curve integral of the associated vector field (cf. [1, Sec. 6.2.3]). Hence, we can add the integral condition

$$\int_{\gamma} *\omega = C_0. \tag{11}$$

However, we have only  $\omega \in \mathring{H}_0^2(d, \delta)$  so  $*\omega \in H^1(\delta)$  so this integral might not be well defined. In order to deal with this, we will again use the operator  $\mathcal{R}$  and  $\bar{I}$  from Sec. 1.1. Instead of using (11) directly we use the condition

$$\int_{\gamma} \mathcal{R} * \omega = C_0.$$

This is equivalent to  $(\bar{I} * \omega)(\gamma) = C_0$ . We know that this is well-defined.

We want to give some justification about why this is a reasonable extension. For any closed  $\eta \in W_p^1(\bar{\Omega})$  we have  $\mathcal{R}\eta = \eta - d\mathcal{A}\eta$  from Thm. 1. For smooth  $\phi \in \Lambda^0(\bar{\Omega})$  we immediately get from the standard Stoke's theorem that  $\int_c d\phi = 0$  for all closed 1-chains  $c$ . If we assume sufficient regularity on  $\mathcal{R}\eta$ ,  $d\mathcal{A}\eta$  and  $\eta$  then we would have indeed

$$\int_{\gamma} \mathcal{R}\eta = \int_{\gamma} \eta - d\mathcal{A}\eta = \int_{\gamma} \eta.$$

**Remark 1.** This justification can be done more rigorously with the help of S-forms (cf. [3]) which correspond to  $W_{\infty,loc}(\overline{\Omega})$  with an additional decaying condition. Then we obtain that the integral is consistent with the integral on these S-forms.

To summarize we obtain the following problem.

**Problem 1.** Find  $\omega \in \dot{H}^2(d; \Omega)$  s.t.

$$\begin{aligned} d\omega &= 0, \\ \delta\omega &= 0 \text{ in } \Omega, \\ (\bar{I}\omega)(\gamma) &= C_0. \end{aligned}$$

We will examine existence and uniqueness of this problem in the next section.

## 2.2 Existence and uniqueness

We start with the following

**Proposition 3.** *Let  $\gamma$  be a closed  $k$ -chain s.t. the homology class  $[\gamma]$  spans the homology space  $H_c^k$ . Then for any  $C_0 \in \mathbb{R} \setminus \{0\}$  if we have closed cochains  $F, G$  s.t.*

$$F(\gamma) = G(\gamma) = C_0$$

*then  $[F] = [G]$  i.e. their cohomology classes are equal.*

*Proof.* From [1, Sec. 2.5] we know that  $\dim H^k(K) = \dim H_c^k$ . Because  $F$  and  $G$  are closed we therefore have  $\lambda_F, \lambda_G \in \mathbb{R}$  and a cohomology class  $[b]$  s.t.  $[F] = \lambda_F[b]$  and  $[G] = \lambda_G[b]$ . This is equivalent to the existence of  $(k-1)$ -cochains  $J_F, J_G$  s.t.

$$F = \lambda_F b + \delta J_F \text{ and } G = \lambda_G b + \delta J_G.$$

so

$$0 \neq \lambda_F b(\gamma) + \delta J_F(\gamma) = F(\gamma) = G(\gamma) = \lambda_G b(\gamma) + \delta J_G(\gamma).$$

Because  $\gamma$  is closed we have for any  $(k-1)$ -chain  $J$ ,  $\delta J(\gamma) = J(\partial\gamma) = 0$  and so we arrive at  $\lambda_F = \lambda_G$  i.e.  $[F] = [G]$ .  $\square$

We are now able to proof existence of a solution. Take a closed cochain  $F \in C_2^1(K)$  s.t.  $F(\gamma) = C_0$  and  $F(\partial d) = 0$  for any 2-chain  $d$  as in Prop. 3. Then we know from Sec. 1.1 that there exists a unique  $[\theta] \in \mathcal{H}_2^1(K)$  s.t.  $[I]([\theta]) = [F]$ . Let us take the 1-form  $\eta := \varphi^{-1}\theta \in W_2^1(\omega)$ . Then  $\int_\gamma \eta = C_0$  holds. Because we need a 2-form we will take a closer look at  $*\eta$ . We have that the Hodge star operator  $* : L_2^1(\Omega) \rightarrow L_2^2(\Omega)$  is a Hilbert space isometry.

First we compute

$$\delta * \eta = (-1)^{n(k-1)+1} * d * \eta = (-1)^{n(k-1)+1} * d\eta = 0.$$

Using the Hodge decomposition for unbounded domains (cf. [5, Lemma 1]) we get therefore that there exists a sequence  $(\phi_i)_{i \in \mathbb{N}} \subseteq H^3(\delta, \Omega)$  and a harmonic  $\omega \in \dot{H}^2(d) \cap H^3(\delta)$  s.t.

$$*\eta = \lim_{i \rightarrow \infty} \delta \phi_i + \omega.$$

where the limit is in  $L_2^1(\Omega)$ .

We will show that  $\omega$  is then a solution. Because  $\omega$  is harmonic we already know that  $d\omega = 0$ ,  $\delta\omega = 0$  and  $\omega \in \dot{H}^2(d)$ . It remains to show that the integral condition (11) is also satisfied.

At first we check that the integral is well-defined. We have  $\delta\omega = 0$  which implies

$$d * \omega = * * d * \omega = (-1)^{n(k-1)+1} * \delta\omega = 0$$

so  $*\omega$  is a closed 1-form. Therefore the integral is well-defined as shown in Sec. ??.

Take a ball of radius  $R > 0$  around the origin with  $R$  large enough s.t.  $\gamma \subseteq B_R$ . Using the fact that the range of  $\delta$  is closed on bounded domains (cf. [5, Lemma 7]) there is some  $\phi_R \in H^3(\delta; B_R)$  s.t.

$$*\eta|_{B_R} = \omega|_{B_R} + \delta\phi_R.$$

We then get for the integral condition (11)

$$\begin{aligned} \int_\gamma *\omega &= \int_\gamma *\omega|_{B_R} = \int_\gamma *(*\eta|_{B_R} - \delta\phi_R) = \int_\gamma \eta|_{B_R} - (-1)^{3 \cdot 2 + 1} d * \phi_R \\ &= \int_\gamma \eta|_{B_R} = C_0 \end{aligned}$$

so the integral condition (11) is fulfilled and  $\omega$  is indeed a solution.

**TBD: Put all this in a theorem.**

**Theorem 2.** *The solution of the problem is unique.*

*Proof.* Let  $\omega, \tilde{\omega}$  both be solutions. Because  $*\omega$  and  $*\tilde{\omega}$  are closed the cochains  $c \mapsto \int_c \mathcal{R} * \omega$  and  $c \mapsto \int_c \mathcal{R} * \tilde{\omega}$  are closed.

Due to  $\int_\gamma \mathcal{R} * \omega = \int_\gamma \mathcal{R} * \tilde{\omega}$  and the assumption that  $[\gamma]$  spans the homology space we have with Prop. 3  $[I(\mathcal{R} * \omega)] = [I(\mathcal{R} * \tilde{\omega})]$  and because  $[I]$  is an isomorphism  $[\mathcal{R} * \omega] = [\mathcal{R} * \tilde{\omega}]$ . Hence,

$$[*\tilde{\omega}] = [\mathcal{R} * \tilde{\omega}] = [\mathcal{R} * \omega] = [* \omega].$$

That is equivalent to the existence of some 0-form  $\phi \in H^0(d)$  s.t.  $*\omega = *\tilde{\omega} + d\phi$ . We continue by applying the Hodge star operator to both sides and use the definition of the codifferential  $\delta$ :

$$\omega = \tilde{\omega} + *d\phi = \tilde{\omega} + *d * *\phi = \tilde{\omega} + (-1)^{(n-k)(k-1)+1} \delta * \phi.$$

Then because  $\omega$  and  $\tilde{\omega}$  are harmonic we have  $\omega, \tilde{\omega} \perp \delta H^3(\delta)$  and therefore

$$\omega = \tilde{\omega}.$$

□

If we now translate this back to standard vector calculus terms we have found the unique solution of the homogeneous magnetostatic on our domain  $\Omega$ .

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