1 Variational formulation of the magnetostatic problem in 2D

For simplicity we will now turn to the 2D case and we assume that our open domain will be bounded and Lipschitz. This involves to introduce a different Hilbert complex with other differential operators.

We define the scalar curl for $\mathbf{v} \in C^1(\Omega, \mathbb{R}^2)$

$$\operatorname{curl} \mathbf{v} = \partial_1 v_2 - \partial_2 v_1.$$

Additionally, we have the vector-valued curl, denoted in bold, defined for $v \in C^1(\Omega)$

$$\mathbf{curl}\, v = \begin{pmatrix} \partial_2 v \\ -\partial_1 v \end{pmatrix}.$$

The cross product for 2D reads for $\mathbf{a}, \mathbf{b} \in \mathbb{R}^2$,

$$\mathbf{a} \times \mathbf{b} := a_1 b_2 - a_2 b_1.$$

A straightforward calculation shows that the following integration-by-parts formula holds for $u \in C^1(\overline{\Omega})$, $\mathbf{v} \in C^1(\overline{\Omega}; \mathbb{R}^2)$, assuming Ω is Lipschitz

$$\int_\Omega \mathbf{curl}\, u \cdot \mathbf{v}\, dx = \int_\Omega u \operatorname{curl} \mathbf{v}\, dx - \int_{\partial\Omega} u\, \mathbf{n} \times \mathbf{v}\, dl$$

where $\mathbf{n}$ is the outward unit normal. Now, analogous to what we did in **??** we can now extend this definition in the weak sense.

First, notice that $\mathbf{curl}\, u = R_{-\pi/2} \operatorname{grad} u$ and thus $\mathbf{curl}$ is well-defined on H^1. For the scalar curl we define

$$H(\operatorname{curl}; \Omega) = \{\mathbf{v} \in L^2 \mid \exists w \in L^2 : \int_\Omega w\phi\, dx = \int_\Omega \mathbf{v} \cdot \mathbf{curl}\, \phi\, dx \quad \forall \phi \in C_0^\infty\}$$

i.e. $(\operatorname{curl}, H(\operatorname{curl})) = (\mathbf{curl}, C_0^\infty)^*$. Analogous to Section **??**, it is then possible to extend the tangential trace to a operator γ_τ defined on $H(\operatorname{curl})$ s.t. the integration by parts formula **??** holds. Also analogous to the 3D case, we can define

$$H_0(\operatorname{curl}) := \{\mathbf{v} \in H(\operatorname{curl}) \mid \gamma_\tau \mathbf{v} = 0\}$$

and then

$$(\operatorname{curl}, H_0(\operatorname{curl})) = (\mathbf{curl}, H^1)^*$$
$$(\operatorname{curl}, H(\operatorname{curl})) = (\mathbf{curl}, H_0^1)^*.$$

Notice that it hold $\div \mathbf{curl} = 0$ and so we thus have the following 2D Hilbert complex

$$H_0^1 \xrightarrow{\mathbf{curl}} H_0(\mathrm{div}) \xrightarrow{\mathrm{div}} L^2.$$

and the dual complex

$$L^2 \xleftarrow{\mathrm{curl}} H(\mathrm{curl}) \xleftarrow{-\,\mathrm{grad}} H^1$$

We use the notation introduced in **??** for general Hilbert complexes i.e. $V^0 = H_0^1$, $V^1 = H_0(\mathrm{div})$, $V_1^* = H(\mathrm{curl})$ etc. Also we remind of the notation $\mathfrak{B}^k$ for the image of a differential operator and $\mathfrak{B}_k^*$ and analogous $\mathfrak{Z}^k$ for the kernel and $\mathfrak{Z}^k$ for the kernel of the adjoint.

We want to look at the following simplified problem. Assume that our current source $\mathbf{J}$ is pointing in z-direction i.e. $\mathbf{J} = J\mathbf{e}_z$. Further assume that there is a $\tilde{\Omega}$ s.t. $\Omega = \tilde{\Omega} \times \mathbb{R}$. Due to symmetry we can then assume further that B_3 does not change in z-direction which implies that

$$0 = \mathrm{div}\,\mathbf{B} = \partial_x B_1 + \partial_y B_2 = \mathrm{div}\,\tilde{\mathbf{B}}.$$

where $\tilde{\mathbf{B}} = (B_1, B_2)^\top$. The last line of the $\mathrm{curl}\,\mathbf{B} = \mathbf{J}$ reads

$$J = \partial_x B_2 - \partial_y B_1 = \mathrm{curl}\,\tilde{\mathbf{B}}$$

For Ω the unit outer normal has zero part in z-direction and thus the boundary condition

$$0 = \mathbf{B} \cdot \mathbf{n} = \tilde{\mathbf{B}} \cdot \tilde{\mathbf{n}}$$

with $\tilde{\mathbf{n}} = (n_1, n_2)^\top$ is the outer unit normal $\tilde{\Omega}$.

Now we will abuse notation and refer to $\tilde{\mathbf{B}}$ as $\mathbf{B}$ $\tilde{\mathbf{n}}$ as $\mathbf{n}$ and $\tilde{\Omega}$ as $\boldsymbol{\Omega}$. Let $J \in L^2$ be given. Then we see that $\mathbf{B}$ must fulfill the following equations

$$\begin{aligned} \mathrm{curl}\,\mathbf{B} &= J, \\ \mathrm{div}\,\mathbf{B} &= 0. \end{aligned}$$

Depending on the domain, this problem is in general not well-posed – just as the problem in 3D – and requires an additional constraint. Let us now make certain restrictions on what type of domain we will consider.

From now on, we assume that the space of harmonic forms $\mathfrak{H}^1$ has dimension one and that our domain is encompassed by two disjoint closed curves (cf. Fig.**??**) i.e. we have curves $\partial\Omega_{in}$ and $\partial\Omega_{out}$ s.t. $\partial\Omega_{in} \cup \partial\Omega_{out}$ is the boundary of Ω. Let now Γ be s.t. it is a closed curve in Ω that goes around

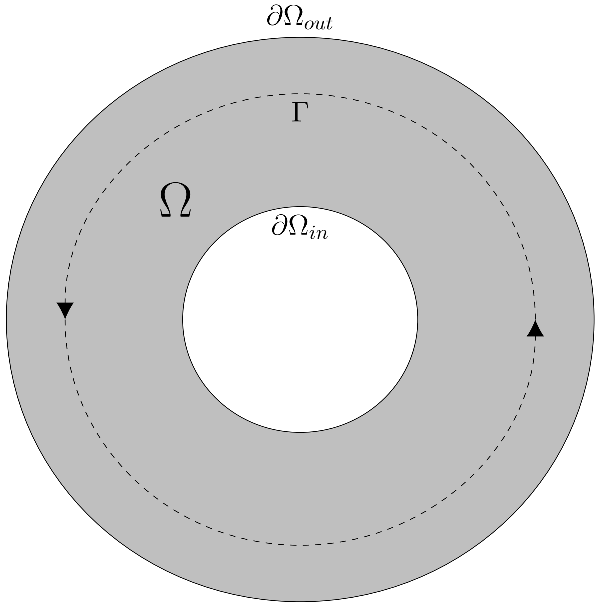

Figure 1: Does this really work? {fig:annulus_dom

the hole in the middle i.e. the area surrounded by Γ contains $\partial\Omega_{in}$. Let Γ be a closed curve with parametrization $\gamma : [0, |\Gamma|]$ s.t. $|\gamma'(t)| = 1$ and assume that γ is bijective i.e. the curve does not intersect itself. We assume that Γ has positive distance from $\partial\Omega_{in}$. We do not assume anything like that for the exterior boundary i.e. Γ can touch or be identical to $\partial\Omega_{out}$ We then denote the area that is enclosed by Γ and $\partial\Omega_{in}$ as Ω_Γ.

From now on, our domain Ω is always assumed to be of that kind. Then we add the curve integral along Γ as an additional constraint, which we assume to be well-defined. So in total, we obtain the following problem.

Problem 1.0.1 (2D magnetostatic problem)**.** Given $J \in L^2$ and $C_0 \in \mathbb{R}$, find $\mathbf{B} \in H_0(\mathrm{div}) \cap H(\mathrm{curl})$ s.t.

$$\begin{aligned} \operatorname{curl} \mathbf{B} &= J, \\ \operatorname{div} \mathbf{B} &= 0, \\ \int_\Gamma \mathbf{B} \cdot dl &= C_0 \end{aligned}$$

The additional constraints are necessary to give a unique solution. We will focus, just as in the first part, on a curve integral as additional constraint. Another option would be an orthogonality constraint as in **??**.

1.1 Mixed formulation

In order to solve this problem numerically using finite elements we have to choose a suitable variational formulation of the problem. Ignoring the curve integral at first, we will use the following. We choose a non-zero harmonic form $\mathbf{p} \in \mathfrak{H}^1$. For any $J \in L^2$ find $\sigma \in H_0^1$, $B \in H_0(\mathrm{div})$ and $\lambda \in \mathbb{R}$ s.t.

$$\langle \sigma, \tau \rangle - \langle u, \mathrm{curl} \tau \rangle = -\langle J, \tau \rangle \quad \forall \tau \in H_0^1, \tag{1.1.1}$$

{eq:first_eq_mix

$$\langle \mathrm{curl} \sigma, \mathbf{v} \rangle + \langle \mathrm{div}\, \mathbf{B}, \mathrm{div}\, \mathbf{v} \rangle + \lambda \langle \mathbf{p}, \mathbf{v} \rangle = 0 \quad \forall \mathbf{v} \in V^k \tag{1.1.2}$$

{eq:second_eq_mi

Here the curve integral condition is missing. It is difficult to include the curve integral condition directly when solving this system numerically. So we will reformulate it below in **??**.

This formulation is of course much more complicated than the first one **??**, but this formulation will turns out to be well-suited for finite element approximations. But it begs the question if the two formulations are equivalent. We will answer it in two propositions. We will treat the harmonic constraint separately.

Proposition 1.1.1. *For any $J \in L^2$, (1.1.1) and (1.1.2) hold i.i.f. $\sigma = 0$, $\lambda = 0$ and $\mathrm{curl}\, \mathbf{B} = J$ and $\mathrm{div}\, \mathbf{B} = 0$ i.e. the magnetostatic* **??** *is fulfilled without the additional constraint.*

Proof. Assume $(\sigma, \mathbf{B}, \mathbf{p})$ is a solution of (1.1.1) and (1.1.2). Then the first equation is

$$\langle \sigma + J, \tau \rangle = \langle \mathbf{B}, \mathrm{curl} \tau \rangle \quad \forall \tau \in H_0^1$$

which is equivalent to $\mathbf{B} \in H(\mathrm{curl})$ and $J + \sigma = \mathrm{curl}\, \mathbf{B}$.

Now assume additionally, that (1.1.2) holds. Then by choosing $\mathbf{v}\mathfrak{p} \in \mathfrak{H}^1$ we get from the definition of the harmonic forms and $\mathfrak{H}^1 \perp \mathrm{curl} H_0^1$ from the Hodge decomposition and thus

$$\langle \mathrm{curl} \sigma, \mathfrak{p} \rangle + \langle du, d\mathfrak{p} \rangle + \lambda \langle p, \mathfrak{p} \rangle = \lambda \langle p, \mathfrak{p} \rangle = 0$$

and so $\lambda = 0$. Then we can choose $\mathbf{v} = \mathrm{curl} \sigma$ to get

$$\langle \mathrm{curl} \sigma, \mathrm{curl} \sigma \rangle + \langle \mathrm{div}\, \mathbf{B}, \mathrm{div}\, \mathrm{curl} \sigma \rangle + \lambda \langle \mathbf{p}, \mathrm{curl} \sigma \rangle = \| \mathrm{curl} \sigma \|^2 .$$

Because $\sigma \in H_0^1$ this gives us $\sigma = 0$. Also we have then $J = \mathrm{curl}\, \mathbf{B}$. At last we choose $\mathbf{v} = \mathbf{B}$ which gives us $\mathrm{div}\, \mathbf{B} = 0$ and thus we proved the first direction.

The other implication is clear i.e. if $\mathbf{B} \in H(\mathrm{curl}) \cap H_0(\mathrm{div})$ with $\mathrm{curl}\, \mathbf{B} = J$ and $\mathrm{div}\, \mathbf{B} = 0$ then the variational formulation clearly holds. $\square$

Notice that the variable λ is not necessary for this variational formulation, but we will need it later, since we will add another equation representing the curve integral constraint and hence we need another variable to have the same number of unknowns and equations. If we now add the same additional constraint to both formulations of the problem then they will remain equivalent.

1.2 Curve integral constraint

We then take ψ s.t. $\psi = 0$ on $\partial\Omega_{in}$ and $\psi = 1$ on Γ and $\psi \equiv 1$ in $\Omega \setminus \Omega_\Gamma$.

Then we know that for $\mathbf{n}$ being the unit outward normal of Ω_Γ we have $\mathbf{n} \perp \gamma'$. If we now take $\mathbf{B}$ that

$$n \times \mathbf{B} = -\mathbf{B} \times n = -(B_1 n_2 - B_2 n_1) = -\mathbf{B} \cdot \begin{pmatrix} n_2 \\ -n_1 \end{pmatrix} = \mathbf{B} \cdot R_{\pi/2}\mathbf{n}$$

then $R_{\pi/2}\mathbf{n}$ is either γ' or $-\gamma'$. Assume w.l.o.g. that $R_{\pi/2}\mathbf{n} = \gamma'$ and thus

$$\mathbf{B} \cdot R_{\pi/2}\mathbf{n} = \mathbf{B} \cdot \gamma'.$$

So we see that we can write

$$n \times \mathbf{B} = \mathbf{B} \cdot \gamma'$$

and so the curve integral becomes

$$\int_\Gamma \mathbf{B} \cdot dl = \int_0^{|\Gamma|} \mathbf{B}(\gamma(t)) \cdot \gamma'(t)\, dt = \int_0^{|\Gamma|} n(\gamma(t)) \times \mathbf{B}(\gamma(t))\, dt = \int_\Gamma n \times \mathbf{B}.$$

Then we observe

$$\int_\Omega \mathrm{curl}\psi \cdot \mathbf{B}\, dx = \int_\Omega \psi\, J\, dx - \int_\partial \Omega n \times \mathbf{B}\, dl = \int_\Omega \psi\, J\, dx - \int_\Gamma \mathbf{B} \cdot dl$$

So we see that

Note that even though right hand side requires some regularity for $\mathbf{B}$ the left hand side makes sense even if $\mathbf{B}$ is only in L^2! So if we are in a situation where the we have curve integral given then we can add this constraint like this without having to compute the curve integral explicitely in our approximation.

Let us assume we are given that the curve integral

$$\int_\Gamma \mathbf{B} \cdot dl = C_0$$

assuming this makes sense. Then we choose ψ and then get the constraint

$$\langle \mathrm{curl}\psi, \mathbf{B} \rangle = \langle J, \psi \rangle - C_0.$$

Note that there are not test functions involved since ψ is fixed. We define $C_1 := \langle J, \psi \rangle - C_0$. Then we have for $\mathbf{B} \in C^1(\Omega; \mathbb{R}^2)$ that

$$\int_\Gamma \mathbf{B} \cdot dl = C_0 \Leftrightarrow \langle \mathrm{curl}\psi, \mathbf{B} \rangle = C_1.$$

This is the motivation why it makes sense to add the right equation as an representation of the curve integral to our system since it is much easier to enforce numerically. Then to get a variational formulation we multiply **??** with an arbitrary $\mu \in \mathbb{R}$. In conclusion, we have the following variational problem: Let $J \in L^2$, $\mathbf{p} \in \mathfrak{H}^1 \setminus \{0\}$. Find $\sigma \in H_0^1$, $\mathbf{B} \in H_0(\mathrm{div})$, $\lambda \in \mathbb{R}$ s.t.

$$\langle \sigma, \tau \rangle - \langle u, \mathrm{curl}\tau \rangle = -\langle J, \tau \rangle \quad \forall \tau \in H_0^1, \tag{1.2.1}$$ {eq:first_eq_mix

$$\langle \mathrm{curl}\sigma, \mathbf{v} \rangle + \langle \mathrm{div}\, \mathbf{B}, \mathrm{div}\, \mathbf{v} \rangle + \langle \lambda \mathbf{p}, \mathbf{v} \rangle = 0 \quad \forall \mathbf{v} \in V^k, \tag{1.2.2}$$ {eq:second_eq_mi

$$\mu \langle \mathbf{curl}\, \psi, \mathbf{B} \rangle = \mu C_1 \quad \forall \mu \in \mathbb{R}. \tag{1.2.3}$$

which gives us the variational formulation of the magnetostatic problem with curve integral constraint. Using the analogous reasoning as in **??** we see that the first two equations are still equivalent to the magnetostatic problem without additional constraint. We will study the well-posedness of this formulation next.

1.3 Well-posedness of the magnetostatic system

The well-posedness is based on the well-known Banach-Nečas-Babuška (BNB) theorem (this formulation is from [**ern˙guermond**] in the real case)

Theorem 1.3.1 (BNB)**.** *Let X be a Banach space and Y be reflexive Banach space. Let $a : X \times Y \to \mathbb{R}$ be a bounded bilinear form and $\ell \in W'$. Then a problem of the form* **??** *is well-posed(?) i.i.f. the following two criteria are fulfilled*

1. $\inf_{x \in X} \sup_{y \in Y} \frac{|a(x,y)|}{\|x\|_X \|y\|_Y} =: \gamma > 0$
2. *if we have $y \in Y$ $a(x, y) = 0$ for every $x \in X$, then $y = 0$.*

Then we also get the stability estimate for a solution x

$$\|x\|_X \leq \frac{1}{\gamma} \|\ell\|_{W'}.$$

Since we are dealing with Hilbert spaces only we can utilize the following proposition to prove it (see [**ern˙guermond**]).

Proposition 1.3.2 (T-coercivity)**.** *Let X and Y be Hilbert spaces. Then* **??** *holds, i.i.f. there exists a bounded bijective operator $T : X \to Y$ s.t.*

$$a(x, T(x)) \geq \eta \|x\|_X^2 \quad \forall x \in X. \tag{1.3.1}$$ {eq:T_coercivity

When we have found such a T then we can choose the γ from **??** as $\eta / \|T\|_{\mathcal{L}(X,Y)}$. Note also that when we have found T s.t. (1.3.1) holds then it must be injective.

Now defining $X := H_0^1 \times H_0(\text{div}) \times \mathbb{R}$ and the bilinear form $a : X \times X \to \mathbb{R}$

$$a(\sigma, \mathbf{B}, \lambda; \tau, \mathbf{v}, \mu) = \langle \sigma, \tau \rangle - \langle u, \text{curl}\tau \rangle + \langle \text{curl}\sigma, \mathbf{v} \rangle + \langle \text{div}\, \mathbf{B}, \text{div}\, \mathbf{v} \rangle + \langle \lambda \mathbf{p}, \mathbf{v} \rangle - \mu \langle \text{curl}\psi, \mathbf{B} \rangle.$$

and

$$\ell(\tau, \mathbf{v}, \mu) = -\langle J, \tau \rangle - \mu C_1.$$

This allows us to rewrite **??** in the standard form: Find $(\sigma, \mathbf{B}, \lambda) \in X$ s.t.

$$a(\sigma, \mathbf{B}, \lambda; \tau, \mathbf{v}, \mu) = \ell(\tau, \mathbf{v}, \mu) \quad \forall (\tau, \mathbf{v}, \mu) \in X.$$

Note that the bilinear form a is not symmetric.

As stated in the proof of the Poincare inequality **??** the $\mathbf{curl}\,|_{\mathfrak{Z}^\perp} : \mathfrak{Z}^\perp \to \mathfrak{B}^1$ is bijective and since it is bounded w.r.t. the V-norm – which is the H^1-norm here – due to Banach inverse theorem it is invertible and we denote this inverse $\mathbf{curl}^{-1}$. This is a slight abuse of notation since it is not really the inverse of the full **curl**.

Usually the last equation **??** is used to enforced the harmonic part of the solution. This implies that we expect that $\mathbf{curl}\,\psi$ must have non-vanishing harmonic part. This is indeed true.

Proposition 1.3.3. *Let $Q^1_{\mathfrak{H}} : L^2 \to \mathfrak{H}^1$ be the orthogonal projection onto the harmonic forms. Then with ψ defined as above we have $Q^1_{\mathfrak{H}}\, \mathbf{curl}\, \psi \neq 0$.*

Proof. Since $\text{div}\, \mathbf{curl}\, \psi = 0$ we know that

$$\mathbf{curl}\, \psi \in \mathfrak{Z}^1 = \mathfrak{B}^1 \overset{\perp}{\oplus} \mathfrak{H}^1$$

using the Hodge decomposition (cf Thm. **??**). Assume for now that $\mathbf{curl}\, \psi \in \mathfrak{B}^1$ i.e. there exists $\psi_0 \in H_0^1$ s.t. $\mathbf{curl}\, \psi_0 = \mathbf{curl}\, \psi$. Since the **curl** is just the rotated gradient we would get that $\text{grad}(\psi - \psi_0) = 0$ and thus $\psi - \psi_0$ is constant almost everywhere. But this is a contradiction since $\text{tr}\, \psi_0$ is zero on $\partial\Omega_{in}$ and $\partial\Omega_{out}$, but $\text{tr}\, \psi = 0$ on $\partial\Omega_{in}$, but $\text{tr}\, \psi = 1$ on $\partial\Omega_{in}$. Thus $\mathbf{curl}\, \psi \notin \mathfrak{B}^1$ and the claim follows. □

Lemma 1.3.4. *Take* $\mathbf{curl}\,\psi = \mathbf{curl}\,\psi_0 + c_\psi \mathbf{p}$. *Assume w.l.o.g.* $c_\psi > 0$ *(otherwise we can choose the representative of* $\mathfrak{H}^1$ *with the opposite sign) and* $\|\mathbf{p}\| = 1$. *Define* $T : X \to X$ *as*

$$T(\sigma, \mathbf{B}, \lambda) = (\sigma - \frac{1}{c_P^2}\,\mathbf{curl}^{-1}, \mathbf{curl}\,\sigma + \mathbf{B} + \lambda\beta\mathbf{p}, \alpha\langle\mathbf{p}, \alpha\langle\mathfrak{p}, \mathbf{B}\rangle + \frac{\lambda}{c_\psi}\rangle).$$

with $\alpha < 0$ *and* $\beta > 0$. *Then* T *is surjective.*

Proof. Take $(\tau, \mathbf{v}, \mu) \in X$ arbitrary. Then Now we choose $\sigma = (1+1/c_P)^{-1}(\tau + (1/c_P^2)\,\mathbf{curl}^{-1}\,\mathbf{v}_{\mathfrak{B}})$ and $\mathbf{B}_{\mathfrak{B}} = \mathbf{v}_{\mathfrak{B}} - \mathbf{curl}\,\sigma$. So

$$\sigma - 1/c_P^2\,\mathbf{curl}^{-1}\,\mathbf{B}_{\mathfrak{B}} = \sigma - 1/c_P^2(\mathbf{curl}^{-1}\,\mathbf{v}_{\mathfrak{B}} - \sigma) = (1 + 1/c_P^2)\sigma - 1/c_P^2\,\mathbf{curl}^{-1}\,\mathbf{v}_{\mathfrak{B}} = \tau.$$

We simply choose $\mathbf{B}_{\mathfrak{B}^*} = \mathbf{v}_{\mathfrak{B}^*}$. For the harmonic part we observe for $\mathbf{v}_{\mathfrak{H}} = c_v p$
Let us look at the system

$$\begin{pmatrix} 1 & \beta \\ \alpha & 1/c_\psi \end{pmatrix} \begin{pmatrix} \kappa_u \\ \lambda \end{pmatrix} = \begin{pmatrix} c_v \\ \mu \end{pmatrix}$$

Now since $c_\psi > 0$ and $\alpha < 0$, $\beta > 0$ we get $1/c_\psi - \alpha\beta \neq 0$ and the system has a solution. Choose $\mathbf{B}_{\mathfrak{H}} = \kappa_u \mathbf{p}$. Then we see

$$\mathbf{v}_{\mathfrak{H}} = c_v p = p(\kappa_u + \beta\lambda) = \mathbf{B}_{\mathfrak{H}} + \beta\lambda p$$

and

$$\mu = \alpha\kappa_u + 1/c_\psi\lambda = \alpha\kappa_u\|p\|^2 + 1/c_\psi\lambda = \alpha\langle\mathbf{B}, \mathbf{p}\rangle + 1/c_\psi\lambda.$$

By combining all that we arrive at $T(\sigma, \mathbf{B}, \lambda) = (\tau, \mathbf{v}, \mu)$. □

We assume from now on that we have chosen $\mathbf{p}$ in a way s.t. c_ψ – as defined in the previous lemma – is positive and $\mathbf{p}$ has norm one. This comes down to choosing $\mathbf{p}$ with the correct sign and normalizing it. Now we can use the T-coercivity **??** to prove the inf-sup condition and thus well-posedness of our formulation.

Theorem 1.3.5. *a satisfies a inf-sup condition with* γ *depending on the Poincaré constant as well as* ψ.

Proof. We will use T-coercivity to prove it. Choose $(\sigma, \mathbf{B}, \lambda) \in X$ arbitrary and define $\rho := \mathbf{curl}^{-1}\,\mathbf{B}_{\mathfrak{B}}$

$$T(\sigma, \mathbf{B}, \lambda) = (\sigma - \frac{1}{c_P^2}\rho, \mathbf{curl}\,\sigma + \mathbf{B} + \beta\lambda\mathbf{p}, \alpha\langle\mathbf{p}, \alpha\langle\mathfrak{p}, \mathbf{B}\rangle + \frac{\lambda}{c_\psi}\rangle)$$

with $\beta = \frac{3c_1^2 c_P^2}{c_\psi^2}$ and $\alpha = -\frac{c_\psi}{4c_1^2 c_P^2}$. Take $c_1 > 0$ s.t. $\|\mathbf{curl}\,\psi_0\| \leq c_1$ (e.g. one could choose $c_1 = \|\mathbf{curl}\,\psi_0\| + 1$). Then T is surjective **??**. Note

$$\langle B, \mathbf{p}\rangle^2 = \|\mathbf{B}_{\mathfrak{H}}\|^2 \langle \frac{\mathbf{B}_{\mathfrak{H}}}{\|\mathbf{B}_{\mathfrak{H}}\|}, \mathbf{p}\rangle^2 = \|\mathbf{B}_{\mathfrak{H}}\|^2$$

where we used in the last equality that $\frac{\mathbf{B}_{\mathfrak{H}}}{\|\mathbf{B}_{\mathfrak{H}}\|}$ is either $+p$ or $-p$ because $\mathfrak{H}^1$ is assumed to be one-dimensional. We split up $d\psi = d\psi_0 + c_\psi \mathfrak{p}$ to get

$$\begin{aligned}
&a(\sigma, \mathbf{B}, \lambda; T(\sigma, \mathbf{B}, \lambda)) \\
&= \langle \sigma, \sigma - \frac{1}{c_P^2}\rho\rangle - \langle \mathbf{B}, \mathbf{curl}\,\sigma - \frac{1}{c_P^2}\mathbf{curl}\,\rho\rangle + \langle \mathbf{curl}\,\sigma, \mathbf{curl}\,\sigma + \mathbf{B} + \beta\lambda\,\mathrm{div}\,\mathbf{p}\rangle \\
&\quad + \langle \mathrm{div}\,\mathbf{B}, \mathrm{div}\,\mathbf{curl} + \sigma\,\mathrm{div}\,\mathbf{B} + \beta\lambda\,\mathrm{div}\,\mathbf{p}\rangle \\
&\quad + \langle \lambda\mathbf{p}, \mathbf{curl}\,\sigma + \mathbf{B} + \beta\lambda\mathfrak{p}\rangle - (\alpha\langle \mathbf{B}, \mathbf{p}\rangle + \frac{\lambda}{c_\psi})\langle \mathbf{B}, \mathbf{curl}\,\psi\rangle \\
&= \|\sigma\|^2 - \frac{1}{c_P^2}\langle \sigma, \rho\rangle + \frac{1}{c_P^2}\|B_{\mathfrak{B}}\|^2 + \|\mathbf{curl}\,\sigma\|^2 + \|\mathrm{div}\,\mathbf{B}\|^2 + \lambda^2\beta - \alpha c_\psi \|\mathbf{B}_{\mathfrak{H}}\|^2 \\
&\quad - \alpha\langle \mathfrak{p}, \mathbf{B}\rangle\langle \mathbf{B}, \mathbf{curl}\,\psi_0\rangle - \frac{\lambda}{c_\psi}\langle B_{\mathfrak{B}}, \mathbf{curl}\,\psi_0\rangle \\
&... \geq \|\sigma\|^2 - \left(\frac{1}{2}\|\sigma\|^2 + \frac{\|B_{\mathfrak{B}}\|^2}{2c_P^2}\right) + \frac{1}{c_P^2}\|B_{\mathfrak{B}}\|^2 + \|\mathbf{curl}\,\sigma\|^2 + \|\mathrm{div}\,\mathbf{B}\|^2 \\
&\quad + \lambda^2\beta - \alpha c_\psi \|\mathbf{B}_{\mathfrak{H}}\|^2 - \left(\frac{\epsilon_1 \alpha^2 \|\mathbf{B}_{\mathfrak{H}}\|^2}{2} + \frac{\|\mathbf{B}_{\mathfrak{B}}\|^2 \|\mathbf{curl}\,\psi_0\|^2}{2\epsilon_1}\right) - \left(\frac{\lambda^2}{2\epsilon_2 c_\psi^2} + \frac{\epsilon_2 \|\mathbf{B}_{\mathfrak{B}}\|^2 \|\mathbf{curl}\,\psi_0\|^2}{2}\right)
\end{aligned}$$

In the last step we used ϵ-Young combined with Cauchy-Schwarz inequality several times. Choose $\epsilon_1 = 4c_1^2 c_P^2$ to get

$$\begin{aligned}
&\frac{1}{2}\|\sigma\|^2 + \frac{1}{2c_P^2}\|B_{\mathfrak{B}}\|^2 + \|\mathbf{curl}\,\sigma\|^2 + \|\mathrm{div}\,\mathbf{B}\|^2 + \lambda^2\left(\beta - \frac{1}{2\epsilon_2 c_\psi^2}\right) \\
&\quad + \|\mathbf{B}_{\mathfrak{H}}\|^2\left(-\alpha c_\psi - \frac{4c_1^2 c_P^2 \alpha^2}{2}\right) - \|B_{\mathfrak{B}}\|^2 \frac{\|\mathbf{curl}\,\psi_0\|^2}{8c_1^2 c_P^2} - \|B_{\mathfrak{B}}\|^2 \frac{\epsilon_2 \|\mathbf{curl}\,\psi_0\|^2}{2}
\end{aligned}$$

Now choose $\epsilon_2 = \frac{1}{4c_1^2 c_P^2}$ and plug in the definition of α to get bound it from below with

$$\begin{aligned}
&\frac{1}{2}\|\sigma\|^2 + \|B_{\mathfrak{B}}\|^2\left(\frac{1}{2c_P^2} - \frac{1}{8c_P^2} - \frac{\|\mathbf{curl}\,\psi_0\|^2}{8c_1^2 c_P^2}\right) + \|\mathbf{curl}\,\sigma\|^2 + \|\mathrm{div}\,\mathbf{B}\|^2 + \lambda^2\left(\beta - \frac{4c_1^2 c_P^2}{2c_\psi^2}\right) \\
&\quad + \|\mathbf{B}_{\mathfrak{H}}\|^2\left(\frac{c_\psi^2}{4c_1^2 c_P^2} - \frac{c_1^2 c_P^2 c_\psi^2}{8c_1^4 c_P^4}\right)
\end{aligned}$$

and finally by using the Poincaré inequality $\|\mathbf{B}_{\mathfrak{B}^*}\| \leq c_P \|\operatorname{div} \mathbf{B}\|$ and $\beta = \frac{3c_1^2 c_P^2}{c_\psi^2}$

$$\frac{1}{2}\|\sigma\|^2 + \frac{1}{4c_P^2}\|B_{\mathfrak{B}}\|^2 + \|\mathbf{curl}\,\sigma\|^2 + \frac{1}{2c_P^2 B_{\mathfrak{B}^*}} + \frac{1}{2}\|\operatorname{div} \mathbf{B}\|^2 + \frac{c_1^2 c_P^2}{c_\psi^2}\lambda^2 + \frac{c_\psi^2}{8c_1^2 c_P^2}\|\mathbf{B}_{\mathfrak{H}}\|^2.$$

We can define now η by taking the minimum of all these coefficients and thus we proved T-coercivity together with **??**. $\square$

Theorem 1.3.6 (Stability)**.** *The system is stable. For solution* $(\sigma, \mathbf{B}, \mathbf{p}) \in X$ *we get*

$$\|\sigma\|_V + \|\mathbf{B}\|_V + |\lambda| \leq \frac{\|J\| + |C_1|}{\gamma}.$$

Proof. The statement follows immediately from **??** and the fact that

$$|\ell(\tau, \mathbf{v}, \mu)| = |-\langle J, \tau\rangle - C_1\mu| \leq (\|J\| + C_1)\|\tau, \mathbf{v}, \mu\|_X$$

and thus $\|\ell\|_{X'} \leq \|J\| + |C_1|$. $\square$

Note that $1/c_\psi$ terms arise stability constant γ. This is not surprising since the term $\langle \mathbf{curl}\,\psi, \mathbf{B}\rangle$ will not enforce the harmonic part if the harmonic part of $\mathbf{curl}\,\psi$ is zero because $\mathbf{B}_{\mathfrak{H}}$ will disappear from the formulation. This also forces us to choose ψ with this in mind to obtain a stable system.

2 Discrete Hilbert complex

In order to approximate the Hodge Laplacian problem we want to use finite elements. We want to use them in a way that we can rebuild the structure of the Hilbert complex in our discretization. This section follows Sec. 5.2 in Arnold's book [1].

Recall situation

Let us assume the we have finite dimensional subspaces $V_h^k \subseteq V^k$. Then we define completely analogous to the continuous case,

$$\begin{aligned}\mathfrak{Z}_h^k &:= \{v \in V_h^k \mid dv = 0\} = \ker d \cap V_h^k \\ \mathfrak{B}_h^k &:= \{dv \mid v \in V_h^{k-1}\}.\end{aligned}$$

We can now also define the discrete harmonic forms. Now the situation is slightly different however. We will not use the continuous adjoint d_k^* to define it. Instead,

$$\mathfrak{H}_h^k := \{v \in \mathfrak{Z}_h^k \mid v \perp \mathfrak{B}_h^k\} = \mathfrak{Z}_h^k \cap \mathfrak{B}_h^{k,\perp}.$$

Notice that we have $\mathfrak{Z}_h^k \subseteq \mathfrak{Z}^k$ and $\mathfrak{B}_h^k \subseteq \mathfrak{b}^k$, but due to $\mathfrak{B}_h^{k,\perp} \supseteq \mathfrak{B}^{k,\perp}$ we have in general

$$\mathfrak{H}^k = \mathfrak{Z}_h^k \cap \mathfrak{B}_h^{k,\perp} \not\subseteq \mathfrak{Z}^k \cap \mathfrak{B}^{k,\perp} = \mathfrak{H}^k.$$

There are three crucial properties that are necessary for stability and convergence of the method. The first one is the common and reasonable assumption that – as usual in finite element theory – we want that the discrete spaces V_h^k approximate the continuous ones V^j. This can be generally summarized that

$$\lim_{h \to 0} \inf_{v_h \in V_h^j} \|w - v_h\| = 0, \quad \forall w \in V^j.$$

This is usually satisfied if we use established finite elements for a given space e.g. if we take Lagragian FE if $V = H^1$ or Raviart-Thomas if $V = H(\mathrm{div})$ [**<empty citation>**].

The next property is more restrictive. We require that $dV_h^{k-1} \subseteq V_h^k$ and $dV_h^j \subseteq V_h^{j+1}$. This shows that the we cannot simply use arbitrary discrete subspaces independent from one another. We say the spaces have to be compatible [**<empty citation>**]. This property has a very nice consequence. It shows that

$$V_h^{k-1} \xrightarrow{d^{k-1}} V_h^k \xrightarrow{d^k} V_h^{k+1}$$

is itself a Hilbert complex and we can apply the general theory from Sec. **??** directly to it. Let us do that.

Denote the restriction of d^j to V_h^j as d_h^j. Then as a linear map between finite spaces the adjoint – denoted as $d_{j,h}^* : V_h^j \to V_h^{j-1}$ – is everywhere defined. It is important to notice that in contrast to d_h the adjoint $d_j^* h$ is not the restriction of the adjoint the continuous adjoint d_j^*. In general, $V_h \not\subseteq V^*$ and so the continuous adjoint might not be well-defined for a given $v_h \in V_h$.

So we obtain the Hilbert complex

$$V_h^{k-1} \xrightarrow{d^{k-1}} V_h^k \xrightarrow{d^k} V_h^{k+1}$$

and its dual complex

$$V_h^{k-1} \xleftarrow{d_k^*} V_h^k \xleftarrow{d_{k+1}^*} V_h^{k+1}$$

From the general Hilbert complex theory (Thm. **??**) we thus obtain the *discrete Hodge decomposition*

$$V_h^j = \mathfrak{B}_h^j \overset{\perp}{\oplus} \mathfrak{H}_h^j \overset{\perp}{\oplus} \mathfrak{B}_{jh}^*.$$

So we achieved our goal of getting a structure like in the continuous case for our discrete approximation. Especially the question how well the discrete harmonic forms approximate the contiinous one will be looked at more closely.

The third crucial assumption is the existence of *bounded cochain projections* π_h. This is a projection that is a cochain map in the sense of cochain complexes **??** i.e. the following diagram commutes: π_h are either bounded in the V or in the W-norm where W-boundedness implies V boundedness. The cochain projection will play an important role in the stablity of the discrete system.

Let us now answer the question about the difference between discrete and continous harmonic forms. In order to do that we need some way to measure the "difference" between two subspaces.

Definition 2.0.1 (Gap between subspaces)**.** For a Banach space W with subspaces Z_1 and Z_2. Let S_1 and S_2 be the unit spheres in Z_1 and Z_2 respectively i.e. $S_1 = \{z \in Z_1 \mid \|z\|_W = 1\}$. Then we define the gap between these subspaces as

$$\mathrm{gap}(Z_1, Z_2) = \max\{\sup_{z_1 \in S_1} \mathrm{dist} z_1, Z_2, \sup_{z_2 \in S_2} \mathrm{dist} z_2, Z_1\}$$

This definition is from [**kato perturbation theory**] and defines a metric on the set of closed subspaces of W (see **??**Remark p.198]) If W is a Hilbert space – as it is throughout this section – and Z_1 and Z_2 are closed then the $\mathrm{gap}(Z_1, Z_2) = \|P_{Z_1} - P_{Z_2}\|$ i.e. the difference in operator norm of the orthogonal projections onto Z_1 and Z_2. This gives us a measure of distance between spaces which we can now apply to the question of the difference of the difference between discrete and continous harmonic forms.

Proposition 2.0.2 (Gap between harmonic forms)**.** *Assume that the discrete complex* **??** *admits a V-bounded cochain projection π_h. Then*

$$\|(I - P_{\mathfrak{H}_h^k})q\|_V \leq \|(I - \pi_h^k)q\|_V, \forall q \in \mathfrak{H}^k$$
$$\|(I - P_{\mathfrak{H}^k})q_h\|_V \leq \|(I - \pi_h^k)P_{\mathfrak{H}^k}q\|_V, \forall q \in \mathfrak{H}^k, \forall q_h \in \mathfrak{H}_h^k$$

and then

$$\mathrm{gap}(\mathfrak{H}, \mathfrak{H}_h) \leq \sup_{q \in \mathfrak{H}, \|q\|=1} \|(I - \pi_h^k)q\|_V$$

Proof. See [1, Thm. 5.2]. □

Do not forget the continuous poincare inequality This then implies that there is a quasi optimal kind of bound of the gap

Proposition 2.0.3 (Discrete Poincare inequality). *Assume that we have a V-bounded cochain projection π_h for the discrete Hilbert complex **??**. Then*

$$\|v\|_V \leq c_P \|\pi_h\|_V \|dv\|, \quad \forall v \in \mathfrak{Z}_h^{k\perp} \cap V_h$$

*with c_P being the Poincare constant from **??**.*

Proof. This indeed is a direct consequence of the existence of bounded cochain projections. Take $v_h \in \mathfrak{Z}_h^{k\perp} \cap V_h$ arbitrary. Since $d(\mathfrak{Z}^{k,\perp} \cap V^k) = \mathfrak{B} \supseteq \mathfrak{B}_h$ we find $z \in \mathfrak{Z}^{k\perp} \cap V_h$ s.t. $dz = dv$. We can apply now the continuous Poincare inequality **??** to get $\|z\|_V \leq c_P \|dz\|_V = c_P \|dv_h\|_V$. Now we can combine the different assumptions about the discrete Hilbert complex teo get $v_h - \pi_h z \in V_h^k$. Now we can use the fact that π_h is a cochain map and the fact that π_h is a projection:

$$d\pi_h^k z = \pi_h^{k+1} dz = \pi_h^{k+1} dv_h = dv_h$$

For the last equality we used also the fact that we have a discrete complex i.e. $d^k V_h^k \subseteq V_h^{k+1}$. That shows that $d(v_h - \pi_h z) = 0$ i.e. $(v_h - \pi_h z) \in \mathfrak{Z}_h^k$. Because $v_h \in \mathfrak{Z}_h^{k,\perp}$ by assumption we have

$$0 = \langle v, v_h - \pi_h z \rangle = \langle v, v_h - \pi_h z \rangle + \langle dv, d(v_h - \pi_h z) \rangle = \langle v, v_h - \pi_h z \rangle_V$$

so $v_h - \pi_h z$ is V orthogonal to v_h. So

$$\|v_h\|_V^2 = \langle v_h, \pi_h^k z \rangle_V + \langle v_h, v_h - \pi_h^k z \rangle_V = \langle v_h, \pi_h^k z \rangle_V \leq \|\pi_h\|_V \|dv\| \overset{Poincare ineq.}{\leq} c_P \|\pi_h\|_V \|dv\|_V$$

□

So we get the inf sup condition with $c_{P,h} = c_P \|\pi_h\|_V$ instead of c_P and obtain well-posedness.

3 Discretized magnetostatic problem

Let us apply this discretized Hilbert complex to the 2D Hilbert complex **??** to get $V_h^0 \subseteq H_0^1$, $V_h^1 \subseteq H_0(\mathrm{div})$ and $V_h^2 \subseteq L^2$ with

$$V_h^0 \xrightarrow{\mathbf{curl}} V_h^1 \xrightarrow{\mathrm{div}} V_h^2$$

and the dual complex

$$V_h^0 \xleftarrow{\widetilde{\mathrm{curl}}_h} V_h^1 \xleftarrow{\widetilde{-\mathrm{grad}}_h} V_h^2$$

where $\widetilde{\mathrm{curl}}_h$ is the adjoint of $\mathbf{curl}_h$ and corresponds thus to weak form of curl and the same for $\widetilde{\mathrm{grad}}_h$. Analogous to the continuous case we assume that $\dim \mathfrak{H}_h^1 = 1$ which is not unreasonable thanks to **??** for $h > 0$ small enough. The discretized version of the magnetostatic problem then states: Find $\mathbf{B}_h \in V_h^1$ s.t.

$$\widetilde{\mathrm{curl}}_h \mathbf{B} = J \text{ and } \operatorname{div} \mathbf{B} = 0.$$

Note that the divergence is enforced strongly while the curl is only enforced weakly. As explained in **??** we will add the curve integral constraint as in **??**. This gives us the following discrete formulation. Find $\sigma_h \in V_h^0$, $\mathbf{B} \in V_h^1$ and $\lambda \in \mathbb{R}$ s.t.

$$\begin{aligned} \langle \sigma_h, \tau_h \rangle - \langle \mathbf{B}_h, \mathrm{curl} \tau_h \rangle &= -\langle J, \tau_h \rangle \quad \forall \tau_h \in V_h^0, \\ \langle \mathrm{curl} \sigma_h, \mathbf{v}_h \rangle + \langle \operatorname{div} \mathbf{B}_h, \operatorname{div} \mathbf{v}_h \rangle + \langle \lambda \mathbf{p}_h, \mathbf{v}_h \rangle &= 0 \quad \forall \mathbf{v}_h \in V_h^1, \\ \mu \langle \mathbf{curl}\, \psi, \mathbf{B}_h \rangle &= \mu C_1 \quad \forall \mu \in \mathbb{R}. \end{aligned}$$

Here we assume for simplicity that $\mathbf{curl}\, \psi \in V_h^1$. Since we can choose ψ this is not unreasonable.

Note that this trial and test space is indeed conforming, but we choose a discrete harmonic form $\mathbf{p}_h \in \mathfrak{H}_h^1$ so the resulting bilinear forms are different. The stability now follows from the exact same arguments as in **??** so we obtain a inf sup condition with a different constant γ_h that involves $c_{P,h}$ from **??**.

4 Implementation in 2D

4.1 Splines

For the discretization we will use the push forward of tensor product splines on a rectangular reference domain $\hat{\Omega}$. We use geometric degrees of freedom which we will introduce below **??**. This section is a recollection of [**broken FEEC framework on mapped multipatch**] since we use the same method as presented in this paper also to fix notation.

We will use two different types knot sequences, non-periodic and periodic ones. We choose a knot sequence $\{\xi_i\}_{i=0}^{n+p}$ with $\xi_0 \leq \xi_1 \leq ... \xi_{n+p}$. We choose two types of sequences.

Let us define $x_0 < x_1 < ... < x_N$ as the physical knots which is our actual grid. We will stick to the equidistant case. Let h be $x_{i+1} - x_i$.

For the non-periodic case we choose an *open* knot sequence by $\xi_0 = ... = \xi_p = x_0$, $\xi_{p+l} = x_l$ for $l = 0, 1, ..., N$ and $\xi_n = \xi_{n+1} = ... = \xi_{n+p} = x_N$

$$\xi_0 = ... = \xi_p < \xi_{p+1} < ...\xi_n = \xi_{n+1} = ... = \xi_{n+p}$$

and for the periodic case $\xi_0 = x_0 - ph, \xi_1 = x_0 - (p-1)h, ..., \xi_p = x_0$, $\xi_{p+l} = x_l$ and $\xi_{n+l} = x_N + lh$ for $l = 0, ..., p$. Thus we always have $N + p = n$.

$$0 = \xi_0 = \xi_1 = ... = \xi_p < \xi_{p+1} < \xi_{p+2} < ... < \xi_{n-1} < \xi_n = \xi_{n+1} = ... = \xi_{n+p}$$

with $n = N + p$ where N is the number of cells. Note that all the knot multiplicities in the interior are one and thus our spline space has maximal regularity. We then define $\mathcal{N}_i^q$ be the normalized B-spline [**multipatch paper**]. We then define the spline spce $\mathbb{S}_q = \mathbb{S}_q(\boldsymbol{\xi}) = span \mathcal{N}_i^q \mid i = 0, ..., n-1$. Since we have maximal regularity we get that

$$\{v \in C^{q-1} \mid v|_{\xi_{q+j,q+j+1}} \in \mathbb{P}_q\}.$$

$\mathcal{N}_0^{p-1}$ vanishes.

We now can take tensor product of spline spaces. We use the notation with $\mathbf{q} \in \{p-1, p\}^2$ and we define with $\mathbf{i} \in [N_0] \times [N_1]$

$$\mathcal{N}_{\mathbf{i}}^{\mathbf{q}} \mathcal{N}_{i_1}^{q_1} \mathcal{N}_{i_2}^{q_2}.$$

We write this as $\mathbb{S}_{\mathbf{q}} = \mathbb{S}_{q_1} \otimes \mathbb{S}_{q_2}$ The spline spaces used in the tensor product can also be periodic or only one of them can be periodic. On $\hat{\Omega}$ we obtain the following discrete Hilbert complex

$$\mathbb{S}_{p,p} \xrightarrow{\mathbf{curl}} \mathbb{S}_{p-1,p} \xrightarrow{\mathrm{div}} \mathbb{S}_{p-1,p-1}$$

and we denote $\hat{V}^0 = \mathbb{S}_{p,p}$, $\hat{V}^1 = \mathbb{S}_{p-1,p}$ and $\hat{V}^2 = \mathbb{S}_{p-1,p-1}$.

It is well-known that if we have a function $\mathbf{v}$ that is piecewise smooth then $\mathbf{v} \in H(\mathrm{div})$ i.i.f. the normal trace across element interfaces agrees almost everywhere. Analogously for $\tau \in H^1$ i.i.f. the values agree on the interfaces almost everywhere. Since we always have $p \geq 2$ we thus know that all our tensor splines are at least continuous globally and thus we get $\hat{V}_h^j \subseteq \hat{V}^j$ for $j = 0, 1, 2$ as desired.

4.2 Basis and degrees of freedom

Let us now investigate the degrees of freedom. We will use geometric degrees of freedom i.e. each degree of freedom can be associated with some geometrical element of our domain. We define Greville points by

$$\zeta_i := \frac{\xi_{i+1} + ... + \xi_{i+p}}{p}$$

i.e. the knot averages for $i = 0, ..., n-1$. Then the spline interpolation at these points is well-defined (see [**I think that was in the isogeometric analysis book**]). Note that in the case periodic some Greville points lie outside of the physical grid. But since the function is periodic it can simply be extended periodically and then interpolated at these points.

This gives us the following geometric elements nodes, edges and cells

$$\hat{\mathtt{n}}_{\mathbf{i}} := (\zeta_{i_1}, \zeta_{i_2}) for \mathbf{i} \in \hat{\mathcal{M}}^0$$
$$\hat{\mathtt{e}}_{d,\mathbf{i}} := [\hat{\mathtt{n}}_{\mathbf{i}}, \hat{\mathtt{n}}_{\mathbf{i}+\mathbf{e}_d}] for (d, \mathbf{i}) \in \hat{\mathcal{M}}^1$$
$$\hat{\mathtt{n}}_{\mathbf{i}} := [\hat{\mathtt{e}}_{1,\mathbf{i}}, \hat{\mathtt{e}}_{1,\mathbf{i}}] = [\zeta_{i_1}, \zeta_{i_1+1}] \times [\zeta_{i_2}, \zeta_{i_2+1}] for \mathbf{i} \in \hat{\mathcal{M}}^2$$

with $[\cdot]$ being the convex hull. As before, $\mathbf{e}_d$ for $d = 1, 2$ is the standard basis vector of $\mathbb{R}^2$. The set of multiindices are defined as

$$\hat{\mathcal{M}}^0 := [n-1]^2$$
$$\hat{\mathcal{M}}^1 := \{(d, \mathbf{i} =) \mid \mathbf{i} \in \hat{\mathcal{M}}^0, d \in \{1, 2\}\}$$
$$\hat{\mathcal{M}}^2 := [n-2]^2$$

Does this stuff go through for periodic case?

Now that we have defined the geometric elements we define the corresponding degrees of freedom as

$$\sigma_{\mathbf{i}}^{\hat{0}}(v) := v(\hat{\mathtt{n}}_{\mathbf{i}}) for \mathbf{i} \in \hat{\mathcal{M}}^0$$
$$\hat{\sigma}_{\mathbf{i}}^2(v) := \int_{\hat{\mathtt{c}}_{\mathbf{i}}} v for (\mathbf{i}) \in \hat{\mathcal{M}}^2$$

We $\mathbf{e}_d^{\perp}$ as the rotation of $\mathbf{e}_d$ by π_2 in counter clockwise direction i.e. $\mathbf{e}_1^{\perp} = \mathbf{e}_2$ and $\mathbf{e}_2^{\perp} = -\mathbf{e}_1$. Something about orientation

These degrees of freedome are unisolvent (explanation) i.e. with some ordering $\mu_0, \mu_1, ...$ of the indices of $\mathcal{M}^l$

$$(\sigma_{\mu_0}^l, \sigma_{\mu_1}^l, ..., \sigma_{\mu_{|\mathcal{M}^l|}}^l) : V_h^l \to \mathbb{R}^{|\mathcal{M}^l}$$

is bijective. And we can thus define our basis functions $\hat{\Lambda}_{\mu}^l$, $\mu \in \mathcal{M}^l$ as the basis which is dual to the degrees of freedom in the sense

$$\hat{\sigma}_{\mu}^l(\hat{\Lambda}_{\nu}^l) = \delta_{\mu,\nu} \quad for \mu, \nu \in \mathcal{M}^l.$$

The question is now on what function spaces these degrees of freedom are defined. We note first that the standard choice as described above with $\hat{V}^0 = H_0^1(\hat{\Omega})$ $\hat{V}^1 = H_0(\mathrm{div})$ and $\hat{V}^2 = L^2(\hat{\Omega})$ can not work because the evaluation at point values is not well-defined for H^1 in 2D since they can not be embedded into the continuous functions. Thus, we need to choose function spaces with higher regularity or integrability.

Let us define the spaces

$$W_{1,2}^1(\hat{\Omega}) := \{v \in L^1(\hat{\Omega}) \mid \partial_1 \partial_2 v \in L^1(\hat{\Omega})\}$$
$$W_d^1(\hat{\Omega}) := \{v \in L^1(\hat{\Omega}) \mid \partial_d v \in L^1(\hat{\Omega})\}$$

Why do we set the sequence to the spaces with L^1 sub instead of the intersection?

This now gives us the discrete setting on the reference domain $\hat{\Omega}$ the idea is now to define the basis functions on the physical domain Ω by a push forward of the basis functions on the reference domain. This is now the inverse of the analogous operations to **??** in 2D.

We will stick to the single patch case meaning that there is a diffeomorphism $F : \hat{\Omega} \to \Omega$. Then we define the pullbacks

$$\mathcal{P}_F^0 : v \mapsto \hat{v} := v \circ F$$
$$\mathcal{P}_F^1 : \mathbf{v} \mapsto \hat{\mathbf{v}} := \det DF \, DF^{-1}\mathbf{v}$$
$$\mathcal{P}_F^2 : v \mapsto \hat{v} := (\det DF)v$$

which map functions on the physical domain Ω to functions on the reference domain $\hat{\Omega}$. Then we have the commuting properties

$$\widehat{\mathbf{curl}}\mathcal{P}_F^0 v = \mathcal{P}_F^1 \, \mathbf{curl}\, v$$
$$\widehat{\mathrm{div}}\mathcal{P}_F^1 \mathbf{v} = \mathcal{P}_F^1 \, \mathrm{div}\, v$$

Using the pullbacks we define the pushforwards as $\mathcal{F}^l := (\mathcal{P}_F^l)^{-1}$ and then we get the basis functions on the physical domain

$$\Lambda_\mu^l := \mathcal{F}^l \hat{\Lambda}_\mu^l$$

and then

$$V_h^l := \mathrm{span}\{\Lambda_\mu^l \mid \mu \in \mathcal{M}^l\}$$

are our discrete spaces. Approximation properties???

Using the geometric degrees from **??** we can now construct the corresponding by

$$\sigma_\mu^l := \hat{\sigma}_\mu^l \circ \mathcal{P}_F^l$$

Then we have by construction that $\sigma^l_\mu(\Lambda^l_\nu) = \delta_{\mu,\nu}$.

We can thus define the projection operators

$$\pi^l_h : U^l \to V^l_h, v \mapsto \sum_{\mu \in \mathcal{M}^l} \sigma^l_\mu \Lambda^l_\mu.$$

Then these degrees of freedom commute with the corresponding differential operators as desired from **??**. They also correspond to geometric elements. σ^0 corresponds to point values in the physical domain, σ^1 to the fluxes through the image of edges and σ^2 to the integral over the mapped cells.

This very simple geometric interpretation gives us the ability to enforce the boundary condition directly by setting the corresponding degrees of freedom to zero.

For $V^0_h \subseteq H^1_0$ we want the trace on the boundary to vanish which means that we set the values at the boundary nodes to zero. Thus, for $\mathtt{n}_{\mathbf{i}}$ on the boundary we want $\sigma^0_{\mathbf{i}}(v) = 0$.

For $V^1_h \subseteq H_0(\mathrm{div})$ we want to have the normal trace zero. So when $\mathtt{e}_{d,\mathbf{i}}(\mathbf{v})$ is boundary edge we require This is then achieved when $\sigma^1_{d,\mathbf{i}}(\mathbf{v}) = 0$.

We now define the spaces $\bar{V}^l_h$ which are the correspoding spaces without any boundary conditions. Then we define the projections $P^l_h : \bar{V}^l_h \to \bar{V}^l_h$ which set the boundary degrees of freedom to zero. They have a very simple matrix representation $\mathbb{P}^l_h$ $(\mathbb{P}^l_h)_{\mu,\nu} = 1$ i.i.f. $\mu = \nu$ and μ does not correspond to a geometric element on the boundary. They are easily constructed by taking the identity matrix and setting the diagonal entries to zero that belong to boundary degrees of freedom.

On the lower level these basis functions are not used explicitely. Instead we use B-splines to compute the corresponding mass matrices etc. We will not go to deep into the details of implementation however. More details about the use of B-splines and the connection with the basis Λ^l_μ can be found in [**multipatch˙paper**]