Materiewellen und Formfaktoren

a)

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b)

$$F(\boldsymbol{q}^{2}) = \int_{\mathbb{R}^{3}} e^{\frac{i\boldsymbol{q}\boldsymbol{x}}{\hbar}} f(\boldsymbol{x}) \, d\boldsymbol{x} = \int_{0}^{2\pi} \int_{0}^{\pi} \int_{0}^{\infty} e^{\frac{i\boldsymbol{q}r\cos(\theta)}{\hbar}} f(r)r^{2} \sin(\theta) \, dr \, d\theta \, d\varphi$$

$$= 2\pi \int_{0}^{\infty} f(r)r^{2} \int_{0}^{\pi} e^{\frac{i\boldsymbol{q}r\cos(\theta)}{\hbar}} \sin(\theta) \, d\theta \, dr = 2\pi \int_{0}^{\infty} f(r)r^{2} \int_{-1}^{1} e^{\frac{i\boldsymbol{q}r\boldsymbol{u}}{\hbar}} \, d\boldsymbol{u} \, dr$$

$$= 4\pi \int_{0}^{\infty} f(r)r^{2} \left(e^{\frac{i\boldsymbol{q}r}{\hbar}} - e^{-\frac{i\boldsymbol{q}r}{\hbar}} \right) \frac{\hbar}{2i\boldsymbol{q}r} \, dr = 4\pi \int_{0}^{\infty} \frac{\sin(\frac{\boldsymbol{q}r}{\hbar})}{\frac{\boldsymbol{q}r}{\hbar}} f(r)r^{2} \, dr$$

c)
$$\langle r^2 \rangle = 4\pi \int_0^\infty r^4 f(r) \, dr; \quad \frac{\sin\left(\frac{qr}{\hbar}\right)}{\frac{qr}{\hbar}} \approx 1 + \frac{r^2 q^2}{6\hbar^2}$$

$$F(q^2) \approx 4\pi \int_0^\infty f(r) \left(r^2 + \frac{r^4 q^2}{6\hbar^2}\right) dr = 4\pi \int_0^\infty f(r) r^2 \, dr + \langle r^2 \rangle \frac{6\hbar^2}{q^2} = 1 + \langle r^2 \rangle \frac{6\hbar^2}{q^2}$$

$$\Rightarrow \langle r^2 \rangle = F(q^2) \frac{6\hbar^2}{r^4 q^2} - 1$$

$$d) f(r) = f_0 e^{-ar}$$

$$\int_{\mathbb{R}^3} f(\mathbf{x}) \, \mathrm{d}\mathbf{x} \stackrel{!}{=} 1 \quad \Longleftrightarrow \quad \int_0^\infty f_0 \, \mathrm{e}^{-ar} r^2 \, \mathrm{d}r = \frac{1}{4\pi} \quad \Longleftrightarrow \quad \underline{f_0 = \frac{a^3}{8\pi}}$$

e)
$$F(q^2) = (1 + \alpha^2)^{-2}$$

$$F(\mathbf{q}^2) = 4\pi \int_0^\infty \frac{\sin\left(\frac{qr}{\hbar}\right)}{\frac{qr}{\hbar}} f(r) r^2 dr = \frac{\hbar a^3}{2q} \int_0^\infty \sin\left(\frac{qr}{\hbar}\right) e^{-ar} r dr = \frac{a^4 \hbar^4}{\left(a^2 \hbar^2 + q^2\right)^2}$$
$$\Rightarrow \alpha(q, a) = \frac{q}{a\hbar}$$

Diese Ladungsverteilung beschreibt Protonen.

f)
$$f(r) = \begin{cases} \frac{3}{4\pi R_0^3} & \text{für } 0 \le r \le R_0 \\ 0 & \text{für } R_0 < r \end{cases}$$

$$F(q^2) = 4\pi \int_0^\infty \frac{\sin\left(\frac{qr}{\hbar}\right)}{\frac{qr}{\hbar}} f(r)r^2 dr = 3\frac{\hbar}{qR_0^3} \int_0^{R_0} \sin\left(\frac{qr}{\hbar}\right)r dr =$$

$$= 3\frac{\sin\left(\frac{qR_0}{\hbar}\right) - \frac{qR_0}{\hbar}\cos\left(\frac{qR_0}{\hbar}\right)}{\frac{q^3R_0^3}{\hbar^3}} = 3\frac{\sin(x) - x\cos(x)}{x^3} , \text{ mit } x(q) = \frac{qR_0}{\hbar}$$

Diese Ladungsverteilung wird in der Natur nicht wiedergefunden, man kann schwerere Kerne aber dadurch approximieren.

g)
$$x(0) = 0$$

$$F(0) = \lim_{x \to 0} \left(3 \frac{\sin(x) - x \cos(x)}{x^3} \right) = 3 \lim_{x \to 0} \left(\frac{1}{x^2} \right)$$

Stabilstes Nuklid einer Isobare

a)

b)

c)

Lum	nosität des LHC	
a)		
b)		