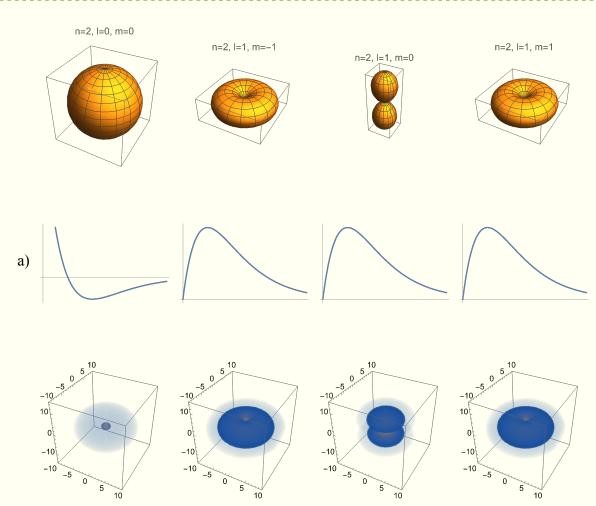
Wasserstoffatom: Aufenthaltswahrscheinlichkeiten



On top you can see the spherical and radial probability of an electron in a hydrogen Atom for different Quantum numbers. Below is a density plot of the square modulus of the whole wave function (as product of the upper 2 plots). If you sum all 4 states up you can kind of recognize a radial symmetry, as the l=0, m=0 state is already symmetric and the other three apparently add up to a sphere.

b)

$$\langle \hat{r} \rangle = \int_{V} r |\psi_{2,1,1}|^{2} dV = \int_{0}^{2\pi} \int_{0}^{\pi} \int_{0}^{\infty} \frac{1}{24a_{0}^{5}} e^{-\frac{r}{a}} r^{2} \frac{3}{8\pi} \sin(\theta)^{2} r r^{2} \sin(\theta) dr d\theta d\phi =$$

$$= \frac{1}{32a_{0}^{5}} \int_{0}^{\infty} r^{5} e^{-\frac{r}{a_{0}}} dr \int_{0}^{\pi} \sin(\theta)^{3} d\theta = \frac{1}{32a_{0}^{5}} 120a_{0}^{6} \frac{4}{3} = \underline{\underline{5a_{0}}}$$

$$\langle \hat{r}^2 \rangle = \int_V r^2 |\psi_{2,1,1}|^2 \, \mathrm{d}V = \frac{1}{32a_0^5} \int_0^\infty r^6 \, \mathrm{e}^{-\frac{r}{a_0}} \, \mathrm{d}r \int_0^\pi \sin(\theta)^3 \, \mathrm{d}\theta = \underline{30a_0^2}$$

$$\langle \Delta \hat{r}^2 \rangle = \langle \hat{r}^2 \rangle - \langle \hat{r} \rangle^2 = 30a_0^2 - 25a_0^2 = \underline{\underline{5a_0^2}}$$

c) d = 1.75 fm

$$P = \int_{V} |\psi_{2,0,0}|^{2} dV = \int_{0}^{2\pi} \int_{0}^{\pi} \int_{0}^{d} \frac{1}{4a_{0}^{3}} e^{-\frac{r}{a}} \left(1 - \frac{r}{a_{0}} + \frac{r^{2}}{4a_{0}^{2}}\right) \frac{1}{4\pi} r^{2} \sin(\theta) dr d\theta d\phi =$$

$$= \frac{1}{2a_{0}^{3}} \int_{0}^{d} e^{-\frac{r}{a}} \left(1 - \frac{r}{a_{0}} + \frac{r^{2}}{4a_{0}^{2}}\right) r^{2} dr = \underline{6.11 \times 10^{-15}}$$

Drehimpulsquantenzahlen

a)
$$\hat{L}_z = -i \hbar \frac{\partial}{\partial \phi}; \quad \psi_{n,l,m} = a_m R_{n,l} \cos(\theta) e^{im\phi} P_l^m$$

$$\langle \hat{L}_z \rangle = \int_V \psi^* \hat{L}_z \psi \, dV = -i \hbar \int_V \psi^* \frac{\partial \psi}{\partial \phi} \, dV = \hbar m \int_V \psi^* \psi \, dV = \underline{\hbar m}$$

b)
$$\hat{\mathbf{L}}^2 Y_l^m(\theta, \phi) = l(l+1)\hbar^2 Y_l^m(\theta, \phi)$$

$$\langle \hat{\mathbf{L}}^2 \rangle = \int_V \psi^* \hat{\mathbf{L}}^2 \psi \, dV = \int_V \psi^* R_{n,l} \hat{\mathbf{L}}^2 Y_l^m \, dV = l(l+1)\hbar^2 \int_V \psi^* \psi \, dV = \underline{l(l+1)\hbar^2}$$

c)

$$\begin{split} \left[\hat{\mathbf{L}}^{2}, \hat{L}_{z}\right] \psi_{n,l,m} &= \hat{\mathbf{L}}^{2} (\hat{L}_{z} \psi_{n,l,m}) - \hat{L}_{z} (\hat{\mathbf{L}}^{2} \psi_{n,l,m}) = \\ &= \hat{\mathbf{L}}^{2} \hbar m \psi_{n,l,m} - \hat{L}_{z} l (l+1) \hbar^{2} \psi_{n,l,m} \end{split}$$

Zeeman-Effekt

a)

b)

c)

d)