

## Polarization of medium

We consider the polarisation density induced linearly by  $\vec{E}$  field.

$\vec{P} := \epsilon_0 \bar{\chi} \vec{E}$ , then  $\bar{\chi}$  is a  $3 \times 3$  Tensor.

$$P_i = \sum_j \chi_{ij} E_j \epsilon_0$$

$$\vec{D} := \vec{P} + \epsilon_0 \vec{E} = \epsilon_0 (\mathbb{I} + \bar{\chi}) \cdot \vec{E} := \epsilon_0 \bar{\epsilon} \cdot \vec{E} \quad \epsilon_{ij} = \delta_{ij} + \chi_{ij}$$

$$D_i = \epsilon_0 \epsilon_{ij} E_j$$

for symmetric  $\epsilon_{ij}$ , we can diagonalise it as.

$$\bar{\epsilon} = \begin{pmatrix} \epsilon_{11} & 0 & 0 \\ 0 & \epsilon_{22} & 0 \\ 0 & 0 & \epsilon_{33} \end{pmatrix} \text{ with a proper choice of coordinates.}$$

Then

$$D_i = \epsilon_0 \epsilon_{ij} E_j$$

The corresponding Maxwell eqns, and wave equation:

$$\textcircled{1} \nabla \cdot \vec{E} = \rho / \epsilon_0 \quad \Rightarrow \quad \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} - \frac{\nabla \cdot \vec{P}}{\epsilon_0} \Rightarrow \nabla \cdot \vec{D} = 0$$

$$\textcircled{2} \nabla \cdot \vec{B} = 0$$

$$\textcircled{3} \nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad \swarrow \text{ignore magnetization}$$

$$\textcircled{4} \nabla \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \Rightarrow \nabla \times \vec{B} = \mu_0 \frac{\partial \vec{D}}{\partial t}$$

use Ansatz  $\vec{E} \Rightarrow \vec{E} e^{-i\omega t + i\vec{k} \cdot \vec{x}}$ , we get from  $\textcircled{3} \textcircled{4}$

$$\begin{cases} i\vec{k} \times \vec{E} = i\omega \vec{B} \\ i\vec{k} \times \vec{B} = -i\omega \mu_0 \epsilon_0 \bar{\epsilon} \cdot \vec{E} \end{cases}$$

$$; \quad c^2 = \frac{1}{\mu_0 \epsilon_0}$$

$$\Rightarrow \vec{k} \times (\vec{k} \times \vec{E}) = - \frac{\omega^2}{c^2} \bar{\epsilon} \cdot \vec{E};$$

$$\Rightarrow \vec{k} (\vec{k} \cdot \vec{E}) - \vec{E} k^2 = - \frac{\omega^2}{c^2} \bar{\epsilon} \cdot \vec{E}; \quad \vec{k} = k \hat{k} \quad ; \quad |\hat{k}| = 1$$

$$\Rightarrow \left( \hat{k} \hat{k} - \mathbb{I} + \frac{\bar{\epsilon}}{n^2} \right) \cdot \vec{E} = 0 \quad \frac{\omega^2}{k^2} \frac{1}{c^2} := n^{-2}$$

given  $\hat{k}$ , the direction of  $\vec{k}$ , we need  $\vec{E} \neq 0$ , so:

$$\det \left| \hat{k}\hat{k} - \mathbb{I} + \frac{\vec{E}}{n^2} \right| = 0 \quad \Leftarrow \text{equation for } n(\hat{k})$$

In general,  $\nabla \cdot \vec{D} = \rho_{\text{free}}$ ; if  $\rho_{\text{free}} = 0$ , we get  $i\vec{k} \cdot \vec{D} = 0 \therefore \vec{k} \perp \vec{D}$

Define  $\eta = \epsilon^{-1}$  where:  $\begin{cases} \vec{D} = \epsilon_0 \vec{E} \cdot \vec{E} \\ \vec{E} = \epsilon_0^{-1} \vec{\eta} \cdot \vec{D} \end{cases}$

choose the coordinates which diagonalize  $\vec{E}$ ; then

$$\vec{E} = \begin{pmatrix} \epsilon_1 & 0 & 0 \\ 0 & \epsilon_2 & 0 \\ 0 & 0 & \epsilon_3 \end{pmatrix} \quad \text{and} \quad \vec{\eta} = \begin{pmatrix} \epsilon_1^{-1} & 0 & 0 \\ 0 & \epsilon_2^{-1} & 0 \\ 0 & 0 & \epsilon_3^{-1} \end{pmatrix} = \begin{pmatrix} \eta_1 & 0 & 0 \\ 0 & \eta_2 & 0 \\ 0 & 0 & \eta_3 \end{pmatrix}$$

$$\hat{k}\hat{k} \cdot \vec{E} - \vec{E} + \frac{\vec{D}}{n^2} = 0 \quad \xrightarrow[\text{* } \epsilon_0]{\text{left dot } \vec{D}} \quad -\vec{D} \cdot \vec{E} \epsilon_0 + \frac{|\vec{D}|^2}{n^2} = 0$$

set  $|\vec{D}|^2 = 1$ ,  $\vec{D} = (D_1, D_2, D_3)$

$$\therefore \vec{D} \cdot \vec{\eta} \cdot \vec{D} = \frac{1}{n^2} \Rightarrow \sum_{i=1}^3 D_i^2 \eta_i = \frac{1}{n^2} \quad \text{or} \quad \sum_{i=1}^3 D_i^2 / \epsilon_i = \frac{1}{n^2}$$

Many people write  $|\vec{D}|=1$  and  $\vec{D} = (x, y, z)$

and:

$$\frac{x^2}{\epsilon_{11}} + \frac{y^2}{\epsilon_{22}} + \frac{z^2}{\epsilon_{33}} = \frac{1}{n^2} \quad \text{or} \quad \frac{x^2}{n_x^2} + \frac{y^2}{n_y^2} + \frac{z^2}{n_z^2} = \frac{1}{n^2}$$

where  $n_i^2 = \epsilon_{ii}$ . This is the "Index Ellipsoid".

With this Ellipsoid, we can calculate  $n$  from the direction of  $\vec{D} = (x, y, z)$  and  $(n_x, n_y, n_z)$  of the material.  
note:  $\vec{D} \perp \vec{k}$

For Example,  $n_x = n_y = n_o$ ;  $n_z = n_e$ .

In this case,  $\vec{z}$  is the symmetry axis. If we apply a static (much slower than optical frequency) field  $\vec{E}$  along  $\vec{z}$ ,  
The leading terms of new  $\tilde{n}_e$  and  $\tilde{n}_o$  are:

$$\begin{cases} \tilde{n}_e = n_e + \frac{1}{2} \chi_{ee} n_e^3 E_z \\ \tilde{n}_o = n_o + \frac{1}{2} \chi_{oe} n_o^3 E_z \end{cases}$$

1st Taylor term

Here we assume no "mixing" induced by  $E_z$ , i.e. no off-diagonal terms in  $\vec{E}$ .

## Further Examples.

Direction-dependened Refractive Index.

$$M := \begin{pmatrix} q_1^2 & q_1 q_2 & q_1 q_3 \\ q_2 q_1 & q_2^2 & q_2 q_3 \\ q_3 q_1 & q_3 q_2 & q_3^2 \end{pmatrix} - \mathbb{I} + \frac{1}{n^2} \begin{pmatrix} \epsilon_{11} & 0 & 0 \\ 0 & \epsilon_{22} & 0 \\ 0 & 0 & \epsilon_{33} \end{pmatrix}; \quad \text{note: } \begin{aligned} \epsilon_{11} &= n_1^2 \\ \epsilon_{22} &= n_2^2 \\ \epsilon_{33} &= n_3^2 \end{aligned}$$

we want,  $\vec{M} \cdot \vec{E} = 0$ , but  $\vec{E} \neq 0$

Consider  $\epsilon_{11} \neq \epsilon_{22} = \epsilon_{33}$

$$\vec{k} \parallel \vec{q} = \begin{pmatrix} \cos \theta \\ \sin \theta \\ 0 \end{pmatrix}$$

$$M = \begin{pmatrix} -\sin^2 \theta & \sin \theta \cos \theta & 0 \\ \sin \theta \cos \theta & -\cos^2 \theta & 0 \\ 0 & 0 & -1 \end{pmatrix} + \frac{1}{n^2} \begin{pmatrix} \epsilon_{11} & & \\ & \epsilon_{22} & \\ & & \epsilon_{22} \end{pmatrix}$$

Case I:  $\vec{E} = \begin{pmatrix} 0 \\ 0 \\ E \end{pmatrix}$   $\begin{aligned} \epsilon_{11} &= n_1^2 \\ \epsilon_{22} &= n_2^2 \end{aligned}$

$$M \cdot \vec{E} = 0 \Rightarrow \frac{\epsilon_{22}}{n^2} = 1 \Rightarrow n = n_2^2$$

Case II:  $\vec{E} = \begin{pmatrix} E_1 \\ E_2 \\ 0 \end{pmatrix}$

$$M \cdot \vec{E} = 0 \Rightarrow \begin{vmatrix} -\sin^2 \theta + \frac{n_1^2}{n^2} & \sin \theta \cos \theta \\ \sin \theta \cos \theta & -\cos^2 \theta + \frac{n_2^2}{n^2} \end{vmatrix} = 0$$

$$\Rightarrow \frac{\sin^2 \theta}{n_1^2} + \frac{\cos^2 \theta}{n_2^2} = \frac{1}{n^2}$$

if  $\theta = 0$   $n^2 = n_2^2$

$\theta = \frac{\pi}{2}$   $n^2 = n_1^2$

