PHYS 325 Final Project: Random Walks

Author: A. J. Ogle Date: May 10th, 2015

PHYS 325 Computational Physics

1 Introduction

The Schrdinger equation describes the motion and energy of quantum particles. It can be used to determine how a system evolves over time, in the case of the time-dependent Schrdinger equation (TDSE). It can also describe the steady state solution for a physical system which does not change over time, in the case of the time-independent Schrdinger equation (TISE).

The TDSE can describe how a particle-wave propagates through space. The TISE can be used for instances in which the wave is stationary, such as a particle trapped in a potential well or box. These "boxes" are simplified representations of actual potential wells, such as the orbitals of an atom. There are instances in which dealing with potentials can be much more complex - so complex, that analytic solutions are difficult, if not impossible, to derive. With this in mind, it is important to develop methods which can numerically solve for the potential field for any geometry in 2 or 3 dimensions. While this paper will not cover the more advanced methods for solving for the wavefunction of a particle in complicated 3 dimensional potential fields, it will provide a method to start from.

2 Background

The starting point will be to grapple with solving for the wavefunction in the context of a 1 dimensional infinite potential square well. This situation produces a square potential well because if applied in the case of 2 dimensions, it produces a square "well" for the particle to be trapped in. To begin, the 1D TISE must be considered:

$$\frac{-\hbar}{2m}\frac{d^2\Psi_x}{dx^2} + V_x\Psi_x = E\Psi_x \tag{1}$$

Since the potential is infinite at the boundaries, it is assumed that there is zero probability of finding the particle outside the potential barrier. It is also assumed that the wavefunction is continuous across the entire dimension. Since the wavefunction is continuous, the limits approaching the points at the boundary (labeled with L, with $L=\pm 1$) must exist. From these assumptions, a solution to the wave equation can be solved for which the waves fulfill the boundary conditions. Analytically, it can be found that:

$$\Psi_x = Acos(\frac{(2n)\pi}{L}x), niseven and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2n-1)\pi}{L}x), nisodd and A = \sqrt{\frac{L}{2}} and \Psi_x = Asin(\frac{(2$$

When solving for stationary solutions numerically, the method of "shooting and matching" will be used. This is a method of assuming an initial value for the energy, then checking whether this gives a solution sufficiently close to the boundary conditions $(Psi_x = 0)$. If a sufficient solution is not achieved, then the simulation "shoots" again, using a guess for the energy that is either higher or lower by some amount, depending on whether the prior guess was too high or too low. This is an algorithm that can be used for finding the solution to some projectile motion, to solve for the projectile hitting a particle value some point away from the projectile's point of launch.

To begin, the initial conditions of $Psi_0 = 1$ are assumed (this is the point right at the middle of the well) for the case of an even energy level (n = 0, 2, 4, 6, ...). This gives a point from which the algorithm can "shoot," using the wave equation and some initial guess for the energy level, to solve for the boundary conditions. In the algorithm, the points of the wave equation can be iterated using:

$$\Psi_{n+1} = 2 * \Psi_n - \Psi_{n-1} - 2(\Delta x)^2 (E - V_n) \Psi_n$$
(3)

The initial conditions for this algorithm are $\Psi_0=1$ and $\frac{d\Psi}{dx}=0$, which implies $\Psi_{n-1}=1$.

An initial guess for E is fed to the algorithm, which then solves for the value of the wavefunction across the space until it reaches the boundary. The solution for the boundary is then compared against the boundary condition. For $\Psi_x > \Psi_{boundary}$, the guess for E is adjusted by some increment -dE, lowering the energy (and hopefully moving Ψ_x closer to $\Psi_{boundary}$. If $\Psi_x < \Psi_{boundary}$, then E is adjusted by some increment dE, moving Ψ_x higher for the next iteration, and hopefully closer to the boundary condition. For each iteration, the value of dE is halved before multiplied by a directional constant, determined by whether Ψ_x is higher or lower than the boundary value.

This process repeats itself until dE has been lowered below some value labeled dE_{min} , after which the algorithm stops and the solution is graphed. This is the method of "shooting and matching."

When graphing the solution, symmetry of the wavefunction is assumed in this case, so the solution to the wavefunction is plotted on one side of the x axis, then this solution is "mirrored" and plotted on the opposite side, presenting the proper waveform expected of the stationary solution for a particle-wave in an infinite square potential well.

3 Results

The aforementioned method of "shooting and matching" was used to determine the values for the first 5 even energy levels (n = 0, 2, 4, 6, 8). The results are listed in

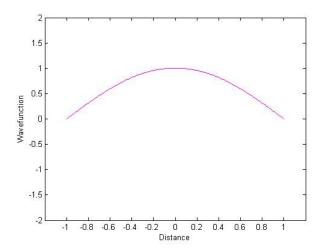


Figure 1: A solution to the wavefunction for initial energy guess of $E_0 = 10$. This solution is the n = 0 state (ground state), with E = 1.2715.

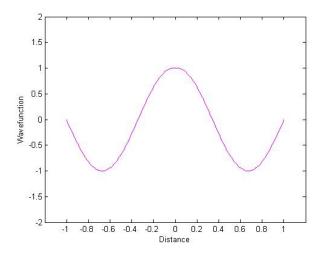


Figure 2: A solution to the wavefunction for initial energy guess of $E_0 = 20$. This solution is the n = 2 state (ground state), with E = 11.4419.

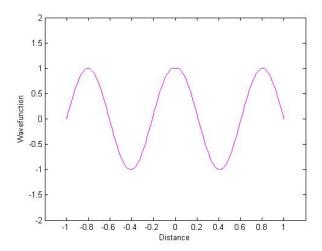


Figure 3: A solution to the wavefunction for initial energy guess of $E_0 = 40$. This solution is the n = 4 state (ground state), with E = 31.7722.

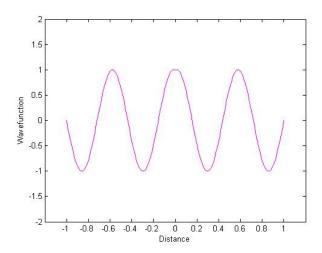


Figure 4: A solution to the wavefunction for initial energy guess of $E_0 = 60$. This solution is the n = 6 state (ground state), with E = 62.2418.

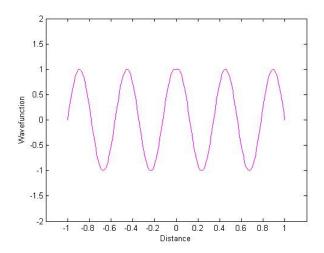


Figure 5: A solution to the wavefunction for initial energy guess of $E_0 = 120$. This solution is the n = 8 state (ground state), with E = 102.8198.

4 Conclusions

References

[1] Giordano, Nicholas J., and Hisao Nakanishi. "10. Quantum" Computational Physics. Upper Saddle River, NJ: Pearson/Prentice Hall, 2006. N. pag. Print.

5 MATLAB code

```
%%Chapter 10.2: One Dimension - Shooting and Matching Methods
clear all:
%Initialize video object
writerObj = VideoWriter('waves_1.avi');
open(writerObj);
prompt = {'Initial Energy Guess:','minimum dE:','speed (<1)'};</pre>
dlg_title = 'Input';
num_lines = 1;
def = {'10','0.00001','0.1'};
input = inputdlg(prompt,dlg_title,num_lines,def)';
V_{max} = 1000;
N = 100; %number of the spatial grids
dx = 0.01; %size of the spatial grids
x = (dx:dx:N*dx);
speed = str2double(input(3)); %the "speed" at which the dE value alters the speed of
%Initial guess for the increment dE
color = 'm';
E = str2double(input(1));
dE = speed*E;
min_dE = str2double(input(2));
%cutoff parameter
b = 1.5;
keep_going = true;
%Wave Amplitude
L = (0.5*length(x)*dx);
A = L^{(-0.5)};
%colors for graphing
c = ['r','g','b','y','m','k'];
"Generate the potential for the "infinite" potential well
V = zeros(1, length(x));
V(1) = V_{max};
V(length(x)) = V_max;
%Settings for an even parity solution
last_diverge = 0;
%Find the solutions for even energy levels
```

```
while (abs(dE) > min_dE && keep_going);
%initial conditions
Psi_x(1) = 1;
Psi_x(2) = 1;
%calculate the wave equation
    for j = 2:N-1
%
      if(mod(n,2) == 0)
%
      k = ((n*pi)/L);
%
      Psi_x(j) = A*sin(k*x(j));
%
      end
%
      if(mod(n,2) == 1)
%
      k = (((2*n-1)*pi)/(2*L));
%
      Psi_x(j) = A*cos(k*x(j));
%
      end
    Psi_x(j+1) = 2*Psi_x(j) - Psi_x(j-1) - 2*(E-V(j))*(dx^2)*Psi_x(j);
        if(abs(Psi_x(j+1)) > b);
            %assume P_x is diverging
            j = (N);
        end
    end
%now we need to mirror the shooting solution for the wave to the other side
%as well
Psi_x_right = Psi_x;
for i = 1:length(Psi_x)
    x_{\text{left}(i)} = -(x(\text{end+1-i}));
    Psi_x_left(i) = Psi_x_right(end+1-i);
end
%Then we put it into a single array
Psi_x_total = [Psi_x_left, Psi_x_right];
x_{total} = [x_{left}, x];
plot(x_total,Psi_x_total,color);
xlabel('Distance')
ylabel('Wavefunction')
axis([-(0.2*N*dx+N*dx) (0.2*N*dx+N*dx) -2 2]);
F(j) = getframe;
writeVideo(writerObj,F(j));
%ask the user if they want to continue the simulation
```

prompt = $\{$ 'Do you like this solution? (1 - ves 0 - no)' $\}$ '.

```
%
      dlg_title = 'Input';
%
      num_lines = 1;
%
      def = {'0'};
%
      prompt_ans = inputdlg(prompt,dlg_title,num_lines,def)';
%
%
      x_input = str2double(prompt_ans);
%
      if(x_iput == 1)
%
          keep_going = false;
%
      end
%
      if(x_iput == 0)
%
          keep_going = true;
%
      end
    if(sign(Psi_x(end))~=sign(last_diverge));
    dE = -0.5*dE;
    end
    E = E + dE;
    last_diverge = sign(Psi_x(end));
    disp(dE);
end
close(writerObj);
disp(E);
```