

**NUCL 510 Nuclear Reactor Theory I**  
**Fall 2011**

**Homework #3**

Due September 15

1. Two mono-energetic neutron beams, each of intensity  $5 \times 10^8$  neutrons/cm<sup>2</sup>-sec, cross at an angle of  $45^\circ$ . Compute the neutron flux and current in the region where the beams intersect.

$$\phi = \int \phi(\Omega) d\Omega = 5 \times 10^8 + 5 \times 10^8 = 10^9 \text{ n/cm}^2 \text{ sec}$$

Taking the direction of one beam as the x-axis, the direction of the other beam can be represented as  $(1/\sqrt{2}, 1/\sqrt{2})$ .

$$\begin{aligned} \vec{J} &= 5 \times 10^8 \vec{e}_x + 5 \times 10^8 (1/\sqrt{2} \vec{e}_x + 1/\sqrt{2} \vec{e}_y) \\ &= 8.536 \times 10^8 \vec{e}_x + 3.536 \times 10^8 \vec{e}_y \text{ n/cm}^2 \text{ sec} \end{aligned}$$

2. The density distribution function for mono-energetic neutrons at a point  $\vec{r}$  is

$$n(\vec{r}, \vec{\Omega}) = n e^{-\lambda/\lambda} (1 + \cos \theta)$$

where  $n$  and  $\lambda$  are constants, and  $\vartheta$  is the angle between  $\vec{\Omega}$  and the z-axis. Compute at  $\vec{r}$  (a) the neutron density, (b) the flux, and (c) the current.

$$\begin{aligned} n(\vec{r}) &= \int_{4\pi} n(\vec{r}, \vec{\Omega}) d\Omega = 2\pi \int_0^\pi n e^{-\lambda/\lambda} (1 + \cos \theta) \sin \theta d\theta \\ &= 2\pi n e^{-\lambda/\lambda} \int_{-1}^1 (1 + \mu) d\mu = 4\pi n e^{-\lambda/\lambda} \end{aligned}$$

$$\phi(\vec{r}) = \int \psi(\vec{r}, \vec{\Omega}) d\Omega = \int n(\vec{r}, \vec{\Omega}) v d\Omega = n(\vec{r}) v = 4\pi n e^{-\lambda/\lambda} v$$

$$\vec{J}(\vec{r}) = \int \psi(\vec{r}, \vec{\Omega}) \vec{\Omega} d\Omega = \vec{J}_x + \vec{J}_y + \vec{J}_z$$

$$\vec{\Omega} = \Omega_x \vec{e}_x + \Omega_y \vec{e}_y + \Omega_z \vec{e}_z = \sin \theta \cos \phi \vec{e}_x + \sin \theta \sin \phi \vec{e}_y + \cos \theta \vec{e}_z$$

$$\vec{J}_x = \int_0^{2\pi} \int_0^\pi n e^{-\lambda/\lambda} v (1 + \cos \theta) \sin \theta \cos \phi (\sin \theta d\theta d\phi) = 0$$

$$\vec{J}_y = \int_0^{2\pi} \int_0^\pi n e^{-\lambda/\lambda} v (1 + \cos \theta) \sin \theta \sin \phi (\sin \theta d\theta d\phi) = 0$$

$$\begin{aligned} \vec{J}_z &= \int_0^{2\pi} \int_0^\pi n e^{-\lambda/\lambda} v (1 + \cos \theta) \cos \theta \sin \theta d\theta d\phi \\ &= 2\pi n e^{-\lambda/\lambda} v \int_{-1}^1 (1 + \mu) \mu d\mu = \frac{4\pi}{3} n e^{-\lambda/\lambda} v \end{aligned}$$

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3. Neutrons are produced uniformly and isotropically throughout a spherical chamber containing a mixture of  $^3\text{H}$  and  $^2\text{H}$  gases at high temperature ( $\sim 10^8$  °K) and low density. (The neutrons originate in the  $^3\text{H}(\text{d},\text{n})^4\text{He}$  and  $^2\text{H}(\text{d},\text{n})^3\text{He}$  fusion reactions.) Show that the neutron flux and current at any point on the surface of the chamber are given by  $\phi = SR/2$  and  $\vec{J} = SR\vec{e}_r/3$ , respectively, where  $S$  is the source density (neutrons/cm<sup>3</sup>-sec),  $R$  is the radius of the chamber, and  $\vec{e}_r$  is a unit radial vector. The neutron mean free path in the medium is essentially infinite.

Since  $\lambda_t = \infty$ , there is no attenuation ( $\Sigma_t = 0$ ).

$$n(\vec{r}, \vec{\omega}) = \frac{S}{4\pi r^2}$$

$$\begin{aligned}\phi &= \int_0^{\pi/2} d\theta \int_0^{2R\cos\theta} dr \int_0^{2\pi} d\phi \frac{S}{4\pi r^2} r^2 \sin\theta \\ &= \frac{S}{2} \int_0^1 d\mu \int_0^{2R\mu} dr = SR \int_0^1 \mu d\mu = \frac{SR}{2}\end{aligned}$$

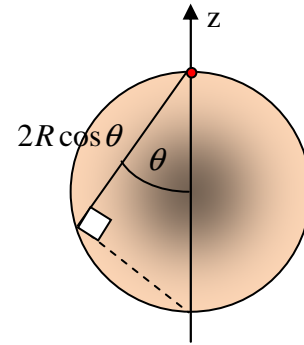
$$\vec{J} = \int_0^{\pi/2} d\theta \int_0^{2R\cos\theta} dr \int_0^{2\pi} d\phi \frac{S}{4\pi r^2} \vec{\Omega} (r^2 \sin\theta)$$

$$\vec{\Omega} = (\sin\theta \cos\phi, \sin\theta \sin\phi, -\cos\theta)$$

$$J_x = J_y = 0 \quad \left( \int_0^{2\pi} \cos\phi d\phi = \int_0^{2\pi} \sin\phi d\phi = 0 \right)$$

$$J_z = - \int_0^{\pi/2} d\theta \int_0^{2R\cos\theta} dr \int_0^{2\pi} d\phi \frac{S}{4\pi} \cos\theta \sin\theta = - \frac{S}{2} \int_0^1 d\mu \int_0^{2R\mu} \mu dr = - \frac{SR}{3}$$

$$\vec{J} = - \frac{SR}{3} \vec{e}_z = \frac{SR}{3} \vec{e}_r$$



4. For an energy-independent s-wave scattering, the energy transfer function is given by

$$\begin{aligned}f_s(E \rightarrow E') &= \frac{\sigma_{s0} \eta^2}{2E} \left\{ \text{erf} \left( \eta \sqrt{\frac{E'}{kT}} - \rho \sqrt{\frac{E}{kT}} \right) \mp \text{erf} \left( \eta \sqrt{\frac{E'}{kT}} + \rho \sqrt{\frac{E}{kT}} \right) \right. \\ &\quad \left. + e^{(E-E')/kT} \left[ \text{erf} \left( \eta \sqrt{\frac{E}{kT}} - \rho \sqrt{\frac{E'}{kT}} \right) \pm \text{erf} \left( \eta \sqrt{\frac{E}{kT}} + \rho \sqrt{\frac{E'}{kT}} \right) \right] \right\}\end{aligned}$$

where

$$\eta = \frac{A+1}{2\sqrt{A}}, \quad \rho = \frac{A-1}{2\sqrt{A}}$$

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and the upper signs are to be used for  $E' > E$  and the lower signs for  $E' < E$ . Plot the energy transfer functions of incident neutrons of 0.025 eV, 10 eV and 100 eV for the following targets: (a) H-1 at 300 °C, (b) O-16 2000 °C and (c) U-238 at 2000 °C.