

NUCL 510 Nuclear Reactor Theory I
Fall 2011

Homework #4

Due September 22

1. A spherical target of radius R is placed in a mono-directional beam of neutrons of intensity I . The area of the beam is larger than the cross sectional area of the sphere. Show that the total reaction rate within the sphere in such a beam is equal to the reaction rate when the sphere is completely immersed in an isotropic flux of magnitude equal to I .

(Hint) For the mono-directional beam, consider the number of neutrons coming into the sphere through an incremental area dA at an angle θ with respect to the beam direction and then the number of interactions made by these neutrons over the traveling path $2R\cos\theta$ within the spherical target. Similarly, for the isotropic flux case, consider the number of neutrons passing through an incremental area dA into an incremental solid angle $d\Omega$ about $\vec{\Omega}$ and then the number of interactions made by these neutrons within the spherical target.

Mono-directional beam

Number of neutrons coming into the sphere through the area dA : $I|\vec{e}_x \cdot \vec{n}|dA$

Number of interactions made by these neutrons over the distance $2R\cos\theta$

$$R_A = \int_0^{2R\cos\theta} \Sigma_t I(x) dx = \int_0^{2R\cos\theta} \Sigma_t I e^{-\Sigma_t x} dx$$

$$= I(1 - e^{-2\Sigma_t R\cos\theta})$$

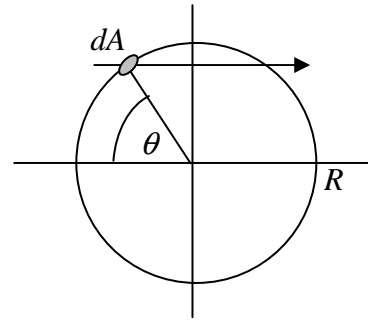
Total reaction rate within the sphere

$$R = \int_A I(1 - e^{-2\Sigma_t R\cos\theta}) \vec{e}_x \cdot \vec{n} dA$$

$$= I \int_{\theta=0}^{\pi/2} \int_{\phi=0}^{2\pi} (1 - e^{-2\Sigma_t R\cos\theta}) \cos\theta R^2 \sin\theta d\phi d\theta$$

$$= 2\pi R^2 I \int_0^1 (1 - e^{-2\Sigma_t R\mu}) \mu d\mu$$

$$= \pi R^2 I \left[1 + \frac{1}{\Sigma_t R} e^{-2\Sigma_t R} + \frac{1}{2(\Sigma_t R)^2} (e^{-2\Sigma_t R} - 1) \right]$$



Isotropic flux

Number of neutrons passing through the area dA into a solid angle $d\Omega$ about $\vec{\Omega}$:

$$\frac{I}{4\pi} |\vec{e}_x \cdot \vec{\Omega}| d\Omega dA$$

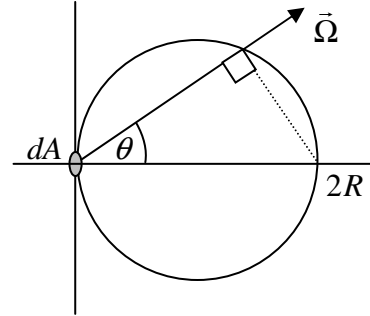
Number of interactions made by these neutrons within the sphere

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$$\begin{aligned}
 R_A &= \frac{I}{4\pi} \int_{\theta=0}^{\pi/2} \int_{\phi=0}^{2\pi} (1 - e^{-2\Sigma_t R \cos \theta}) \cos \theta \sin \theta d\phi d\theta \\
 &= \frac{I}{2} \int_0^1 (1 - e^{-2\Sigma_t R \mu}) \mu d\mu \\
 &= \frac{I}{4} \left[1 + \frac{1}{\Sigma_t R} e^{-2\Sigma_t R} + \frac{1}{2(\Sigma_t R)^2} (e^{-2\Sigma_t R} - 1) \right]
 \end{aligned}$$

Total reaction rate within the sphere

$$\begin{aligned}
 R &= \int_A R_A dA = R_A (4\pi R^2) \\
 &= \pi R^2 I \left[1 + \frac{1}{\Sigma_t R} e^{-2\Sigma_t R} + \frac{1}{2(\Sigma_t R)^2} (e^{-2\Sigma_t R} - 1) \right]
 \end{aligned}$$



2. Let f_k be an eigenfunction of an operator A corresponding to eigenvalue a_k , and let g_l be an eigenfunction of the adjoint operator A^* with eigenvalue b_l . Show that either b_l is the complex conjugate of a_k the eigenvector g_l is orthogonal to f_k .

$$\begin{aligned}
 Af_k &= a_k f_k \Rightarrow (g_l, Af_k) = (g_l, a_k f_k) = a_k (g_l, f_k) \\
 A^* g_l &= b_l g_l \Rightarrow (A^* g_l, f_k) = (g_l, Af_k) = (b_l g_l, f_k) = \bar{b}_l (g_l, f_k) \\
 \Rightarrow (a_k - \bar{b}_l)(g_l, f_k) &= 0 \Rightarrow a_k = \bar{b}_l \text{ or } (g_l, f_k) = \delta_{kl}
 \end{aligned}$$

3. Represent $x^5 - x^3$ in terms of Legendre polynomials $P_n(x)$.

$$\begin{aligned}
 f(x) &= x^5 - x^3 = \sum_{n=0}^5 a_n P_n(x); \quad a_n = \frac{2n+1}{2} \int_{-1}^1 f(x) P_n(x) dx = \frac{2n+1}{2} \int_{-1}^1 (x^5 - x^3) P_n(x) dx \\
 a_0 &= \frac{1}{2} \int_{-1}^1 (x^5 - x^3) dx = 0; \\
 a_1 &= \frac{3}{2} \int_{-1}^1 (x^5 - x^3) x dx = -\frac{6}{35}; \\
 a_2 &= \frac{5}{2} \int_{-1}^1 (x^5 - x^3) \frac{1}{2} (3x^2 - 1) dx = 0; \\
 a_3 &= \frac{7}{2} \int_{-1}^1 (x^5 - x^3) \frac{1}{2} (5x^3 - 3x) dx = \frac{4}{315} \\
 a_4 &= \frac{9}{2} \int_{-1}^1 (x^5 - x^3) \frac{1}{8} (35x^4 - 30x^2 + 3) dx = 0; \\
 a_5 &= \frac{11}{2} \int_{-1}^1 (x^5 - x^3) \frac{1}{8} (63x^5 - 70x^3 - 15x) dx = \frac{16}{693}
 \end{aligned}$$

4. Represent the following functions in terms of spherical harmonics functions $Y_{lk}(\theta, \varphi)$, where

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θ and φ denote the polar and azimuthal angles, respectively. (a) $\sin \theta \cos \varphi$, (b) $\sin \theta \sin \varphi$ and (c) $\cos \theta$.

$$Y_{0,0}(\theta, \varphi) = \frac{1}{2\sqrt{\pi}}; \quad Y_{1,-1}(\theta, \varphi) = \frac{1}{2} e^{-i\varphi} \sqrt{\frac{3}{2\pi}} \sin \theta$$

$$Y_{1,0}(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{3}{\pi}} \cos \theta; \quad Y_{1,1}(\theta, \varphi) = -\frac{1}{2} e^{i\varphi} \sqrt{\frac{3}{2\pi}} \sin \theta$$

$$\Omega_x = \sin \theta \cos \varphi = \sqrt{\frac{2\pi}{3}} [Y_{1,-1}(\theta, \varphi) - Y_{1,1}(\theta, \varphi)]$$

$$\Omega_y = \sin \theta \sin \varphi = -i \sqrt{\frac{2\pi}{3}} [Y_{1,-1}(\theta, \varphi) + Y_{1,1}(\theta, \varphi)]$$

$$\Omega_z = \cos \theta = 2 \sqrt{\frac{\pi}{3}} Y_{1,0}(\theta, \varphi)$$

5. The angular flux for mono-energetic neutrons at a point \vec{r} is given by $\psi(\vec{r}, \vec{\Omega}) = a + b \cos \theta$

where a and b are constants, and θ is the angle between $\vec{\Omega}$ and the z-axis. Compute at \vec{r} (a) the flux, (b) the current, (c) the partial current in the positive z direction, and (d) the partial current in the negative z direction.

$$\phi(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta (a + b \cos \theta) = 2\pi \int_{-1}^1 d\mu (a + b\mu) = 4\pi a$$

$$\vec{J}(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta \vec{\Omega} (a + b \cos \theta)$$

$$J_x(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta \Omega_x (a + b \cos \theta) = \int_0^{2\pi} d\varphi \cos \varphi \int_0^\pi d\theta \sin^2 \theta (a + b \cos \theta) = 0$$

$$J_y(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta \Omega_y (a + b \cos \theta) = \int_0^{2\pi} d\varphi \sin \varphi \int_0^\pi d\theta \sin^2 \theta (a + b \cos \theta) = 0$$

$$J_z(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta \Omega_z (a + b \cos \theta) = \int_0^{2\pi} d\varphi \int_0^\pi d\theta \sin \theta \cos \theta (a + b \cos \theta)$$

$$= 2\pi \int_{-1}^1 d\mu \mu (a + b\mu) = \frac{4\pi}{3} b$$

$$J_z^+(\vec{r}) = \int_0^{2\pi} d\varphi \int_0^{\pi/2} d\theta \sin \theta \Omega_z (a + b \cos \theta) = \int_0^{2\pi} d\varphi \int_0^{\pi/2} d\theta \sin \theta \cos \theta (a + b \cos \theta)$$

$$= 2\pi \int_0^1 d\mu \mu (a + b\mu) = 2\pi \left(\frac{1}{2} a + \frac{1}{3} b \right) = \frac{1}{4} \phi(\vec{r}) + \frac{1}{2} J_z(\vec{r})$$

$$J_z^-(\vec{r}) = \int_0^{2\pi} d\varphi \int_{\pi/2}^\pi d\theta \sin \theta |-\Omega_z| (a + b \cos \theta) = \int_0^{2\pi} d\varphi \int_{\pi/2}^\pi d\theta \sin \theta \cos \theta (a + b \cos \theta)$$

$$= 2\pi \int_0^{-1} d\mu \mu (a + b\mu) = 2\pi \left(\frac{1}{2} a - \frac{1}{3} b \right) = \frac{1}{4} \phi(\vec{r}) - \frac{1}{2} J_z(\vec{r})$$