

# SHEARS

## 1. EQUIDISTRIBUTION OF $H$ -TRANSLATES OF A CLOSED GEODESIC

Let  $\Gamma < G = \mathrm{SL}(2, \mathbb{R})$  be a lattice, say co-compact for now, and assume after conjugation that  $\Gamma$  contains the hyperbolic element  $\mathrm{diag}(e^{\ell/2}, e^{-\ell/2})$ . Then the unit tangent bundle  $T_1(\Gamma \backslash \mathbb{H})$  has a closed geodesic of length  $\ell$  from the point  $(i, \uparrow)$ . Assume for simplicity that  $L^2(\Gamma \backslash G)$  has no exceptional spectrum. Let  $w = k_{\pi/4}$  denote the element which maps  $(i, \uparrow)$  to  $(i, \rightarrow)$ , and  $H = w^{-1}Aw$  be the geodesic flow orthogonal to  $A$ . (Acting on  $(i, \uparrow)$  by  $H$  and  $A$ , we can move to any point in  $\mathbb{H}$  with tangent vector pointing away from the origin, so this is the same as polar coords,  $z = re^{i\theta}$ , with  $r$  determined by  $A$  and  $\theta$  by  $H$ . Here's the equidistribution theorem.

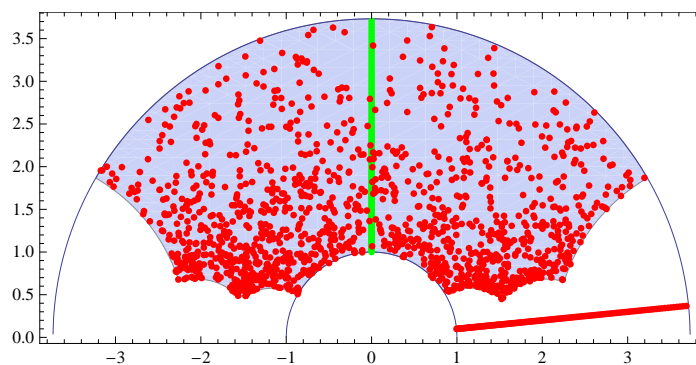
**Theorem 1.** *For a nice test function  $f$  on  $\Gamma \backslash G$ ,*

$$\int_{(A\Gamma) \backslash A} f(ah_t) da = \frac{1}{\mathrm{vol}} \int_{\Gamma \backslash G} f + O_f(e^{(-1/8+\varepsilon)t}),$$

as  $t \rightarrow \infty$ .

The optimal bound here would replace  $-1/8$  by  $-1/2$ , and is what Valentin did in his notes by spectral methods. The point of this note is to do something very soft and easy to get *some* power savings.

Geometrically, this is the following equidistribution: (green is the starting closed geodesic, red is its  $h$ -flow, shown in  $\mathbb{H}$  and  $\Gamma \backslash \mathbb{H}$ ; I can send the Mathematica code if anyone wants to play with it)



Let  $\bar{N}$  be the group  $w^{-1}Nw$ , where  $N = \left\{ \begin{pmatrix} 1 & x \\ 0 & 1 \end{pmatrix} \right\}$ . We have the decomposition

$$G = AH\bar{N}$$

(this is easy to see algebraically or from geometry in  $T^1(\mathbb{H})$ ). For  $\varepsilon > 0$ , let  $\psi = \psi_\varepsilon$  be a function on  $A \backslash G$ , defined in the above coordinates by

$$\psi(ah\bar{n}) = \eta_\varepsilon(h)\phi_\varepsilon(\bar{n}),$$

where  $\eta$  and  $\phi$  are smooth normalized bump functions of a ball of radius  $\varepsilon$  in  $H$  and  $\bar{N}$ , respectively. Normalized means

$$\int_{(\Gamma \cap A) \backslash G} \psi = 1.$$

Average  $\psi$  over  $\Gamma$  to  $\Psi$ , that is, let

$$\Psi(g) := \sum_{\gamma \in (\Gamma \cap A) \backslash \Gamma} \psi(\gamma g).$$

Consider the matrix coefficient:

$$\langle \pi(h_t) \cdot f, \Psi \rangle = \int_{\Gamma \backslash G} f(gh_t) \Psi(g) dg = \int_{(\Gamma \cap A) \backslash G} f(gh_t) \psi(g) dg,$$

where we unfolded. Parametrize  $G = AH\bar{N}$ :

$$= \int_{(\Gamma \cap A) \backslash A} \int_H \int_{\bar{N}} f(ah\bar{n}h_t) \psi(ah\bar{n}) d\bar{n} dh da = \int_{(\Gamma \cap A) \backslash A} \int_H \int_{\bar{N}} f(ah\bar{n}h_t) \eta_\varepsilon(h) \phi_\varepsilon(\bar{n}) d\bar{n} dh da.$$

Now for the “wavefront lemma”: note that if  $h_t = w^{-1}a_t w$  and  $\bar{n}_x = w^{-1}n_x w$ , then

$$h_t^{-1} \bar{n}_x h_t = w_{-1} a_t^{-1} n_x a_t w = \bar{n}_{x e^{-t}}.$$

So if  $\bar{n} \in \text{supp } \phi_\varepsilon$ , then so is its  $h_t$  conjugate. (If we wanted the limit the other way,  $t \rightarrow -\infty$ , then just choose the lower-triangular  $N$ .) Assume  $f$  is uniformly Lipschitz, so

$$f(ah\bar{n}h_t) = f(ah_t h \bar{n}') = f(ah_t) + O_f(\varepsilon),$$

where  $\bar{n}'$  is the conjugate. So we have

$$\langle \pi(h_t) \cdot f, \Psi \rangle = \int_{(\Gamma \cap A) \backslash A} f(ah_t) da \cdot \int_H \int_{\bar{N}} \eta_\varepsilon(h) \phi_\varepsilon(\bar{n}) d\bar{n} dh + O(\varepsilon) = \int_{(\Gamma \cap A) \backslash A} f(ah_t) da + O(\varepsilon).$$

On the other hand, expand spectrally and apply Howe-Moore:

$$\langle \pi(h_t) \cdot f, \Psi \rangle = \frac{1}{\text{vol}} \langle f, \mathbf{1} \rangle \langle \mathbf{1}, \Psi \rangle + O(e^{-t/2} \mathcal{S} f \mathcal{S} \Psi),$$

where  $\mathcal{S}$ 's are Sobolev 1-norms. (I'm ignoring an epsilon here; in general the error term is  $e^{-(1/2-\theta)t}$ , where  $\theta$  controls the exceptional spectrum,  $\theta = 0$  being Ramanujan, which we assumed.) By construction of  $\psi$ , we have  $\mathcal{S} \Psi \ll \varepsilon^{-3}$  (I did this calculation quickly, so check if you agree).

Putting it together, we have

$$\int_{(\Gamma \cap A) \backslash A} f(ah_t) da = \frac{1}{\text{vol}} \int_{\Gamma \backslash G} f + O(e^{-t/2} \mathcal{S} f \varepsilon^{-3} + \varepsilon),$$

and optimizing  $\varepsilon$  gives the claim.