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Modules

The following modules refer to the algebra

$$\mathcal{L}_n = \{h_1, \dots, h_n\},$$

that fulfils the commutation rule

$$[h_i, h_j] = i \hbar \sum_{k=1}^n c_{i,j,k} h_k,$$

characterized by the structure constants $c_{i,k,j}$.

LieGetMa

LieGetMa[**c**, **J**, **vα**] generates the M transformations of the for $M_k = e^{-Q_k}$ for $k = 1, \dots,$

n and $M_{n+1} = M_1 \dots M_n$ is an $n \times n$ matrix and M is an $n \times n \times n$ tensor corresponding to the structure constants c .

- **c** : $n \times n \times n$ tensor containing the structure constants,
- **J** : $n \times n$ matrix, $h' = Jh$ is a new representation of the \mathcal{L}_n ,
- **vα**: dimension n list $v\alpha = \{\alpha_1, \alpha_2, \dots, \alpha_n\}$ containing the transformation parameters for the U_A transformation.

```

LieGetMa[c_, J_, va_] := Module[{dim, k, Q, M},
  dim = Dimensions[c][[1]];
  k = Dimensions[J][[1]];
  M = Table[0, {k1, 1, k + 1}, {k2, 1, dim},
    {k3, 1, dim}];
  M[[k + 1]] = IdentityMatrix[k];
  Do[

    Q =
      Table[Sum[c[[k2, k3, k4]] J[[k1, k2]] va[[k1]],
        {k2, 1, dim}], {k3, 1, dim}, {k4, 1, dim}];
    M[[k1]] = MatrixExp[-Q];
    M[[k + 1]] = M[[k + 1]].M[[k1]];
    , {k1, 1, k}];
  M
]

```

LieGetNu

LieGetNu[M,J] generates the v matrix where v is an $n \times n$ matrix.

- **M** : $n \times n \times n$ tensor containing the transformation matrices,
- **J** : $n \times n$ matrix, $h' = Jh$ is a new representation of the \mathcal{L}_n .

```

LieGetNu[M_, J_] := Module[{dim, Mk, Ik, vt, vt1},
  dim = Dimensions[M[[1]]][[1]];
  Ik = Normal[SparseArray[{{1, 1} → 1}, dim]];
  vt1 = Ik.J;
  Do[
    Mk = M[[k1]];
    Ik = Normal[SparseArray[{{k1, k1} → 1}, dim]];
    vt = vt1.Mk + Ik.J;
    vt1 = vt;
    , {k1, 2, dim}];
  Transpose[vt]
]

```

LieTrans

LieTrans[**M**, **J**, **va**, **va'**, **k**, **t**] transforms the coefficients h into va' under the k 'th transformation U_k corresponding to the M_k matrix. Under this transformation the original Floquet operator $H - p_t = va^T h - p_t$ is transformed into $H' - p_t = U_k (H - p_t) U_k = va^T M_k h - \alpha_k h_k - p_t = (va')^T h - p_t$ where va' is the new set of coefficients.

- **M** : $n \times n \times n$ tensor containing the transformation matrices,
- **J** : $n \times n$ matrix, $h' = Jh$ is a new representation of the \mathcal{L}_n ,
- **va** : dimension n list $va = \{a_1, a_2, \dots, a_n\}$ containing the coefficients of the original Floquet operator,

- **$\mathbf{v}\alpha$** : dimension n list $\mathbf{v}\alpha = \{\alpha_1, \alpha_2, \dots, \alpha_n\}$ containing the transformation parameters for the U_A transformation,
- **\mathbf{k}** : integer that tags the number of transformation to be used,
- **\mathbf{t}** : time parameter.

LieTrans $[M_ , J_ , \mathbf{va_} , \mathbf{v}\alpha_ , k_ , t_] :=$
 $\mathbf{va} . M[[k]] - D[\mathbf{v}\alpha[[k]], t] J[[k]]$

LieGetu

LieGetu $[M, J, \mathbf{va}, \mathbf{v}\alpha, \mathbf{t}]$ transforms the original coefficients \mathbf{va} into u under the complete transformation U_A corresponding to $M_a = M_1 M_2 \dots M_n$. Under this transformation the original Floquet operator $H - p_t = \mathbf{va} . h - p_t$ is transformed into $H' - p_t = U_A (H - p_t) U_A = \mathbf{va}^T M_a h - \dot{\alpha}^T v^T h_k - p_t = u^T h - p_t$.

- **\mathbf{M}** : is an $n \times n \times n$ tensor containing the transformation matrices,
- **\mathbf{J}** : $n \times n$ matrix, $h' = Jh$ is a new representation of the \mathcal{L}_n ,
- **\mathbf{va}** : dimension n list $\mathbf{va} = \{a_1, a_2, \dots, a_n\}$ containing the coefficients of the original Floquet operator,
- **$\mathbf{v}\alpha$** : dimension n list containing the transformation parameters for the U_A transformation. The α parameters must be functions of the time parameter t .
- **\mathbf{t}** : time parameter.

```

LieGetu[M_, J_, va_, va_, t_] := Module[{k, vu, vw},
  k = Dimensions[va][[1]];
  vu = va;
  Do[
    vw = LieTrans[M, J, vu, va, k1, t];
    vu = vw;
    , {k1, 1, k}];
  vu
]

```

LieGetDifEqLambda

LieGetDifEqLambda[J,vi,v α ,v β , λ ,ci] calculates a list containing the differential equations with respect to the auxiliary parameter λ that connects the α and β parameters.

- **J** : $n \times n$ matrix, $h' = J.h$ is a new representation of the \mathcal{L}_n ,
- **vi** : inverse of the $n \times n$ matrix v calculated with LieGetNu[M,J],
- **v α** : dimension n list $v\alpha = \{\alpha_1, \alpha_2, \dots, \alpha_n\}$ containing the transformation parameters for the U_A . The α parameters must be functions of the auxiliary parameter λ ,
- **v β** : dimension n list $v\beta = \{\beta_1, \beta_2, \dots, \beta_n\}$ containing the transformation parameters for the U_B . The β parameters are NOT functions of the auxiliary parameter λ ,
- **λ** : is the auxiliary parameter that helps relate the α and β transformation parameters.

- ci** : is a Boolean variable. If $ci == \text{True}$,
the initial conditions $\alpha_1[0] == 0$, $\alpha_2[0] == 0$, ...,
 $\alpha_n[0] == 0$ is appended to the list of differential equations. If on $ci == \text{False}$ then the output is just the list of differential equations,
- **λ** : is the auxiliary parameter that helps relate the α and β transformation parameters.

```
LieGetDifEqLambda[J_, v i_, v α_, v β_, λ_, ci_] :=
Module[{dim, v},
  dim = Dimensions[v α][[1]];
  v = v i.v β;
  If[ci == True,
    Join[Table[D[v α[[k1]], λ] == v[[k1]], {k1, 1, dim}],
      Table[(v α[[k1]] /. λ → 0) == 0, {k1, 1, dim}]],
    Table[D[v α[[k1]], λ] == v[[k1]], {k1, 1, dim}]
  ]
]
```

Main Program

Definition of the structure constants $c_{i,j,k}$

The elements of this algebra are given by the operators $h_1 = 1$, $h_2 = x$, $h_3 = p$, $h_4 = m \omega_0^2 x^2 + p^2$. The following lines define the algebra dimension and the structure constants.

```

n = 4;
d = Table[0, {k1, 1, n}, {k2, 1, n}, {k3, 1, n}];
d[[2, 3, 1]] = 1;
d[[3, 2, 1]] = -1;
d[[4, 2, 3]] = -2;
d[[2, 4, 3]] = 2;
d[[4, 3, 2]] = 2 m2 ω02;
d[[3, 4, 2]] = -2 m2 ω02;

R = IdentityMatrix[n];
Ri = Inverse[R];
c =
  Table[Sum[R[[k1, m1]] R[[k2, m2]] Ri[[m3, k3]]
    d[[m1, m2, m3]], {m1, 1, n}, {m2, 1, n},
    {m3, 1, n}], {k1, 1, n}, {k2, 1, n}, {k3, 1, n}];

```

Since $J=\mathcal{I}$, the structure of the algebra elements is preserved.

```
J = IdentityMatrix[n];
```

Derivation of the time differential equations for $\alpha_i(t)$

We calculate \mathcal{U} using Eq. (13). Notice that in this case $v\alpha = \{\alpha_1(t), \dots, \alpha_n(t)\}$ is a function of time.

```

vα = Table[Subscript[α, k1][t], {k1, 1, n}];
va = Table[Subscript[a, k1], {k1, 1, n}];
Ma = LieGetMa[c, J, vα];
vu = LieGetu[Ma, J, va, vα, t];
MatrixForm[Simplify[ExpToTrig[vu]]]

```

$$\begin{pmatrix} a_1 - a_3 \alpha_2[t] + a_4 \alpha_2[t]^2 + a_2 \alpha_3[t] + m^2 a_4 \omega_0^2 \alpha_3[t]^2 \\ \cos[2 m \omega_0 \alpha_4[t]] (a_2 + 2 m^2 a_4 \omega_0^2 \alpha_3[t] - \alpha_2'[t]) - m \sin[2 m \omega_0 \alpha_4[t]] \\ \frac{\sin[2 m \omega_0 \alpha_4[t]] (a_2 + 2 m^2 a_4 \omega_0^2 \alpha_3[t] - \alpha_2'[t])}{m \omega_0} + \cos[2 m \omega_0 \alpha_4[t]] \\ a_4 - \alpha_4'[t] \end{pmatrix}$$

Compare these results with the ones in Eqs. (82)-(85).

The simplified differential equations for $\alpha_i(t)$ are obtained from Eq. (15)

```

v = LieGetNu[Ma, J];
vi = Inverse[v];
ε = Simplify[vi.vu];
MatrixForm[ε]

```

$$\begin{pmatrix} a_1 - a_3 \alpha_2[t] + a_4 \alpha_2[t]^2 - m^2 a_4 \omega_0^2 \alpha_3[t]^2 - \alpha_1'[t] \\ a_2 + 2 m^2 a_4 \omega_0^2 \alpha_3[t] - \alpha_2'[t] \\ a_3 - 2 a_4 \alpha_2[t] - \alpha_3'[t] \\ a_4 - \alpha_4'[t] \end{pmatrix}$$

Compare these results with the ones in Eqs. (87)-(90).

Relation between $\alpha(t)$ and $\beta(t)$ via the solution of the λ differential equations

Using Eq. (16) we workout the λ differential equations for the $\alpha_i(\lambda, t)$ parameters. Note that in this case $v\alpha = \{\alpha_1(\lambda), \dots, \alpha_n(\lambda)\}$ is a

function of λ therefore, M_a and v have to be recalculated. In **LieGetDifEqLamda** the condition **ci** is set to **True** in order to include the initial conditions.

```

vα = Table[Subscript[α, k1][λ], {k1, 1, n}];
Ma = LieGetMa[c, J, vα];
v = LieGetNu[Ma, J];
vi = Inverse[v];
vβ = Table[Subscript[β, k1], {k1, 1, n}];
difeqsλ =
  Simplify[
    ExpToTrig[LieGetDifEqLambda[J, vi, vα, vβ, λ,
      True]]];
MatrixForm[difeqsλ]

```

$$\left(\begin{array}{l}
 (\cos[2 m \omega_0 \alpha_4[\lambda]] \beta_2 + m \sin[2 m \omega_0 \alpha_4[\lambda]] \beta_3 \omega_0) \alpha_3[\lambda] + \alpha_1'[\lambda] \\
 \cos[2 m \omega_0 \alpha_4[\lambda]] \beta_2 + m \sin[2 m \omega_0 \alpha_4[\lambda]] \beta_3 \omega_0 = \alpha_2'[\lambda] \\
 \cos[2 m \omega_0 \alpha_4[\lambda]] \beta_3 = \frac{\sin[2 m \omega_0 \alpha_4[\lambda]] \beta_2}{m \omega_0} + \alpha_3'[\lambda] \\
 \beta_4 = \alpha_4'[\lambda] \\
 \alpha_1[0] = 0 \\
 \alpha_2[0] = 0 \\
 \alpha_3[0] = 0 \\
 \alpha_4[0] = 0
 \end{array} \right)$$

Compare these results with Eqs. (95)-(98).

These equations are simple enough that we can attempt to solve them with **DSolve**.

```
sol = Simplify[DSolve[difeqs $\lambda$ , va,  $\lambda$ ][[1]]]
```

$$\left\{ \begin{aligned} \alpha_4[\lambda] &\rightarrow \lambda \beta_4, \alpha_2[\lambda] \rightarrow \frac{\sin[m \lambda \beta_4 \omega_0] (\cos[m \lambda \beta_4 \omega_0] \beta_2 + m \sin[m \lambda \beta_4 \omega_0] \beta_3 \omega_0)}{m \beta_4 \omega_0}, \\ \alpha_3[\lambda] &\rightarrow \frac{-2 \sin[m \lambda \beta_4 \omega_0]^2 \beta_2 + m \sin[2 m \lambda \beta_4 \omega_0] \beta_3 \omega_0}{2 m^2 \beta_4 \omega_0^2}, \\ \alpha_1[\lambda] &\rightarrow \frac{1}{16 m^3 \beta_4^2 \omega_0^3} \left(-8 m \cos[2 m \lambda \beta_4 \omega_0] \sin[m \lambda \beta_4 \omega_0]^2 \beta_2 \beta_3 \omega_0 - \beta_2^2 \right. \\ &\quad \left. (-4 \sin[2 m \lambda \beta_4 \omega_0] + \sin[4 m \lambda \beta_4 \omega_0] + 4 m \lambda \beta_4 \omega_0) + m^2 \right. \\ &\quad \left. \omega_0^2 (16 m \lambda \beta_1 \beta_4^2 \omega_0 + \beta_3^2 (\sin[4 m \lambda \beta_4 \omega_0] - 4 m \lambda \beta_4 \omega_0)) \right) \end{aligned} \right\}$$

Setting $\lambda=1$ we can obtain a relation between $\alpha(t)$ and $\beta(t)$ of the form (18).

```
va1 = Table[Subscript[a, k1], {k1, 1, n]];
eqs = Table[va1[[k1]] == ((va[[k1]] /. sol) /. { $\lambda \rightarrow 1$ }),
  {k1, 1, n}]
```

$$\left\{ \begin{aligned} \alpha_1 &== \frac{1}{16 m^3 \beta_4^2 \omega_0^3} \left(-8 m \cos[2 m \beta_4 \omega_0] \sin[m \beta_4 \omega_0]^2 \beta_2 \beta_3 \omega_0 - \right. \\ &\quad \left. \beta_2^2 (-4 \sin[2 m \beta_4 \omega_0] + \sin[4 m \beta_4 \omega_0] + 4 m \beta_4 \omega_0) + \right. \\ &\quad \left. m^2 \omega_0^2 (16 m \beta_1 \beta_4^2 \omega_0 + \beta_3^2 (\sin[4 m \beta_4 \omega_0] - 4 m \beta_4 \omega_0)) \right), \\ \alpha_2 &== \frac{\sin[m \beta_4 \omega_0] (\cos[m \beta_4 \omega_0] \beta_2 + m \sin[m \beta_4 \omega_0] \beta_3 \omega_0)}{m \beta_4 \omega_0}, \\ \alpha_3 &== \frac{-2 \sin[m \beta_4 \omega_0]^2 \beta_2 + m \sin[2 m \beta_4 \omega_0] \beta_3 \omega_0}{2 m^2 \beta_4 \omega_0^2}, \\ \alpha_4 &== \beta_4 \end{aligned} \right\}$$

Relation between $\alpha(t)$ and $\beta(t)$ via the eigenvalue one

eigenvectors of M_a^T

Working out the inverse relation between α and β might be difficult in this case given the complexity of the previous equations.

Therefore, it is convenient to obtain a relation between $\alpha(t)$ and $\beta(t)$ by obtaining the eigenvalue one eigenvectors of M_a^T . There are two eigenvalue one eigenvectors.

```

vα = Table[Subscript[α, k1], {k1, 1, n}];
Ma = LieGetMa[c, J, vα];
Mat = Transpose[Ma[[n + 1]]];
eval = Simplify[Eigenvalues[Mat]];
evec = Simplify[Eigenvectors[Mat]];
eval[[2]]
ρ1 = Simplify[ExpToTrig[evec[[2]]]]
eval[[3]]
ρ2 = Simplify[evec[[3]]]
vβ = γ1 ρ1 + γ2 ρ2

```

1

$$\{0, -m \omega_0 (-\cot[m \alpha_4 \omega_0] \alpha_2 + m \alpha_3 \omega_0), \alpha_2 + m \cot[m \alpha_4 \omega_0] \alpha_3 \omega_0, 1\}$$

1

$$\{1, 0, 0, 0\}$$

$$\{\gamma_2, -m \gamma_1 \omega_0 (-\cot[m \alpha_4 \omega_0] \alpha_2 + m \alpha_3 \omega_0), \gamma_1 (\alpha_2 + m \cot[m \alpha_4 \omega_0] \alpha_3 \omega_0), \gamma_1\}$$

Compare this last result with the one in Eq. (100).