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# Template for writing LHCb papers

The LHCb collaboration<sup>†</sup>

#### Abstract

Guidelines for the preparation of LHCb documents are given. This is a "living" document that should reflect our current practice. It is expected that these guidelines are implemented for papers before they go into the first collaboration wide review. Please contact the Editorial Board chair if you have suggestions for modifications. This is the title page for journal publications (PAPER). For a CONF note or ANA note, switch to the appropriate template by uncommenting the corresponding line in the file main.tex.

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#### 1 Introduction

<sup>2</sup> CP violation enters the Standard Model (SM) via the single irreducible phase in the

- <sup>3</sup> Cabbibo-Kobayashi-Maskawa (CKM) matrix. The unitarity requirement placed on this
- 4 matrix allows us to construct the Unitray Triangle (UT) in the complex plane from the
- 5 complex equation:

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0 (1)$$

This is of particular interest as all angles and sides are of measurable size and represent the weak phases and decay amplitudes respectively, of inter- and intra-generational quark transitions. The least well known angle,  $\phi_3 = arg(-V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$ , is accessible via tree-level processes, therefore, as well as being a fundamental parameter of the SM, a precise measurement of this angle at tree level can be compared to higher order processes in order to probe physics beyond the SM (BSM) that may enter decays via loops.

The LHCb experiment has previously probed  $\phi_3$  by studying the interference between  $b \to c\bar{u}s$  and, the colour and CKM suppressed,  $b \to u\bar{c}s$  transitions in  $B^\pm \to DK$  decays, where the D is an admixture of  $D^0$  and  $\bar{D}^0$  mesons decaying to a common final state. These could be two-body, non-self-conjugate multibody, or self-conjugate multibody final states. Similar studies have also been performed on neutral  $B^0$  and  $B^0_s$  mesons. This analysis makes use of self-conjugate multibody final states,  $D \to K^0_s h^+ h^-$ , in decays of the form  $B^\pm \to D^* K^\pm$ , where the charmed meson is produced in an excited  $D^*$  state and decays strongly via  $D^{*0} \to D\pi^0/\gamma$ . Sensitivity to  $\phi_3$  is then obtained by studying the Dalitz distribution of  $D \to K^0_s h^+ h^-$  events comparatively for  $B^+$  and  $B^-$  mesons. In order to achieve this, the strong phase variation of the D meson decay over the bins of the Dalitz plot must be known.

It is possible to determine this distribution using an amplitude model, derived from flavour-tagged  $D \to K_s^0 h^+ h^-$  decays. This has been successfully applied in previous analyses by the BaBar, Belle and LHCb collaborations. The alternative approach is a model-independent method, as has been demonstrated by Belle and LHCb, making use of direct measurements of the strong phase variation measured at CLEO-c in decays of  $\psi(3770)$  mesons to quantum-correlated  $D^0\bar{D}^0$  pairs. This data-driven approach is attractive as it avoids making the assumptions required by the model-dependant case.

This paper presents a model-independent study of  $B^{\pm} \to D^*K^{\pm}$  decays, with  $D^* \to D\pi^0$ ,  $D^* \to D\gamma$  and  $D \to K_s^0\pi^+\pi^-$ ,  $D \to K_s^0K^+K^-$ . The data used corresponds to a total integrated luminosity of 3.295 fb<sup>-1</sup> of pp collision acquired with the LHCb detector, at center-of-mass energies of  $\sqrt{s} = 7, 8, 13$  TeV over the years 2011, 2012 and 2015, respectively.

 $[1] \quad [2]$ 

## 2 Analysis Overview

The neutral  $D^*$  meson produced in the decay of the  $B^-$  is a superposition of particle, anti-particle states and can be expressed as:

$$D^* = D^{*0} + r_B e^{i(\delta_B - \phi_3)} \overline{D}^{*0}$$
 (2)

Here,  $r_B = \frac{|A(B^- \to \bar{D}^*K^-)|}{|A(B^- \to D^{*0}K^-)|}$ , the ratio of the strong decay amplitudes, and  $\delta_B$  is the strong phase difference between them. Writing this in terms of odd and even CP eigenstates,

where  $D_+^* = \frac{D^{*0} + \overline{D}^{*0}}{2}$  and  $D_-^* = \frac{D^{*0} - \overline{D}^{*0}}{2}$ :

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$$D^* = \frac{D^{*+} + D^{*-}}{\sqrt{2}} + r_B e^{i(\delta_B - \phi_3)} \frac{D^{*+} - D^{*-}}{\sqrt{2}}$$
 (3)

The CP eigenvalue of the  $D^*$  state is given by the product  $\lambda_{D^*} = \lambda_D \times \lambda_{\pi^0/\gamma} \times (-1)^l$ . In the case of the strong  $D^*$  decay via  $\pi^0$  emission,  $\lambda_{\pi^0} = -1$  and l = 1 therefore  $\lambda_{D^*} = \lambda_D$  and  $D^{*\pm} \to D^{\pm}\pi^0$ . Where as, for  $\gamma$  emission,  $\lambda_{\gamma} = +1$  and l = 1 therefore  $\lambda_{D^*} = -\lambda_D$  and  $D^{*\pm} \to D^{\mp}\gamma$ . This introduces a strong phase shift of  $\pi$  between the two neutral D mesons produced in the decay  $B^- \to D^*K^-$ . Defining  $\delta = \delta_B$  for the  $D^* \to D\pi^0$  decay and  $\delta = \delta_B + \pi$  for  $D^{*\pm} \to D_{\mp}\gamma$ , we can write the equivalent of Eq. (2) for the D meson produced in  $B^-$  decay:

$$D = D^0 + r_B e^{i(\delta - \phi_3)} \overline{D}^0 \tag{4}$$

This can be translated into an amplitude in Dalitz space,

$$A_B(m_-^2, m_+^2) \propto A + r_B e^{i\delta - \phi_3} \bar{A} \tag{5}$$

where  $m_{\pm}^2$  are the invariant masses squared of the  $K_s^0h^{\pm}$  combinations.  $A=A(m_-^2,m_+^2)$  is the Dalitz amplitude for the  $D^0$  decay and analogously for the  $\overline{D}^0$  decay,  $\overline{A}=\overline{A}(m_-^2,m_+^2)$ . Making the substitutions  $\phi_3 \to -\phi_3$  and  $A \leftrightarrow \overline{A}$  gives the equivalent expression for the amplitude of the D meson produced in the charge-conjugated decay,  $B^+ \to DK^+$ . The exent to which CP is violated in  $D^0 - \overline{D}^0$  mixing and Cabibbo-favoured D meson decays is known to be small, therefore, neglecting these second order effects, the conjugate amplitudes can be related by  $A_B(m_-^2, m_+^2) = \overline{A}_B(m_+^2, m_-^2)$ .

The Dalitz plot is constructed in 2N bins labelled -N to +N (excluding 0), symmetric about the line  $m_{-}^2 = m_{+}^2$ , where N> 0 corresponds to  $m_{-}^2 > m_{+}^2$ . It can be seen from Eq. 5 that the square of the amplitude is dependant on the strong phase difference between the  $D^0$  and  $\overline{D}^0$  decay:

$$\Delta \delta_D \equiv \delta_D(m_+^2, m_-^2) - \delta_D(m_-^2, m_+^2)$$
 (6)

where  $\delta_D(m_+^2, m_-^2) \equiv \arg A$  -  $\arg \bar{A}$  is the phase of  $A(m_+^2, m_-^2)$ . The equation for the cosine of the average strong phase difference, weighted by the decay rate, within each bin i of area  $\mathcal{D}_i$ , is given by:

$$c_i \equiv \frac{\int_{\mathcal{D}_i} (|A||\bar{A}|\cos \delta_D) d\mathcal{D}}{\sqrt{\int_{\mathcal{D}_i} |A|^2 d\mathcal{D}} \sqrt{\int_{\mathcal{D}_i} |\bar{A}|^2 d\mathcal{D}}}$$
(7)

An analogous expression can be constructed for the sine of the weighted strong phase difference within bin  $i, s_i$ . Both parameters have been measured by the CLEO-c experiment, who observed  $D^0 \overline{D}{}^0$  pairs produced at the  $\psi(3770)$  resonance. One D meson is reconstructed in  $K_{\rm S}^0$   $h^+h^-$  or  $K_{\rm L}^0$   $h^+h^-$  final state, and the other in a decay to  $K_{\rm S}^0$   $h^+h^-$  or a CP eigenstate. Combining the efficiency-corrected events yields in each bin with flavour-tagged information allows  $s_i$  and  $c_i$  to be determined, where their finite precision introduces a systematic uncertainty into the final result. The benefit of this method is that it avoids making assumptions about the nature of the intermediate resonances contributing to the  $K_{\rm S}^0$   $h^+h^-$  final state, which are required in the model-dependant approach when fitting an amplitude model to flavour-tagged  $D^0$  decays, in which an assumed functional form for |A|,  $|\bar{A}|$  and  $\delta_D$  are constructed.

\*\*\* Insert paragraph describing choice of binning scheme \*\*\*

The CP observables that are measured in this analysis are defined as:

$$x_{\pm} \equiv r_B \cos \left(\delta \pm \phi_3\right) \text{ and } y_{\pm} \equiv r_B \sin \left(\delta \pm \phi_3\right)$$
 (8)

and from these parameters, it is possible to extract  $r_B$ ,  $\delta_B$  and  $\phi_3$ . As described in Sec. 1,  $\delta = \delta_B$  and  $\delta = \delta_B + \pi$  for  $D^{*0} \to D\pi^0$  and  $D^* \to D\gamma$ , respectively. The extra strong phase difference of  $\pi$  introduced by the latter decay simply results in a negative sign in front of  $x_{\pm}$  and  $y_{\pm}$ , therefore does not provide any extra information.

Selections placed on the data in order to extract signal events result in nonuniformities within the Dalitz phase space, and the associated effect on each position in relation to the others must be accounted for. This is achieved by constructing a relative selection and reconstruction efficiency profile as a function of Dalitz plot position,  $\epsilon = \epsilon(m_-^2, m_+^2)$ . The fraction of events in bin i is therefore given by:

$$F_i = \frac{\int_{\mathcal{D}_i} |A|^2 \epsilon d\mathcal{D}}{\sum_i \int_{\mathcal{D}_i} |A|^2 \epsilon d\mathcal{D}}$$

$$\tag{9}$$

Denoting the number of events in positive (negative) bins due to  $B^+$  decays as  $N_{+i}^+$   $(N_{-i}^+)$ , and for  $B^-$  decays  $N_{+i}^ (N_{-i}^-)$ , and considering Eq. 5, it then follows that

$$N_{\pm i}^{+} = h_{B^{+}} \left[ F_{\mp i} + (x_{+}^{2} + y_{+}^{2}) F_{\pm i} + 2\sqrt{F_{i}F_{-i}} (x_{+}c_{\pm i} - y_{+}s_{\pm i}) \right]$$

$$N_{\pm i}^{-} = h_{B^{-}} \left[ F_{\pm i} + (x_{-}^{2} + y_{-}^{2}) F_{\mp i} + 2\sqrt{F_{i}F_{-i}} (x_{-}c_{\pm i} + y_{-}s_{\pm i}) \right]$$

$$(10)$$

where  $h_{B^{\pm}}$  accounts for the relative normalisation of  $B^+$  and  $B^-$  decays that arise from asymmetries in the production rates of bottom and anti-bottom mesons.

The decay  $(\overline{B}) \to D^{*\pm}\mu_{\mp}\nu_{\mu}X$ , where  $D^{*\pm} \to (\overline{D})^0\pi^{\pm}$  and  $(\overline{D})^0 \to K_s^0h^+h^-$  (X denotes other particles that could be produced in the  $(\overline{B})$  decay), is used as a control mode to determine  $F_i$ . Comparions are made between simulated  $(\overline{B}) \to D^{*\pm}\mu_{\mp}\nu_{\mu}X$  and  $B^{\pm} \to D^*K^{\pm}$  decays in order to correct for differences in their reconstruction and selection efficiencies. Decays of the higher statistics mode,  $B^{\pm} \to D^*\pi^{\pm}$ , are also studied in order to develop the signal selection procedure and determine the yield of  $B^{\pm} \to D^*\pi^{\pm}$  missidentified as  $B^{\pm} \to D^*K^{\pm}$  candidates.

The effect of  $D^0$ - $\overline{D}^0$  mixing was neglected in the CLEO-c measurements of  $c_i$  and  $s_i$ , as it is in the current analysis. This introduces a bias of 0.2° in the derived measurement of  $\phi_3$ . The CP violating effect in  $K_s^0$  decays was also not considered, and is expected to introduce  $\mathcal{O}(1^\circ)$  uncertainty to the final result. An uncertainty of similar magnitude enters due to the different interaction cross sections of the  $K^0$  and  $\overline{K}^0$  mesons. The results of these neglected effects are negligible in the current analysis, considering the expected precision of the  $\phi_3$  measurement that can be obtained.

\*\*\* Insert paragraph describing ANA note layout \*\*\*

#### 3 Detector and simulation

The LHCb detector [?,1] is a single-arm forward spectrometer covering the pseudorapidity range  $2 < \eta < 5$ , designed for the study of particles containing b or c quarks. The detector

includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the pp interaction region [?]\*, a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes [?]\* placed downstream of the magnet. The tracking system provides a measurement of momentum, p, of charged particles with a relative uncertainty that varies from 0.5% at low momentum to 1.0% at  $200 \,\text{GeV}/c$ . The minimum distance of a track to a primary vertex (PV), the impact parameter (IP), is measured with a resolution of  $(15 + 29/p_T) \mu m$ , where  $p_T$  is the component of the momentum transverse to the beam, in GeV/c. Different types of charged hadrons are distinguished using information from two ring-imaging Cherenkov detectors [?]\*. Photons, electrons and hadrons are identified by a calorimeter system consisting of scintillating-pad and preshower detectors, an electromagnetic calorimeter and a hadronic calorimeter. Muons are identified by a system composed of alternating layers of iron and multiwire proportional chambers [?]\*. The online event selection is performed by a trigger [?]\*, which consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage, which applies a full event reconstruction.

At the hardware trigger stage, events are required to have a muon with high  $p_{\rm T}$  or a hadron, photon or electron with high transverse energy in the calorimeters. For hadrons, the transverse energy threshold is 3.5 GeV. The software trigger requires a two-, three- or four-track secondary vertex with a significant displacement from the primary pp interaction vertices. At least one charged particle must have a transverse momentum  $p_{\rm T} > 1.7 \, {\rm GeV}/c$  and be inconsistent with originating from a PV. A multivariate algorithm [3] is used for the identification of secondary vertices consistent with the decay of a b hadron.

Decays of  $K_s^0 \to \pi^+\pi^-$  are reconstructed in two different categories: the first involving  $K_s^0$  mesons that decay early enough for the daughter pions to be reconstructed in the vertex detector; and the second containing  $K_s^0$  that decay later such that track segments of the pions cannot be formed in the vertex detector. These categories are referred to as long and downstream, respectively. The long category has better mass, momentum and vertex resolution than the downstream category.

In the simulation, pp collisions are generated using Pythia [4] with a specific LHCb configuration [5]. Decays of hadronic particles are described by EvtGen [6], in which final-state radiation is generated using Photos [7]. The interaction of the generated particles with the detector, and its response, are implemented using the Geant4 toolkit [8] as described in Ref. [9].

## 142 Appendices

### A Standard References

Below is a list of common references, as well as a list of all LHCb publications. As they are already in prepared bib files, they can be used as simply as \cite{Alves:2008zz} to get the LHCb detector paper. The references are defined in the files main.bib, LHCb-PAPER.bib, LHCb-CONF.bib, LHCb-DP.bib LHCb-TDR.bib files, with obvious contents. Each of these have their LHCb-ZZZ-20XX-0YY number as their cite code. If you believe there is a problem with the formatting or content of one of the entries, then get in contact with the Editorial Board rather than just editing it in your local file, since you are likely to need the latest version just before submiting the article.

Description	cite code	Reference
LHCb detector	Alves:2008zz	[1]
LHCb simulation	LHCb-PROC-2011-006	[9]
PDG 2014	PDG2014	[10]
HFAG	HFAG	[11]
Рутніа	Sjostrand:2006za, *Sjostrand:2007gs	[4]
LHCb Pythia tuning	LHCb-PROC-2010-056	[5]
Geant4	Allison:2006ve, *Agostinelli:2002hh	[8]
EVTGEN	Lange: 2001uf	[6]
PHOTOS	Golonka:2005pn	[7]
DIRAC	Tsaregorodtsev:2010zz, *BelleDIRACAmazon	[12]
Crystal Ball function <sup>1</sup>	Skwarnicki:1986xj	[13]
Wilks' theorem	Wilks:1938dza	[14]
BDT	Breiman	[15]
BDT training	AdaBoost	[16]
HLT2 topo	BBDT	[3]
DecayTreeFitter	Hulsbergen:2005pu	[17]
sPlot	Pivk:2004ty	[18]
Punzi's optimization	Punzi:2003bu	[19]
$f_s/f_d$	fsfd	[20]

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<sup>&</sup>lt;sup>1</sup>A valid alternative for most papers where the normalisation is not critical is to use the expression "Gaussian function with a low-mass power-law tail" or "Gaussian function with power-law tails". In that case, no citation is needed

LHCb-DP number	Title
LHCb-DP-2016-001 [?]	TESLA project
LHCb-DP-2014-002 [?]	LHCb detector performance
LHCb-DP-2014-001 [?]	Performance of the LHCb Vertex Locator
LHCb-DP-2013-004 [?]	Performance of the LHCb calorimeters
LHCb-DP-2013-003 [?]	Performance of the LHCb Outer Tracker
LHCb-DP-2013-002 [?]	Measurement of the track reconstruction efficiency at LHCb
LHCb-DP-2013-001 [?]	Performance of the muon identification at LHCb
LHCb-DP-2012-005 [?]	Radiation damage in the LHCb Vertex Locator
LHCb-DP-2012-004 [?]	The LHCb trigger and its performance in 2011
LHCb-DP-2012-003 [?]	Performance of the LHCb RICH detector at the LHC
LHCb-DP-2012-002 [?]	Performance of the LHCb muon system
LHCb-DP-2012-001 [?]	Radiation hardness of the LHCb Outer Tracker
LHCb-DP-2011-002 [?]	Simulation of machine induced background
LHCb-DP-2011-001 [?]	Performance of the LHCb muon system with cosmic rays
LHCb-DP-2010-001 [?]	First spatial alignment of the LHCb VELO

LHCb-TDR number	Title
LHCb-TDR-016 [?]	Trigger and online upgrade
LHCb-TDR-015 [?]	Tracker upgrade
LHCb-TDR-014 [?]	PID upgrade
LHCb-TDR-013 [?]	VELO upgrade
LHCb-TDR-012 [?]	Framework TDR for the upgrade
LHCb-TDR-011 [?]	Computing
LHCb-TDR-010 [?]	Trigger
LHCb-TDR-009 [?]	Reoptimized detector
LHCb-TDR-008 [?]	Inner Tracker
LHCb-TDR-007 [?]	Online, DAQ, ECS
LHCb-TDR-006 [?]	Outer Tracker
LHCb-TDR-005 [?]	VELO
LHCb-TDR-004 [?]	Muon system
LHCb-TDR-003 [?]	RICH
LHCb-TDR-002 [?]	Calorimeters
LHCb-TDR-001 [?]	Magnet

Table 2: LHCb-PAPERs (which have their identifier as their cite code). Note that LHCb-PAPER-2011-039 does not exist.

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Table 3: LHCb-CONFs (which have their identifier as their cite code). Note that LHCb-CONF-2011-032 does not exist.

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LHCb-CONF-2012-032 [?]	LHCb-CONF-2012-031 [?]

<sup>&</sup>lt;sup>2</sup>If you cite the gamma combination, always also cite the latest gamma paper as \cite{LHCb-PAPER-2013-020,\*LHCb-CONF-2014-004} (unless you cite LHCb-PAPER-2013-020 separately too).

- continued from previous page. LHCb-CONF-2012-030 [?] LHCb-CONF-2012-029 [? LHCb-CONF-2012-027 [?] LHCb-CONF-2012-028 LHCb-CONF-2012-026 |?| LHCb-CONF-2012-025 |?| LHCb-CONF-2012-024 [?] LHCb-CONF-2012-023 [?] LHCb-CONF-2012-022 [?] LHCb-CONF-2012-021 [?] LHCb-CONF-2012-020 [?] LHCb-CONF-2012-019 [?] LHCb-CONF-2012-018 [?] LHCb-CONF-2012-017 [?] LHCb-CONF-2012-016 [?] LHCb-CONF-2012-015 [?] LHCb-CONF-2012-014 [?] LHCb-CONF-2012-013 [?] LHCb-CONF-2012-012 [?] LHCb-CONF-2012-011 [?] LHCb-CONF-2012-010 [?] LHCb-CONF-2012-009 [?] LHCb-CONF-2012-008 |?| LHCb-CONF-2012-007 |?| LHCb-CONF-2012-006 [?] LHCb-CONF-2012-005 [?] LHCb-CONF-2012-004 [?] LHCb-CONF-2012-003 [?] LHCb-CONF-2012-002 [? LHCb-CONF-2012-001 |?| LHCb-CONF-2011-062 [? LHCb-CONF-2011-061 [?] LHCb-CONF-2011-060 [?] LHCb-CONF-2011-059 [?] LHCb-CONF-2011-058 [?] LHCb-CONF-2011-057 [?] LHCb-CONF-2011-055 [?] LHCb-CONF-2011-056 [?] LHCb-CONF-2011-054 [?] LHCb-CONF-2011-053 [?] LHCb-CONF-2011-052 |?| LHCb-CONF-2011-051 |?| LHCb-CONF-2011-050 [?] LHCb-CONF-2011-049 [?] LHCb-CONF-2011-048 [?] LHCb-CONF-2011-047 [?] LHCb-CONF-2011-045 [?] LHCb-CONF-2011-046 |? LHCb-CONF-2011-044 [?] LHCb-CONF-2011-043 [?] LHCb-CONF-2011-042 [?] LHCb-CONF-2011-041 [?] LHCb-CONF-2011-040 [?] LHCb-CONF-2011-039 [?] LHCb-CONF-2011-038 [?] LHCb-CONF-2011-037 [?] LHCb-CONF-2011-036 [?] LHCb-CONF-2011-035 [?] LHCb-CONF-2011-034 [?] LHCb-CONF-2011-033 [?] LHCb-CONF-2011-031 |?| LHCb-CONF-2011-030 [?] LHCb-CONF-2011-029 [?] LHCb-CONF-2011-027 [?] LHCb-CONF-2011-028 [?] LHCb-CONF-2011-026 [?] LHCb-CONF-2011-025 [?] LHCb-CONF-2011-024 [?] LHCb-CONF-2011-023 [?] LHCb-CONF-2011-023 [?] LHCb-CONF-2011-021 |?| LHCb-CONF-2011-020 [?] LHCb-CONF-2011-019 [?] LHCb-CONF-2011-018 [?] LHCb-CONF-2011-017 [?] LHCb-CONF-2011-016 [?] LHCb-CONF-2011-015 [?] LHCb-CONF-2011-014 [?] LHCb-CONF-2011-013 [?] LHCb-CONF-2011-012 [?] LHCb-CONF-2011-011 [?] LHCb-CONF-2011-010 [?] LHCb-CONF-2011-009 [?] LHCb-CONF-2011-008 |?| LHCb-CONF-2011-007 |?| LHCb-CONF-2011-006 [?] LHCb-CONF-2011-005 [?]

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LHCb-CONF-2010-008 [?]	

Some LHCb papers quoted together will look like [344–348]. The combination of CMS and LHCb results on  $B_{(s)}^0 \to \mu^+\mu^-$  should be cited like [?].

## 159 B Standard symbols

As explained in Sect. ?? this appendix contains standard typesetting of symbols, particle names, units etc. in LHCb documents.

In the file lhcb-symbols-def.tex, which is included, a large number of symbols is defined. While they can lead to quicker typing, the main reason is to ensure a uniform notation within a document and between different LHCb documents. If a symbol like \CP to typeset CP violation is available for a unit, particle name, process or whatever, it should be used. If you do not agree with the notation you should ask to get the definition in lhcb-symbols-def.tex changed rather than just ignoring it.

All the main particles have been given symbols. The B mesons are thus named  $B^+$ ,  $B^0$ ,  $B_s^0$ , and  $B_c^+$ . There is no need to go into math mode to use particle names, thus saving the typing of many \$ signs. By default particle names are typeset in italic type to agree with the PDG preference. To get roman particle names you can just change \setboolean{uprightparticles}{false} to true at the top of this template.

There is a large number of units typeset that ensures the correct use of fonts, capitals and spacing. As an example we have  $m_{B_s^0} = 5366.3 \pm 0.6 \,\mathrm{MeV}/c^2$ . Note that  $\,\mu\mathrm{m}$  is typeset with an upright  $\,\mu$ , even if the particle names have slanted greek letters.

A set of useful symbols are defined for working groups. More of these symbols can be included later. As an example in the Rare Decay group we have several different analyses looking for a measurement of  $C_7^{'(\text{eff})}$  and  $O_7^{'}$ .

## C List of all symbols

#### 30 C.1 Experiments

	ackslashlhcb	LHCb	ackslash	ATLAS	$\backslash \mathtt{cms}$	CMS
\	$ar{alice}$	ALICE	ackslashbabar	BaBar	\belle	Belle
	\cleo	CLEO	$\backslash \mathtt{cdf}$	CDF	\dzero	D0
181	$ackslash  ext{aleph}$	ALEPH	$\backslash \mathtt{delphi}$	DELPHI	$\setminus$ opal	OPAL
	$ackslash  ext{lthree}$	L3	$\backslash \mathtt{sld}$	SLD	$\backslash \mathtt{cern}$	CERN
	$\backslash  ext{lhc}$	LHC	\lep	LEP	ackslashtevatron	Tevatron

#### $^{182}$ C.1.1 LHCb sub-detectors and sub-systems

	\velo	VELO	$\$ rich	RICH	$\backslash { t richone}$	RICH1
,	$\backslash \mathtt{richtwo}$	RICH2	ackslash ttracker	$\operatorname{TT}$	$\setminus$ intr	$\operatorname{IT}$
	\st	ST	\ot	OT	\spd	SPD
	$\backslash \mathtt{presh}$	PS	\ecal	ECAL	$\$ hcal	HCAL
183	$\setminus$ MagUp	MagUp	$\backslash { t MagDown}$	MagDown	$\setminus$ ode	ODE
	\daq	DAQ	\tfc	TFC	\ecs	ECS
,	$\backslash  exttt{lone}$	L0	$ackslash  ext{hlt}$	HLT	$\hline$	HLT1
	$ackslash  ext{hlttwo}$	HLT2				

#### 184 C.2 Particles

#### 185 C.2.1 Leptons

	$\setminus$ electron	e	\en	$e^-$	\ep	$e^+$
	$\operatorname{\mathtt{f epm}}$	$e^{\pm}$	$\backslash \mathtt{epem}$	$e^+e^-$	$\backslash \mathtt{muon}$	$\mu$
	\mup	$\mu^+$	$\backslash \mathtt{mun}$	$\mu^-$	$\backslash$ mumu	$\mu^+\mu^-$
	ackslashtauon	au	ackslashtaup	$ au^+$	$\setminus$ taum	$ au^-$
186	ackslashtautau	$ au^+ au^-$	$\setminus$ lepton	$\ell$	$\backslash \mathtt{ellm}$	$\ell^-$
	\ellp	$\ell^+$	ackslashellell	$\ell^+\ell^-$	\neu	$\nu$
	\neub	$\overline{ u}$	$\setminus$ neue	$ u_e$	\neueb	$\overline{ u}_e$
	$\nextriangle$	$ u_{\mu}$	$\setminus$ neumb	$\overline{ u}_{\mu}$	$\setminus$ neut	$ u_{ au}$
	$\setminus$ neutb	$\overline{ u}_{ au}$	$\setminus \mathtt{neul}$	$ u_{\ell}$	$\ne $	$\overline{ u}_\ell$

## $_{187}$ C.2.2 Gauge bosons and scalars

	\g	$\gamma$	$\backslash H$	$H^0$	$\backslash \mathtt{Hp}$	$H^+$
	$\backslash Hm$	$H^-$	$\backslash \texttt{Hpm}$	$H^\pm$	$\backslash W$	W
188	\Wp	$W^+$	\Wm	$W^-$	\Wpm	$W^{\pm}$
	$\setminus Z$	Z				

#### 189 C.2.3 Quarks

	$\quark$	q	$\setminus$ quarkbar	$\overline{q}$	$\q$	$q\overline{q}$
	$\setminus$ uquark	u	$ackslash  ext{uquarkbar}$	$\overline{u}$	$\setminus$ uubar	$u\overline{u}$
	$\backslash \mathtt{dquark}$	d	$\setminus$ dquarkbar	$\overline{d}$	$\backslash \mathtt{ddbar}$	$d\overline{d}$
190	$\setminus \mathtt{squark}$	s	$ackslash  ext{squarkbar}$	$\overline{s}$	$\backslash \mathtt{ssbar}$	$s\overline{s}$
	$\backslash \mathtt{cquark}$	c	$\setminus$ cquarkbar	$\overline{c}$	$\backslash \mathtt{ccbar}$	$c\overline{c}$
	ackslashbquark	b	ackslashbquarkbar	$\overline{b}$	ackslashbbbar	$b\overline{b}$
	$\t$ tquark	t	$\$ tquarkbar	$\overline{t}$	\ttbar	$t\overline{t}$

## 191 C.2.4 Light mesons

	$\backslash \texttt{hadron}$	h	\pion	$\pi$	\piz	$\pi^0$
	$ackslash  ext{pizs}$	$\pi^0\mathrm{s}$	\pip	$\pi^+$	\pim	$\pi^-$
	\pipm	$\pi^\pm$	$\neq$	$\pi^{\mp}$	$\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ $	$\rho$
	$\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ $	$ ho^0$	$\r$ hop	$ ho^+$	$\r$ hom	$ ho^-$
	$\rdown$	$ ho^\pm$	$\rdot rhomp$	$ ho^{\mp}$	$\setminus$ kaon	K
	\Kb	$\overline{K}$	$\backslash \texttt{KorKbar}$	$(\overline{K})$	$\backslash \mathtt{Kz}$	$K^0$
192	$\backslash \texttt{Kzb}$	$\overline{K}{}^0$	$\backslash Kp$	$K^+$	$\backslash \mathtt{Km}$	$K^{-}$
	$\backslash \mathtt{Kpm}$	$K^{\pm}$	$\backslash \mathtt{Kmp}$	$K^{\mp}$	\KS	$K_{\mathrm{S}}^{0}$ $\overline{K}^{*0}$
	$\backslash KL$	$K_{\scriptscriptstyle m L}^0$	$\setminus \texttt{Kstarz}$	$K^{*0}$	$\setminus$ Kstarzb	$\overline{K}^{*0}$
	$\setminus$ Kstar	$K^*$	$\setminus$ Kstarb	$\overline{K}^*$	$\setminus$ Kstarp	$K^{*+}$
	$\setminus \texttt{Kstarm}$	$K^{*-}$	$\setminus \texttt{Kstarpm}$	$K^{*\pm}$	$\setminus \texttt{Kstarmp}$	$K^{*\mp}$
	ackslashetaz	$\eta$	$\backslash \mathtt{etapr}$	$\eta'$	$ackslash  ext{phiz}$	$\phi$
	$\backslash \mathtt{omegaz}$	$\omega$				

## 193 C.2.5 Heavy mesons

\D	D	\Db	$\overline{D}$	$\backslash { t Dor Dbar}$	$(\overline{D})$
\Dz	$D^0$	\Dzb	$\overline{D}{}^0$	\Dp	$D^+$
\Dm	$D^-$	$\backslash \mathtt{Dpm}$	$D^{\pm}$	$\backslash \mathtt{Dmp}$	$D^{\mp}$
\Dstar	$D^*$	$ackslash  exttt{Dstarb}$	$\overline{D}^*$	$ackslash \mathtt{Dstarz}$	$D^{*0}$
\Dstarzb	$\overline{D}^{*0}$	$ackslash  exttt{Dstarp}$	$D^{*+}$	$\backslash \mathtt{Dstarm}$	$D^{*-}$
$\backslash \mathtt{Dstarpm}$	$D^{*\pm}$	$\backslash \mathtt{Dstarmp}$	$D^{*\mp}$	\Ds	$D_s^+$
$\backslash \mathtt{Dsp}$	$D_s^+$	$\backslash \mathtt{Dsm}$	$D_s^-$	$\backslash \mathtt{Dspm}$	$D_s^{\pm}$
$\backslash \mathtt{Dsmp}$	$D_s^{\mp}$	$ackslash  extsf{Dss}$	$D_s^{*+}$	$\backslash \mathtt{Dssp}$	$D_s^{*+}$
$\backslash \mathtt{Dssm}$	$D_s^{*-}$	$\backslash \mathtt{Dsspm}$	$D_s^{*\pm}$	$\backslash \mathtt{Dssmp}$	$\overline{B}^{*\mp}$
\B	B	\Bbar	$\overline{B}$	\Bb	$\overline{B}$
\BorBbar	$(\overline{B})$	\Bz	$B^0$	\Bzb	$\overline{B}{}^0$
\Bu	$B^+$	Bub	$B^-$	\Bp	$B^+$
\Bm	$B^-$	$\backslash \mathtt{Bpm}$	$B^{\pm}$	$\backslash \mathtt{Bmp}$	$B^{\mp}$
\Bd	$B^0$	\Bs	$B_s^0$	\Bsb	$\overline{B}_s^0$
\Bdb	$\overline{B}{}^0$	\Bc	$B_c^+$	$\backslash \mathtt{Bcp}$	$B_c^+$
\Bcm	$B_c^-$	$\backslash \texttt{Bcpm}$	$B_c^{\pm}$		
	\Dz \Dm \Dstar \Dstarzb \Dstarpm \Dsp \Dsmp \Dsmp \Dssm \B \BorBbar \Bu \Bu \Bd	$egin{array}{lll} egin{array}{lll} egin{array}{lll} Dz & D^0 \\ egin{array}{lll} DDm & D^- \\ egin{array}{lll} DStar & D^* \\ egin{array}{lll} DStarpm & D^{*\pm} \\ egin{array}{lll} DSp & D_s^+ \\ egin{array}{lll} DSmp & D_s^{\mp} \\ egin{array}{lll} DSSm & D_s^{*-} \\ egin{array}{lll} B & B \\ egin{array}{lll} BOrBbar & \overline{B}^0 \\ egin{array}{lll} BD & B^- \\ egin{array}{lll} BD & B^0 \\ BD & BD \\ egin{array}{lll} BD & B^0 \\ BD & BD $	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$

#### 195 C.2.6 Onia

	\jpsi	$J\!/\psi$	$ackslash  exttt{psitwos}$	$\psi(2S)$	$ackslash  ext{psiprpr}$	$\psi(3770)$
	ackslashetac	$\eta_c$	ackslashchiczero	$\chi_{c0}$	ackslashchicone	$\chi_{c1}$
196	$\backslash \mathtt{chictwo}$	$\chi_{c2}$	$\setminus \mathtt{OneS}$	$\Upsilon(1S)$	\TwoS	$\Upsilon(2S)$
	ThreeS	$\Upsilon(3S)$	FourS	$\Upsilon(4S)$	\FiveS	$\Upsilon(5S)$
	$\backslash \mathtt{chic}$	$\chi_c$				

#### 197 C.2.7 Baryons

	$\proton$	p	$\antiproton$	$\overline{p}$	$\new red$	n
	$\setminus$ antineutron	$\overline{n}$	$\backslash \mathtt{Deltares}$	$\Delta$	$\backslash \mathtt{Deltaresbar}$	$\overline{\Delta}$
	\Xires	arvarepsilon	$\backslash \mathtt{Xiresbar}$	$\overline{\Xi}$	\Lz	$\Lambda$
	\Lbar	$\overline{\varLambda}$	$\backslash  ext{LorLbar}$	$(\overline{\Lambda})$	$\backslash \texttt{Lambdares}$	$\Lambda$
	$\backslash \texttt{Lambdaresbar}$	$\overline{\varLambda}$	$\backslash \texttt{Sigmares}$	$\Sigma$	$\backslash \texttt{Sigmaresbar}$	$\overline{\Sigma}$
	$\backslash \mathtt{Omegares}$	$\Omega$	$\backslash {\tt Omegaresbar}$	$\overline{\varOmega}$	\Lb	$A_b^0$
198	\Lbbar	$\overline{\varLambda}^0_b$	\Lc	$\Lambda_c^+$	$ackslash  ext{Lcbar}$	$\overline{\Lambda}_c^-$
	\Xib	$\Xi_b$	\Xibz	$\Xi_b^0$	$\backslash \mathtt{Xibm}$	$\Xi_b^-$
	\Xibbar	$\overline{\Xi}_b$	$\setminus \texttt{Xibbarz}$	$\frac{\Xi_b^0}{\Xi_b^0}$	$\setminus \texttt{Xibbarp}$	$\overline{\Xi}_b^+$
	\Xic	$\frac{\Xi_c}{\Xi_c}$	\Xicz	$\frac{\Xi_c^0}{\Xi_c^0}$	\Xicp	$\begin{array}{ccc} \Xi_b^- \\ \overline{\Xi}_b^+ \\ \overline{\Xi}_c^- \\ \overline{\Xi}_c^- \end{array}$
	\Xicbar	$\overline{\Xi}_c$	$\setminus \texttt{Xicbarz}$	$\overline{\Xi}_c^0$	$\setminus \mathtt{Xicbarm}$	$\overline{\Xi}_c^-$
	$\backslash \mathtt{Omegac}$	$\Omega_c^0$	$\backslash {\tt Omegacbar}$	$\overline{\Omega}_c^0$	$\backslash \mathtt{Omegab}$	$\Omega_b^-$
	$\backslash {\tt Omegabbar}$	$\overline{\varOmega}_b^+$				

## 199 C.3 Physics symbols

#### 200 C.3.1 Decays

	\BF	$\mathcal B$	\BRvis $\mathcal{B}_{ ext{vis}}$	$\backslash \mathtt{BR}$	${\cal B}$
201	$\decay[2] \decay\{a \}\{b c \}$	$a \rightarrow bc$	$\$ $\rightarrow$	$\backslash  exttt{to}$	$\rightarrow$

## 202 C.3.2 Lifetimes

	$\setminus \mathtt{tauBs}$	$ au_{B^0_s}$	$\backslash { t tauBd}$	$ au_{B^0}$	$\backslash { t tauBz}$	$ au_{B^0}$
203	$\setminus$ tau $Bu$	$ au_{B^+}$	$\setminus \mathtt{tauDp}$	$ au_{D^+}$	$\setminus \mathtt{tauDz}$	$ au_{D^0}$
	$\setminus \mathtt{tauL}$	$ au_{ m L}$	$\setminus \mathtt{tauH}$	$ au_{ m H}$		

#### 204 C.3.3 Masses

	$\backslash mBd$		$\backslash \mathtt{mBp}$	$m_{B^+}$	$\backslash \mathtt{mBs}$	$m_{B_s^0}$
205	$\backslash mBc$	$m_{B_c^+}$	$\mbox{mLb}$	$m_{A_b^0}$		

## 206 C.3.4 EW theory, groups

	\grpsuthree	SU(3)	$\grpsutw$	SU(2)	$\grpu$ one	U(1)
207	\ssqtw	$\sin^2\! heta_{ m W}$	$\backslash \mathtt{csqtw}$	$\cos^2 \theta_{ m W}$	ackslashstw	$\sin \theta_{ m W}$
	\ctw	$\cos  heta_{ m W}$	$\setminus \mathtt{ssqtwef}$	$\sin^2\! heta_{ m W}^{ m eff}$	$\backslash \mathtt{csqtwef}$	$\cos^2 \theta_{ m W}^{ m eff}$
	\stwef	$\sin heta_{ m W}^{ m eff}$	$\backslash \mathtt{ctwef}$	$\cos heta_{ m W}^{ m eff}$	\gv	$g_{ m \scriptscriptstyle V}$
	\ga	$g_{ m A}$	$\backslash \mathtt{order}$	$\mathcal O$	$\backslash \mathtt{ordalph}$	$\mathcal{O}(\alpha)$
	\ordalsq	$\mathcal{O}(lpha^2)$	$\backslash \mathtt{ordalcb}$	$\mathcal{O}(lpha^3)$		

## $^{208}$ C.3.5 QCD parameters

#### 210 C.3.6 CKM, CP violation

	\eps	arepsilon	\epsK	$arepsilon_K$	\epsB	$\varepsilon_B$
	\epsp	$arepsilon_K'$	\CP	CP	\CPT	CPT
	$\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ $	$\overline{ ho}$	\etabar	$\overline{\eta}$	\Vud	$V_{ud}$
	$\backslash \mathtt{Vcd}$	$V_{cd}$	\Vtd	$V_{td}$	\ <b>Vus</b>	$V_{us}$
211	\Vcs	$V_{cs}$	\Vts	$V_{ts}$	\Vub	$V_{ub}$
	Vcb	$V_{cb}$	\\Vtb	$V_{tb}$	\Vuds	$V_{ud}^*$
	Vcds	$V_{cd}^*$	\Vtds	$V_{td}^*$	Vuss	$V_{us}^*$
	Vcss	$V_{cs}^*$	Vtss	$V_{ts}^*$	\\\Vubs	$V_{ub}^*$
	Vcbs	$V_{ab}^*$	\Vtbs	$V_{\iota\iota}^*$	·	200

#### 212 C.3.7 Oscillations

	$\backslash dm$	$\Delta m$	$\backslash \mathtt{dms}$	$\Delta m_s$	$\backslash \mathtt{dmd}$	$\Delta m_d$
\	\DG	$\Delta\Gamma$	$ackslash  exttt{DGs}$	$\Delta\Gamma_s$	\DGd	$\Delta\Gamma_d$
	\Gs	$\Gamma_s$	\Gd	$\Gamma_d$	$\backslash \mathtt{MBq}$	$M_{B_q}$
\	\DGq	$\Delta\Gamma_q$	$\backslash \mathtt{Gq}$	$\Gamma_q$	$\backslash \mathtt{dmq}$	$\Delta m_q$
	\GL	$\Gamma_{ m L}$	$\backslash \mathtt{GH}$	$\Gamma_{ m H}$	$\backslash \mathtt{DGsGs}$	$\Delta\Gamma_s/\Gamma_s$
213	$\backslash \mathtt{Delm}$	$\Delta m$	$\backslash \texttt{ACP}$	$\mathcal{A}^{CP}$	$\setminus \texttt{Adir}$	$\mathcal{A}^{\mathrm{dir}}$
	$\backslash \texttt{Amix}$	$\mathcal{A}^{ ext{mix}}$	$ackslash  exttt{ADelta}$	$\mathcal{A}^{\Delta}$	$\backslash  exttt{phid}$	$\phi_d$
\	$\backslash \mathtt{sinphid}$	$\sin \phi_d$	$ackslash  ext{phis}$	$\phi_s$	ackslash	$\beta_s$
	$\backslash \mathtt{sbetas}$	$\sigma(\beta_s)$	ackslashstbetas	$\sigma(2\beta_s)$	$\backslash \mathtt{stphis}$	$\sigma(\phi_s)$
	$\backslash \mathtt{sinphis}$	$\sin \phi_s$				

#### 214 C.3.8 Tagging

	\edet	$arepsilon_{ m det}$	\erec	$\varepsilon_{ m rec/det}$	\esel	$\varepsilon_{ m sel/rec}$
	ackslashetrg	$\varepsilon_{\mathrm{trg/sel}}$	\etot	$arepsilon_{ ext{tot}}$	\mistag	$\omega$
215	$\backslash \mathtt{wcomb}$	$\omega^{\mathrm{comb}}$	\etag	$\varepsilon_{ m tag}$	ackslashetagcomb	$arepsilon_{ ext{tag}}^{ ext{comb}}$
	\effeff	$arepsilon_{ ext{eff}}$	\effeffcomb	$arepsilon_{ ext{eff}}^{ ext{comb}}$	\	$\varepsilon_{\rm tag}(1-2\omega)^2$
	\effD	$\varepsilon_{\mathrm{tag}} D^2$	\etagprompt	$arepsilon_{ ext{tag}}^{ ext{Pr}}$	$\backslash \mathtt{etagLL}$	$arepsilon_{ ext{tag}}^{ ext{LL}}$

#### 216 C.3.9 Key decay channels

	$\backslash \texttt{BdToKstmm}$	$B^0 \to K^{*0} \mu^+ \mu^-$	$^-ackslash BdbToKstmm$	$\overline{B}{}^0 \to \overline{K}{}^{*0} \mu^+ \mu^-$	$^-ackslash BsToJPsiPhi$	$B_s^0 \to J/\psi  \phi$
	$\BdToJPsiKst$	$B^0 \rightarrow J/\psi K^{*0}$	$\BdbToJPsiKst$	$\overline B{}^0\! o J\!/\!\psi\overline K^{*0}$	ackslash BsPhiGam	$B_s^0 \to \phi \gamma$
217	$\backslash \texttt{BdKstGam}$	$B^0\! o K^{*0}\gamma$	$\backslash \mathtt{BTohh}$	$B \rightarrow h^+ h^{\prime -}$	$\backslash \texttt{BdTopipi}$	$B^0 \rightarrow \pi^+\pi^-$
	\BdToKpi	$B^0 \rightarrow K^+\pi^-$	$\backslash \mathtt{BsToKK}$	$B_s^0 \rightarrow K^+K^-$	$\backslash  exttt{BsTopiK}$	$B_s^0 \rightarrow \pi^+ K^-$

#### 218 C.3.10 Rare decays

#### 220 C.3.11 Wilson coefficients and operators

#### 222 C.3.12 Charm

#### 224 C.3.13 QM

#### 226 C.4 Units

 $_{227}$  \unit[1] \unit{kg} kg

#### 228 C.4.1 Energy and momentum

	\tev	TeV	\gev	$\operatorname{GeV}$	$\backslash \mathtt{mev}$	MeV
	kev	keV	\ev	eV	\gevc	GeV/c
229	$\backslash \mathtt{mevc}$	MeV/c	\gevcc	$\text{GeV}/c^2$	\gevgevcccc	$\text{GeV}^2/c^4$
	\mevcc	$MeV/c^2$				

#### 230 C.4.2 Distance and area

	$\backslash \mathtt{km}$	km	$\backslash m$	m	$\backslash \mathtt{ma}$	$\mathrm{m}^2$
	$\backslash \mathtt{cm}$	cm	$\backslash \mathtt{cma}$	$\mathrm{cm}^2$	$\backslash mm$	mm
,	$\backslash \mathtt{mma}$	$\mathrm{mm}^2$	$\backslash \mathtt{mum}$	μm	$\backslash \mathtt{muma}$	$\mu\mathrm{m}^2$
	$\backslash \mathtt{nm}$	nm	$\backslash \mathtt{fm}$	fm	\barn	b
231	\mbarn	mb	\mub	μb	\nb	nb
	\invnb	$nb^{-1}$	\pb	pb	\invpb	$pb^{-1}$
	\fb	fb	\invfb	$fb^{-1}$	\ab	ab
	\ \invab	$ab^{-1}$	•		•	

#### 232 C.4.3 Time

	\sec	s	$\backslash \mathtt{ms}$	ms	$\backslash \mathtt{mus}$	$\mu s$
	\ns	ns	\ps	ps	\fs	fs
233	mhz	MHz	\khz	kHz	\hz	Hz
	\ \invps	$ps^{-1}$	\ \invns	$\mathrm{ns}^{-1}$	\yr	yr
	\hr	hr				

#### 234 C.4.4 Temperature

$$^{\rm 235}$$
 \degc  $^{\circ}{\rm C}$  \degk  ${\rm K}$ 

#### 236 C.4.5 Material lengths, radiation

	\Xrad	$X_0$	\NIL	$\lambda_{int}$	\mip	MIP
237	\neutroneq	$n_{eq}$	$\neq cmcm$	$n_{eq}/cm^2$	\kRad	kRad
	\MRad	MRad	\ci	Ci	\mci	mСi

#### 238 C.4.6 Uncertainties

#### 240 C.4.7 Maths

	\order	$\mathcal{O}$	\chisq	$\chi^2$	$\backslash \mathtt{chisqndf}$	$\chi^2/\mathrm{ndf}$
١	ackslashchisqip	$\chi^2_{ m IP}$	\chisqvs	$\chi^2_{ m VS}$	$\backslash \mathtt{chisqvtx}$	$\chi^2_{ m vtx}$
	ackslashchisqvtxndf	$\chi^2_{ m vtx}/{ m ndf}$	\deriv	d	$\backslash \mathtt{gsim}$	$\gtrsim$
241	$\backslash  exttt{lsim}$	$\lesssim$	$\mathbb{1} \operatorname{mean}[1]$	$\langle x \rangle$	$\abs[1] \abs[x]$	x
,	\Real	$\mathcal{R}e$	\Imag	$\mathcal{I}m$	\PDF	PDF
	\sPlot	sPlot				

#### 242 C.5 Kinematics

#### 243 C.5.1 Energy, Momenta

	$\backslash \mathtt{Ebeam}$	$E_{\scriptscriptstyle  m BEAM}$	\sqs	$\sqrt{S}$	\ptot	p
244	\pt	$p_{ m T}$	\et	$E_{ m T}$	$\backslash \mathtt{mt}$	$M_{ m T}$
	$\backslash \mathtt{dpp}$	$\Delta p/p$	$\backslash \mathtt{msq}$	$m^2$	\dedx	dE/dx

#### 245 C.5.2 PID

$$\label{eq:local_ppi} $$ \dlkpi \ DLL_{K\pi} $$ \dllppi \ DLL_{p\pi} $$ \dllepi \ DLL_{e\pi} $$$$

#### <sup>247</sup> C.5.3 Geometry

#### 249 C.5.4 Accelerator

$$_{250}$$
 \betastar  $eta^*$  \lum  ${\cal L}$  \intlum[1] \intlum[2 fb $^{-1}$  }  $\int {\cal L} = 2 ext{ fb}^{-1}$ 

## 251 C.6 Software

#### 252 C.6.1 Programs

	\bcvegpy	BCVEGPY	\boole	BOOLE	\brunel	Brunel
	\davinci	DAVINCI	$ackslash  ext{dirac}$	DIRAC	\evtgen	EVTGEN
	\fewz	Fewz	\fluka	FLUKA	\ganga	Ganga
,	\gaudi	Gaudi	\gauss	Gauss	$\backslash \mathtt{geant}$	Geant4
253	$\backslash \texttt{hepmc}$	НЕРМС	ackslashherwig	HERWIG	$\backslash \mathtt{moore}$	Moore
		NeuroBayes	$ackslash  exttt{photos}$	Photos	$\setminus$ powheg	Powheg
	$\protect\pro$	Рутніа	$ackslash{ ext{resbos}}$	ResBos	\roofit	RooFit
	\root	Root	\spice	SPICE	$\backslash \mathtt{urania}$	Urania

#### 254 C.6.2 Languages

	$\backslash cpp$	C++ SVN	$\ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ \ $	Ruby	$ackslash  extsf{fortran}$	FORTRAN
255	$\setminus$ svn	SVN				

#### 256 C.6.3 Data processing

	$ackslash  ext{kbytes}$	kbytes	ackslashkbsps	kbits/s	ackslashkbits	kbits
	ackslash kbsps	kbits/s	$ackslash { t mbsps}$	Mbytes/s	$ackslash { t mbytes}$	Mbytes
257	$\backslash \mathtt{mbps}$	Mbyte/s	\mbsps	Mbytes/s	\gbsps	Gbytes/s
	gbytes	Gbytes	\gbsps	Gbytes/s	\tbytes	Tbytes
	\tbpy	Tbytes/yr	ackslash dst	DST		

#### 258 C.7 Detector related

#### 259 C.7.1 Detector technologies

١	$\setminus$ nonn	$n^+$ -on- $n$	\ponn	$p^+$ -on- $n$	$\setminus nonp$	$n^+$ -on- $p$
260	\cvd	CVD	\mwpc	MWPC	\gem	$\operatorname{GEM}$

#### 261 C.7.2 Detector components, electronics

\	tell1	TELL1	\ukl1	UKL1	\beetle	Beetle
\	otis	OTIS	\croc	CROC	carioca	CARIOCA
\	dialog	DIALOG	\sync	SYNC	\cardiac	CARDIAC
\	\gol	GOL	ackslash vcsel	VCSEL	\ttc	TTC
\	ttcrx	TTCrx	\hpd	HPD	$\backslash \mathtt{pmt}$	PMT
262	specs	SPECS	$\backslash \mathtt{elmb}$	ELMB	\fpga	FPGA
\	plc	PLC	rasnik	RASNIK	\elmb	ELMB
\	can	CAN	\lvds	LVDS	\ntc	NTC
\	\adc	ADC	\led	LED	\ccd	CCD
\	hv	HV	\lv	LV	\pvss	PVSS
\	cmos	CMOS	\fifo	FIFO	\ccpc	CCPC

#### 263 C.7.3 Chemical symbols

	\cfourften	$C_4F_{10}$	ackslash cffour	$CF_4$	\cotwo	$CO_2$
264	\csixffouteen	$C_6F_{14}$	$\backslash \mathtt{mgftwo}$	$MgF_2$	\siotwo	$SiO_2$

## 265 C.8 Special Text

# D Supplementary material for LHCb-PAPER-20XX-YYY

This appendix contains supplementary material that will posted on the public cds record but will not appear in the paper.

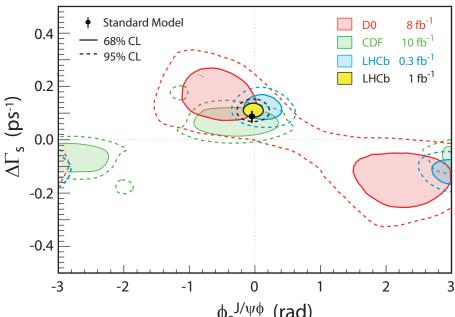
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Please leave the above sentence in your draft for first and second circulation and replace what follows by your actual supplementary material. For more information about other types of supplementary material, see Section ??. Plots and tables that follow should be well described, either with captions or with additional explanatory text.



 $\varphi_s^{J/\psi\varphi}$  (rad) Figure 1: Comparison of our result to those from other experiments. Note that the style of this figure differs slightly from that of Figure ??

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