Environmental Fluid Dynamics: Lecture 19

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Overview



- In the broadest sense, stability is a dividing line between laminar and turbulent flows.
- If a flow is stable, it can become or remain laminar.
- If a flow is unstable, it can become or remain turbulent.
- There are many stabilizing and destabilizing factors to consider.



- These factors are often terms in the TKE balance equation
- To simplify and understand the problem, a destabilizing term is often paired with a stabilizing term to form a dimensionless ratio.
- This ratio is used to determine which factor "wins" and whether the flow becomes turbulent or not.
- Examples include the Reynolds number, Richardson number, Rossby number, Rayleigh number, and Froude number.



- Recall that we previously discussed static stability.
- "Static" means no motion, so it described stability that is independent of the wind.
- Static stability determined whether a flow was capable of buoyant convection.
- Hence, we framed our discussion of static stability around vertical profiles of potential temperature.
- · Today we will discuss dynamic instability.



- "Dynamic" means having motion
- · Unlike static stability, dynamic stability depends on the wind
- Even if the atmosphere is statically stable in some layer, wind shear may be sufficient to generate turbulence
- One example is Kelvin-Helmholtz instability









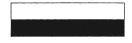






- Consider low density fluid on top of high density fluid (statically stable)
- Imagine that shear (ΔU) exists at their interface.
- If shear becomes large enough, the flow is dynamically unstable,
- The amplitude of the waves grows until they break.
- The breaking wave is called a Klevin-Helmholtz (KH) wave.
- The physics are different than a breaking wave on the ocean's surface, for instance.













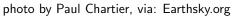


- Within each wave roll there are "packets" of statically unstable fluid (locally heavy above light fluid)
- The static and dynamic instabilities act to generate turbulence.
- The turbulence expands and leads to diffusion (mixing), which transfers momentum and reduces the shear.
- If the shear falls below the critical value, then the dynamic instability ends.
- In the absence of the shear, turbulence decays and the flow becomes laminar.



- KH waves likely occur often in statically-stable shear layers, although they are rarely observed visually.
- If sufficient moisture is present in the atmosphere, clouds can form in the ascending portions of the wave.
- These are called billow clouds.





- Interestingly, both static and dynamic instabilities cause the flow to react in a way to remove the instability.
- So, turbulence is an "effort" by the flow to undo the cause of the instability.
- Static: convection moves buoyant air upward, which stabilizes the flow.
- Dynamic: turbulence reduces wind shears, which stabilizes the flow.



• In other words, turbulence exists to end its own existance



- Since observations show prolonged periods of turbulence in the atmosphere, there must be external forces that act to destabilize the PBL.
- Static: solar heating of ground.
- Dynamic: pressure gradients from synoptic-scale features causes winds to fight dissipation.



- To understand when a flow might become dynamically unstable, we can compare the relative magnitudes of the shear and buoyancy terms in our TKE balance equation.
- Shear leads to mechanical production of turbulence a destabilizing effect.
- Buoyancy can either enhance or suppress turbulence so it can be stabilizing or destabilizing.
- One such ratio is the Richardson number Ri.



Richardson Number

- Consider a statically-stable environment.
- turbulent motions have to fight gravity, so buoyancy suppresses turbulence here.
- Conversely, wind shear acts to generate turbulence.
- The buoyant production/consumption term in the TKE balance is negative and the shear production term is positive.
- It is useful to examine their ratio.



 The ratio of buoyant production/consumption to shear production is called the Flux Richardson Number:

$$\operatorname{Ri}_{\mathrm{f}} = \frac{\left(\frac{g}{\overline{\theta_{v}}}\right)(\overline{w'\theta_{v}'})}{(\overline{u_{i}'u_{j}'})\frac{\partial \overline{u_{i}}}{\partial x_{j}}} = \frac{(\overline{w'b'})}{(\overline{u_{i}'u_{j}'})\frac{\partial \overline{u_{i}}}{\partial x_{j}}}$$

 The Richardson number is dimensionless and contains nine terms int he denominator!



• If we neglect subsidence ($\overline{w}=0$) and assume horizontal homogeneity, we arrive:

$$\operatorname{Ri}_{\mathrm{f}} = \frac{\left(\frac{g}{\overline{\theta_{v}}}\right)(\overline{w'\theta_{v}'})}{(\overline{u'w'})\frac{\partial \overline{u}}{\partial z} + (\overline{v'w'})\frac{\partial \overline{v}}{\partial z}} = \frac{(\overline{w'b'})}{(\overline{u'w'})\frac{\partial \overline{u}}{\partial z} + (\overline{v'w'})\frac{\partial \overline{v}}{\partial z}}$$

- \bullet For statically stable, unstable, neutral flows, ${\rm Ri}_{\rm f}$ is positive, negative, 0 respectively.
- Richardson suggested ${
 m Ri}_{
 m f}=1$ is a critical value since mechanical and buoyant terms are in balance.
- At any value less than $\mathrm{Ri_f}=1$, static stability is too weak to prevent mechanical generation of turbulence.
- For Ri_f < 1, buoyancy contributes to the generation of turbulence.

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\begin{cases} \mbox{Flow is dynamically unstable (turbulent)} & Ri_f < 1 \\ \mbox{Flow is dynamically stable (laminar)} & Ri_f > 1 \end{cases}
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 Note: statically unstable flow is by definition dynamically unstable.



- ullet With Ri_f , notice that our terms involve turbulent fluxes.
- So, it determines when a turbulent flow becomes laminar, but not when a laminar flow becomes turbulent.
- Using the idea of K-theory, we can say that $-\overline{w'\theta_v'}$ is proportional to $\partial \overline{\theta_v}/\partial z$, $-\overline{u'w'}$ is proportional to $\partial \overline{u}/\partial z$, and $-\overline{v'w'}$ is proportional to $\partial \overline{v}/\partial z$.
- We can use these assumptions to create a new Richardson number.



• The **Gradient Richardson Number** is given by:

$$\operatorname{Ri} = \frac{\frac{g}{\overline{\theta_v}} \left(\frac{\partial \overline{\theta_v}}{\partial z} \right)}{\left(\frac{\partial \overline{u}}{\partial z} \right)^2 + \left(\frac{\partial \overline{v}}{\partial z} \right)^2}$$

 \bullet If you see ${\rm R}i$ without any subscript - the authors probably mean the gradient Richardson number.



• Theoretical and experimental data suggest that laminar flow becomes turbulent when Ri is below a critical value Ri_c , while turbulent flow ceases at some termination value Ri_t .

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\begin{cases} \text{Laminar flow becomes turbulent} & \mathrm{Ri} < \mathrm{Ri_c} \\ \text{Turbulent flow becomes laminar} & \mathrm{Ri} > \mathrm{Ri_t} \end{cases}
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- There is still debate about what values should be used, but typically ${\rm Ri}_c=0.25$ and ${\rm Ri}_t=1.$
- \bullet There is a hysteresis since $\mathrm{Ri}_{\mathrm{t}}>\mathrm{Ri}_{\mathrm{c}}$



Why the hysteresis?

- The idea is that we need two conditions for turbulence: instability and a trigger mechanism.
- Consider KH waves as a trigger mechanism.
- In the absence of existing turbulence, Ri must fall well below Ri_{t} in order for KH waves to form.
- \bullet Experimental data suggests KH waves form when $\mathrm{Ri} < \mathrm{Ri}_c.$



Why the hysteresis?

- Thus, we have a hysteresis.
- In other words, the Richardson number of a non-turbulent flow must be lowered to Ri_{c} , while a turbulent flow can remain so until the Richardson number surpasses Ri_{t} .



- ${\rm Ri_c}\approxeq 0.25$ is theoretically based on local gradients of wind and temperature.
- In the real world, we don't really know those local gradients.
- However, we can approximate them using discrete height layers.
- For example, we can approximate $\partial \overline{\theta_v}/\partial z$ as $\Delta \overline{\theta_v}/\Delta z$.
- We can apply these approximations to the gradient Richardson number to create a new ratio.



• The Bulk Richardson Number is given by:

$$\mathrm{Ri_B} = \frac{\frac{g}{\overline{\theta_v}} \frac{\Delta \overline{\theta_v}}{\Delta z}}{\left(\frac{\Delta \overline{u}}{\Delta z}\right)^2 + \left(\frac{\Delta \overline{v}}{\Delta z}\right)^2} = \frac{g\Delta \overline{\theta_v} \Delta z}{\overline{\theta_v} \left[(\Delta \overline{u})^2 + (\Delta \overline{v})^2\right]}$$

- This is the form most often used by meteorologists.
- Observational and NWP data give wind and temperature at discrete points in the vertical.
- Note: $\Delta \overline{u} = \overline{u}(\mathsf{top}) \overline{u}(\mathsf{bottom})$



- ullet A few words of caution with ${
 m Ri}_{
 m B}$ are warranted.
- The critical value is based on local gradients and not finite differences across layers.
- Gradients likely get "washed out" as the layer grows in thickness.
- Uncertainty arises in our ability to determine whether turbulence might form.
- \bullet Since some observations are spread far apart in the vertical, we must take care in interpreting the computed Ri_B when trying to characterize the considered flow.

