Geometric and Combinatorial Properties of Self-similar Multifractal Measures

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ABSTRACT. For any self-similar measure μ in $\mathbb{R}$, we show that the distribution of μ is controlled by products of non-negative matrices governed by a finite or countable graph depending only on the IFS. This generalizes the net interval construction of Feng from the equicontractive finite type case. When the measure satisfies the weak separation condition, we prove that this directed graph has a unique attractor. This allows us to verify the multifractal formalism for restrictions of μ to certain compact subsets of $\mathbb{R}$, determined by the directed graph. When the measure satisfies the generalized finite type condition with respect to an open interval, the directed graph is finite and we prove that if the multifractal formalism fails at some $q \in \mathbb{R}$, there must be a cycle with no vertices in the attractor. As a direct application, we verify the complete multifractal formalism for an uncountable family of IFSs with exact overlaps and without logarithmically commensurable contraction ratios.

CONTENTS

This paper is in final form and no version of it will be submitted for publication elsewhere.

1. INTRODUCTION

Self-similar measures in $\mathbb{R}$ are perhaps the simplest examples of measures which exhibit complex local structure. These measures are associated with finite sets of similarity maps in $\mathbb{R}$. To be precise, by an *iterated function system of similarities* (IFS) we mean a finite set of maps $\{S_i\}_{i\in\mathcal{I}}$ where each $S_i(x) = r_i x + d_i$ and $0 < |r_i| < 1$. The *attractor*, or *self-similar set*, of this system is the unique compact set K satisfying $\bigcup_{i\in\mathcal{I}} S_i(K) = K$. Given a probability vector $\boldsymbol{p} = (p_i)_{i\in\mathcal{I}}$ where each $p_i > 0$ and $\sum_i p_i = 1$, the associated *self-similar measure* is the unique Borel probability measure satisfying

$$\mu_{\boldsymbol{p}}(E) = \sum_{i\in\mathcal{I}} p_i \mu_{\boldsymbol{p}} \circ S_i^{-1}(E)$$

for any Borel set $E \subseteq \mathbb{R}$. For a more through discussion of the background and basic properties of self-similar sets and measures, we refer the reader to Falconer's book [6].

In order to understand the general structure of the measure $\mu_{\boldsymbol{p}}$ or the self-similar set K, one often considers basic dimensional quantities such as the Hausdorff dimension $\dim_H K$ and analgous statements for measures, or other notions of dimension. Computing these values can be highly non-trivial for general iterated function systems of similarities and there is significant literature on this matter (see, for example, [2, 11, 15, 22, 25, 28, 31, 34]). In this paper, we focus on a more fine-grained notion of dimension known as the local dimension. Given a point $x \in K = \operatorname{supp} \mu_{\boldsymbol{p}}$, the *local dimension* is given by

$$\dim_{\operatorname{loc}} \mu_{\boldsymbol{p}}(x) = \lim_{t\to 0} \frac{\log \mu_{\boldsymbol{p}}(B(x,t))}{\log t},$$

when the limit exists. From the perspective of multifractal analysis, one is interested in determining geometric properties of the sets $K(\alpha) := \{x \in K : \dim_{\operatorname{loc}} \mu_{\boldsymbol{p}}(x) = \alpha\}$.

On the other hand, the L^q-*spectrum* of $\mu_{\boldsymbol{p}}$ is given by

$$\tau(\mu_{\boldsymbol{p}}, q) = \tau(q) := \liminf_{t \to 0} \frac{\log \sup \sum_i \mu_{\boldsymbol{p}}(B(x_i, t))^q}{\log t}$$

for each $q \in \mathbb{R}$, where the supremum is over disjoint families of closed balls with centres $x_i \in K$.

An important objective of multifractal analysis is to understand the relationship between the L^q-spectrum of the measure $\mu_{\boldsymbol{p}}$, and the *dimension spectrum* $\dim_H K(\alpha)$. A heuristic relationship between $\tau(q)$ and $\dim_H K(\alpha)$, known as the multifractal formalism, was introduced by Halsey *et al.* [16]. The *multifractal formalism* states, roughly speaking, that the dimension spectrum can be computed as the concave conjugate of $\tau(q)$, i.e.

$$\dim_H K(\alpha) = \tau^*(\alpha) := \inf_{q \in \mathbb{R}} \{q\alpha - \tau(q)\}$$

for any α in the domain of $\tau^*(\alpha)$; see Definition 4.2 for a complete definition in our setting. This concave conjugate relationship has been studied by many authors (see, for example, [3, 7, 10, 12, 13, 16, 26, 27, 32, 33, 36]). As a particularly elegant example, it has been verified in general for iterated function systems satisfying the strong separation condition ($S_i(K) \cap S_j(K) \neq \emptyset$ if and only if $i = j$) [3]. This separation requirement has been relaxed to the open set condition [24] and the concave conjugate relationship has been verified [1, 32, 33]. In both cases, $\tau(q)$ is differentiable for all $q \in \mathbb{R}$ and is determined uniquely by the implicit formula $\sum_{i \in \mathcal{I}} p_i^q r_i^{-\tau(q)} = 1$.

However, when the open set condition fails, outside specialized analysis of some families of examples (for example, Bernoulli convolutions associated with the unique positive root of the polynomial $x^k - x^{k-1} - \cdots - x - 1$ [9]), there has been much less progress on verifying the multifractal formalism at all $q \in \mathbb{R}$. For $q > 0$, the function $x \mapsto x^q$ is increasing so the summation in the definition of $\tau(q)$ is dominated by closed balls with large measure. On the other hand, for $q < 0$, the summation is dominated by closed balls of small measure. Generally speaking, understanding the multifractal analysis of measures when $q < 0$ is substantially more challenging than the case $q \geq 0$. Gaining more information about this case is our focus in this document.

1.1. **The weak separation condition.** Notably, neither the strong separation condition nor the open set condition allows for the existence of exact overlaps. We introduce some notation: let $\mathcal{I}^*$ denote the set of all finite words on $\mathcal{I}$. For $\sigma = (i_1, \ldots, i_n) \in \mathcal{I}^*$, write $S_\sigma = S_{i_1} \circ \cdots \circ S_{i_n}$, $r_\sigma = r_{i_1} \cdots r_{i_n}$ and, if $n \geq 1$, $\sigma^- = (i_1, \ldots, i_{n-1})$. By *exact overlaps* we mean the existence of words $\sigma \neq \tau \in \mathcal{I}^*$ such that $S_\sigma = S_\tau$. To study examples allowing exact overlaps while still maintaining separation of non-overlapping words, Lau and Ngai introduced the weak separation condition and studied basic conditions under which the multifractal formalism holds [27]. For any $t > 0$ and Borel set $E \subseteq \mathbb{R}$, define

$$\Lambda_t(E) = \{\sigma \in \mathcal{I}^* : r_\sigma < t \leq r_{\sigma^-}, S_\sigma(K) \cap E \neq \emptyset\}.$$

Then the *weak separation condition* (WSC) is equivalent to requiring that

$$(1.1) \qquad \sup_{x \in \mathbb{R}, t > 0} \#\{S_\sigma : \sigma \in \Lambda_t(U(x, t))\} < \infty$$

where $\#X$ denotes the cardinality of a set X and $U(x, t)$ is the open ball about x with radius t. Note that the definition only considers functions S_σ rather than the words σ so as to allow exact overlaps. To see an equivalent formulation with respect to exact overlaps or the equivalence with the original definition of Lau and Ngai, see [38, Thm. 1].

Under the weak separation condition, verification of the multifractal formalism is subtle. One of the earliest examples of exceptional behaviour is with respect to self-similar measures of the system of Bernoulli convolutions $\{x \mapsto \rho x, x \mapsto \rho x + (1 - \rho)\}$ where the contraction ratio ρ is the reciprocal of the golden mean. In this case, the L^q-spectrum $\tau(q)$ has a *phase transition*, or a point where $\tau(q)$ is not differentiable. Nevertheless, the multifractal formalism still holds and $\tau(q)$ is analytic for other values of q [9]. Another example of exceptional behaviour is the 3-fold convolution of the uniform Cantor measure. In this case, it was observed that the set of attainable local dimensions is not an interval and the multifractal formalism fails [23]. The problem here is, in some sense, that the measure $\mu_{\boldsymbol{p}}$ is too small at certain points in K. This measure, and other related measures, were studied in detail [13, 20, 30, 36] and a modified multifractal formalism was proven therein. In these cases, the failure occurs at some point $q < 0$.

In an important paper, Feng and Lau [12] obtain deep results about the multifractal formalism under the weak separation condition. Using a subtle Moran construction [14], they prove that the multifractal formalism holds for any value $q \geq 0$, and for $q < 0$, they give a modified multifractal formalism by considering suitable restrictions to an open ball U_0 which attains the supremum in the definition of the WSC (1.1). Unfortunately, this result does not directly give information on the validity of the multifractal formalism for values $q < 0$. In some sense, the restriction avoids the breakdown of the multifractal formalism by avoiding points in K where the measure is too small.

To extend this perspective, we develop some new ideas. Even in regions where the overlap is not dense (i.e. away from any maximal open ball U_0), through a general graph construction, we will show that the measure may be "combinatorially linked" to regions with high density where the multifractal formalism holds. For example, consider the IFS given by the maps

$$(1.2) \qquad S_1(x) = \rho x \qquad S_2(x) = rx + \rho(1 - r) \qquad S_3(x) = rx + 1 - r$$

where $\rho > 0$, $r > 0$ satisfy $\rho + 2r - \rho r \leq 1$. This IFS was first studied by Lau and Wang [29] and satisfies the weak separation condition. In Section 5.3.3, we show that the maximal open sets U_0 can never contain the point 1 in the self-similar set, which is a phenomenon similar to the situation of the Cantor convolution. Despite this, we can prove (as a consequence of our more general results) that the multifractal

formalism still holds for the measure $\mu_{\boldsymbol{p}}$, without restriction to a subset and with any probabilities. Our main goal in this paper is to provide a new, natural perspective for understanding the failure of the multifractal formalism, and to provide combinatorial conditions under which the multifractal formalism holds or in which one might suspect that fails.

Our starting point is the net interval construction of Feng [8]. In that document, for iterated function systems of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ with $0 < r < 1$ satisfying a combinatorial overlap condition known as the finite type condition [31], he obtains formulas for the values of $\mu_{\boldsymbol{p}}(\Delta)$ on families of intervals $\mathcal{F}_n$ as products of non-negative matrices. He then uses properties of matrix products to verify differentiability of the L^q-spectrum (and thus the multifractal formalism by the prior work of Lau and Ngai [27]) for values $q > 0$. Using some different perspectives but with the same underlying approach, he proves a modified multifractal formalism for values of $q < 0$ [10].

In recent work, following the techniques of Feng and operating in the same setting, Hare, Hare, and other collaborators [17, 18] define a finite graph called the transition graph corresponding to the IFS. Then they determine that the set of local dimensions at special points in K called interior essential points form a closed interval, and show that the failure for the set of local dimensions to be a closed interval is determined by the existence of certain combinatorial structures in the transition graph called non-essential loop classes.

However, as observed by Testud [37], when the IFS does not have a common contraction ratio or a similar property (for example, $\log r_i / \log r_j \in \mathbb{Q}$ for all i, j [21]), one cannot apply Feng's net interval construction in a natural way.

1.2. **Summary of main results.** Our first contribution is a generalization of the net interval construction to apply to any IFS of similarities. We determine that the distribution of $\mu_{\boldsymbol{p}}$ on certain intervals which we call *net intervals* is determined by a local overlap structure which we call the *neighbour set* of the net interval (see [19] for the first appearance of this construction). Our first key observations, Lemma 2.3 and Theorem 2.8, are that the neighbour set completely determines the local geometry of the attractor K and the distribution of the measure $\mu_{\boldsymbol{p}}$ (up to fixed constants of comparability). This allows us in Section 2.4 to construct a countable directed graph which we call the *transition graph* of the IFS, where the vertices are the distinct neighbour sets. Then in Section 2.5, we associate to each edges of the transition graph a non-negative matrices called a *transition matrix* such that the distribution of $\mu_{\boldsymbol{p}}$ on net intervals is given by products of these non-negative matrices. Since we do not make any assumptions on the contraction ratios, we introduce two simple but important ideas: the notion of the *transition generation* (Definition 2.4), and the notion of the *length of an edge* (Definition 2.9). These definitions resolve the issues with the original net interval construction recognized above.

In Section 3, we turn our attention to the IFSs satisfying the WSC. In particular, we prove the existence of a relatively open subset $K_{\mathrm{ess}} \subseteq K$ called the *set of interior*

essential points, and a corresponding subgraph of the transition graph called the *essential class* on which the self-similar measure has certain important regularity properties (Lemma 3.9). We call a net interval *essential* if its neighbour set is a vertex in the essential class. We determine that the set of interior essential points is large in two different senses:

Theorem 1.1. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition.*
(i) If U_0 is any open set which attains the maximality in (1.1), then U_0 is contained in a finite union of essential net intervals. In particular, $U_0 \subseteq K_{\mathrm{ess}}$.
(ii) If $\mu_{\boldsymbol{p}}$ is any associated self-similar measure, then $\mu_{\boldsymbol{p}}(K \setminus K_{\mathrm{ess}}) = 0$.

See Proposition 3.7 and Theorem 3.10 for proofs of these facts.

We also obtain dimensional results at certain points in K called *periodic points*, an idea introduced by Hare, Hare, and Matthews. In Proposition 3.15, we prove that an elegant formula holds for the local dimensions at such points, and in Theorem 4.1 we show that the sets of local dimensions at periodic points are dense in the sets of upper and lower local dimensions at points in K_{ess}. This generalizes a pre-existing result [18, Cor. 3.15] to the weak separation case.

We then focus on understanding the multifractal formalism from the perspective of the essential class. Our main result in this section is the following (see Theorem 4.5 for a complete statement and proof):

Theorem 1.2. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition and let $\mu_{\boldsymbol{p}}$ be an associated self-similar measure. Let $\Delta_1, \ldots, \Delta_n$ be any essential net intervals and let $\nu = \mu_{\boldsymbol{p}}|_{\Delta_1 \cup \cdots \cup \Delta_n}$. Then ν satisfies the multifractal formalism and*

$$(1.3) \qquad \{\dim_{\mathrm{loc}} \nu(x) : x \in \mathrm{supp}\, \nu\} = \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}.$$

Moreover, the values of $\tau(\nu, q)$ do not depend on the choice of $\Delta_1, \ldots, \Delta_n$ and for $q \geq 0$, $\tau(\mu_{\boldsymbol{p}}, q) = \tau(\nu, q)$.

Our verification of this modified multifractal formalism begins with [12, Thm 1.2], but then uses the matrix product structure of the transition graph to move the weight of the measure from the sets U_0 to any net interval in the essential class. We note some minor improvements: rather than considering restrictions of the L^q-spectrum to an open set, we obtain the results as a restriction to a compact subset $\Delta_1 \cup \cdots \cup \Delta_n$, where this subset can strictly contain any open set U_0 attaining the maximum in (1.1). In some sense, this says that the measure $\mu_{\boldsymbol{p}}$ is also sufficiently regular near the boundary of U_0.

In fact, our matrix product structure provides a more general perspective for understanding the quasi-product property of Feng and Lau [12]; a natural analogue holds in our setting where their set Ω is replaced by a set of net intervals which have the same neighbour set and determined by a fixed essential net interval. As a result, a more direct proof of Theorem 1.2 is possible. However, many details of this proof overlap with the approach of Feng and Lau, so we do not include this approach.

Combining this result with Theorem 1.1, we prove the following modified multifractal formalism for any IFS satisfying the weak separation condition:

Corollary 1.3. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition with associated self-similar measure $\mu_{\boldsymbol{p}}$. Then there exists a sequence of compact sets $(K_m)_{m=1}^{\infty}$ with $K_m \subseteq K_{m+1} \subseteq K$ for each $m \in \mathbb{N}$ such that*
  (i) $\lim_{m \to \infty} \mu_{\boldsymbol{p}}(K_m) = 1$,
  (ii) each $\mu_m := \mu_{\boldsymbol{p}}|_{K_m}$ satisfies the multifractal formalism, and
  (iii) $\tau(\mu_m, q)$ and $D(\mu_m)$ do not depend on the index m.

We note the similarity of this result to a result of Feng [10, Thm. 1.2], which follows from general results about the multifractal formalism of certain matrix-valued functions satisfying an irreducibilty condition. However, the techniques used therein only apply naturally in the finite type case for IFSs of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$.

We also obtain the following important corollary:

Corollary 1.4. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition with transition graph $\mathcal{G}$. Suppose there is a bound on the maximum length of a path with no vertices in the essential class. Then if $\mu_{\boldsymbol{p}}$ is any associated measure, $\mu_{\boldsymbol{p}}$ satisfies the multifractal formalism.*

In particular, suppose $\mathcal{G}$ is finite. In this situation, the only mechanism for the failure of the multifractal formalism is the existence of a cycle (a path in the transition graph which begins and ends at the same vertex) which is not contained in the essential class. This gives a combinatorial condition which guarantees that the multifractal formalism holds. In this situation, it is possible to write a finite algorithm to determine whether such a cycle exists.

In particular, in Theorem 5.7, we apply this to the family of IFS defined in (1.2):

Corollary 1.5. *Let $\{S_i\}_{i=1}^{3}$ be the IFS defined in (1.2). Then for any probability weights $\boldsymbol{p} = (p_i)_{i=1}^{3}$, the associated self-similar measure $\mu_{\boldsymbol{p}}$ satisfies the complete multifractal formalism.*

To the best knowledge of the author, this is the first example of an IFS with exact overlaps and without logarithmically commensurable contraction ratios for which the complete multifractal formalism is proven to hold. Understanding failure of the multifractal formalism is based critically on understanding the properties of cycles in the transition graph outside the essential class.

By combining our results with the work of Deng and Ngai [4], we can also gain information about differentiability of the L^q-spectrum in this case. In a slightly specialized case, [4, Thm. 1.2] states that, for probabilities $p_2 > p_3$,

$$f(\alpha) := \dim_H \{x \in K : \dim_{\text{loc}} \mu_{\boldsymbol{p}}(x) = \alpha\}$$

is the concave conjugate of a differentiable function. Combining this with Corollary 1.4 and involutivity of concave conjugation, we obtain the following result:

Corollary 1.6. *Let $\{S_i\}_{i=1}^{3}$ be the IFS defined in (1.2). Then if $p_2 > p_3$, the L^q-spectrum $\tau(\mu_{\boldsymbol{p}}, q)$ is differentiable for any $q \in \mathbb{R}$.*

This answers some of the questions raised in [4].

Finally, in Section 5, we investigate some specific families of IFSs to illustrate these results; notably, we give an in-depth analysis of the IFS given in (1.2). In fact, every example in that section has a finite transition graph: this is equivalent to the generalized finite condition of Lau and Ngai [28] holding with respect to an open interval (see [19, Thm. 3.4] and Remark 5.2 for a proof). Moreover, when K is a convex set, a recent result gives that the weak separation condition is equivalent to the finiteness of the transition graph [19, Thm. 4.4]. In general, the author believes this to be true without any convexity assumption on K:

Conjecture 1.7. *Let $\{S_i\}_{i\in\mathcal{I}}$ be an IFS in $\mathbb{R}$ with transition graph $\mathcal{G}$. Then $\{S_i\}_{i\in\mathcal{I}}$ satisfies the weak separation condition if and only if $\mathcal{G}$ is finite.*

The results obtained in this paper under the weak separation condition, and the similar strength to results proven under various finite type conditions, provides some more evidence towards this equivalence in general.

1.3. **Limitations and future work.** We note here that the result Corollary 1.4 is not a dichotomy. While the non-existence of cycles outside the transition graph guarantees that the multifractal formalism holds, the converse need not hold. We have examples of measures satisfying the open set condition (with respect to an open set that is not an open interval) with cycles outside the essential class, while the open set condition guarantees that the multifractal formalism does hold. This situation is likely a by-product of the net interval construction, since our perspective is always with respect to images of the entire interval $[0,1]$. However, there are also cases such as the Bernoulli measure associated with the IFS $\{x \mapsto \rho x, x \mapsto \rho x + (1-\rho)\}$ where $1/\rho$ is the Golden mean. In this situation, the attractor is the entire interval $[0,1]$ so that the net interval construction is a natural choice. Here, even though the L^q-spectrum contains a point of non-differentiability at some $q_0 < 0$ and contains a cycle not contained in the essential class, the measure still satisfies the multifractal formalism [9]. This phenomenon, and other related special cases, is studied in recent work of Hare, Hare, and Shen [20].

More work is needed to address the general case. In a work in preparation, the author investigates the multifractal analysis of measures when the transition graph is finite to provide a more detailed understanding of such examples. In particular, we obtain a greater understanding of the multifractal formalism outside the essential class as a continuation of our analysis here.

1.4. **Notational conventions.** We briefly mention here some of the conventions we use through out the document. Given any set X, we write $\#X$ to denote the cardinality of X. The set $\mathbb{R}$ is always the metric space equipped with the usual Euclidean metric. The set $\mathbb{N}$ is the set of natural numbers beginning at 1. The set $B(x,t)$ is always a closed ball about x with radius t, and $U(x,t)$ denotes the open ball. Let

$E, F \subseteq \mathbb{R}$ be Borel sets. We denote by $\mathrm{diam}(E) = \sup\{|x - y| : x, y \in E\}$ and $\mathrm{dist}(E, F) = \inf\{|x - y| : x \in E, y \in F\}$. By E°, we mean the topological interior of E.

Boldface quantities are typically vectors. If M is a square matrix, we denote by $\mathrm{sp}(M)$ the spectral radius of M. If $\boldsymbol{v}, \boldsymbol{w}$ are vectors with the same dimension, we write $\boldsymbol{v} \preccurlyeq \boldsymbol{w}$ if $\boldsymbol{v}_i \leq \boldsymbol{w}_i$ for each i. All matrices in this document are non-negative.

Given families of real numbers $(a_i)_{i \in I}$ and $(b_i)_{i \in I}$, we write $a_i \asymp b_i$ if there exist constants $c_1, c_2 > 0$ such that $c_1 a_i \leq b_i \leq c_2 a_i$ for all $i \in I$.

The maps $\{S_i\}_{i \in \mathcal{I}}$ always denotes an iterated function system. We assume that $\#\mathcal{I} \geq 2$ and its attractor K is not a singleton. Sets denoted by Δ are closed intervals and often net intervals. Indices s, t are used to refer to generations and radii of open and closed balls. Greek letters $\sigma, \tau, \omega, \phi, \xi$ typically refer to words in $\mathcal{I}^*$. The Greek η typically refers to a path in the transition graph. The character T refers to either a transition matrix or, more occasionally, a similarity map, depending on context.

1.5. **Acknowledgements.**

2. ITERATED FUNCTION SYSTEMS THROUGH NET INTERVALS

2.1. **Iterated function systems of similarities in** $\mathbb{R}$.

Let $\mathcal{I}$ be a non-empty finite index set. By an iterated function system of similarities (IFS) $\{S_i\}_{i \in \mathcal{I}}$ we mean a finite set of similarities

$$(2.1) \qquad S_i(x) = r_i x + d_i : \mathbb{R} \to \mathbb{R} \text{ for each } i \in \mathcal{I}$$

with $0 < |r_i| < 1$.

Each IFS generates a unique non-empty compact set K satisfying

$$K = \bigcup_{i \in \mathcal{I}} S_i(K).$$

This set K is known as the associated *self-similar set*. Throughout, we will assume K is not a singleton. By rescaling and translating the d_i if necessary, without loss of generality we may assume the convex hull of K is $[0, 1]$.

Given a probability vector $\boldsymbol{p} = (p_i)_{i \in \mathcal{I}}$ where $p_i > 0$ and $\sum_{i \in \mathcal{I}} p_i = 1$, there exists a unique Borel measure $\mu_{\boldsymbol{p}}$ with $\mathrm{supp}\,\mu_{\boldsymbol{p}} = K$ satisfying

$$(2.2) \qquad \mu_{\boldsymbol{p}}(E) = \sum_{i \in \mathcal{I}} p_i \mu_{\boldsymbol{p}}(S_i^{-1}(E))$$

for any Borel set $E \subseteq K$. This measure $\mu_{\boldsymbol{p}}$ as known as an associated self-similar measure.

Let $\mathcal{I}^*$ denote the set of all finite words on $\mathcal{I}$. Given $\sigma = (\sigma_1, \ldots, \sigma_j) \in \mathcal{I}^*$, we denote

$$\sigma^- = (\sigma_1, \ldots, \sigma_{j-1}), S_\sigma = S_{\sigma_1} \circ \cdots \circ S_{\sigma_j} \text{ and } r_\sigma = r_{\sigma_1} \cdots r_{\sigma_j}.$$

Given $t > 0$, put

$$\Lambda_t = \{\sigma \in \mathcal{I}^* : |r_\sigma| < t \le |r_{\sigma^-}|\}.$$

We refer to the set of $\sigma \in \Lambda_t$ as the *words of generation t*. We remark that in the literature it is more common to see this defined by the rule $|r_\sigma| \le t < |r_{\sigma^-}|$. The two choices are essentially equivalent, but this choice is more convenient for our purposes.

2.2. Neighbour sets.

The notions of net intervals and neighbour sets were introduced in [8] and [21]. In [19], these notions were extended to an arbitary IFS, and we present those definitions here. We then continue the discussion to define the children of a net interval, and show in Theorem 2.8 that the children depend only on the neighbour set of the parent.

Let $h_1, \ldots, h_{s(t)}$ be the collection of distinct elements of the set $\{S_\sigma(0), S_\sigma(1) : \sigma \in \Lambda_t\}$ listed in strictly ascending order; we refer to this set as the *endpoints of generation t*. Set

$$\mathcal{F}_t = \{[h_j, h_{j+1}] : 1 \le j < s(t) \text{ and } (h_j, h_{j+1}) \cap K \neq \emptyset\}.$$

Elements of $\mathcal{F}_t$ are called *net intervals of generation t*. Write $\mathcal{F} = \bigcup_{t>0} \mathcal{F}_t$ to denote the set of all possible net intervals.

Suppose $\Delta \in \mathcal{F}$. We denote by T_Δ the unique contraction $T_\Delta(x) = rx + a$ with $r > 0$ such that

$$T_\Delta([0, 1]) = \Delta.$$

Of course, $r = \mathrm{diam}(\Delta)$ and a is the left endpoint of Δ.

Definition 2.1. *We will say that a similarity $f(x) = Rx + a$ is a* neighbour *of $\Delta \in \mathcal{F}_t$ if there exists some $\sigma \in \Lambda_t$ such that $S_\sigma(K) \cap \Delta^\circ \neq \emptyset$ and $f = T_\Delta^{-1} \circ S_\sigma$. In this case, we also say that S_σ* generates *the neighbour f. The* neighbour set *of Δ is the maximal set*

$$\mathcal{V}_t(\Delta) = \{f_1, \ldots, f_m\}$$

where each $f_i = T_\Delta^{-1} \circ S_{\sigma_i}$ is a distinct neighbour of Δ.

Since $K = \bigcup_{\sigma \in \Lambda_t} S_\sigma(K)$, we easy to see that every net interval has a non-empty neighbour set.

Note that if σ generates a neighbour of Δ, then $S_\sigma([0, 1]) \supseteq \Delta$. When the generation of Δ is implicit, we will simply write $\mathcal{V}(\Delta)$. For notational convenience, we define the quantity $R_{\max}(\Delta) = \max\{|R| : \{x \mapsto Rx + a\} \in \mathcal{V}(\Delta)\}$, which depends only on $\mathcal{V}(\Delta)$.

Remark 2.2. *For an IFS of the form $\{S_i(x) = rx + d_i\}_{i \in \mathcal{I}}$ where $0 < r < 1$ is fixed, the notion of a neighbour set is related to the characteristic vector of Feng [8]. We describe the equivalence here.*

Let $\Delta = [a, b] \in \mathcal{F}_t$ be some net interval and let n be such that $r^n < t \le r^{n-1}$. Let $\sigma_1, \ldots, \sigma_m$ generate distinct neighbours of Δ, so that $r_{\sigma_i} = r^n$ for each $1 \le i \le m$. Then the (reduced) characteristic vector of Δ (see [8, Sec. 2] for notation) is determined by

$$\ell_n(\Delta) = r^{-n} \mathrm{diam}(\Delta) \qquad V_n(\Delta) = \{r^{-n}(a - S_{\sigma_i}(0)) : 1 \le i \le m\}.$$

whereas the neighbour set of Δ is given by

$$\mathcal{V}(\Delta) = \{T_\Delta^{-1} \circ S_{\sigma_i}\} = \{x \mapsto \frac{S_{\sigma_i}(x) - a}{\operatorname{diam}(\Delta)}\}$$

$$= \{x \mapsto \frac{x}{r^{-n} \operatorname{diam}(\Delta)} + \frac{S_{\sigma_i}(0) - a}{\operatorname{diam}(\Delta)}\}.$$

Thus, when the IFS has a common positive contraction ratio, our neighbour set construction can be interpreted directly as a normalized version of Feng's characteristic vector.

When the IFS has arbitrary contraction ratios, there is no clear choice of normalization factor analgous to $\ell_n(\Delta)$ that is uniform across all net intervals $\Delta \in \mathcal{F}_t$. This issue is resolved by normalizing directly by $\operatorname{diam}(\Delta)$, but now it is no longer clear how to define the children of a net interval in a global way. Instead, a local definition for the children of net intervals, and the analogue of [8, Lem. 2.1], are given in Section 2.3.

Neighbour sets of net intervals are relevant in the sense that they completely determine the local geometry of K in the net interval, as well as the behaviour of associated self-similar measures on Borel subsets of the net interval. To be precise, we have the following lemma:

Lemma 2.3. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS as in (2.1) with attractor K and associated self-similar measure $\mu_{\boldsymbol{p}}$. Suppose Δ_1, Δ_2 are net intervals with $\mathcal{V}(\Delta_1) = \mathcal{V}(\Delta_2)$. Then there exists a surjective similarity $g : \Delta_1 \cap K \to \Delta_2 \cap K$ and constants $c_1, c_2 > 0$ such that if $E \subseteq \Delta_1$ is any Borel set,*

$$c_1 \mu_{\boldsymbol{p}}(E) \leq \mu_{\boldsymbol{p}}(g(E)) \leq c_2 \mu_{\boldsymbol{p}}(E).$$

Proof. By definition of the neighbour set, if Δ is any net interval, we have

$$\Delta \cap K = \bigcup_{f \in \mathcal{V}(\Delta)} (T_\Delta \circ f(K)) \cap \Delta.$$

Set $g = T_{\Delta_2} \circ T_{\Delta_1}^{-1}$ so that g is clearly a similarity, and applying this observation to Δ_1 and Δ_2, we have

$$g(\Delta_1 \cap K) = \bigcup_{f \in \mathcal{V}(\Delta_1)} g(T_{\Delta_1} \circ f(K) \cap \Delta_1) = \bigcup_{f \in \mathcal{V}(\Delta_1)} (g \circ T_{\Delta_1} \circ f(K)) \cap g(\Delta_1)$$

$$= \bigcup_{f \in \mathcal{V}(\Delta_2)} (T_{\Delta_2} \circ f(K)) \cap \Delta_2 = \Delta_2 \cap K.$$

Thus g is a surjective with the correct image.

We now verify the measure property. Now by the invariant property of the self-similar measure (2.2), if $\Delta \in \mathcal{F}_t$ is any net interval and $E \subseteq \Delta$ is any Borel set,

$$\mu_{\boldsymbol{p}}(E) = \sum_{\sigma \in \Lambda_t} p_\sigma \mu_{\boldsymbol{p}} \circ S_\sigma^{-1}(E) = \sum_{f \in \mathcal{V}(\Delta)} \mu_{\boldsymbol{p}}\big(f^{-1} \circ T_\Delta^{-1}(E)\big) \sum_{\substack{\sigma \in \Lambda_t \\ \sigma \text{ generates } f}} p_\sigma.$$

Since f is a neighbour of Δ, there is at least one σ generating f. In particular, say $\Delta_1 \in \mathcal{F}_{t_1}$ and $\Delta_2 \in \mathcal{F}_{t_2}$, write $\mathcal{V}(\Delta_1) = \mathcal{V}(\Delta_2) = \{f_1, \ldots, f_m\}$, and set for each $1 \leq i \leq m$ and $j = 1, 2$

$$q_{i,j} := \sum_{\substack{\sigma \in \Lambda_{t_j} \\ \sigma \text{ generates } f_i}} p_\sigma > 0.$$

Set $c_1 = \min\{q_{i,2}/q_{i,1} : 1 \leq i \leq m\}$. We then have for $E \subseteq \Delta_1$ that $g(E) \subseteq \Delta_2$ so that

$$\mu_{\boldsymbol{p}}(g(E)) = \sum_{i=1}^m \mu_{\boldsymbol{p}}\big(f_i^{-1} \circ T_{\Delta_2}^{-1} \circ g(E)\big)q_{i,2}$$

$$\geq c_1 \sum_{i=1}^m \mu_{\boldsymbol{p}}\big(f_i^{-1} \circ T_{\Delta_1}^{-1}(E)\big)q_{i,1} = c_1\mu_{\boldsymbol{p}}(E).$$

Similarly, we have $\mu_{\boldsymbol{p}}(g(E)) \leq c_2\mu_{\boldsymbol{p}}(E)$ where $c_2 = \min\{q_{i,1}/q_{i,2} : 1 \leq i \leq m\}$. $\qquad\square$

We will revisit these ideas in Section 2.5.

2.3. **Children of net intervals.** Let $\Delta \in \mathcal{F}$ have neighbour set $\{f_1, \ldots, f_m\}$, and for each i, let S_{σ_i} generate the neighbour f_i (recall that this means that $S_{\sigma_i}(K) \cap \Delta^\circ \neq \emptyset$ and $f_i = T_\Delta^{-1} \circ S_{\sigma_i}$).

Definition 2.4. *We define the* ancestoral generation *of Δ, denoted* $\mathrm{ag}(\Delta)$, *and the* transition generation *of Δ, denoted* $\mathrm{tg}(\Delta)$, *to be positive real values such that*

$$\bigcap_{i=1}^j (|r_{\sigma_i}|, |r_{\sigma_i^-}|] = (\mathrm{tg}(\Delta), \mathrm{ag}(\Delta)].$$

Note that $0 < \mathrm{tg}(\Delta) \leq 1$; if $\Delta = [0, 1]$, we say $\mathrm{ag}(\Delta) = \infty$. It is straightforward to verify that
- $\mathrm{tg}(\Delta) = R_{\max}(\Delta) \cdot \mathrm{diam}(\Delta)$,
- $t \in (\mathrm{tg}(\Delta), \mathrm{ag}(\Delta)]$,
- for any $s \in (\mathrm{tg}(\Delta), \mathrm{ag}(\Delta)]$, $\Delta \in \mathcal{F}_s$ and $\mathcal{V}_s(\Delta) = \mathcal{V}_t(\Delta)$, and
- if $s \notin (\mathrm{tg}(\Delta), \mathrm{ag}(\Delta)]$, either $\Delta \notin \mathcal{F}_s$ or $\mathcal{V}_s(\Delta) \neq \mathcal{V}_t(\Delta)$.

Let $t > 0$ and $\Delta \in \mathcal{F}_t$. Let $(\Delta_1, \ldots, \Delta_n) \in \mathcal{F}_{\mathrm{tg}(\Delta)}$ be the distinct net intervals, ordered from left to right, of generation $\mathrm{tg}(\Delta)$ contained in Δ. Note that either $n > 1$ or if $n = 1$, then $\mathcal{V}(\Delta) \neq \mathcal{V}(\Delta_1)$. Then we call the tuple $(\Delta_1, \ldots, \Delta_n)$ the *children* of $\Delta \in \mathcal{F}_t$. Note that for any child Δ_i of Δ, $\mathrm{ag}(\Delta_i) = \mathrm{tg}(\Delta)$.

Similarly, we define the *parent* of $\Delta \in \mathcal{F}_t$ to be the net interval $\widehat{\Delta} \in \mathcal{F}_s$ with $s > t$, so that Δ is a child of $\widehat{\Delta}$.

Remark 2.5. *One way to think about the children of a net interval is as follows. Enumerate the points $\big\{\prod_{i\in\mathcal{I}} |r_i^{a_i}| : a_i \in \{0\} \cup \mathbb{N}\big\}$ in decreasing order $(t_i)_{i=1}^\infty$. Since $\mathrm{tg}(\Delta) = |r_\sigma|$ for some $\sigma \in \mathcal{I}^*$, the transitions to new generations must happen at some t_i. However, if $\Delta \in \mathcal{F}_{t_k}$,*

it may not hold that $\mathrm{tg}(\Delta) = t_{k+1}$. *The children are the net intervals in generation* t_m *where* $m \geq k+1$ *is minimal such that either* $\Delta \notin \mathcal{F}_{t_m}$ *or* $\mathcal{V}_{t_m}(\Delta) \neq \mathcal{V}_{t_k}(\Delta)$.

If the IFS is of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ *for some fixed* $0 < r < 1$ *and* $\Delta \in \mathcal{F}_{r^n}$, *then* $\mathrm{tg}(\Delta) = r^{n+1}$.

Example 2.6. *For a worked example of neighbour set and children computations of a non-commensurable IFS, see Section 5.3.*

A key feature of the preceding definitions is that, in a sense that will be made precise, the neighbour set of some net interval $\Delta \in \mathcal{F}_\alpha$ completely determines the placement and the neighbour set of each child of the net interval.

Definition 2.7. *Suppose* $\Delta = [a, b] \in \mathcal{F}$ *has children* $(\Delta_1, \ldots, \Delta_n)$ *in generation* $\mathrm{tg}(\Delta)$. *For some fixed child* $\Delta_i = [a_i, b_i]$, *we define the* position index $q(\Delta, \Delta_i) = (a_i - a)/\mathrm{diam}(\Delta)$.

One purpose of the position index is to distinguish the children of Δ which have the same neighbour set.

We have the following basic result. The insight behind this result is straightforward. The children of a net interval are determined precisely by the words which generate the neighbours of maximal length. Up to normalization by the position of Δ, these correspond uniquely to the neighbours of Δ with maximal contraction factor.

Theorem 2.8. *Let* $\{S_i\}_{i \in \mathcal{I}}$ *be an arbitrary IFS. Let* $\Delta \in \mathcal{F}_t$ *be arbitrary, with children* $(\Delta_1, \ldots, \Delta_n)$ *in* $\mathcal{F}_{\mathrm{tg}(\Delta)}$. *Then for any* $\Delta' \in \mathcal{F}_s$ *with* $\mathcal{V}(\Delta) = \mathcal{V}(\Delta')$ *and children* $(\Delta'_1, \ldots, \Delta'_{n'})$ *in* $\mathcal{F}_{\mathrm{tg}(\Delta')}$, *we have that* $n = n'$ *and for any* $1 \leq i \leq n$,

(i) $\mathcal{V}(\Delta'_i) = \mathcal{V}(\Delta_i)$,

(ii) $q(\Delta', \Delta'_i) = q(\Delta, \Delta_i)$,

(iii) $\frac{\mathrm{diam}(\Delta'_i)}{\mathrm{diam}(\Delta')} = \frac{\mathrm{diam}(\Delta_i)}{\mathrm{diam}(\Delta)}$, *and*

(iv) $\frac{\mathrm{tg}(\Delta_i)}{\mathrm{tg}(\Delta)} = \frac{\mathrm{tg}(\Delta'_i)}{\mathrm{tg}(\Delta_i)}$.

Proof. Given a map $f(x) = rx + d$, we set $R(f) = |r|$.

Write $\mathcal{V}(\Delta') = \mathcal{V}(\Delta) = \{f_1, \ldots, f_m\}$, and let

$$\mathcal{W}' = \{T_{\Delta'} \circ f_i : R(f_i) = R_{\max}(\Delta'), 1 \leq i \leq m\}$$
$$\mathcal{W} = \{T_\Delta \circ f_i : R(f_i) = R_{\max}(\Delta), 1 \leq i \leq m\}$$

denote the corresponding sets of neighbours corresponding to functions with maximal contraction factor, where $R_{\max}(\Delta') = R_{\max}(\Delta)$. Then let

$$\mathcal{C}' = \{S_\tau : \tau \in \Lambda_{\mathrm{tg}(\Delta')}, S_\tau(K) \cap (\Delta')^\circ \neq \emptyset\}$$
$$\mathcal{C} = \{S_\tau : \tau \in \Lambda_{\mathrm{tg}(\Delta)}, S_\tau(K) \cap \Delta^\circ \neq \emptyset\}.$$

In other words, $\mathcal{C}$ is the set of words of generation $\mathrm{tg}(\Delta)$ which contribute to some child of Δ, and similarly for Δ'. Using the observation that the only new words are those which are one-level descendents of those which generate neighbours of

maximal length, we have

$$\mathcal{C} = \{f \circ S_j : f \in \mathcal{W}, f \circ S_j(K) \cap \Delta^\circ \neq \emptyset\} \cup \{T_\Delta^{-1} \circ f_i : R(f_i) \neq R_{\max}(\Delta)\}$$

$$(2.3) \qquad = \{T_\Delta \circ T_{\Delta'}^{-1} \circ f : f \in \mathcal{C}'\}.$$

Note that, in the above set of equalities, we use the fact that for $f \in \mathcal{W}$

$$f \circ S_j(K) \cap \Delta^\circ \neq \emptyset \iff T_\Delta^{-1} \circ f \circ S_j(K) \cap (0,1) \neq \emptyset$$
$$\iff T_{\Delta'} \circ T_\Delta^{-1} \circ f \circ S_j(K) \cap (\Delta')^\circ \neq \emptyset$$

where $T_{\Delta'} \circ T_\Delta^{-1} \circ f \in \mathcal{W}'$.

Write $\Delta = [a,b]$ and $\Delta' = [a',b']$. Now consider the set $H = \{a,b\} \cup \{f(0), f(1) : f \in \mathcal{C}\} \cap \Delta$ so that H is the set of all endpoints of generation $\mathrm{tg}(\Delta)$ contained in Δ. Then if $H' = \{a',b'\} \cup \{f(0), f(1) : f \in \mathcal{C}'\} \cap \Delta'$, we observe by (2.3) that $T_{\Delta'}^{-1}(H') = T_\Delta^{-1}(H)$. Let $a = h_1 < \cdots < h_{k+1} = b$ denote the ordered elements of H and $a' = h_1' < \cdots < h_{k+1}' = b'$ the ordered elements of H' where $k = |H| - 1 = |H'| - 1$. By Lemma 2.3, we observe that $(h_i, h_{i+1}) \cap K \neq \emptyset$ if and only if $(h_i', h_{i+1}') \cap K \neq \emptyset$. Thus the children of Δ are $\{[h_i, h_{i+1}] : (h_i, h_{i+1}) \cap K \neq \emptyset\}$ and the children of Δ' are $\{T_{\Delta'} \circ T_\Delta^{-1}([h_i, h_{i+1}]) : (h_i, h_{i+1}) \cap K \neq \emptyset\}$, so $k = n = n'$.

Now fix some $1 \leq i \leq n$. Note that $T_\Delta \circ T_{\Delta'}^{-1}(\Delta_i') = \Delta_i$ so that $T_{\Delta_i}^{-1} \circ T_\Delta \circ T_{\Delta'}^{-1} = T_{\Delta_i'}^{-1}$.

(i) By direct computation,

$$\mathcal{V}(\Delta_i) = \{T_{\Delta_i}^{-1} \circ f : f \in \mathcal{C}, f(K) \cap \Delta_i^\circ \neq \emptyset\}$$
$$= \{T_{\Delta_i}^{-1} \circ T_\Delta \circ T_{\Delta'}^{-1} \circ f : f \in \mathcal{C}', T_\Delta \circ T_{\Delta'}^{-1} \circ f(K) \cap (T_\Delta \circ T_{\Delta'}^{-1}(\Delta_i'))^\circ \neq \emptyset\}$$
$$= \{T_{\Delta_i'}^{-1} \circ f : f \in \mathcal{C}', f(K) \cap (\Delta_i')^\circ \neq \emptyset\} = \mathcal{V}(\Delta_i')$$

(ii) Since the T_Δ are isometries, $q(\Delta, \Delta_i) = \frac{h_i - h_1}{\mathrm{diam}(\Delta)} = T_\Delta^{-1}(h_i)$ since $T_\Delta^{-1}(h_1) = 0$. Then the result follows since $T_\Delta^{-1}(h_i) = T_{\Delta'}^{-1}(h_i')$.

(iii) We have

$$\frac{\mathrm{diam}(\Delta_i)}{\mathrm{diam}(\Delta)} = \mathrm{diam}(T_\Delta^{-1}(\Delta_i)) = \mathrm{diam}(T_{\Delta'}^{-1}(\Delta_i')) = \frac{\mathrm{diam}(\Delta_i')}{\mathrm{diam}(\Delta')}$$

(iv) Recall that for an arbitrary net interval, $\mathrm{tg}(\Delta_0) = R_{\max}(\Delta_0) \cdot \mathrm{diam}(\Delta_0)$ where $R_{\max}(\Delta_0)$ depends only on $\mathcal{V}(\Delta_0)$. Apply (i) and (iii).

We thus have the desired result. □

2.4. **The transition graph of an iterated function system.** In the context of Theorem 2.8, to understand the behaviour of the IFS, it is in a sense sufficient to track the behaviour of the neighbour sets. Thus, we construct the *transition graph* of the IFS. The transition graph is a directed graph $\mathcal{G}(\{S_i\}_{i \in \mathcal{I}})$, possibly with loops and multiple edges, (denoted by $\mathcal{G}$ when the IFS is clear from the context) defined as follows. The vertex set of $\mathcal{G}$, denoted $V(\mathcal{G})$, is $\{\mathcal{V}(\Delta) : \Delta \in \mathcal{F}\}$, the set of distinct neighbour sets. The edge set of $\mathcal{G}$, denoted $E(\mathcal{G})$, is a set of triples (v_1, v_2, q) where v_1 is the source vertex, v_2 is the target vertex, and q is the edge label to distinguish multiple edges. The

edges are given as follows: for each net interval $\Delta \in \mathcal{F}_t$ with children $(\Delta_1, \ldots, \Delta_m)$ and for each i, we introduce an edge $e = (\mathcal{V}_t(\Delta), \mathcal{V}_{\mathrm{tg}(\Delta)}(\Delta_i), q(\Delta, \Delta_i))$. By Theorem 2.8, this construction is well-defined since it depends only on the neighbour set of Δ.

An *(admissible) path* η in $\mathcal{G}$ is a sequence of edges $\eta = (e_1, \ldots, e_n)$ in $\mathcal{G}$ where the target of e_i is the source of e_{i+1}. A path in $\mathcal{G}$ is a *cycle* if the path begins and ends at the same vertex.

We can encode the behaviour of the IFS symbolically using the transition graph. Given $\Delta \in \mathcal{F}_t$, consider the sequence $(\Delta_0, \ldots, \Delta_n)$ where $\Delta_0 = [0, 1]$, $\Delta_n = \Delta$, and each Δ_i is a child of Δ_{i-1}. Then the *symbolic representation* of Δ is the path $\eta = (e_1, \ldots, e_n)$ of G where for each $1 \leq i \leq n$

$$ e_i = \big(\mathcal{V}(\Delta_{i-1}), \mathcal{V}(\Delta_i), q(\Delta_{i-1}, \Delta_i) \big). $$

Conversely, if η is any admissible path, we say that $(\Delta_i)_{i=0}^k$ is a *(net interval) realization* of η if

- each Δ_i is a child of Δ_{i-1}, and
- each $e_i = (\mathcal{V}(\Delta_{i-1}), \mathcal{V}(\Delta_i), q(\Delta_{i-1}, \Delta_i))$.

By construction, every admissible path has a net interval realization.

Now let $x \in K$ be arbitrary and let $(\Delta_i)_{i=0}^\infty$ be a sequence of nested intervals where $\Delta_0 = [0, 1]$ and Δ_{i+1} a child of Δ_i and $\{x\} = \bigcap_{i=1}^\infty \Delta_i$. The symbolic representation of x corresponding to sequence $(\Delta_i)_{i=0}^\infty$ is the infinite path $(e_i)_{i=1}^\infty$ where for each n, $(e_1, \ldots, e_n)$ is the symbolic representation of Δ_n. The symbolic representation uniquely determines x, but if x is an endpoint of some net interval, it can happen that there are two distinct symbolic representations.

Suppose $\{S_i\}_{i \in \mathcal{I}}$ is of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ where $0 < r < 1$. Then if $\Delta \in \mathcal{F}_t$ is any net interval with symbolic representation $\eta = (e_1, \ldots, e_n)$, $\mathrm{tg}(\Delta) = r^n$ and $r^n < t \leq r^n / r_{\min}$. In other words, given the symbolic representation, we can approximate the generation of Δ.

However, when the IFS is not of this form, paths with the same length can result in net intervals in substantially different generations, and if the contraction ratios are not logarithmically commensurable (i.e. $\log r_i / \log r_j \in \mathbb{Q}$ for any $i, j \in \mathcal{I}$), there is no way to resolve this in a uniform way. Thus in order to approximate the change in generation along a path in the transition graph, it is necessary to assign distinct values to the edges in the transition graph.

Definition 2.9. *Let $\mathcal{G}$ be the transition graph of an IFS. We define the* edge length function $L : E(\mathcal{G}) \to (0, 1)$ *as follows. For a particular edge e, let the source and target be given by v_1 and v_2, where $v_i = \mathcal{V}(\Delta_i)$ for some Δ_1 the parent of Δ_2, and define $L(e) = \mathrm{tg}(\Delta_2) / \mathrm{tg}(\Delta_1)$.*

This function is well-defined by Theorem 2.8. When $\eta = (e_1, \ldots, e_n)$ is an admissible path, we say $L(\eta) = L(e_1) \cdots L(e_n)$.

Remark 2.10. *If $\{S_i\}_{i \in \mathcal{I}}$ is of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ where $0 < r < 1$, then $L(e) = r$ for any edge $e \in E(\mathcal{G})$.*

The main point here is that if $\Delta \in \mathcal{F}_t$ is any net interval with symbolic representation η, then $L(\eta) \asymp t$ with constants of comparability not depending on Δ. While the above choice of the length for an edge is not unique with this property, a straightforward argument shows that any such function must agree with L on any cycle.

2.5. **Encoding the invariant measure by the transition graph.** Given an IFS $\{S_i\}_{i \in \mathcal{I}}$ with a corresponding invariant measure $\mu_{\boldsymbol{p}}$, we are interested in formulas for computing or approximating $\mu_{\boldsymbol{p}}(E)$ where $E \subseteq K$ is an arbitrary Borel set. When $\{S_i\}_{i \in \mathcal{I}}$ satisfies the strong separation condition (that is, for $i \neq j$, $S_i(K)$ and $S_j(K)$ are disjoint), this is straightforward since $\mu_{\boldsymbol{p}}(S_\sigma(K)) = p_\sigma$. However, when images of K overlap, such a formula no longer holds.

The net interval construction can be thought of as a way of converting the behaviour of the IFS on overlapping images of K into behaviour on net intervals, which are disjoint except on a countable set (which has $\mu_{\boldsymbol{p}}$-measure 0). It turns out that one may also encode the dynamics of the invariant measure $\mu_{\boldsymbol{p}}$ using products of matrices. This technique was developed in the equicontractive case for IFS of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ with $0 < r < 1$ by Feng [8], and extended to IFS which satisfy the finite type condition [21]. Using similar techniques, we describe here how to generalize this construction to an arbitrary IFS.

Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS and $\mu_{\boldsymbol{p}}$ the self similar measure associated to probabilities $\{p_i\}_{i \in \mathcal{I}}$. The main mechanism to compute the approximate measure of net intervals is through transition matrices. Recall that $\mathcal{G}$ has vertex set $V(\mathcal{G}) = \{\mathcal{V}(\Delta) : \Delta \in \mathcal{F}\}$. Fix some total ordering on the set of all neighbours $\{f : f \in \mathcal{V}(\Delta), \Delta \in \mathcal{F}\}$.

Let $e \in E(\mathcal{G})$ be a fixed edge with source v_1 and target v_2. Suppose $\Delta_1 \supseteq \Delta_2$ are net intervals such that Δ_1 is the parent of Δ_2 and $e = (\mathcal{V}(\Delta_1), \mathcal{V}(\Delta_2), q(\Delta_1, \Delta_2))$. Suppose the neighbour sets are given by $\mathcal{V}(\Delta_1) = \{f_1, \ldots, f_m\}$ and $\mathcal{V}(\Delta_2) = \{g_1, \ldots, g_n\}$ where $f_1 < \cdots < f_m$ and $g_1 < \cdots < g_n$. We then define the *transition matrix* $T(e)$ as the non-negative $m \times n$ matrix given by

$$(2.4) \qquad\qquad T(e)_{i,j} = \frac{\mu_{\boldsymbol{p}}(g_j^{-1}((0,1)))}{\mu_{\boldsymbol{p}}(f_i^{-1}((0,1)))} \cdot p_\ell$$

if there exists an index $\ell \in \mathcal{I}$ such that f_i is generated by σ and g_j is generated by $\sigma\ell$; otherwise, set $T(e)_{i,j} = 0$. This is well-defined since a neighbour f has $f^{-1}((0,1)) \cap K \neq \emptyset$ by definition. Recall that if σ' generates any neighbour of Δ_2, then necessarily $\sigma' = \sigma\ell$ for some σ which generates a neighbour of Δ_1; thus, every column of $T(e)$ has a positive entry. However, it may not hold that each row of $T(e)$ has a positive entry.

It is clear from Theorem 2.8 that this definition depends only on the edge e. If $\eta = (e_1, \ldots, e_n)$ is an admissible path, we define $T(\eta) = T(e_1) \cdots T(e_n)$. Note that it need not hold that every row of $T(\eta)$ has a positive entry.

Example 2.11. *See Section 5.3 and Figure 1 for a complete transition graph example.*

Throughout, we will denote by $\|T\| = \sum_{i,j} |T_{ij}|$ to denote the matrix 1-norm. Suppose $\Delta \in \mathcal{F}_t$ is an arbitrary net interval. From the defining identity of the self-similar measure,

$$\mu_{\boldsymbol{p}}(\Delta) = \sum_{\sigma \in \Lambda_t} p_\sigma \mu_{\boldsymbol{p}}(S_\sigma^{-1}(\Delta))$$

where, since $\mu_{\boldsymbol{p}}$ is non-atomic, the summation may be taken over σ such that $S_\sigma^{-1}(\Delta^\circ) \cap K$ is non-empty. Note that $S_\sigma^{-1}(\Delta^\circ) = S_\sigma^{-1} \circ T_\Delta((0,1)) = f^{-1}((0,1))$ where $f \in \mathcal{V}(\Delta)$. We thus have

(2.5)
$$\mu_{\boldsymbol{p}}(\Delta) = \sum_{f \in \mathcal{V}(\Delta)} \mu_{\boldsymbol{p}}(f^{-1}((0,1))) \sum_{\substack{\sigma \in \Lambda_t \\ \sigma \text{ generates } f}} p_\sigma.$$

Let $\mathcal{V}(\Delta) = \{f_1, \dots, f_m\}$ with $f_1 < \dots < f_m$; then, we denote the *vector form* of $\mu_{\boldsymbol{p}}$ by $\boldsymbol{Q}_{\boldsymbol{p}}(\Delta) = (q_1, \dots, q_m)$ where

$$q_i = \mu_{\boldsymbol{p}}(f_i^{-1}((0,1))) \sum_{\substack{\sigma \in \Lambda_t \\ \sigma \text{ generates } f_i}} p_\sigma.$$

In particular, $\boldsymbol{Q}_{\boldsymbol{p}}(\Delta)$ is a strictly positive vector for any Δ, and $\mu_{\boldsymbol{p}}(\Delta) = \|\boldsymbol{Q}_{\boldsymbol{p}}(\Delta)\|$.

With this notation, we have the following theorem:

Theorem 2.12. *Let $\{S_i\}_{i \in \mathcal{I}}$ have associated self-similar measure $\mu_{\boldsymbol{p}}$. If η is any admissible path realized by $(\Delta_i)_{i=0}^m$,*

$$\boldsymbol{Q}_{\boldsymbol{p}}(\Delta_0)T(\eta) = \boldsymbol{Q}_{\boldsymbol{p}}(\Delta_m).$$

Proof. Suppose $\Delta_0 \in \mathcal{F}_t$ and $\Delta_m \in \mathcal{F}_s$. Say $\mathcal{V}(\Delta_0) = \{f_1, \dots, f_\ell\}$ with $f_1 < \dots < f_\ell$ and $\mathcal{V}(\Delta_m) = \{g_1, \dots, g_m\}$ with $g_1 < \dots < g_m$. For each i, assume τ_i generates the neighbour f_i, and set $\mathcal{A}_{ij} = \{\omega : \tau_i\omega \in \Lambda_s, \tau_i\omega \text{ generates } g_j\}$. Then for any $1 \le j \le m$, we have

$$\left(\boldsymbol{Q}_{\boldsymbol{p}}(\Delta_0)T(\eta)\right)_j = \sum_{i=1}^{\ell} \mu_{\boldsymbol{p}}(f_i^{-1}((0,1))) \left(\sum_{\substack{\sigma \in \Lambda_t \\ \sigma \text{ generates } f_i}} p_\sigma \right) \cdot \left(\sum_{\omega \in \mathcal{A}_{ij}} \frac{\mu_{\boldsymbol{p}}(g_j^{-1}((0,1)))}{\mu_{\boldsymbol{p}}(f_i^{-1}((0,1)))} p_\omega \right)$$

$$= \mu_{\boldsymbol{p}}(g_j^{-1}((0,1))) \sum_{i=1}^{\ell} \left(\sum_{\substack{\sigma \in \Lambda_t \\ \sigma \text{ generates } f_i}} p_\sigma \right) \cdot \left(\sum_{\omega \in \mathcal{A}_{ij}} p_\omega \right)$$

$$= \mu_{\boldsymbol{p}}(g_j^{-1}((0,1))) \sum_{\substack{\omega \in \Lambda_s \\ \omega \text{ generates } g_j}} p_\omega$$

so that $\boldsymbol{Q}_{\boldsymbol{p}}(\Delta_0)T(\eta) = \boldsymbol{Q}_{\boldsymbol{p}}(\Delta_m)$. $\qquad\square$

3. ITERATED FUNCTION SYSTEMS SATISFYING THE WEAK SEPARATION CONDITION

We now focus our attention on self-similar measures associated with IFSs satisfying the weak separation condition. We give a definition which is slightly different than the original [27], but is known to be equivalent when K is not a singleton [38]. Given a Borel set $E \subset K$ and $t > 0$, we define

$$\Lambda_t(E) = \{\sigma \in \Lambda_t : S_\sigma(K) \cap E \neq \emptyset\}$$
$$\mathcal{S}_t(E) = \{S_\sigma : \sigma \in \Lambda_t(E)\}$$

Let $U(x, t)$ denote the open ball about x with radius t.

Definition 3.1. *We say that the IFS $\{S_i\}_{i \in \mathcal{I}}$ satisfies the* weak separation condition *if*

$$(3.1) \qquad\qquad \sup_{x \in \mathbb{R}, t > 0} \#\mathcal{S}_t(U(x, t)) < \infty.$$

We can obtain an equivalent formulation of the weak separation condition in terms of a variant of the neighbour set which we call the *covering neighbour set*. Given a net interval $\Delta \in \mathcal{F}_t$, we write $\mathcal{V}^c(\Delta) = \{T_\Delta^{-1} \circ S_\sigma : \sigma \in \Lambda_t, S_\sigma([0,1]) \supseteq \Delta\}$. We refer to elements of $\mathcal{V}^c(\Delta)$ as *covering neighbours*. Notably, we omit the requirement that a neighbour $f \in \mathcal{V}^c(\Delta)$ has $f(K) \cap (0, 1) \neq \emptyset$.

Remark 3.2. *We always have $\mathcal{V}(\Delta) \subseteq \mathcal{V}^c(\Delta)$ with strict inequality possible. Moreover, we note that if Δ and Δ' are any net intervals with $\mathcal{V}^c(\Delta) = \mathcal{V}^c(\Delta')$, then necessarily $\mathcal{V}(\Delta) = \mathcal{V}(\Delta)$ following similar arguments to Lemma 2.3 and Theorem 2.8. Note that the covering neighbour set is taken as the definition of neighbour set in [19].*

We have the following characterization, which is [19, Prop. 4.3]:

Proposition 3.3 ([19]). *The IFS $\{S_i\}_{i \in \mathcal{I}}$ satisfies the weak separation condition if and only if*

$$\sup_{\Delta \in \mathcal{F}} \#\mathcal{V}^c(\Delta) < \infty.$$

Net intervals for which $\#\mathcal{V}^c(\Delta)$ attain the supremum in Proposition 3.3 will play an important role in our analysis in this section.

3.1. **The essential class of the transition graph.**
Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS with associated transition graph $\mathcal{G}$. Recall that in a directed graph $\mathcal{G}$, an induced subgraph $\mathcal{G}'$ is a subgraph for which there exists some set of vertices $H \subseteq V(\mathcal{G})$ such that $\mathcal{G}'$ has vertex set H and edge set composed of every outgoing edge from a vertex in H which connects to another vertex in H.

Definition 3.4. *An* essential class *of $\mathcal{G}$ is an induced subgraph $\mathcal{G}'$ of $\mathcal{G}$ such that*
   (i) for any $v, v' \in \mathcal{G}'$, there exists a path from v to v', and
   (ii) if $v \in \mathcal{G}$ and $v' \in \mathcal{G}'$ and there is a path from v' to v, then $v \in \mathcal{G}'$.

In a finite graph, there is always at least one essential class [35, Lem. 1.1]. In an infinite graph, there need not be an essential class; moreover, the essential class, if it exists, need not be finite. When $\mathcal{G}$ has exactly one essential class, we denote it by $\mathcal{G}_{\mathrm{ess}}$.

We have the following basic observation. The proof of this result is similar to the idea in [19, Lemma 4.2], but we reiterate the aspects of the proof that we need here for clarity.

Proposition 3.5. *Let $\{S_i\}_{i\in\mathcal{I}}$ be an IFS satisfying the weak separation condition. Then its transition graph $\mathcal{G}$ has a unique essential class.*

Proof. It suffices to show that there exists some vertex v such that if w is any other vertex, there exists an admissible path from w to v. Then the essential class is the set of all vertices v' for which there is a path from v to v'. By Proposition 3.3, there exists some $t > 0$ and net interval $\Delta_0 \in \mathcal{F}_t$ such that $\#\mathcal{V}^c(\Delta_0)$ is maximal; let $v := \mathcal{V}(\Delta_0)$.

Now, let $w \in V(\mathcal{G})$ be arbitrary and $\Delta \in \mathcal{F}$ such that $\mathcal{V}(\Delta) = w$. Since $\Delta^\circ \cap K \neq \emptyset$, there exists some $\sigma \in \mathcal{I}^*$ such that $S_\sigma(K) \subseteq \Delta$ and $r_\sigma > 0$. Set $\gamma = r_\sigma \cdot t$ and let $\Delta_1 := S_\sigma(\Delta_0)$

Let $\Delta_0 = [a, b]$ have covering neighbours generated by words $\{\omega_1, \ldots, \omega_m\}$ with $\omega_i \in \Lambda_t$. By definition of γ, $\{\sigma\omega_1, \ldots, \sigma\omega_m\}$ are words of generation Λ_γ. Note that $(\Delta_1)^\circ \cap K \neq \emptyset$ and that the endpoints of Δ_1 are of the form $S_{\sigma\zeta}(z)$ where $z \in \{0, 1\}$ and $\zeta \in \Lambda_s$, so that $\sigma\zeta \in \Lambda_\gamma$. In particular, if $\Delta_1 \notin \mathcal{F}_\gamma$, then there exists some $\tau \in \Lambda_\gamma$ such that $S_\tau \notin \{S_{\sigma\omega_1}, \ldots, S_{\sigma\omega_m}\}$ and $S_\tau([0,1]) \supseteq \Delta_1$. But then there exists some $\Delta_2 \in \mathcal{F}_\gamma$ with $\Delta_2 \subseteq \Delta_1 \cap S_\tau([0,1])$, where Δ_2 has distinct covering neighbours generated by $\{\omega_1, \ldots, \omega_m\} \cup \{\tau\}$, contradicting the maximality of $\#\mathcal{V}^c(\Delta_0)$.

Thus Δ_1 is in fact a net interval of generation γ. Moreover, since $r_\sigma > 0$, we have $T_{\Delta_1} = S_\sigma \circ T_{\Delta_0}$, so that

$$\mathcal{V}^c(\Delta_1) = \{T_{\Delta_1}^{-1} \circ S_{\sigma\omega_i}\}_{i=1}^m = \{T_{\Delta_0}^{-1} \circ S_\sigma^{-1} \circ S_\sigma \circ S_{\omega_i}\}_{i=1}^m = \mathcal{V}^c(\Delta_0).$$

Thus by Remark 3.2, we have $\mathcal{V}(\Delta_1) = v$ and $\Delta_1 \subseteq \Delta$, so that there exists a path from $\mathcal{V}(\Delta)$ to $\mathcal{V}(\Delta_1)$, as claimed. $\qquad\square$

Definition 3.6. *We say that a point $x \in K$ is an* essential point *if for some symbolic representation $(e_j)_{j=1}^\infty$ of x, there exists some $N \in \mathbb{N}$ so that for all $k \geq N$, $e_k \in E(\mathcal{G}_{\mathrm{ess}})$. We say that a point $x \in K$ is an* interior essential point *if every symbolic representation has this property. We denote the set of all interior essential points by K_{ess}. We say a net interval $\Delta \in \mathcal{F}$ is an* essential net interval *if $\mathcal{V}(\Delta) \in V(\mathcal{G}_{\mathrm{ess}})$.*

If Δ is an essential net interval, then $\Delta^\circ \cap K \subseteq K_{\mathrm{ess}}$. Of course, a given path $(e_j)_{j=1}^\infty$ is eventually in the essential class if and only if a single edge is in the essential class. One may verify that the set of interior essential points is the topological interior of the set of essential points; in particular, the essential points form an open set in K. Interior essential points play an important role in the multifractal analysis of self-similar measures under the weak separation condition.

In the next proposition, we observe that interior essential points are abundant.

Proposition 3.7. *Let $\{S_i\}_{i\in\mathcal{I}}$ be an IFS satisfying the weak separation condition. Let $U(x_0, t_0)$ be any open ball which attains the maximal value in (3.1). Then the following hold:*

(i) If $\sigma \in \mathcal{I}^*$ is arbitrary, then $S_\sigma(U(x_0, t_0))$ also attains the maximal value in (3.1).

(ii) $U(x_0, t_0) \cap K$ is contained in a finite union of essential net intervals. In particular, $U(x_0, t_0) \cap K \subseteq K_{\text{ess}}$.

Remark 3.8. *In Section 5.3.3, we show that the converse of (ii) need not hold: there exists some IFS $\{S_i\}_{i \in \mathcal{I}}$ satisfying the weak separation condition and an essential net interval Δ such that $\Delta \cap K$ is not contained a finite union of balls $U(x_0, t_0)$. In the same example, we show that if W is the union of all balls $U(x_0, t_0)$ which attain the maximal value in (3.1), then $W \subsetneq K_{\text{ess}}$.*

Proof. To see that $S_\sigma(U(x_0, t_0))$ also attains the maximal value in (3.1), if

$$\mathcal{S}_{t_0}(U(x_0, t_0)) = \{S_{\phi_1}, \ldots, S_{\phi_m}\},$$

then $S_{\sigma\phi_i} \in \mathcal{S}_{|r_\sigma|t_0}(S_\sigma(U(x_0, t_0)))$ for each i and $\#\mathcal{S}_{|r_\sigma|t_0}(S_\sigma(U(x_0, t_0))) \geq m$. Then equality holds by maximality of m.

We now see (ii). By definition of net intervals, we know that for any $t > 0$, $U(x_0, t_0) \cap K$ is contained in a finite union of net intervals of generation t. In particular, it suffices to show that there is some $t_1 > 0$ such the set

$$\{\Delta \in \mathcal{F}_{t_1} : \Delta \cap U(x_0, t_0) \neq \emptyset\}$$

is composed only of essential net intervals. Let Δ_0 be a fixed essential net interval and let $\sigma_0 \in \mathcal{I}^*$ have $r_{\sigma_0} > 0$ and $S_{\sigma_0}([0, 1]) \subseteq \Delta_0$. As argued above, $S_{\sigma_0}(U(x_0, t_0))$ also attains the maximal value in (3.1). Let

$$H = \{S_\sigma : \sigma \in \Lambda_{r_{\sigma_0}t_0}, S_\sigma(K) \cap S_{\sigma_0}(U(x_0, t_0)) = \emptyset\}.$$

Since $S_{\sigma_0}(U(x_0, t_0))$ is open, there exists some $\epsilon_0 > 0$ such that for any ϵ with $|\epsilon| < \epsilon_0$, $S_{\sigma_0}(U(x_0 + \epsilon, t_0))$ also attains the maximal value in (3.1). In particular, if $S_\sigma \in H$ is arbitrary, we in fact have $S_\sigma(K) \cap S_{\sigma_0}(B(x_0, t_0)) = \emptyset$. Since H is a finite set, take

$$t_1 = \min\{\min\{\text{dist}(f(K), S_{\sigma_0}(B(x_0, t_0))) : f \in H\}, t_0\} > 0.$$

It remains to show that such a t_1 works.

Write $\mathcal{S}_{t_0}(U(x_0, t_0)) = \{S_{\phi_1}, \ldots, S_{\phi_m}\}$ and set

$$F = \{\Delta \in \mathcal{F}_{t_1} : \Delta \cap U(x_0, t_0) \neq \emptyset\}.$$

Suppose for contradiction there is some $\Delta \in F$ that is not an essential net interval, and let Δ have neighbours generated by distinct functions $\{S_{\omega_1}, \ldots, S_{\omega_k}\}$ with $\omega_i \in \Lambda_{t_1}$. As argued in Proposition 3.5, since $\Delta_1 := S_{\sigma_0}(\Delta)$ is not a net interval with neighbour set $\mathcal{V}(\Delta)$ (or Δ_1 would be a descendent of Δ_0, and hence essential), there exists some $\tau \in \Lambda_{r_{\sigma_0}t_1}$ such that $S_\tau(K) \cap \Delta^\circ \neq \emptyset$ and $S_\tau \neq S_{\sigma_0\omega_i}$ for each $1 \leq i \leq k$. We also observe that

$$(3.2) \qquad \{S_{\sigma_0\omega_1}, \ldots, S_{\sigma_0\omega_k}\} = \{S_{\sigma_0\xi} : \xi \in \Lambda_{t_1}, S_{\sigma_0\xi}(K) \cap \Delta_1^\circ \neq \emptyset\}.$$

Since $t_1 \leq t_0$, let $\tau_1 \preccurlyeq \tau$ be the unique prefix in $\Lambda_{r_{\sigma_0}t_0}$. Suppose for contradiction $S_{\tau_1}(K) \cap S_{\sigma_0}(U(x_0, t_0)) \neq \emptyset$. Since $S_{\sigma_0}(U(x_0, t_0))$ attains the maximal value in (3.1), we have $S_{\tau_1} = S_{\sigma_0} \circ S_\omega$ for some $S_\omega \in \mathcal{S}_{r_{\sigma_0}t_0}(S_\sigma(U(x_0, t_0)))$. Thus there exists some

word ξ such that $S_\tau = S_\sigma \circ S_\xi$, which contradicts (3.2). We thus have that $S_{\tau_1}(K) \cap S_{\sigma_0}(U(x_0, t_0)) = \emptyset$ so that $S_{\tau_1} \in H$.

But by definition of Δ_1, we have that $\Delta_1 \cap S_{\sigma_0}(U(x_0, t_0)) \neq \emptyset$ and $\Delta_1^\circ \cap S_{\tau_1}(K) \neq \emptyset$, so

$$\operatorname{dist}(S_{\tau_1}(K), S_{\sigma_0}(U(x_0, t_0))) < \operatorname{diam}(\Delta_1) \leq t_1,$$

contradicting the choice of t_1. Thus every $\Delta \in F$ is in fact essential, as claimed. $\square$

3.2. **An important measure approximation lemma.** The following technical lemma is a key approximation property for measures satisfying the weak separation condition, and the main factor behind the regularity of the measure on the essential class. Note the similarity of the result to the weak separation "counting" results; see, for example, Feng and Lau [12, Prop. 4.1].

Lemma 3.9. *Suppose the IFS $\{S_i\}_{i \in \mathcal{I}}$ satisfies the weak separation condition, and let $v \in V(\mathcal{G}_{\mathrm{ess}})$ be fixed. Then there exist constants $c, C > 0$ (depending on v) such that for any ball $B(x, t)$ with $\mu_{\boldsymbol{p}}(B(x, t)) > 0$, there exists $t \geq s \geq ct$ and $\Delta \in \mathcal{F}_s$ such that $\Delta \subseteq B(x, 2t)$, $\mathcal{V}(\Delta) = v$, and $\boldsymbol{Q}_{\boldsymbol{p}}(\Delta)_j \geq C \cdot \mu_{\boldsymbol{p}}(B(x, t))$ for each $1 \leq j \leq \#v$.*

Proof. Since $\mu_{\boldsymbol{p}}(B(x, t)) > 0$ and $\mu_{\boldsymbol{p}}$ is non-atomic, $U(x, t) \cap K \neq \emptyset$. From the weak separation condition, there exists some $\ell \in \mathbb{N}$ such that $\#\mathcal{S}_t(B(x, t)) \leq \ell$ for any $x \in \mathbb{R}$ and $t > 0$. By the invariant property of $\mu_{\boldsymbol{p}}$ and since $\mu_{\boldsymbol{p}}$ is a probability measure, we have

$$\mu_{\boldsymbol{p}}(B(x, t)) = \sum_{\sigma \in \Lambda_t(B(x,t))} p_\sigma \mu_{\boldsymbol{p}} \circ S_\sigma^{-1}((B(x, t))) \leq \sum_{\sigma \in \Lambda_t(B(x,t))} p_\sigma$$

$$= \sum_{\substack{S_\omega \in \mathcal{S}_t(B(x,t))}} \sum_{\substack{\sigma \in \Lambda_t(B(x,t)) \\ S_\sigma = S_\omega}} p_\sigma.$$

In particular, since $\#\mathcal{S}_t(B(x, t)) \leq \ell$, get ω_0 such that

$$(3.3) \qquad \sum_{\substack{\sigma \in \Lambda_t(B(x,t)) \\ S_\sigma = S_{\omega_0}}} p_\sigma \geq \mu_{\boldsymbol{p}}(B(x, t))/\ell.$$

Note that $S_{\omega_0}(K) \cap B(x, t) \neq \emptyset$, so that $S_{\omega_0}([0, 1]) \subseteq B(x, 2t)$. If $r_{\omega_0} < 0$, get $k \in \mathcal{I}$ with $r_k < 0$ and set $\omega_1 = \omega_0 k$; otherwise, take $\omega_1 = \omega_0$. Now, let $\Delta_0 \in \mathcal{F}_{s_0}$ be such that $\#\mathcal{V}^c(\Delta_0)$ is maximal. Exactly as argued in Proposition 3.5, $\Delta_1 := S_{\omega_1}(\Delta_0)$ is a net interval in generation $r_{\omega_1} \cdot s_0$ with $\mathcal{V}(\Delta_1) = \mathcal{V}(\Delta_0)$. Moreover, we know that if σ generates some neighbour f of Δ_0, then $\omega_1 \sigma$ generates the same neighbour f of Δ_1. Fix some $1 \leq j \leq \#\mathcal{V}(\Delta_1)$ and let f_j be the neighbour of Δ_1 corresponding to the

index j. We then have by using (3.3) and the above observation that

$$(\boldsymbol{Q_p}(\Delta_1))_j = \mu(f_j^{-1}((0,1))) \sum_{\substack{\sigma \in \Lambda_{s_0 r \omega_1} \\ \sigma \text{ generates } f_j}} p_\sigma$$

$$\geq p_k \Big(\sum_{\substack{\sigma \in \Lambda_t(B(x,t)) \\ S_\sigma = S_{\omega_1}}} p_\sigma \Big) \cdot \mu(f_j^{-1}((0,1))) \cdot \sum_{\substack{\sigma \in \Lambda_{s_0} \\ \sigma \text{ generates } f_j}} p_\sigma$$

$$\geq \mu_{\boldsymbol{p}}(B(x,t)) \cdot \frac{p_k \cdot (\boldsymbol{Q_p}(\Delta_0))_j}{\ell} \geq \mu_{\boldsymbol{p}}(B(x,t)) \cdot C_1$$

where $C_1 := p_k \cdot \min_j(\boldsymbol{Q_p}(\Delta_0))_j/\ell$, which depends only on the IFS and choice of probabilities.

Now let η be any fixed path from $\mathcal{V}(\Delta_0)$ to v and let ϵ be the smallest strictly positive entry of $T(\eta)$. Let Δ be the unique net interval with symbolic $\gamma\eta$ where γ is the symbolic representation of Δ_0. Since $T(\eta)$ is non-negative and $\boldsymbol{Q_p}(\Delta) = \boldsymbol{Q_p}(\Delta_1)T(\eta)$ is a positive vector, we have that $(\boldsymbol{Q_p}(\Delta))_j \geq \mu_{\boldsymbol{p}}(B(x,t)) \cdot C_1 \cdot \epsilon$. Taking $C := C_1 \epsilon$, we see that C satisfies the requirements. Moreover, since $\Delta_0 \in \mathcal{F}_{r_{\omega_0} s_0}$, taking $c = s_0 L(\eta) \cdot r_{\min}^2$ and noting that $t \cdot r_{\min} \leq |r_{\omega_0}| \leq t$, we have that $\Delta \in \mathcal{F}_s$ where $s \geq ct$. Finally, $\Delta \subseteq \Delta_1 \subseteq S_{\omega_0}([0,1]) \subseteq B(x,2t)$ as required. $\square$

3.3. **Measure properties of the essential class.** As our first consequence of this lemma, we establish that the interior essential points form a large subset of K.

Theorem 3.10. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition with attractor K and let $v \in V(\mathcal{G}_{\mathrm{ess}})$ be arbitrary. Let*

$$E = \bigcup_{\substack{\Delta \in \mathcal{F} \\ \mathcal{V}(\Delta) = v}} \Delta \cap K.$$

Then if $\mu_{\boldsymbol{p}}$ is any associated self-similar measure, $\mu_{\boldsymbol{p}}(K \setminus E) = 0$. In particular, $\mu_{\boldsymbol{p}}(K \setminus K_{\mathrm{ess}}) = 0$.

Proof. By Lemma 3.9, there exist constants $c, C > 0$ such that for any $t > 0$ and ball $B(x,t)$ with $\mu_{\boldsymbol{p}}(B(x,t)) > 0$, there exists some net interval $\Delta \in \mathcal{F}$ with $\Delta \subseteq B(x,2t)$, $\mathcal{V}(\Delta) = v$, and $\mu_{\boldsymbol{p}}(\Delta) \geq C\mu_{\boldsymbol{p}}(B(x,r))$. We will construct a nested family of sets $E_1 \supseteq E_2 \supseteq \cdots$ such that each E_n is a finite union of intervals, $\mu_{\boldsymbol{p}}(E_n) \leq (1 - C/3)^n$, and $K \setminus E \subseteq \bigcap_{n=1}^\infty E_n$. From this, the result clearly follows.

First consider the ball $B_1 = B(0,1)$. Get $\Delta_1 \subseteq B(0,2)$ with $\mathcal{V}(\Delta_1) = v$, set $E_1 = [0,1] \setminus \Delta_1$ so that $\mu_{\boldsymbol{p}}(E_1) \leq 1 - C \leq 1 - C/3$. Since Δ_1 is an interval, E_1 is a finite union of intervals and clearly $K \setminus E \subseteq E_1$. Inductively, suppose E_n is a finite union of intervals with $\mu_{\boldsymbol{p}}(E_n) \leq (1 - \lambda)^n$. Since each E_n is a finite union of intervals, there is a family of balls $\{B(x_i, t_i)\}_{i=1}^m$ such that the balls only overlap pairwise on endpoints, $E_n = \bigcup_{i=1}^m B(x_i, t_i)$, and for any distinct i_1, i_2, i_3,

$$(3.4) \qquad B(x_{i_1}, 2t_{i_1}) \cap B(x_{i_2}, 2t_{i_2}) \cap B(x_{i_3}, 2t_{i_3})$$

is either a singleton or the empty set and hence has measure 0, as $\mu_{\boldsymbol{p}}$ has no atoms. Now for each $1 \leq i \leq m$, apply Lemma 3.9 to get $\Delta_n^{(i)} \subseteq B(x_i, 2t_i)$ with $\mu_{\boldsymbol{p}}(\Delta_n^{(i)}) \geq C\mu_{\boldsymbol{p}}(B(x_i, t_i))$. While the $\Delta_n^{(i)}$ need not be disjoint, by (3.4), there exists a subcollection labelled without loss of generality $\{\Delta_n^{(i)}\}_{i=1}^{m'}$ such that

$$\sum_{i=1}^{m'} \mu_{\boldsymbol{p}}(\Delta_n^{(i)}) \geq \frac{1}{3} \sum_{i=1}^{m} \mu_{\boldsymbol{p}}(\Delta_n^{(i)})$$

and $\Delta_n^{(i)} \cap \Delta_n^{(j)}$ is at most a singleton for $i \neq j$. (To do this, pick the interval $\Delta_n^{(i)}$ with the largest measure and remove any net intervals $\Delta_n^{(j)}$ where $\Delta_n^{(j)} \cap \Delta_n^{(i)}$ is not a singleton. By (3.4) and the geometry in $\mathbb{R}$, there are at most 2 such indices j. Then repeat until the set is exhausted.)

Set $E_{n+1} = E_n \setminus \bigcup_{i=1}^{m'} \Delta_n^{(i)}$. Each $\Delta_n^{(i)}$ is an interval with $\mathcal{V}(\Delta_n^{(i)}) = v$, so that E_{n+1} is a finite union of intervals with $K \setminus E \subseteq E_{n+1}$, and

$$\mu_{\boldsymbol{p}}(E_{n+1}) = \mu_{\boldsymbol{p}}(E_n) - \sum_{i=1}^{m'} \mu_{\boldsymbol{p}}(\Delta_n^{(i)}) \leq \mu_{\boldsymbol{p}}(E_n) - \frac{C}{3} \sum_{i=1}^{m} \mu_{\boldsymbol{p}}(B(x_i, t))$$

$$\leq (1 - C/3)\mu_{\boldsymbol{p}}(E_n) \leq (1 - C/3)^{n+1}$$

as claimed. $\qquad\qquad\square$

Remark 3.11. *It can also be shown, using similar techniques, that if $s = \dim_H K$, then $\mathcal{H}^s(K \setminus K_{\mathrm{ess}}) = 0$ where $\mathcal{H}^s$ is the s-dimensional Hausdorff measure. This follows from Ahlfors regularity of self-similar sets under the weak separation condition [15, Thm. 2.1] along with Lemma 2.3, in place of Lemma 3.9.*

3.4. **Local dimensions and periodic points.** The notion of a periodic point was introduced by Hare, Hare and Matthews for IFS of the form $\{x \mapsto rx + d_i\}_{i \in \mathcal{I}}$ with $0 < r < 1$ satisfying the finite type condition [17]. In this section, we take advantage of the general matrix product formula, Theorem 2.12, to establish symbolic formulas for the local dimensions at certain points which we call periodic.

Definition 3.12. *Given a Borel probability measure μ, by the lower local dimension of μ at $x \in \operatorname{supp} \mu$, we mean the number*

$$\underline{\dim}_{\mathrm{loc}}\mu(x) = \liminf_{t \downarrow 0} \frac{\log \mu(B(x,t))}{\log t}.$$

The upper local dimension is defined analgously; when the upper and lower local dimensions coincide, we call the shared value the local dimension of μ at x, denoted by $\dim_{\mathrm{loc}} \mu(x)$.

Definition 3.13. *A periodic point is a point $x \in K$ where every symbolic representation of x is of the form*

$$[x] = (e_1, \ldots, e_n, \theta, \theta, \ldots)$$

where n is minimal and $\theta = (\theta_1, \ldots, \theta_m)$ is a cycle of $\mathcal{G}$ with minimal length. In this case, we call θ a period of the symbolic representation.

Intuitively, periodic points are the natural analogue of the rational numbers; for example, with respect to the IFS $\{x \mapsto x/2, x \mapsto x/2 + 1/2\}$, the periodic points of thie IFS are precisely the rational numbers in $[0, 1]$. Under the weak separation condition, it is straightforward to see that the periodic points form a countable dense subset of K: if $x, y \in K$ have symbolic representations of the form $\gamma\eta_1$ and $\gamma\eta_2$, then both x and y are in the net interval with symbolic representation γ.

The proofs of Lemma 3.14 and Proposition 3.15 are motivated by the proofs [18, Thm. 2.6 and Prop. 2.7].

Fix some $x \in K$. Enumerate $\{h_j : j = 1, \ldots, n\} = \{S_\sigma(0), S_\sigma(1) : \sigma \in \Lambda_t\}$ with $h_1 < \cdots < h_n$. If $x \neq h_j$ for each $1 \leq j \leq n$, then there is a unique net interval $\Delta_t(x) = [h_i, h_{i+1}]$ of generation t containing x. We then say $\Delta_t^-(x)$ is the empty set if $i = 1$ or $(h_{i-1}, h_i) \cap K = \emptyset$, and $\Delta_t^-(x) = [h_{i-1}, h_i]$ otherwise, and we define Δ_t^+ similarly. Then set

$$M_t(x) = \Delta_t(x)^- \cup \Delta_t(x) \cup \Delta_t(x)^+$$

Otherwise, $x = h_m$ for some m, and we write $\Delta_t^{(1)}(x) = [h_{m-1}, h_m]$ if $m \neq 1$ and $(h_{m-1}, h_m) \cap K$ is non-empty, and similarly for $\Delta_t^{(2)}(x)$, and set

$$M_t(x) = \Delta_t^{(1)}(x) \cup \Delta_t^{(2)}(x).$$

We have the following basic estimation:

Lemma 3.14. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS as in (2.1) and let $x \in K$ be such that $\sup\{R_{\max}(\Delta) : x \in \Delta, \Delta \in \mathcal{F}\} < \infty$. Then if $\mu_{\boldsymbol{p}}$ is any associated self-similar measure,*

$$\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) = \lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(M_t(x))}{\log t}$$

provided the limit on the right exists. Similar statements hold with respect to the limit supremum and limit infimum for the upper and lower local dimensions respectively.

Proof. Suppose the local dimension exists and equals D. Recall that if $\Delta \in \mathcal{F}_t$, then $t \geq \mathrm{tg}(\Delta) = R_{\max}(\Delta) \operatorname{diam}(\Delta)$. Thus there exists some constant $0 < \epsilon$ such that for any $t > 0$ and $\Delta \in \mathcal{F}_t$ with $x \in \Delta$, $\epsilon t < \operatorname{diam}(\Delta)$. Moreover, $\operatorname{diam}(\Delta) \leq t$ always holds by the net interval construction.

If x is a boundary point, get s such that x is an endpoint of $\Delta_s(x)$ and

$$B(x, \epsilon s) \subseteq \Delta_s^{(1)}(x) \cup \Delta_s^{(2)}(x) \subseteq B(x, 2s)$$

where the notation is as above. Otherwise if x is not a boundary point, then

$$B(x, \epsilon s) \subseteq \Delta_s^-(x) \cup \Delta_s(x) \cup \Delta_s^+(x) \subseteq B(x, 2s)$$

In either case, $B(x, \epsilon s) \subseteq M_s(x) \subseteq B(x, 2s)$ so that

$$\left(\frac{\log \epsilon + \log s}{\log s} \right) \left(\frac{\log \mu_{\boldsymbol{p}}(B(x, \epsilon s))}{\log \epsilon s} \right)$$

$$\leq \frac{\log \mu_{\boldsymbol{p}}(M_s(x))}{\log s} \leq \left(\frac{\log s + \log 2}{\log s} \right) \left(\frac{\log \mu_{\boldsymbol{p}}(B(x, 2s))}{\log 2s} \right).$$

The limit of the left and right both exist and are equal to D; hence, the limit of the middle expression exists and equals D. The arguments for the upper and lower dimension follow similarly. $\qquad\square$

In the following proposition, recall that for a path θ, $L(\theta)$ is the length of the path defined in Definition 2.9.

Proposition 3.15. *Let $\{S_i\}_{i \in \mathcal{I}}$ be any IFS and suppose x is a periodic point with period $\theta = (e_1, \ldots, e_s)$. Then the local dimension of μ at x exists and is given by*

$$\dim_{loc} \mu(x) = \frac{\log \mathrm{sp}(T(\theta))}{\log L(\theta)}$$

where if x is a boundary point of a net interval with two different symbolic representations given by periods θ and ϕ, then θ is chosen to satisfy

$$\frac{\log \mathrm{sp}(T(\theta))}{\log L(\theta)} \leq \frac{\log \mathrm{sp}(T(\phi))}{\log L(\phi)}.$$

Proof. First, suppose x is a periodic point with two distinct symbolic representations with periods $\theta = (\theta_1, \ldots, \theta_\ell)$ and $\phi = (\phi_1, \ldots, \phi_{\ell'})$, so that x is an endpoint of some net interval $\Delta \in \mathcal{F}$. We first note that

$$\mu_{\boldsymbol{p}}(\Delta_t^{(1)}(x)) = \left\| T(e_1, \ldots, e_j, \underbrace{\theta, \ldots, \theta}_{m}, \theta_1, \ldots, \theta_t) \right\|$$

$$\mu_{\boldsymbol{p}}(\Delta_t^{(2)}(x)) = \left\| T(e'_1, \ldots, e'_{j'}, \underbrace{\phi, \ldots, \phi}_{m'}, \phi_1, \ldots, \phi_{t'}) \right\|$$

for t sufficiently small, $t < \ell$, and $t' < \ell'$. Now, get constants c_i which do not depend on t such that

$$\left\| (T(\theta))^{m+1} \right\| \leq \left\| T(\underbrace{\theta, \ldots, \theta}_{m}, \theta_1, \ldots, \theta_t) \right\| \cdot \left\| T(\theta_{t+1}, \ldots, \theta_\ell) \right\|$$

(3.5) $$\leq c_1 \left\| T(e_1, \ldots, e_j, \underbrace{\theta, \ldots, \theta}_{m}, \theta_1, \ldots, \theta_t) \right\| \leq c_2 \left\| T(\theta)^m \right\|.$$

Moreover, since

$$L(e_1, \ldots, e_j) L(\theta)^m L(\theta_1, \ldots, \theta_t) r_{\min} \leq t \leq L(e_1, \ldots, e_j) L(\theta)^m L(\theta_1, \ldots, \theta_t),$$

we have $L(\theta)^m \asymp t$ with constants of comparability not depending on t. Thus, there exist k_i not depending on t so that

$$\frac{\log k_1 \, \|T(\theta)^{m+1}\|^{1/(m+1)}}{\log k_3 \cdot L(\theta)} \geq \frac{\log \mu(\Delta_t^{(1)}(x))}{\log t} \geq \frac{\log k_2 \, \|(T(\theta))^m\|^{1/m}}{\log k_4 \cdot L(\theta)}$$

and taking the limit as t goes to 0 yields

$$\lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(\Delta_t^{(1)}(x))}{\log t} = \frac{\log \operatorname{sp}(T(\theta))}{\log L(\theta)}$$

In the exact same way, we get

$$\lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(\Delta_t^{(2)}(x))}{\log t} = \frac{\log \operatorname{sp}(T(\phi))}{\log L(\phi)}.$$

Now, since x is a periodic point, the set $\{\mathcal{V}(\Delta) : x \in \Delta, \Delta \in \mathcal{F}\}$ is finite. Since $R_{\max}(\Delta)$ depends only on $\mathcal{V}(\Delta)$, $\sup\{R_{\max}(\Delta) : x \in \Delta, \Delta \in \mathcal{F}\} < \infty$ and the assumptions for Lemma 3.14 hold. Then by the power mean inequality, we have

$$\begin{aligned}
\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) &= \lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(\Delta_t^{(1)}(x)) + \mu_{\boldsymbol{p}}(\Delta_t^{(2)}(x))}{\log t} \\
&= \min\left(\lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(\Delta_t^{(1)}(x))}{\log t}, \lim_{t \to 0} \frac{\log \mu_{\boldsymbol{p}}(\Delta_t^{(2)}(x))}{\log t} \right) \\
&= \min\left(\lim_{t \to 0} \frac{\log \operatorname{sp} T(\theta)}{\log L(\theta)}, \lim_{t \to 0} \frac{\log \operatorname{sp} T(\phi)}{\log L(\phi)} \right)
\end{aligned}$$

since the final two limits in the maximum exist, as claimed.

If x is an endpoint of some net interval but has only one symbolic representation, then either $\Delta_t^{(1)}(x)$ or $\Delta_t^{(2)}(x)$ is empty for sufficiently small t and the argument is identical, but easier.

Finally, suppose x is not an endpoint of any net interval, and thus has unique symbolic representation $[x] = (e_1, \ldots, e_j, \theta, \theta, \ldots)$ where $\theta = (\theta_1, \ldots, \theta_\ell)$. In this situation, Δ_1 has symbolic representation $(e_1, \ldots, e_j, \theta^n)$ and Δ_2 has symbolic representation $(e_1, \ldots, e_j, \theta^{n+1})$ for any $n \in \mathbb{N}$, we have $\Delta_2 \subseteq \Delta_1^\circ$. Thus for any t sufficiently small, there exists some $m \in \mathbb{N}$, such that $\Delta_1 \subseteq \Delta_t(x) \subseteq M_t(x) \subseteq \Delta_2$ where Δ_1 has symbolic representation $(e_1, \ldots, e_j, \theta^m)$ and Δ_2 has symbolic representation $(e_1, \ldots, e_j, \theta^{m+2})$. Similarly as argued in (3.5), there exist constants c_1, c_2 such that $\|T(\theta)^{m+2}\| \leq c_1 \mu(\Delta_t(x)) \leq c_2 \|T(\theta)^m\|$. In addition, since $M_t(x) \subseteq \Delta_1$, we have $\mu(M_t(x)) \leq \mu(\Delta_1)$ and there exist constants c_1', c_2' such that $\|T(\theta)^{m+2}\| \leq c_1' \mu(M_t(x)) \leq c_2' \|T(\theta)^m\|$.

The argument proceeds identically as before. $\qquad\square$

4. Multifractal formalism under the weak separation condition

In this section, we prove the multifractal formalism results under the weak separation condition.

4.1. **Density of local dimensions at periodic points.** We first show that under the weak separation condition periodic points are abundant, in that the set of local dimensions at periodic points is dense in the set of local dimensions in the essential class. This generalizes a result of Hare, Hare and Ng on local dimensions [18, Cor. 3.15] for IFS satisfying substantially stricter conditions. This property can be useful in computing the exact set of possible local dimensions; see, for example, Section 5.3.2 or the discussions of examples in [17, 18, 21].

Theorem 4.1. *Let $\{S_i\}_{i\in\mathcal{I}}$ be an IFS satisfying the weak separation condition and $\mu_{\boldsymbol{p}}$ an associated self-similar measure. Then the set of local dimensions at periodic points is dense in $\{\overline{\dim}_{\mathrm{loc}}(x) : x \in K_{\mathrm{ess}}\}$ and $\{\underline{\dim}_{\mathrm{loc}}(x) : x \in K_{\mathrm{ess}}\}$.*

Proof. Let x be an interior essential point. Either there exists some s_0 such that there is a unique essential net interval $\Delta_0 \in \mathcal{F}_{s_0}$ containing x, or there exists essential net intervals $\Delta_0^{(1)}, \Delta_0^{(2)}$ such that $\{x\} = \Delta_0^{(1)} \cap \Delta_0^{(1)}$. The cases are similar, but the latter is slightly harder, so we treat that here.

Let $t_0 > 0$ be such that $B(x, 2t_0) \subseteq \Delta_0^{(1)} \cup \Delta_0^{(2)}$. Arguing similarly to Lemma 3.9, there exists constants $c, C > 0$ such that for any $0 < t \le t_0$, there exists $\Delta_t^{(1)} \subseteq \Delta_t^{(2)} \subseteq B(x, 2t)$ and for each $k = 1, 2$, we have $\Delta_t^{(k)} \in \mathcal{F}_s$ where $t \ge s \ge ct$,

$$\min\{\boldsymbol{Q}_{\boldsymbol{p}}(\Delta_t^{(k)})_j : 1 \le j \le \#\mathcal{V}(\Delta_t^{(k)})\} \ge C\mu_{\boldsymbol{p}}(B(x,t)),$$

and $\mathcal{V}(\Delta_i^{(k)}) = \mathcal{V}(\Delta_0^{(k)})$. We may also assume that $\Delta_t^{(1)}$ and $\Delta_t^{(2)}$ do not contain x as an endpoint. In particular, for each $0 < t \le t_0$, there exists some $k \in \{1, 2\}$ such that $\Delta_t^{(k)} \subseteq (\Delta_0^{(k)})^\circ$. Set $\Delta_t = \Delta_t^{(k)}$ and let η_t be the path in the transition graph corresponding to $\Delta_t^{(k)} \subseteq \Delta_0^{(k)}$, which is a cycle since the two net intervals have the same neighbour set. Let γ_1 be the symbolic representation of $\Delta_0^{(1)}$ and γ_2 the symbolic representation of $\gamma_0^{(2)}$

For each $0 < t \le t_0$, let x_t be any periodic point with period η_t. We note that since x_t is not the boundary point of any net interval, we have by Proposition 3.15

$$\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x_t) = \frac{\log \operatorname{sp} T(\eta_t)}{\log L(\eta_t)}.$$

Fix t as above, and let $\Delta_0 \in \{\Delta_0^{(1)}, \Delta_0^{(2)}\}$ be such that $x_0 \in \Delta_0^\circ$. Let Δ_0 have symbolic representation γ. By definition of c, we observe that $t \ge \operatorname{tg}(\Delta_t) \ge cr_{\min}t$. Since $\operatorname{tg}(\Delta_t) = L(\gamma)L(\eta_t)$, there exist constants $c_1, c_2 > 0$ (not depending on t) such that

$$c_1 2t \le L(\eta_t) \le c_2 t.$$

We also bound $\operatorname{sp} T(\eta_t)$. Since $\Delta_t \subseteq B(x, 2t)$ has symbolic representation $\gamma\eta_t$, we have $\|T(\gamma\eta_t)\| \leq \mu_{\boldsymbol{p}}(B(x, 2t))$ and since $T(\gamma)$ is a transition matrix, there exists some $C_1 > 0$ such that

$$\operatorname{sp} T(\eta_t) \leq \|T(\eta_t)\| \leq C_1 \mu_{\boldsymbol{p}}(B(x, 2t))$$

(just take C_1 to be the smallest strictly positive entry of $T(\gamma_1)$ and $T(\gamma_2)$). On the other hand, since $\boldsymbol{Q_p}(\Delta_t) = \boldsymbol{Q_p}(\Delta_0)T(\eta_t)$, we have

$$\operatorname{sp} T(\eta_i) \geq \frac{\min\{\boldsymbol{Q_p}(\Delta_t)_j : 1 \leq j \leq \#v\}}{\max\{\boldsymbol{Q_p}(\Delta_0)_j : 1 \leq j \leq \#v\}} \geq \frac{C\mu_{\boldsymbol{p}}(B(x, t))}{\max\{\boldsymbol{Q_p}(\Delta_0^{(k)})_j : 1 \leq j \leq \#v, 1 \leq k \leq 2\}}$$
$$= C_2 \mu_{\boldsymbol{p}}(B(x, t)).$$

To summarize, we have shown that

$$\frac{\log C_2 + \log \mu_{\boldsymbol{p}}(B(x, t))}{\log c_2 + \log t} \geq \dim_{\operatorname{loc}} \mu_{\boldsymbol{p}}(x_t) = \frac{\log \operatorname{sp} T(\eta_t)}{\log L(\eta_t)}$$
$$\geq \frac{\log C_1 + \log \mu_{\boldsymbol{p}}(B(x, 2t))}{\log c_1 + \log 2t}.$$

Let $\alpha = \overline{\dim}_{\operatorname{loc}}\mu_{\boldsymbol{p}}(x)$ and let $\epsilon > 0$ be arbitrary. Get some $t_1 > 0$ such that for all $0 < t \leq t_1$,

$$\frac{\log C_2 + \log \mu_{\boldsymbol{p}}(B(x, t))}{\log c_2 + \log t} \leq \alpha + \epsilon$$

and then choose $0 < t \leq \min\{t_0, t_1\}$ such that

$$\frac{\log C_1 + \log \mu_{\boldsymbol{p}}(B(x, 2t))}{\log c_1 + \log 2t} \geq \alpha - \epsilon.$$

Since $\epsilon > 0$ was arbitrary, it follows that the set of local dimensions at periodic points is dense in $\{\overline{\dim}_{\operatorname{loc}}(x) : x \in K_{\operatorname{ess}}\}$. The result for lower local dimensions holds identically. $\qquad\square$

4.2. **The L^q-spectrum, dimension spectrum, and multifractal formalism.** In this section, we show how to extend a result of Feng and Lau [12] to hold with respect to a larger, more natural class of intervals.

Let μ be a compactly supported finite Borel measure and let $V \subseteq \mathbb{R}$ be any open set with $\mu(V) > 0$. Then the L^q-spectrum of μ on V, denoted by $\tau_V(\mu, q)$, is given by

$$\tau_V(\mu, q) = \liminf_{t \downarrow 0} \frac{\log \sup \sum_i \mu(B(x_i, t))^q}{\log t}$$

where the supremum is over families of disjoint closed balls $\{B(x_i, t)\}_i$ with $x_i \in K$ and $B(x_i, t) \subseteq V$. A direct application of Hölder's inequality shows that $\tau_V(q)$ is a concave function. When $V = \mathbb{R}$, we write $\tau(\mu, q) = \tau_{\mathbb{R}}(\mu, q)$.

Since $\tau_V(\mu, q)$ is a concave function in q, its *concave conjugate* is given by

$$\tau_V^*(\mu, \alpha) := \inf\{\alpha q - \tau_V(q) : q \in \mathbb{R}\}.$$

We set

$$D_V(\mu) = \{\alpha \in \mathbb{R} : \dim_{\mathrm{loc}} \mu(x) = \alpha \text{ for some } x \in K \cap V\}$$

and

$$K_V(\mu, \alpha) = \{x \in K \cap V : \dim_{\mathrm{loc}} \mu(x) = \alpha\}.$$

Understanding the geometric properties of the sets $K_V(\mu, \alpha)$ is a natural way to understand the structure of μ.

A heuristic relationship between the values if $\dim_H K_V(\mu, \alpha)$ and the concave conjugate of L^q-spectrum, known as the multifractal formalism, has been studied by many authors (see, for example, [3, 7, 10, 12, 13, 16, 26, 27, 32, 33, 36]).

Definition 4.2. *Let μ be a compactly supported finite Borel measure and let $V \subset \mathbb{R}$ have $\mu(V) > 0$. We say that the measure μ satisfies the* complete multifractal formalism *with respect to V if*

(i) $D_V(\mu) = [\alpha_{\min}, \alpha_{\max}]$ where

$$\alpha_{\min} = \lim_{q \to +\infty} \frac{\tau_V(q)}{q} \qquad\qquad \alpha_{\max} = \lim_{q \to -\infty} \frac{\tau_V(q)}{q}.$$

(ii) For any $\alpha \in [\alpha_{\min}, \alpha_{\max}]$, $\tau_V^*(\alpha) = \dim_H K_V(\alpha)$.

Note here that we do not comment on differentiability of $\tau_V(q)$. The following result is contained in [12, Thm. 5.4]:

Proposition 4.3 ([12])**.** *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition and let $\mu_{\boldsymbol{p}}$ be an associated self-similar measure. Let U_0 be any open ball which attains the maximal value in (3.1). Then $\mu_{\boldsymbol{p}}$ satisfies the complete multifractal formalism with respect to U_0.*

4.3. **Multifractal formalism for the essential class.** Using the notion of the essential class, we can obtain a strictly stronger extension of this proposition. We first note the following straightforward lemma:

Lemma 4.4 ([12])**.** *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition and let $\mu_{\boldsymbol{p}}$ be an associated self-similar measure. Let U_0 be any open ball which attains the maximal value in (3.1). Then if $\sigma \in \mathcal{I}^*$ is arbitrary,*

(i) $\tau_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}, q) = \tau_{U_0}(q)$,
(ii) $D_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}) = D_{U_0}(\mu_{\boldsymbol{p}})$, and
(iii) $\dim_H K_{U_0}(\mu_{\boldsymbol{p}}, \alpha) = \dim_H K_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}, \alpha)$.

Proof. Statement (i) is [12, Cor. 5.6]. Statements (ii) and (iii) are implicit in the usage of [12, Lem. 2.5]. $\qquad\square$

We obtain the following extension of Proposition 4.3. In light of Proposition 3.7 and the remark which follows it, our result is strictly stronger.

Theorem 4.5. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition and let $\mu_{\boldsymbol{p}}$ be a self-similar measure. Let $\Delta_1, \ldots, \Delta_n$ be any essential net intervals and let $\nu = \mu_{\boldsymbol{p}}|_{\Delta_1 \cup \cdots \cup \Delta_n}$. Then*

(i) ν satisfies the complete multifractal formalism,
(ii) the set

$$P(\mu_{\boldsymbol{p}}) := \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}, x \text{ periodic}\}$$

is dense in $D(\nu)$, and
(iii) the sets of local dimensions satisfy

$$D(\nu) = \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}, \dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) \text{ exists}\}$$
$$= \{\underline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\} = \{\overline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}.$$

Moreover, the values of $\tau(\nu, q)$ do not depend on the choice of $\Delta_1, \ldots, \Delta_n$ and for $q \geq 0$, $\tau(\mu_{\boldsymbol{p}}, q) = \tau(\nu, q)$.

Proof. We split the proof into two parts, for clarity.

Part 1. *The statement (i) holds, the values of $\tau(\nu, q)$ do not depend on the choice of $\Delta_1, \ldots, \Delta_n$, and for $q \geq 0$, $\tau(\mu_{\boldsymbol{p}}, q) = \tau(\nu, q)$.*

Let U_0 be an open ball which attains the maximal value in (3.1).

To verify (i), by Proposition 4.3, it suffices to show that

$$\tau_{U_0}(\mu_{\boldsymbol{p}}, q) = \tau(\nu, q) \qquad D_{U_0}(\mu_{\boldsymbol{p}}) = D(\nu) \qquad \dim_H K_{U_0}(\mu_{\boldsymbol{p}}, \alpha) = \dim_H K(\nu, \alpha).$$

Set $E = \Delta_1 \cup \cdots \cup \Delta_n$. Let σ be such that $S_\sigma(U_0) \subseteq E$, and we see directly from the definitions and Lemma 4.4 that

$$\tau_{U_0}(\mu_{\boldsymbol{p}}, q) = \tau_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}, q) \geq \tau(\nu, q)$$
$$D_{U_0}(\mu_{\boldsymbol{p}}) = D_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}) \subseteq D(\nu)$$
$$\dim_H K(\nu, \alpha) \geq \dim_H K_{S_\sigma(U_0)}(\mu_{\boldsymbol{p}}, \alpha) = \dim_H K_{U_0}(\mu_{\boldsymbol{p}}, \alpha).$$

We now establish the reverse inequalities.

That $\tau(\mu_{\boldsymbol{p}}, q) = \tau_{U_0}(q) = \tau(\nu, q)$ for $q \geq 0$ is straightforward; see, for example, [12, Prop. 3.1].

Since the Δ_i are net intervals, for any i, j we must have one of $\Delta_i \subseteq \Delta_j$, $\Delta_j \subseteq \Delta_i$, or $\Delta_i^\circ \cap \Delta_j^\circ = \emptyset$. We may thus assume that $\Delta_i^\circ \cap \Delta_j^\circ = \emptyset$ for $i \neq j$. Since U_0 is open and the Δ_i are essential, there exist net intervals $\Delta_1^*, \ldots, \Delta_n^*$ such that $\mathcal{V}(\Delta_i) = \mathcal{V}(\Delta_i^*)$ for each $1 \leq i \leq n$, the Δ_i^* are pairwise disjoint, and $E^* := \Delta_1^* \cup \cdots \cup \Delta_n^* \subseteq U_0$. Let $\nu^* := \mu_{\boldsymbol{p}}|_{E^*}$. Again, it is clear from the definitions that $\tau_{U_0}(\mu_{\boldsymbol{p}}, q) \leq \tau(\nu^*, q)$. It remains to show that $\tau(\nu^*, q) \leq \tau(\nu, q)$ for $q < 0$. By Lemma 2.3, get similarities $g_i : \Delta_i \cap K \to \Delta_i^* \cap K$ and some $c > 0$ such that if $E \subseteq \Delta_i$ is an arbitrary Borel set,

$$(4.1) \qquad\qquad c_1 \nu^*(g_i(E)) \leq \nu(E) \leq c_2 \nu^*(g_i(E)).$$

Let each g_i have contraction ratio ρ_i. Set

$$t_0 = \frac{\min\{\mathrm{dist}(\Delta_i^*, \Delta_j^*) : i \neq j\}}{2 \cdot \max\{\rho_i : 1 \leq i \leq n\}}$$
$$r_0 = \min\{\rho_i : 1 \leq i \leq n\}$$

and let $0 < t < t_0$ be arbitrary. Suppose $\{B(x_j, t)\}_{j=1}^m$ is an arbitrary family of disjoint closed balls where $x_j \in E \cap K$. For each j, there is some $i(j)$ such that $x_j \in \Delta_{i(j)} \cap K$ so that $g_{i(j)}(x_j) \in \Delta_{i(j)}^* \cap K$. By choice of $t < t_0$, we have for $j \neq k$ that $g_{i(j)}(B(x_j, t)) \cap g_{i(k)}(B(x_k, t)) = \emptyset$ and $g_{i(j)}(B(x_k, t)) \cap E^* \subseteq \Delta_{i(j)} \cap K$. In particular, we see for each $1 \leq j \leq m$

$$\nu(B(x_j, t)) \geq \nu(B(x_j, t) \cap \Delta_{i(j)}) \geq c_1 \nu^*(B(g_{i(j)}(x_j), \rho_{i(j)} t))$$
$$\geq c_1 \nu^*(B(g_{i(j)}(x_j), r_0 t))$$

so that $\nu(B(x_j, t))^q \leq c_1^q \nu^*(B(x_j^*, r_0 t))^q$ where $x_j^* = g_{i(j)}(x_j)$. But $\{B(x_j, t)\}_{j=1}^m$ was an arbitrary cover, so that

$$\frac{\log \sup \sum_j \nu(B(x_j, t))^q}{\log t} \geq \frac{\log c_1^q + \log \sup \sum_j \nu^*(B(x_j^*, r_0 t))^q}{\log r_0^{-1} + \log r_0 t}.$$

Taking limits, it follows that $\tau(\nu, q) \geq \tau(\nu^*, q)$ for $q < 0$.

We now see that $D(\nu) \subseteq D_{U_0}(\mu_{\boldsymbol{p}})$. First note that $D_{U_0}(\mu_{\boldsymbol{p}}) = [\alpha_{\min}, \alpha_{\max}]$ where

$$\alpha_{\min} = \lim_{q \to +\infty} \frac{\tau(\nu, q)}{q} = \lim_{q \to +\infty} \frac{\tau_{U_0}(\mu_{\boldsymbol{p}}, q)}{q} \quad \alpha_{\max} = \lim_{q \to -\infty} \frac{\tau(\nu, q)}{q} = \lim_{q \to -\infty} \frac{\tau_{U_0}(\mu_{\boldsymbol{p}}, q)}{q},$$

since $\tau(\nu, q) = \tau_{U_0}(\mu_{\boldsymbol{p}}, q)$. Let $x \in \operatorname{supp} \nu$ be arbitrary with $\alpha = \dim_{\mathrm{loc}} \nu(x)$. Then for any $q \in \mathbb{R}$ and $t > 0$, we have

$$\log \sup \sum_i \nu(B(x_i, t))^q \geq \log \nu(B(x, t))^q$$

where the supremum is over disjoint balls $B(x_i, t)$ with $x_i \in \operatorname{supp} \nu$, and therefore $\tau(\nu, q) \leq q\alpha$. Since $\tau(\nu, q)$ is concave, it follows that $\alpha \in [\alpha_{\min}, \alpha_{\max}] = D_{U_0}(\mu_{\boldsymbol{p}})$. Note that this argument also shows that $D(\nu) = D(\nu^*)$.

Finally, to verify that $\dim_H K(\nu, \alpha) \leq \dim_H K_{U_0}(\alpha)$, it suffices to show that

$$\dim_H K(\nu, \alpha) \leq \dim_H K(\nu^*, \alpha).$$

First note by (4.1) that if $x \in \Delta_i^\circ \cap K$ for some i, then $g_i(x) \in \Delta_i^* \cap K$ has $\dim_{\mathrm{loc}} \nu(x) = \dim_{\mathrm{loc}} \nu^*(g_i(x))$. Thus $g_i(K(\nu, \alpha) \cap \Delta_i^\circ) \subseteq K(\nu^*, \alpha)$ and

$$\dim_H K(\nu, \alpha) \cap \bigcup_{i=1}^n \Delta_i^\circ \leq \dim_H K(\nu^*, \alpha).$$

Since $D(\nu) = D(\nu^*)$ and $E \setminus \bigcup_{i=1}^n \Delta_i^\circ$ is a finite set (and hence has Hausdorff dimension 0), the result follows.

Thus the complete multifractal formalism holds.

Since U_0 was arbitrary, it is clear that $\tau(\nu, q)$ does not depend on the choice of $\Delta_1, \ldots, \Delta_n$.

Part 2. *Statements (ii) and (iii) hold.*

To see that

(4.2) $D(\nu) = \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}, \dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) \text{ exists}\},$

if $x \in K_{\mathrm{ess}}$ is arbitrary, there exists essential net intervals Δ_1, Δ_2 with $x \in (\Delta_1 \cup \Delta_2)^{\circ}$ and

$$\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) = \dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}|_{\Delta_1 \cup \Delta_2} \in [\alpha_{\min}, \alpha_{\max}].$$

Conversely, if $x \in D(\nu)$, then there exists some $y \in U_0$ such that $\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(y) = \dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x)$. But $U_0 \subseteq K_{\mathrm{ess}}$ by Proposition 3.7, so that (4.2) follows.

By Theorem 4.1, we have that

$$P(\mu_{\boldsymbol{p}}) = \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}, x \text{ periodic}\}$$

is dense in the set of upper and lower local dimensions in K_{ess}. Now $P(\mu_{\boldsymbol{p}}) \subseteq D(\nu)$ from (4.2) and $D(\nu) = [\alpha_{\min}, \alpha_{\max}]$ is a closed set with $D(\nu) \subseteq \{\overline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}$. But again, Theorem 4.1 shows that $P(\mu_{\boldsymbol{p}})$ is a dense subset of $\{\overline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}$, forcing

$$D(\nu) = \{\overline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}.$$

Of course, we also have $D(\nu) = \{\underline{\dim}_{\mathrm{loc}}\mu_{\boldsymbol{p}}(x) : x \in K_{\mathrm{ess}}\}$ by the same argument, finishing the proof of the theorem. □

Corollary 4.6. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the weak separation condition with associated self-similar measure $\mu_{\boldsymbol{p}}$. Then there exists a sequence of non-empty compact sets $(K_m)_{m=1}^{\infty}$ with $K_m \subseteq K_{m+1} \subseteq K$ for each $m \in \mathbb{N}$ such that*
   (i) $\lim_{m \to \infty} \mu_{\boldsymbol{p}}(K_m) = 1$,
   (ii) each $\mu_m := \mu_{\boldsymbol{p}}|_{K_m}$ satisfies the complete multifractal formalism, and
   (iii) $\tau(\mu_m, q)$ and $D(\mu_m)$ do not depend on the index m.

Proof. Let $(t_m)_{m=1}^{\infty}$ be any sequence monotonically decreasing to 0 where t_1 is the first generation such that there is some essential net interval $\Delta \in \mathcal{F}_m$. For each m, set

$$K_m := \bigcup_{\substack{\Delta \in \mathcal{F}_{t_m} \\ \mathcal{V}(\Delta) \in V(\mathcal{G}_{\mathrm{ess}})}} \Delta \cap K,$$

i.e. each K_m is the set of points in K contained in some essential net interval of generation m. Since any descendent of an essential net interval is also essential, we necessarily have $K_m \subseteq K_{m+1}$ for each $m \in \mathbb{N}$. By Theorem 3.10, since $K_{\mathrm{ess}} \subseteq \bigcup_{m=1}^{\infty} K_m$ and $\mu_{\boldsymbol{p}}(K_{\mathrm{ess}}) = 1$, we have (i). By Theorem 4.5, each $\mu_m := \mu_{\boldsymbol{p}}|_{K_m}$ satisfies the complete multifractal formalism and $\tau(\mu_m, q)$ and $D(\mu_m)$ do not depend in the index m, giving (ii) and (iii). □

In some situations, the above theorem can also be used to verify that the complete multifractal formalism holds with respect to the invariant measure $\mu_{\boldsymbol{p}}$.

Corollary 4.7. *Suppose $\{S_i\}_{i \in \mathcal{I}}$ is an IFS satisfying the weak separation condition with transition graph $\mathcal{G}$. Suppose there is a bound on the maximum length of a path with no vertices in the essential class. Then if $\mu_{\boldsymbol{p}}$ is any associated self-similar measure, $\mu_{\boldsymbol{p}}$ satisfies*

the complete multifractal formalism and the local dimensions at periodic points are dense in the set of all local dimensions in K.

Proof. If M is the bound on the maximum length of a path, since $L(e) \geq r_{\min}$ for any $e \in E(\mathcal{G})$, we have that any net interval in $\mathcal{F}_{r_{\min}^M}$ is an essential net interval. In particular, $\operatorname{supp} \mu_{\boldsymbol{p}}$ is contained in a finite union of essential net intervals. Apply Theorem 4.5. $\square$

Remark 4.8. *For example, if the neighbour set $\mathcal{V}([0,1]) = \{x \mapsto x\}$ is contained in the essential class, then $\mathcal{G} = \mathcal{G}_{\mathrm{ess}}$ and the conditions for the Corollary 4.7 are satisfied.*

Corollary 4.9. *Suppose $\{S_i\}_{i \in \mathcal{I}}$ is an IFS such that the associated transition graph $\mathcal{G}$ is finite. Suppose that any cycle in $\mathcal{G}$ is contained in the essential class. Then if $\mu_{\boldsymbol{p}}$ is any associated self-similar measure, $\mu_{\boldsymbol{p}}$ satisfies the complete multifractal formalism and the local dimensions at periodic points are dense in the set of all local dimensions in K.*

Proof. When $\mathcal{G}$ is finite, it is clear that the assumption in Corollary 4.7 is equivalent to the assumption that any cycle is contained in the essential class. $\square$

5. THE FINITE NEIGHBOUR CONDITION AND EXAMPLES

5.1. **The finite neighbour condition.** Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS as in (2.1). The finite neighbour condition was defined in [19] in a way following naturally from the finite type conditions studied in the literature [28, 31].

Definition 5.1. *We say that $\{S_i\}_{i \in \mathcal{I}}$ satisfies the* finite neighbour condition *if there are only finitely many neighbour sets. Equivalently, its transition graph $\mathcal{G}$ is finite.*

Remark 5.2. *The definition of neighbour in Definition 2.1 differs slightly from [19, Def'n. 2.7]. Namely, for a net interval $\Delta \in \mathcal{F}$, we require $T(K) \cap (0,1) \neq \emptyset$ rather than $T([0,1]) \supseteq [0,1]$. However, using [5, Cor. 3.4] with respect to the generation $k_0 := r_{\min}/M$ where $M = \sup_{\Delta \in \mathcal{F}} R_{\max}(\Delta)$ and the characterization [19, Thm 3.4.], one can verify that the finiteness assumptions are in fact equivalent.*

It is shown in [19] that the finite neighbour condition is equivalent to the generalized finite type condition [28] holding with respect to the invariant open set $(0,1)$. Moreover, under the assumption that the attractor K is an interval, it is proven in [19] that the finite neighbour condition is in fact equivalent to the weak separation condition. The author is not aware of any IFS of similarities in $\mathbb{R}$ which satisfies the weak separation condition but not the finite neighbour condition.

Of course, when an IFS satisfies the finite neighbour condition, it also satisfies the weak separation condition (see, for example, [28, Thm. 1.1] or [19, Thm. 3.7] and thus has a unique finite essential class $\mathcal{G}_{\mathrm{ess}}$. Interestingly, the converse also holds:

Theorem 5.3. *The IFS $\{S_i\}_{i \in \mathcal{I}}$ satisfies the finite neighbour condition if and only if $\mathcal{G}(\{S_i\}_{i \in \mathcal{I}})$ has a finite essential class.*

Proof. ($\Longrightarrow$) Since the finite neighbour condition implies the weak separation condition, this follows immediately from Proposition 3.5 since $\mathcal{G}$ is a finite graph.

($\Longleftarrow$) We first define a construction on neighbour sets. Let $v_1 = \{f_1, \dots, f_{\ell_1}\}$ and $v_2 = \{g_1, \dots, g_{\ell_2}\}$ be a pair of neighbour sets. We denote by $J(v_1, v_2)$ the set of all subsets $w = \{h_1, \dots, h_m\}$ such that there exist indices i, j and $T = f_i \circ g_j^{-1}$ such that

$$\{T_\Delta \circ h_1, \dots, T_\Delta \circ h_m\} \subset \{f_1, \dots, f_{\ell_1}\}$$

where $\Delta = \big[\min\{0, T(0), T(1)\}, \max\{1, T(0), T(1)\}\big]$ and $T_\Delta(x) = rx + d$ with $r > 0$ where $T_\Delta([0,1]) = \Delta$. Clearly there are only finitely many functions T, so that $J(v_1, v_2)$ is a finite set. When F is a finite set, we denote by $J(F) = \bigcup_{v_1, v_2 \in F} J(v_1, v_2)$, which is also finite.

Now, by assumption, $\mathcal{G}$ has a finite essential class $\mathcal{G}_{\mathrm{ess}}$ so that $J_0 := J(V(\mathcal{G}_{\mathrm{ess}}))$ is finite. Let $\Delta_0 \in \mathcal{F}_\alpha$ be an arbitrary net interval; we will see that $\mathcal{V}(\Delta_0) \in J_0$, from which it follows that $\{S_i\}_{i \in \mathcal{I}}$ satisfies the finite neighbour condition.

First, let σ be such that $r_\sigma > 0$ and $S_\sigma([0,1])$ is a finite union of essential net intervals (just take σ such that $S_\sigma([0,1])$ is contained in some essential net interval; if $r_\sigma < 0$, append some $i \in \mathcal{I}$ with $r_i < 0$). Let $\mathcal{V}(\Delta_0)$ have neighbours generated by words $\{\omega_1, \dots, \omega_m\}$ in Λ_α; note that each $\sigma\omega_i \in \Lambda_{r_\sigma \alpha}$. Let $\Delta_1 = S_\sigma(\Delta_0)$ and write $\Delta_1 = [a, b]$. Then there exist essential net intervals $\Delta_a, \Delta_b \in \mathcal{F}_{r_\sigma \alpha}$ such that $\Delta_a = [a, a_0]$ and $\Delta_b = [b_0, b]$; perhaps $\Delta_a = \Delta_b$. Note that $\sigma\omega_1$ has $\Delta_a, \Delta_b \subseteq S_{\sigma\omega_1}([0,1])$ since $\Delta_a, \Delta_b \subseteq \Delta_1$ so that $\sigma\omega_1$ generates a neighbour f_a of Δ_a and f_b of Δ_b.

We see that $\mathcal{V}(\Delta_0)$ is a join of $(\mathcal{V}(\Delta_a), \mathcal{V}(\Delta_b))$. Set $T = f_a \circ f_b^{-1}$. We first note that

- $T_{\Delta_a} \circ f_a = S_{\sigma\omega_1} = T_{\Delta_b} \circ f_b$, so that $T := f_a \circ f_b^{-1} = T_{\Delta_b} \circ T_{\Delta_a}^{-1}$ and
- $\Delta := \big[\min\{0, T(0), T(1)\}, \max\{1, T(0), T(1)\}\big] = T_{\Delta_a}^{-1}(\Delta_1)$ so that $T_\Delta = T_{\Delta_a}^{-1} \circ T_{\Delta_1}$.

Now let $h \in \mathcal{V}(\Delta_0)$ be arbitrary. Since $r_\sigma > 0$, $T_{\Delta_1} = S_\sigma \circ T_{\Delta_0}$. Then if $h = T_{\Delta_0}^{-1} \circ S_{\omega_i}$, we have

$$T_\Delta \circ h = (T_{\Delta_a}^{-1} \circ T_{\Delta_1}) \circ (T_{\Delta_1}^{-1} \circ S_\sigma \circ S_{\omega_i}) = T_{\Delta_a}^{-1} \circ S_{\sigma\omega_i},$$

where $\sigma\omega_i$ generates a neighbour of Δ_a, and thus $T_\Delta \circ h \in \mathcal{V}(\Delta_a)$, as required. $\square$

Remark 5.4. *If $v, w \in V(\mathcal{G}_{\mathrm{ess}})$, then there are at most $\#v \cdot \#w$ distinct functions T, so that $\#J(v, w) \leq \#v \cdot \#w \cdot 2^{\#v}$. Moreover, there are at most $(\#V(\mathcal{G}_{\mathrm{ess}}))^2$ pairs (v, w). In particular, if there are m distinct neighbours in $\mathcal{G}_{\mathrm{ess}}$, then $\#V(\mathcal{G}_{\mathrm{ess}}) \leq 2^m$ and $\#v \leq m$ for any $v \in V(\mathcal{G}_{\mathrm{ess}})$, so that*

$$\#V(\mathcal{G}) \leq (\#V(\mathcal{G}_{\mathrm{ess}}))^2 \cdot m^2 \cdot 2^m \leq m^2 8^m.$$

Thus the above proof gives a quantitative bound on the size of $\mathcal{G}$ as a function of the number of distinct neighbours in $\mathcal{G}_{\mathrm{ess}}$.

5.2. **Approximate transition matrices.** Under the finite neighbour condition, we may approximate the transition matrix $T(e)$ by the matrix $T^*(e)$ given by $T^*(e)_{ij} = p_\ell$ in the same context as (2.4). Since there are only finitely many values $\frac{\mu_{\boldsymbol{p}}(f_i^{-1}((0,1)))}{\mu_{\boldsymbol{p}}(g_j^{-1}((0,1)))}$, there exist constants $c_1, c_2 > 0$ such that $c_1 T^*(\eta) \leq T(\eta) \leq c_2 T^*(\eta)$ element-wise for any

admissible path η. Moreover, since $\mu_{\boldsymbol{p}}$ is a probability measure, direct computation shows that $\|T^*(\eta)\|_1 \leq \mu_{\boldsymbol{p}}(\Delta)$. Applying Theorem 2.12, we have:

Corollary 5.5. *Let $\{S_i\}_{i \in \mathcal{I}}$ be an IFS satisfying the finite neighbour condition with associated self-similar measure $\mu_{\boldsymbol{p}}$.*

- *There exist constants $c_1, c_2 > 0$ such that for any path η realized by $(\Delta_i)_{i=0}^n$,*

$$c_1 \boldsymbol{Q_p}(\Delta_m) \preccurlyeq T^*(\eta) \boldsymbol{Q_p}(\Delta_0) \preccurlyeq c_2 \boldsymbol{Q_p}(\Delta_m)$$

  where the inequalities hold pointwise.
- *There exists a constant $c > 0$ such that for any $\Delta \in \mathcal{F}$ with symbolic representation η,*

$$c\mu_{\boldsymbol{p}}(\Delta) \leq \|T^*(\eta)\|_1 \leq \mu_{\boldsymbol{p}}(\Delta).$$

One may also observe that the same principle works for periodic points. We have the natural analogue of Proposition 3.15:

Corollary 5.6. *Let $\{S_i\}_{i \in \mathcal{I}}$ be any IFS and suppose x is a periodic point with period $\theta = (e_1, \ldots, e_s)$. Then the local dimension of μ at x exists and is given by*

$$\dim_{loc} \mu(x) = \frac{\log \mathrm{sp}(T^*(\theta))}{\log L(\theta)}$$

where if x is a boundary point of a net interval with two different symbolic representations given by periods θ and ϕ, then θ is chosen to satisfy

$$\frac{\log \mathrm{sp}(T^*(\theta))}{\log L(\theta)} \geq \frac{\log \mathrm{sp}(T^*(\phi))}{\log L(\phi)}.$$

Proof. The proof is identical to the proof of Proposition 3.15, noting that the analogue of Corollary 5.5 holds since the set $\{\mathcal{V}(\Delta) : x \in \Delta, \Delta \in \mathcal{F}\}$ is finite. $\qquad\square$

5.3. **An overlapping IFS with non-commensurable contraction ratios.** Consider the IFS given by the maps

$$S_1(x) = \rho \cdot x \qquad S_2(x) = r \cdot x + \rho(1 - r) \qquad S_3(x) = r \cdot x + 1 - r$$

where $0 < \rho, r < 1$ satisfy $\rho + 2r - \rho r \leq 1$, i.e. $S_2(1) \leq S_3(0)$. This IFS was initially studied by [29] and was the first example of an iterated function system with overlaps and satisfying the weak separation condition without commensurable contraction ratios. It is known that the Hausdorff dimension of the attractor K is the unique solution to the equation $\rho^s + 2r^s - (\rho r)^s = 1$ (see [29, Prop. 4.9] or [28, Ex. 5.1]).

Under the assumption that $\rho > r > \rho^2$, we will compute the neighbour sets and the transition graph. We also give formulas to compute the range of local dimensions. We will also show (for any valid parameters r, ρ) that any associated self-similar measure satisfies the complete multifractal formalism.

5.3.1. *Neighbour sets and the transition graph.* We first compute the neighbour sets and children in complete detail. The net interval Δ_0 has $\mathcal{V}(\Delta_0) = \{x \mapsto x\}$ and $\mathrm{tg}(\Delta_0) = 1 = m(\Delta_0) \cdot 1$ since 1 is the maximal contraction ratio of any of its neighbours. Thus Δ_0 has children

$$(\Delta_1 = [0, \rho(1-r)], \Delta_2 = [\rho(1-r), \rho], \Delta_3 = [\rho, \rho + r - \rho r], \Delta_4 = [1-r, r])$$

in $\mathcal{F}_1$. Note that when $\rho + 2r - \rho r < 1$, $[\rho + r - \rho r, 1 - r]$ is not a net interval since its interior does not intersect K. One may compute

$$\mathcal{V}(\Delta_1) = \{x \mapsto x/(1-r)\} \qquad \mathcal{V}(\Delta_2) = \{x \mapsto x/\rho, x \mapsto x/r + \frac{1}{r} - 1\}$$

$$\mathcal{V}(\Delta_3) = \{x \mapsto \frac{x}{1-\rho} + \frac{\rho}{1-\rho}\} \qquad \mathcal{V}(\Delta_4) = \{x \mapsto x\}.$$

Since $\mathcal{V}(\Delta_4) = \mathcal{V}(\Delta_0)$, the children of Δ_4 are scaled versions of the children of Δ_0 and have the same neighbour sets by Theorem 2.8.

- Since $\rho > r$, Δ_1 has $\mathrm{tg}(\Delta_1) = m(\Delta_1) \cdot (1/(1-r)) = \rho$, so Δ_1 has children

  $$(\Delta_5 = [0, \rho^2(1-r)], \Delta_6 = [\rho^2(1-r), \rho^2], \Delta_7 = [\rho^2, \rho(\rho + r - \rho r)])$$

  where $\mathcal{V}(\Delta_5) = \mathcal{V}(\Delta_1)$, $\mathcal{V}(\Delta_6) = \mathcal{V}(\Delta_2)$, and $\mathcal{V}(\Delta_7) = \mathcal{V}(\Delta_3)$.
- Δ_2 has $\mathrm{tg}(\Delta_2) = \rho$ and one child $\Delta_8 = [\rho - \rho r, \rho]$ with $\mathcal{V}(\Delta_8) = \{x \mapsto x, x \mapsto x/\rho\}$. Note that $\Delta_8 = \Delta_2$, but $\mathcal{V}(\Delta_8) \neq \mathcal{V}(\Delta_2)$.
- Δ_3 has $\mathrm{tg}(\Delta_3) = r$ and two children $(\Delta_9 = [\rho, \rho + r^2 - \rho r^2], \Delta_{10} = [r - r^2, r])$ with $\mathcal{V}(\Delta_9) = \mathcal{V}(\Delta_3)$ and $\mathcal{V}(\Delta_{10}) = \mathcal{V}(\Delta_0)$.
- Δ_8 has children $\Delta_{11} = [\rho - \rho r, \rho - \rho r^2], \Delta_{12} = [\rho - \rho r^2, \rho]$ with $\mathcal{V}(\Delta_{11}) = \mathcal{V}(\Delta_1)$ and $\mathcal{V}(\Delta_{12}) = \mathcal{V}(\Delta_2)$.

Thus by Theorem 2.8, there are no new neighbour sets and the IFS satisfies the finite neighbour condition.

For simplicity, fix $v_0 = \mathcal{V}(\Delta_0)$, $v_1 = \mathcal{V}(\Delta_1)$, $v_2 = \mathcal{V}(\Delta_2)$, $v_3 = \mathcal{V}(\Delta_3)$ and $v_4 = \mathcal{V}(\Delta_8)$. Let $\mu_{\boldsymbol{p}}$ be a self-similar measure associated with the IFS, where $\boldsymbol{p} = (p_1, p_2, p_3)$. Observing that v_4 has exactly one child, we can construct an equivalent transition graph by removing v_4, concatenating the incoming edges with the outgoing edge, and multiplying the corresponding edge lengths and transition matrices. This results in the modified transition graphs and edge lengths, described in Figure 1.

5.3.2. *The attainable local dimensions.* We see that the conditions for Corollary 4.7 are satisfied, so that the measure $\mu_{\boldsymbol{p}}$ satisfies the complete multifractal formalism and that the local dimensions at periodic points are dense in the set of upper and lower local dimensions.

We now compute the range of local dimensions at periodic points. We first make note of the following obvious inequality: if $0 < a, b, c, d$ and $\log a / \log b \leq \log c / \log d$, then

(5.1) $$\frac{\log a}{\log b} \leq \frac{\log ac}{\log bd} \leq \frac{\log c}{\log d}.$$

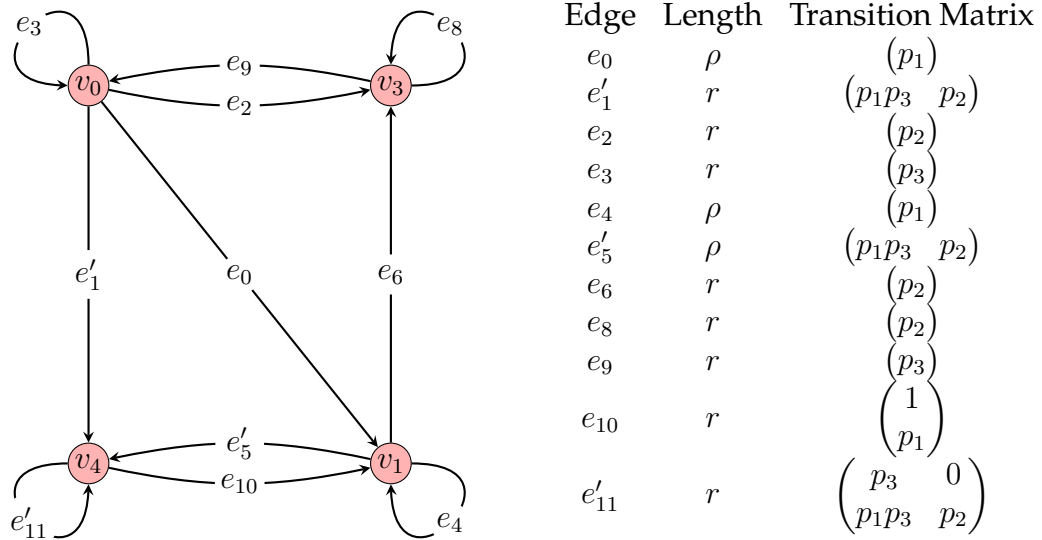| Edge	Length	Transition Matrix
e_0	ρ	(p_1)
e_1'	r	$(p_1 p_3 \quad p_2)$
e_2	r	(p_2)
e_3	r	(p_3)
e_4	ρ	(p_1)
e_5'	ρ	$(p_1 p_3 \quad p_2)$
e_6	r	(p_2)
e_8	r	(p_2)
e_9	r	(p_3)
e_{10}	r	$\begin{pmatrix} 1 \\ p_1 \end{pmatrix}$
e_{11}'	r	$\begin{pmatrix} p_3 & 0 \\ p_1 p_3 & p_2 \end{pmatrix}$

FIGURE 1. Modified transition graph with edge lengths and transition matrices

Now let η be any cycle contained in $\mathcal{G}$. If η only passes through v_4, since $\operatorname{sp} T^*(e_{11}') = \max\{p_2, p_3\}$, the local dimension corresponding to the cycle (e_{11}') is $\frac{\log \max\{p_2, p_3\}}{\log r}$. Otherwise, η passes through some vertex other than v_4. Thus without loss of generality, η begins and ends and some vertex $v \neq v_4$. Suppose η visits some vertex $w \neq v_4$ twice, i.e. $\eta = \eta_1 \eta_2 \eta_3$ where η_1 is a path from v to w, η_2 is a cycle from w to w, and η_3 is a path from w to v. Then η can be written as a concatenation of cycles η_2 and $\eta_3 \eta_1$, where $T(\eta_2)$ and $T(\eta_3 \eta_1)$ are singletons, and by (5.1), we have that

$$\min\left\{\frac{\log \operatorname{sp} T(\eta_2)}{\log L(\eta_2)}, \frac{\log \operatorname{sp} T(\eta_3 \eta_1)}{\log L(\eta_3 \eta_1)}\right\} \leq \frac{\log \operatorname{sp} T(\eta)}{\log L(\eta)}$$
$$\leq \max\left\{\frac{\log \operatorname{sp} T(\eta_2)}{\log L(\eta_2)}, \frac{\log \operatorname{sp} T(\eta_3 \eta_1)}{\log L(\eta_3 \eta_1)}\right\}.$$

In other words, the minimum and maximum local dimensions on cycles are attained at cycles which do not repeat any vertex other than v_4. Thus it suffices to consider all such families of cycles.

If η does not pass through v_4, the only non-repeating cycles are (e_3), (e_4), (e_8), and (e_2, e_9). We thus see that the maximum and minimum possible local dimensions are attained at the points in

$$S = \left\{\frac{\log p_1}{\log \rho}, \frac{\log p_2}{\log r}, \frac{\log p_3}{\log r}\right\}.$$

Otherwise, η passes through v_4. A straightforward induction argument shows that

$$T^*(e'_{11})^n = \begin{cases} \begin{pmatrix} p_3^n & 0 \\ \frac{p_1 p_3(p_2^n - p_3^n)}{p_2 - p_3} & p_2^n \end{pmatrix} & : p_2 \neq p_3 \\ \begin{pmatrix} p^n & 0 \\ np^n p_1 & p^n \end{pmatrix} & : p_2 = p_3 =: p \end{cases}.$$

Now, let

$$\eta_{1,n} = (e_6, e_9, e'_1, \underbrace{e'_{11}, \ldots, e'_{11}}_{n}, e_{10}) \qquad\qquad \eta_{2,n} = (e'_5, \underbrace{e'_{11}, \ldots, e'_{11}}_{n}, e_{10})$$

denote the two possible families of cycles which go through v_4 and do not repeat a vertex not in v_4. We then have that

$$a_n := \mathrm{sp}\, T^*(\eta_{1,n}) = \begin{cases} \frac{p_1 p_2 p_3 (p_2^{n+2} - p_3^{n+2})}{p_2 - p_3} & : p_2 \neq p_3 \\ (2+n)p^{n+2}(1-2p)^2 & : p_2 = p_3 =: p \end{cases}$$

$$b_n := \mathrm{sp}\, T^*(\eta_{2,n}) = \begin{cases} \frac{p_1 (p_2^{n+2} - p_3^{n+2})}{p_2 - p_3} & : p_2 \neq p_3 \\ (2+n)p^{n+1}(1-2p)^2 & : p_2 = p_3 =: p \end{cases}$$

$$L(\eta_{1,n}) = r^{n+4}$$

$$L(\eta_{2,n}) = \rho r^{n+1}.$$

Let

$$a_{\min} = \inf_n \frac{\log a_n}{(n+4)\log r} \qquad\qquad a_{\max} = \sup_n \frac{\log a_n}{(n+4)\log r}$$

$$b_{\min} = \inf_n \frac{\log b_n}{(n+1)\log r + \log \rho} \qquad\qquad b_{\max} = \sup_n \frac{\log b_n}{(n+1)\log r + \log \rho}.$$

Then the minimal local dimension is equal to

$$\alpha_{\min} := \min\Big\{ \frac{\log p_1}{\log \rho}, \frac{\log p_2}{\log r}, \frac{\log p_3}{\log r}, a_{\min}, b_{\min} \Big\}.$$

and the maximal local dimension is equal to

$$\alpha_{\max} := \max\Big\{ \frac{\log p_1}{\log \rho}, \frac{\log p_2}{\log r}, \frac{\log p_3}{\log r}, a_{\max}, b_{\max} \Big\}.$$

The parameters $\alpha_{\min}$ and $\alpha_{\max}$ can be determined exactly in many situations, but generic solutions are tedious. Additional details are left to the reader.

5.3.3. *The maximal open sets of the weak separation condition.* Here we show that the essential net interval $[0,1] \cap K$ is not contained in a union of open balls U_0 satisfying the maximal value in (3.1). In fact, we show that for any $\epsilon > 0$, the open set $(1-\epsilon, 1) \cap K$ is not contained a finite union of such open balls for any $\epsilon > 0$. In addition, this shows

that for any $U(x,t)$ with $\#\mathcal{S}_t(U(x,t))$ maximal, we must have $1 \notin U(x,t)$, whereas $1 \in K = K_{\mathrm{ess}}$.

We first note that $\sup_{x \in \mathbb{R}, t > 0} \#\mathcal{S}_t(U(x,t)) \geq 5$. To see this, take $t = 1/4$ and $U_0 := U(1/4, 1/4)$. Then for each $\sigma \in \{11, 12, 13, 22, 23\}$, we have $S_\sigma(K) \cap U_0 \neq \emptyset$ (since $S_{13} = S_{21}$, we exclude the word 21).

To show that $(1 - \epsilon, 1) \cap K$ is not contained in a finite union of maximal open balls for each $\epsilon > 0$, since 1 is an accumulation point for K it suffices to show that if $t > 0$ and $U(x,t)$ is any open ball such that $x + t = 1$, $\#\mathcal{S}_t(U(x,t)) < 5$. A direct check shows that for $t > 1/4$, $\#\mathcal{S}_t(U(x,t)) < 5$. Otherwise, let $m \geq 1$ be such that $1/4^{m+1} < t \leq 1/4^m$. Since the rightmost child of $[0,1]$ is the net interval $[3/4, 1] \in \Lambda_{1/4}$ with $\mathcal{V}([3/4, 1]) = \mathcal{V}([0,1])$, the net interval in generation t containing 1 is the interval $\Delta = [1 - 1/4^m, 1]$ which has $\mathcal{V}(\Delta) = \mathcal{V}([0,1])$, and thus $U(x,t) \subseteq \Delta' = [1 - 1/4^{m-1}]$ where $\mathcal{V}(\Delta') = \mathcal{V}([0,1])$. But then up to normalization, we know that the net intervals contained in Δ' are the same as the net intervals contained in $[0,1]$ so the case for general t reduces to the case $t > 1/4$.

5.3.4. *On the multifractal formalism.* We now dispense with the assumptions on the parameters ρ, r and establish the following result.

Theorem 5.7. *Any invariant measure $\mu_{\boldsymbol{p}}$ associated with the IFS*

$$S_1(x) = \rho \cdot x \qquad S_2(x) = r \cdot x + \rho(1 - r) \qquad S_3(x) = r \cdot x + 1 - r$$

where $0 < \rho, r < 1$ satisfy $\rho + 2r - \rho r \leq 1$ satisfies the complete multifractal formalism.

Proof. By Corollary 4.7 and the following remark, since the IFS satisfies the weak separation condition, it suffices to show that the vertex $v_0 := \{x \mapsto x\}$ is contained in the essential class. As argued in Section 5.3.1, the net interval $[0,1]$ has children

$$(\Delta_1 = [0, \rho(1 - r)], \Delta_2 = [\rho(1 - r), \rho], \Delta_3 = [\rho, \rho + r - \rho r], \Delta_4 = [1 - r, r])$$

in $\mathcal{F}_1$ with neighbour sets

$$\mathcal{V}(\Delta_1) = \{x \mapsto x/(1 - r)\} \qquad \mathcal{V}(\Delta_2) = \{x \mapsto x/\rho, x \mapsto x/r + \frac{1}{r} - 1\}$$

$$\mathcal{V}(\Delta_3) = \{x \mapsto \frac{x}{1 - \rho} + \frac{\rho}{1 - \rho}\} \qquad \mathcal{V}(\Delta_4) = \{x \mapsto x\}.$$

In particular, there is an edge from v_0 to v_0. Moreover, since the word 23 is in Λ_r, where $S_{23}([0,1])$ is disjoint from $S_3([0,1])$, $S_{22}([0,1])$, and $S_1([0,1])$ by the assumptions on ρ and r, it follows that $S_{23}([0,1])$ is a net interval with neighbour set v_0. Thus there is an edge from $\mathcal{V}(\Delta_3)$ to v_0. Similarly, the words 11 and 12 are in Λ_ρ, where $S_{12}([0,1])$ is disjoint from $S_2([0,1])$, so as computed in Section 5.3.1, the children of Δ_1 have neighbour sets $\mathcal{V}(\Delta_1)$, $\mathcal{V}(\Delta_2)$, and $\mathcal{V}(\Delta_3)$. Since there is an edge from $\mathcal{V}(\Delta_3)$ to v_0, there is a path from $\mathcal{V}(\Delta_1)$ to v_0.

It remains to consider the offspring of $v_2 := \mathcal{V}(\Delta_2)$. We will treat the case where $r > \rho$; the case where $r \leq \rho$ follows by an analogous argument. Let m be maximal such that $r^m > \rho$. We will compute the net intervals in generation Λ_{r^m}.

For $0 \le k \le m$ write

$$\sigma_k = \underbrace{2\dots 2}_{k \text{ times}} 1 \qquad\qquad\qquad \tau_k = \underbrace{2\dots 2}_{k \text{ times}}.$$

For simplicity, given $t > 0$, write $\Gamma_t = \{S_\omega : \omega \in \Lambda_t, S_\omega((0,1)) \cap \Delta_2 \ne \emptyset\}$. Note that $S_2(S_1(1)) = S_1(1)$ where $S_1(1)$ is the right endpoint of Δ_2, so that $S_{\sigma_k}(S_1(1)) = S_1(1)$. Thus by choice of m, we have for $k \le m$

$$\Gamma_{r^k} = \{S_{\tau_{k+1}}, S_{\sigma_0}, S_{\sigma_1}, \dots, S_{\sigma_k}\}.$$

First assume $r^{m+1} < \rho$. Since $r^{m+1} < \rho$ and $S_1([0,1]) \supseteq \Delta_2 \supseteq \Delta^{(i)}$, $\mathrm{tg}(\Delta^{(i)}) = \rho$. Thus since $S_{12}(1) \le S_2(0)$ and $S_{\sigma_0 2} = S_{\sigma_1}$, we have

$$\Gamma_\rho = \{\tau_{m+1}, S_{\sigma_1}, \dots, S_{\sigma_m}\}.$$

Since $S_\tau(0) > S_{\sigma_m}(1)$, the net intervals in $\mathcal{F}_\rho$ contained in Δ_2 are given, ordered from left to right,

$$\Delta^{(i)} = [S_{\sigma_i}(0), S_{\sigma_{i+1}}(0)] \qquad \Delta^{(m)} = [S_{\sigma_m}(0), S_\tau(0)] \qquad \Delta^{(m+1)} = [S_\tau(0), S_{\sigma_m}(1)]$$

for $1 \le i < m$. Since $\Gamma_{r^{m-1}} = \{S_2^{-1} \circ g : g \in \Gamma_\rho\}$, for each $1 \le i \le m+1$, $S_2^{-1}(\Delta^{(i)}) \in \mathcal{F}_{r^{m-1}}$ with $\mathcal{V}(S_2^{-1}(\Delta^{(i)})) = \mathcal{V}(\Delta^{(i)})$. But again S_2 fixes the right endpoint of Δ_2, so that

$$\Delta_2 \supseteq S_2^{-1}(\Delta^{(i)}) \supseteq \Delta^{(j)}$$

for each $2 \le i \le m+1$. In particular, every child of $\mathcal{V}(\Delta^{(i)})$ is of the form $\mathcal{V}(\Delta^{(j)})$, and there is a path from $\mathcal{V}(\Delta^{(i)})$ to $\mathcal{V}(\Delta^{(1)})$ for each $i \ge 2$. Moreover, a direct computation shows that $\mathcal{V}(\Delta^{(1)} = \mathcal{V}(\Delta_1)$, so there is a path from $\mathcal{V}(\Delta^{(1)})$ to v_0. Thus there are no new net intervals, and there is a path to $v_0 = \{x \mapsto x\}$ from any vertex in the transition graph, as required.

In the case $r^{m+1} = \rho$, we get

$$\Gamma_\rho = \{\tau_{m+2}, S_{\sigma_1}, \dots, S_{\sigma_{m+1}}\}$$

so that Γ_ρ is a rescaled version of Γ_{r^m}, and the argument follows similarly. $\qquad\square$

5.4. **On an example of Deng and Ngai.**

In [4, Ex. 8.5], Deng and Ngai introduced the following IFS similar in structure to Section 5.3 but with an additional overlap. Consider IFS defined by following four maps

$$S_1(x) = \rho x \qquad\qquad\qquad S_2(x) = rx + \rho(1-r)$$
$$S_3(x) = \rho^{-1}r^2 x + (1-r)(\rho + r) \qquad\qquad S_4(x) = rx + (1-r)$$

where $0 < \rho, r \in< 1$ satisfy $r^2 < \rho$ and $\rho(r-1)(\rho + r - 1) > r^2$. The constraints on ρ and r ensure that $S_3((0,1)) \cap S_4((0,1)) = \emptyset$.

The parameters of this IFS are chosen so that $S_{14} = S_{21}$ and $S_{24} = S_{31}$. On can verify, arguing similarly to Theorem 5.7, that $\mathcal{G} = \mathcal{G}_{\mathrm{ess}}$ and hence any associated self-similar measure satisfies the complete multifractal formalism.

5.5. **A modified multifractal formalism for Cantor-like measures.** Consider the family of IFS given by maps

$$\Big\{S_j(x) = \frac{x}{r} + \frac{j}{mr}(r-1) : 0 \le j \le m\Big\}$$

where $m \ge r \ge 2$ and m, r are integers. This family of IFS, with appropriate probabilities, contains rescaled versions of measures such as convolutions of the usual Cantor measure. In particular, certain self-similar measures in this family were among the first recognized for which the multifractal formalism can fail [23]. The set of local dimensions is known to consist of a closed interval and, with appropriate probabilities, an isolated point. The L^q-spectra have also been computed, as well as a modified multifractal formalism [12, 13, 30, 36]; our results here are minor improvements of existing results and are primarily useful as illustrations of the theorems.

Fix any IFS $\{S_i\}_{i \in \mathcal{I}}$ in this family with attractor K and associated self-similar measure μ. Arguing similarly to [17, Prop. 7.1], one may verify that $K = [0,1]$, $K_{\mathrm{ess}} = (0,1)$, and for each $t_m = 1/r^{m-1}$, we have

$$K_m := \bigcup_{\substack{\Delta \in \mathcal{F}_{t_m} \\ \mathcal{V}(\Delta) \in V(\mathcal{G}_{\mathrm{ess}})}} = \Big[\frac{r-1}{kr^m}, 1 - \frac{(r-1)}{kr^m}\Big].$$

Then Corollary 4.6 states that each $\mu_m := \mu_{\boldsymbol{p}}|_{K_m}$ satisfies the complete multifractal formalism and

$$D(\mu_m) = \{\dim_{\mathrm{loc}} \mu_{\boldsymbol{p}}(x) : x \in (0,1)\}.$$

This provides an alternative proof of some of the results contained in [13, 36] (without the assumption $k < 2r - 2$) and a variation of [12, Example 6.2].

From the perspective of Corollary 4.7, the obstruction to the multifractal formalism is combinatorial: there exists a cycle outside the essential class which contributes a point with local dimension not contained in the closed interval $\{\dim_{\mathrm{loc}} \mu(x) : x \in K_{\mathrm{ess}}\}$.

REFERENCES

1. Matthias Arbeiter and Norbert Patzschke, *Random Self-Similar Multifractals*, Math. Nachr. **181** (1996), no. 1, 5–42. 3
2. Christoph Bandt and Siegfried Graf, *Self-Similar Sets 7. A Characterization of Self-Similar Fractals with Positive Hausdorff Measure*, Proc. Amer. Math. Soc. **114** (1992), no. 4, 995. 2
3. Robert Cawley and R. Daniel Mauldin, *Multifractal decompositions of Moran fractals*, Adv. Math. **92** (1992), no. 2, 196–236. 3, 29
4. Guotai Deng and Sze-Man Ngai, *Differentiability of L^q-spectrum and multifractal decomposition by using infinite graph-directed IFSs*, Adv. Math. **311** (2017), 190–237. 7, 8, 40
5. Qi-Rong Deng, Ka-Sing Lau, and Sze-Man Ngai, *Separation Conditions for Iterated Function Systems with Overlaps*, Contemporary Mathematics (David Carfì, Michel Lapidus, Erin Pearse, and Machiel van Frankenhuijsen, eds.), vol. 600, American Mathematical Society, Providence, Rhode Island, 2013, pp. 1–20. 33
6. Kenneth J. Falconer, *Techniques in fractal geometry*, Wiley, Chichester ; New York, 1997. 2

7. De-Jun Feng, *Lyapunov exponents for products of matrices and multifractal analysis. Part I: Positive matrices*, Isr. J. Math. **138** (2003), no. 1, 353–376. 3, 29

8. _____ , *Smoothness of the L^q-spectrum of self-similar measures with overlaps*, J. Lond. Math. Soc. **68** (2003), no. 01, 102–118. 5, 10, 11, 16

9. _____ , *The limited Rademacher functions and Bernoulli convolutions associated with Pisot numbers*, Adv. Math. **195** (2005), no. 1, 24–101. 3, 4, 8

10. _____ , *Lyapunov exponents for products of matrices and multifractal analysis. Part II: General matrices*, Isr. J. Math. **170** (2009), no. 1, 355–394. 3, 5, 7, 29

11. De-Jun Feng and Huyi Hu, *Dimension theory of iterated function systems*, Comm. Pure Appl. Math. **62** (2009), no. 11, 1435–1500. 2

12. De-Jun Feng and Ka-Sing Lau, *Multifractal formalism for self-similar measures with weak separation condition*, J. Math. Pures Appl. **92** (2009), no. 4, 407–428. 3, 4, 6, 21, 28, 29, 30, 41

13. De-Jun Feng, Ka-Sing Lau, and Xiang-Yang Wang, *Some exceptional phenomena in multifractal formalism: Part II*, Asian J. Math. **9** (2005), no. 4, 473–488. 3, 4, 29, 41

14. De-Jun Feng, Ka-Sing Lau, and Jun Wu, *Ergodic Limits on the Conformal Repellers*, Adv. Math. **169** (2002), no. 1, 58–91. 4

15. Jonathan M. Fraser, Alexander M. Henderson, Eric J. Olson, and James C. Robinson, *On the Assouad dimension of self-similar sets with overlaps*, Adv. Math. **273** (2015), 188–214. 2, 23

16. Thomas C. Halsey, Mogens H. Jensen, Leo P. Kadanoff, Itamar Procaccia, and Boris I. Shraiman, *Fractal measures and their singularities: The characterization of strange sets*, Phys. Rev. A **33** (1986), no. 2, 1141–1151. 3, 29

17. Kathryn E. Hare, Kevin G. Hare, and Kevin R. Matthews, *Local dimensions of measures of finite type*, J. Fractal Geom. **3** (2016), no. 4, 331–376. 5, 23, 27, 41

18. Kathryn E. Hare, Kevin G. Hare, and Michael K.S. Ng, *Local dimensions of measures of finite type II: Measures without full support and with non-regular probabilities*, Can. J. Math. **70** (2018), no. 4, 824–867. 5, 6, 24, 27

19. Kathryn E. Hare, Kevin G. Hare, and Alex Rutar, *When the Weak Separation Condition implies the Generalized Finite Type Condition*, Proc. Amer. Math. Soc. **149** (2021), no. 4, 1555–1568 (en). 5, 8, 10, 18, 19, 33

20. Kathryn E. Hare, Kevin G. Hare, and Wanchun Shen, *The L^q-spectrum for a class of self-similar measures with overlap*, Asian J. Math. (to appear). 4, 8

21. Kathryn E. Hare, Kevin G. Hare, and Grant Simms, *Local dimensions of measures of finite type III — Measures that are not equicontractive*, J. Math. Anal. Appl **458** (2018), no. 2, 1653–1677. 5, 10, 16, 27

22. Michael Hochman, *On self-similar sets with overlaps and inverse theorems for entropy*, Ann. Math. **180** (2014), no. 2, 773–822. 2

23. Tian-You Hu and Ka-Sing Lau, *Multifractal Structure of Convolution of the Cantor Measure*, Adv. Appl. Math. **27** (2001), no. 1, 1–16. 4, 41

24. John E. Hutchinson, *Fractals and Self Similarity*, Indiana Univ. Math. J. **30** (1981), no. 5, 713–747. 3

25. Thomas Jordan and Ariel Rapaport, *Dimension of ergodic measures projected onto self-similar sets with overlaps*, arXiv:1908.00271 [math] (2020). 2

26. Ka-Sing Lau, *Self-Similarity, L^p-Spectrum and Multifractal Formalism*, Fractal Geometry and Stochastics (Christoph Bandt, Siegfried Graf, and Martina Zähle, eds.), Birkhäuser Basel, Basel, 1995, pp. 55–90. 3, 29

27. Ka-Sing Lau and Sze-Man Ngai, *Multifractal Measures and a Weak Separation Condition*, Adv. Math. **141** (1999), no. 1, 45–96. 3, 5, 18, 29

28. _____ , *A generalized finite type condition for iterated function systems*, Adv. Math. **208** (2007), no. 2, 647–671. 2, 8, 33, 35

29. Ka-Sing Lau and Xiang-Yang Wang, *Iterated function systems with a weak separation condition*, Studia Math. **161** (2004), no. 3, 249–268. 4, 35

30. _____ , *Some exceptional phenomena in multifractal formalism: Part I*, Asian J. Math. **9** (2005), no. 2, 275–294. 4, 41

31. Sze-Man Ngai and Yang Wang, *Hausdorff dimension of self-similar sets with overlaps*, J. Lond. Math. Soc. **63** (2001), no. 3, 655–672. 2, 5, 33

32. Norbert Patzschke, *Self-Conformal Multifractal Measures*, Adv. Appl. Math. **19** (1997), no. 4, 486–513. 3, 29

33. Yakov Pesin and Howard Weiss, *A multifractal analysis of equilibrium measures for conformal expanding maps and Moran-like geometric constructions*, J Stat Phys **86** (1997), no. 1-2, 233–275. 3, 29

34. Andreas Schief, *Separation properties for self-similar sets*, Proc. Amer. Math. Soc. **122** (1994), no. 1, 111–111. 2

35. Eugene Seneta, *Non-negative matrices and Markov chains*, second ed., Springer Series in Statistics, Springer, New York, 1981. 18

36. Pablo Shmerkin, *A Modified Multifractal Formalism for a Class of Self-similar Measures with Overlap*, Asian J. Math. **9** (2005), no. 3, 323–348. 3, 4, 29, 41

37. Benoît Testud, *Mesures quasi-Bernoulli au sens faible : Résultats et exemples*, Ann. Inst. Henri Poincaré B **42** (2006), no. 1, 1–35. 5

38. Martin P.W. Zerner, *Weak Separation Properties for Self-Similar Sets*, Proc. Amer. Math. Soc. **124** (1996), no. 11, 3529–3539. 4, 18

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