

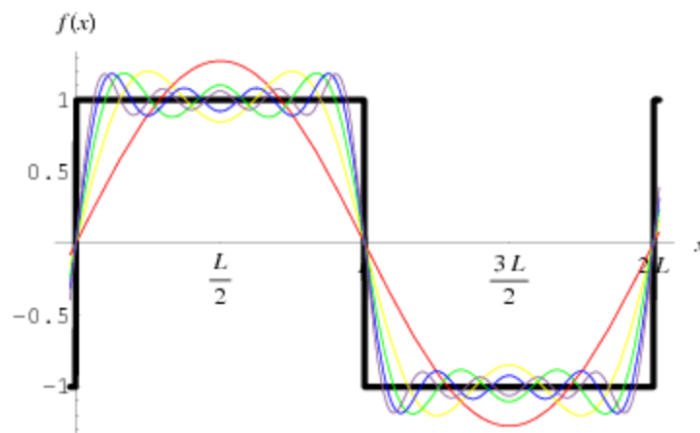
6.058 Fourier Examples

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1 Fourier Series

We are going to take the Fourier Series of a square wave over the range 0 to $2L$. A square wave, as is noted by the sine waves



on it, is a periodic function which is odd.

Since the function is odd, we only have to consider the odd part of the Fourier Series. We will calculate the odd Fourier series coefficients, b_n using

$$b_n = \frac{1}{L} \int_0^{2L} f(t) \sin\left(\frac{n\pi t}{L}\right) dt$$

This equation reduces to

$$b_n = \frac{2}{L} \int_0^L f(t) \sin\left(\frac{n\pi t}{L}\right) dt$$

$$b_n = \frac{4}{n\pi} \sin^2\left(\frac{1}{2}n\pi\right)$$

$$b_n = \frac{2}{n\pi} (1 - (-1)^n)$$

The Fourier series is therefore

$$f(t) = \frac{4}{\pi} \sum_{n=1,3,5,\dots}^{\infty} \frac{1}{n} \sin\left(\frac{n\pi t}{L}\right)$$

2 Continuous Fourier Transform

We shall consider the Fourier Transform of several very important signals.

For the signal $\delta(t)$

$$X(j\omega) = \int_{-\infty}^{\infty} \delta(t) e^{-j\omega t} dt = 1$$

Since a δ in time domain maps to a constant in frequency domain, a constant in time maps to a δ in frequency. In particular, the Fourier transform of 1 is $2\pi\delta(t)$.

For the signal $\delta(t - t_0)$

$$X(j\omega) = \int_{-\infty}^{\infty} \delta(t - t_0) e^{-j\omega t} dt = \int_{-\infty}^{\infty} \delta(0) e^{-j\omega t_0} dt = e^{-j\omega t_0}$$

The signal $e^{-at}u(t)$ for $Re(a) > 0$

$$X(j\omega) = \int_{-\infty}^{\infty} e^{-at}u(t)e^{-j\omega t}dt = \int_0^{\infty} e^{(-a-j\omega)t}dt = \left. \frac{e^{(-a-j\omega)t}}{(-a-j\omega)} \right|_0^{\infty} = \frac{1}{a+j\omega}$$

The signal $\cos(\omega_c t)$

$$X(j\omega) = \int_{-\infty}^{\infty} \frac{e^{j\omega_c t} + e^{-j\omega_c t}}{2} e^{-j\omega t} dt = \int_{-\infty}^{\infty} \frac{e^{j(\omega-\omega_c)t} + e^{j(\omega+\omega_c)t}}{2} dt = 2\pi\delta(\omega - \omega_c) + 2\pi\delta(\omega + \omega_c)$$

3 Discrete Fourier Transform

Now let's flex our muscles by working through a few simple examples in the discrete case. Let's start with the unit impulse:

$$X(j\Omega) = \sum_{n=-\infty}^{\infty} \delta[n]e^{-j\Omega n} = 1$$

We see that the delta function gets smeared out over the whole space, just as in the continuous case. What about a delta function that's been translated?

$$X(j\Omega) = \sum_{n=-\infty}^{\infty} \delta[n-k]e^{-j\Omega n} = e^{-j\Omega k}$$

Notice that both $1 = e^{j\Omega 0}$ and $e^{-j\Omega k}$ are both eigenfunctions. Thus, it seems that shifting the delta function changes the eigenfunction that is spread across the frequency domain. Now for something slightly harder: what is the Fourier transform of the function equal to 1 for all n ($x[n] = 1$)?

$$X(j\Omega) = \sum_{n=-\infty}^{\infty} e^{-j\Omega n}$$

This is a strange summation. For any non-zero Ω , the sum converges to zero, since this amounts to summing together uniformly distributed vectors on the unit circle (if you don't quite see this spend some time thinking about it; it's a neat exercise). But when $\Omega = 0$, we get an infinite summation. So $X(j\Omega) = 0$ for all $\Omega \neq 0$ but $X(j\Omega) = \infty$ when $\Omega = 0$; this smells like a Dirac delta function. To confirm this suspicion, let's see what happens when we integrate over it:

$$\int_{-\pi}^{\pi} \sum_{n=-\infty}^{\infty} e^{-j\Omega n} d\Omega = \sum_{n=-\infty}^{\infty} \int_{-\pi}^{\pi} e^{-j\Omega n} d\Omega = 2\pi$$

Note that in the last step we employed the fact that

$$\int_{-\pi}^{\pi} e^{-j\Omega n} d\Omega = 0$$

in all cases except for $n = 0$, in which case

$$\int_{-\pi}^{\pi} e^{-j\Omega 0} d\Omega = \int_{-\pi}^{\pi} 1 d\Omega = 2\pi$$

Thus, we have a Dirac delta centered at the origin of the frequency domain yielding 2π when integrated over:

$$X(j\Omega) = 2\pi\delta(\Omega)$$

Above we noted that, when we are transforming delta functions, shifted delta functions return different eigenfunctions to be spread across the Fourier domain. Naturally, we might ask: what is the Fourier transform of an arbitrary eigenfunction spread across the time domain (say, $x[n] = e^{j\Omega_0 k}$)?

$$X(j\Omega) = \int_{-\pi}^{\pi} \sum_{n=-\infty}^{\infty} e^{j\Omega_0 k} e^{-j\Omega n} d\Omega = \sum_{n=-\infty}^{\infty} \int_{-\pi}^{\pi} e^{j\Omega_0 k} e^{-j\Omega n} d\Omega = 2\pi\delta(\Omega - \Omega_0)$$

4 System Functional

Suppose we have the RC circuit system from PSET 1 with the known impulse response of $h(t) = \frac{1}{RC}e^{-\frac{1}{RC}t}u(t)$. We wish to find the output of this system to arbitrary input. In order to do this we must calculate $H(j\omega)$. From the above continuous identity that $e^{-at}u(t)$ with $Re(a) > 0$ has fourier transform $\frac{1}{a+j\omega}$ we find that

$$H(j\omega) = \frac{1}{RC} \frac{1}{\frac{1}{RC} + j\omega}$$

If we were to put $x(t) = \cos(6t)$ input the system. In order to do this we must first take the Fourier transform. We note that the Fourier series is $X(j\omega) = 2\pi\delta(\omega - 6) + 2\pi\delta(\omega + 6)$. Thus, $Y(j\omega) = \frac{2\pi\delta(\omega-6)}{RC} \frac{1}{\frac{1}{RC} + j\omega} + \frac{2\pi\delta(\omega+6)}{RC} \frac{1}{\frac{1}{RC} + j\omega}$

5 Real World Systems: RLC Circuit

The differential equation modeling an RLC circuit is

$$\frac{d^2}{dt^2}i(t) + \frac{R}{L} \frac{d}{dt}i(t) + \frac{i(t)}{LC} = \frac{1}{L} \frac{d}{dt}V(t)$$

where $i(t)$ is the current, R is the resistance, C is the capacitance, and L is the inductance, and $V(t)$ is the voltage supplied to the circuit.

Suppose we want to find the transfer function, $H(j\omega)$. We know that $Y(j\omega) = X(j\omega)H(j\omega)$ so finding $H(j\omega)$ reduces to finding $\frac{Y(j\omega)}{X(j\omega)}$. We note that a property of the Fourier Transform is that the Fourier transform of a derivative corresponds to pulling $j\omega$ out in front of a Fourier Transform. Thus, this equation, by the Fourier transform, reduces to

$$(j\omega)^2 Y(j\omega) + \frac{R}{L} (j\omega) Y(j\omega) + \frac{1}{LC} Y(j\omega) = \frac{1}{L} (j\omega) X(j\omega)$$

where $Y(j\omega)$ is the Fourier transform of the current and $X(j\omega)$ is the Fourier Transform of the voltage.

Thus, we find that

$$\frac{Y(j\omega)}{X(j\omega)} = \frac{\frac{j\omega}{L}}{(j\omega)^2 + \frac{R}{L}(j\omega) + \frac{1}{LC}} = H(j\omega)$$

6 Real World System: Low Pass Filter

Suppose we have a system in the real world that obeys the equation

$$y(t) = k \frac{dx(t)}{dt} + x(t)$$

Let us find the transfer function.

$$\frac{Y(j\omega)}{X(j\omega)} = \frac{1}{j\omega k + 1} = H(j\omega)$$

We shall plot the magnitude of this Fourier transform as we increase ω . We get an image similar to this:

