

growpacity: A computationally efficient dust opacity model suitable for coagulation models

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DOI: [10.xxxxxx/draft](https://doi.org/10.xxxxxx/draft)

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Submitted: 01 January 1970

Published: unpublished

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Summary

growpacity is a Python and C toolkit for calculating dust opacities in astrophysical environments. It provides a framework for generating, storing, and interpolating Rosseland and Planck mean opacities over customizable grids of grain size distributions, temperature, and powerlaw exponents. This package is designed as a wrapper around [OpTool](#), allowing users to specify a grain composition and obtain tabulated temperature-, maximum grain size-, and dust size distribution-dependent mean opacities. The resulting tables are lightweight and optimized for use in radiation hydrodynamics and dust coagulation models, where grain size distributions evolve dynamically.

Statement of need

Radiation hydrodynamics simulations often require a dust opacity model under the assumption of a particular grain size or size distribution. It is however not straightforward to implement a dust opacity model that accounts for the dynamical evolution of this grain size distribution. With growpacity, we outline a method to construct such a model and provide a python package that interfaces to OpTool to compute temperature-, maximum grain size-, and dust size distribution-dependent mean opacities, tabulated in a lightweight format and complemented with efficient interpolation methods for usage in radiation hydrodynamics simulations.

Method

For a given grain composition and assuming a grain size distribution with number density $n(a) \propto a^q$ for grains with size a between minimum and maximum grain sizes a_{\min} and a_{\max} , the [OpTool](#) package ([Dominik et al., 2021](#)) can compute the absorption and scattering opacities $\kappa_{\text{abs}}(\nu)$ and $\kappa_{\text{sca}}(\nu)$ (in cm^2/g) as well as the asymmetry factor $g(\nu)$ ([Henyey & Greenstein, 1941](#)) over a frequency grid ν . This calculation is done using the Distribution of Hollow Spheres method (DHS, [Min et al., 2005](#)). The Rosseland and Planck mean opacities κ_{R} and κ_{P} can then be computed as

$$\kappa_{\text{P}}(T) = \frac{\int_0^\infty \kappa_{\text{abs}} B_\nu(T) d\nu}{\int_0^\infty B_\nu(T) d\nu}, \quad \kappa_{\text{R}}(T) = \frac{\int_0^\infty u_\nu(T) d\nu}{\int_0^\infty [\kappa_{\text{abs}} + (1 - g) \kappa_{\text{sca}}]^{-1} u_\nu(T) d\nu}, \quad u_\nu(T) = \left. \frac{dB_\nu}{dT} \right|_T,$$

where $B_\nu(T)$ is the Planck function at temperature T . By fixing the grain composition and a_{\min} , OpTool can be used to calculate the absorption and scattering opacities over a grid of q and a_{\max} . The above equation can then be used to compute and tabulate κ_{R} and κ_{P} over

35 q , a_{\max} , and T . The resulting 3D tables can be used in any context where q and a_{\max} are
 36 dynamically evolved according to a dust coagulation model (Birnstiel et al., 2017; Pfeil et al.,
 37 2024; Robinson et al., 2024; Stammer & Birnstiel, 2022).

38 Interpolation algorithm

39 Once the 3D tables of $\kappa_R(q, a_{\max}, T)$ and $\kappa_P(q, a_{\max}, T)$ have been computed, they can be
 40 interpolated within the range of tabulated values. We interpolate for $\log \kappa$ as a function of q ,
 41 $\log a_{\max}$, and $\log T$, with regular sampling in this space (i.e., logarithmic spacing for a_{\max} and
 42 T). This works best when the mean opacities follow a powerlaw with respect to temperature
 43 $\kappa \propto T^b \Rightarrow \log \kappa \propto b \log T$, which is a reasonable approximation and especially holds for small
 44 grains (Bell & Lin, 1994; Semenov et al., 2003). It also ensures that the interpolated opacities
 45 are always positive in case extrapolation is needed.

46 We use a trilinear interpolation scheme, which is fast and simple to implement, and takes
 47 advantage of the fact that the arrays q , $\log a_{\max}$, and $\log T$ are sorted and regularly spaced
 48 to efficiently locate the indices of the grid points that surround the point of interest. For
 49 each array $x \in \{q, \log a_{\max}, \log T\}$ and for a target value x_t , we first find the index i such
 50 that $x_i \leq x_t < x_{i+1}$ as $i = \lfloor (x_t - x_0) \Delta x^{-1} \rfloor$, where x_0 is the first (smallest) value in the
 51 sampling space and $\Delta x = x_{i+1} - x_i$ is the (constant) sampling spacing. As this information
 52 is known *a priori*, this reduces the complexity of finding the required indices from $\mathcal{O}(\log N)$ to
 53 $\mathcal{O}(1)$, where N is the number of grid points in x , and is especially efficient for larger grids. Of
 54 course, care must be taken to ensure that the indices are within the bounds of the grid.

55 Implementation

56 The above method is implemented in the `growcapacity` package. The package provides a simple
 57 python interface to `OpTool` to compute accurate dust opacities for a given grain composition
 58 (provided in `OpTool` format) and for a fixed a_{\min} , over a grid of a_{\max} , q , and T . The resulting
 59 3D tables of κ_R and κ_P are saved in binary files that can be easily loaded with python. We
 60 also provide an implementation of the interpolating function in C, that can be readily used in
 61 radiation hydrodynamics codes. These files are especially lightweight: a typical calculation
 62 involving $q \in [-4.5, -2.5]$ with $\Delta q = 0.25$, $a_{\max} \in [0.1 \mu\text{m}, 1 \text{m}]$ sampled twice per decade,
 63 and $T \in [1, 2000] \text{ K}$ with 100 points results in two arrays of $9 \times 15 \times 100$ elements, or
 64 about 220 kB of memory. This choice of spacing means that opacities are exactly evaluated
 65 for $q \in [-3.75, -3.5, -3]$, corresponding to dust size distributions in equilibrium due to
 66 small/large-scale turbulence, radial drift, or in a non-equilibrium growth regime (for more
 67 information, see Pfeil et al., 2024).

68 The code and documentation are available at
 69 <https://github.com/alexziab/growcapacity>.

70 Limitations and extensions

71 We underscore that our intent is not to provide a new or more realistic dust opacity model,
 72 but rather one suitable for use in coagulation models, where dust densities and distributions
 73 can vary as a function of position and time—the applicability of the model depends on entirely
 74 user-defined choices. The method can be easily extended to include gas opacities (Malygin
 75 et al., 2014; e.g., Semenov et al., 2003) and prescriptions for the sublimation of dust species
 76 (e.g., Isella & Natta, 2005), for a more complete opacity model in regimes where gas opacities
 77 are significant.

Acknowledgements

AZ and TB acknowledge funding from the European Union under the Horizon Europe Research and Innovation Programme 101124282 (EARLYBIRD). Views and opinions expressed are those of the authors only.

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