Clase 22

Manuel Garcia.

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Conexión de levi-civita 1

Compatible con $g: \nabla_{\alpha} g_{\mu\nu} = 0$

 $\Gamma^{\alpha}_{\mu\nu} = \Gamma^{\alpha}_{\nu\mu} \to \text{Simetrica}$.

$$\begin{split} \bar{\Gamma}^{\alpha}_{\mu\nu} &= \begin{bmatrix} \alpha \\ \mu\nu \end{bmatrix} + K^{\alpha}_{\mu\nu} \\ \Delta\Gamma^{\alpha}_{\mu\nu} &= t^{\alpha}_{\mu\nu} \quad \rightarrow \quad \Gamma^{\alpha}_{\mu\nu} = \hat{\Gamma}^{\alpha}_{\mu\nu} + t^{\alpha}_{\mu\nu} \\ t^{\alpha}_{\mu\nu} &= -K^{\alpha}_{\mu\nu} \quad \rightarrow \quad \Gamma^{\alpha}_{\mu\nu} = \begin{bmatrix} \alpha \\ \mu\nu \end{bmatrix} \end{split}$$

Componentes independientes del tensor de Riemann 2

Podemos determinar la cantidad de componentes independientes del tensor de Riemann en una variedad M, dim(M) = m, teniendo en cuenta sus simetrías:

$$(1) \quad R_{\sigma\rho\mu\nu} = -R_{\sigma\rho\nu\mu}$$

(2)
$$R_{\sigma \alpha \mu \nu} = -R_{\sigma \alpha \mu}$$

$$(3) \quad R_{\sigma\rho\mu\nu} = R_{\mu\nu\sigma\rho}$$

1)
$$R_{\sigma\rho\mu\nu} = -R_{\sigma\rho\nu\mu}$$
 (2) $R_{\sigma\rho\mu\nu} = -R_{\rho\rho\mu\nu}$ (3) $R_{\sigma\rho\mu\nu} = R_{\mu\nu\sigma\rho}$ (4) $R_{\alpha\mu\nu\sigma} + R_{\alpha\nu\sigma\mu} + R_{\alpha\sigma\mu\nu} = 0$

Solo vale para la conexion de levi-civita.

$$R^{\alpha}_{\mu\alpha\nu} = R_{\mu\nu} \quad \to \quad G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

2.1 Ecuaciones de campo relatividad general

$$\boldsymbol{\nabla}^{\mu}R_{\mu\nu}=\kappa\boldsymbol{\nabla}^{\mu}T_{\mu\nu}=0\quad\rightarrow$$

$$R_{\mu\nu}-\frac{1}{2}g_{\mu\nu}R=\kappa T_{\mu\nu}\quad\rightarrow\quad\text{Ec. de campo Einstein}$$

$$\mathbb{L} = \sqrt{-g}R$$
, \rightarrow $S_{EH} = \int \sqrt{-g}Rd^4x$ \rightarrow Accion de Einstein-Hilbert

¿Por qué este lagrangiano? Hasta en la teoria mas basica tenemos $g_{\mu\nu}$, $\Gamma^{\alpha}_{\mu\nu}=g(\partial g)$. Podemos obtener $R^{\alpha}_{\mu\nu\sigma}$. Con los objetos $R_{\alpha\mu\nu\sigma}$, $g_{\mu\nu}$ debemos construir un escalar. Escalar de ricci $R_{\nu\beta}g^{\nu\beta}=R$.