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MMMGRUBS: MOLECULAR MOTION MOVIES AND  
GEOMETRY RECONSTRUCTION USING BAYESIAN  
STATISTICS



# MMMGRUBS: MOLECULAR MOTION MOVIES AND GEOMETRY RECONSTRUCTION USING BAYESIAN STATISTICS

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Hakuna matata?

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## TECHNICAL ABSTRACT

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Ever since the early days of Coulomb explosion imaging, the direct imaging of the structure and dynamics of individual molecules has been promised; however, in practice no one has been able to accurately retrieve this structure and all imaging relies on plotting the momentum vectors of the fragments in different ways to infer crude changes in geometry. The momentum vectors tell a large part of the story but do not provide a clear picture of the molecular dynamics everyone seeks so an accurate method of retrieving the structure is highly desirable. I use Bayesian inference and Markov chain Monte Carlo methods to elucidate the molecular geometries, allowing for the study of larger systems than before and for the inclusion of measurement error in these studies for the first time. Using this method, the ultrafast isomerization of carbonyl sulfide and acetylene is imaged in position-space and showcased as a molecular movie. The method, however, is much larger in scope and is generalized to any problem involving a physical system described by a system of ordinary differential equations, where the final conditions are experimentally measured and inference of the initial conditions is desirable.



## PLAIN LANGUAGE ABSTRACT

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Insert better abstract.





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## Part I

### HOW TO EXPLODE A MOLECULE

To image the dynamics of a molecule by destroying it seems paradoxical. As we shall see, the molecule's structure is encoded in the atomic shrapnel left behind after an explosion. However, to destroy one of nature's simplest creations is no easy task. Molecules are held together by strong chemical bonds. Our best line of attack is to shoot them with a short laser pulse—engulfing the molecule in an intense oscillating electric field will strip away its electrons and cause it to break up quickly. In these first few chapters, I will discuss how to create a short laser pulse and how these pulses interact with matter. Then I will introduce the technique of pump-probe Coulomb explosion imaging, which we will use to probe the dynamic structure of small molecules by studying the atomic fragments resulting from the explosion.



## THE MAKING OF A SHORT LASER PULSE

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### 1.1 SPONTANEOUS AND STIMULATED EMISSION

### 1.2 PROPERTIES OF A LASER

### 1.3 PRODUCING LASER LIGHT

### 1.4 MAKING LASER PULSES

### 1.5 MAKING SHORT LASER PULSES

### 1.6 SHORT PULSE AMPLIFICATION

#### 1.6.1 *Chirped pulse amplification*

#### 1.6.2 *Optical parametric amplification*

### 1.7 LASER AMPLIFIERS





## WHAT DOES A SHORT PULSE DO?

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### 2.1 SHORT LASER PULSE PROPERTIES

### 2.2 INTERACTION WITH SINGLE ATOMS

#### 2.2.1 *Multiphoton ionization*

#### 2.2.2 *Tunneling ionization*

#### 2.2.3 *Keldysh parameter*

### 2.3 INTERACTION WITH SINGLE ELECTRONS

#### 2.3.1 *Ponderomotive force*

#### 2.3.2 *Relativistic Thomson scattering*

#### 2.3.3 *Nonlinear Compton scattering*

### 2.4 DIATOMIC MOLECULES IN INTENSE LASER FIELDS

### 2.5 SMALL POLYATOMIC MOLECULES IN INTENSE LASER FIELDS



## COULOMB EXPLOSION IMAGING

## 3.1 EXPERIMENTAL SCHEME

In pump-probe Coulomb explosion imaging (CEI) one ultrashort laser pulse is split into two pulses through the use of an asymmetric beam-splitter. One of the pulses, the pump pulse, is usually much weaker than the other, the probe pulse. A time delay  $\tau$  between the pulses is then created such that the pump pulse goes first and the probe pulse second. The job of the pump pulse will be to initiate some change in the molecule. One example could include an isomerization of the molecule. Thus the pump pulse “pumps” the molecule into some excited state. The job of the powerful probe pulse is to engulf the molecule in an intense enough laser field such that multiple electrons are stripped off of it. The molecule’s individual atoms are left in a highly-charged state and begin to behave as individual point charges in a purely Coulombic potential. The entire process occurs in the presence of a constant electric field and so the positively-charged ions all accelerate upwards towards a time- and position-sensitive detector. Thus the probe pulse allows for the “probing” of the excited state.

Suppose  $U$  is a simply connected open subset of the complex plane, and  $a_1, \dots, a_n$  are finitely many points of  $U$  and  $f$  is a function which is defined and holomorphic  $U \setminus \{a_1, \dots, a_n\}$ . If  $\gamma$  is a closed rectifiable curve in  $U$  which does not meet any of the  $a_k$ ,

$$\oint_{\gamma} f(z) dz = 2\pi i \sum_{k=1}^n I(\gamma, a_k) \text{Res}(f, a_k). \quad (3.1)$$

## 3.1.1 Femtosecond Multiple Pulse Length Spectroscopy

## 3.1.2 Pump-probe Coulomb explosion imaging

## 3.2 EXPERIMENTAL DETAILS

## 3.3 HISTORY AND ACCOMPLISHMENTS

CEI was first performed by VAGER, NAAMAN, and KANTER (1989) whereby the Coulomb explosion was initiated by passing the molecule through a thin carbon film at high velocities.

### 3.4 OTHER WAYS OF INITIATING COULOMB EXPLOSIONS

There exist other methods of initiating CEI, among them thin foils, free-electron lasers, highly-charged ion impact, single X-ray photons from a synchrotron source.

### 3.5 TIME AND POSITION MEASUREMENT

### 3.6 CALCULATING ATOMIC FRAGMENT MOMENTA

### 3.7 UNCERTAINTY IN MOMENTUM MEASUREMENTS

## Part II

### HOW TO IMAGE A MOLECULE

Having destroyed a molecule and measured the momentum of each of its atomic fragments, we are left with the inverse problem of inferring its structure. While explosions proceed in a deterministic fashion, that is structures map bijectively to momentum measurements, the converse is not true. Two very different structures may produce the same momentum measurements. To make matters worse, there is no analytic solution to the problem and as the molecule grows, the problem of finding its structure becomes increasingly high dimensional. To combat this problem we will require the use of various mathematical and statistical methods. I first discuss some results from the theory of inverse problems to shed some general insight on these problems. I then follow with a discussion of optimization methods which may be used to tackle the problem for very small molecules. However, for full imaging of larger polyatomic molecules with an analysis of measurement error, Bayesian inference using Markov chain Monte Carlo methods is the way to go, which I discuss in the last chapter of this part.



### Part III

## SHOOTING THE MOVIE

Equipped with the means to explode a molecule and the necessary mathematical tools, we can now image the molecule's structure and shoot our molecular movie.





## MOLECULAR MOVIES

## 4.1 HISTORY AND ACCOMPLISHMENTS

The idea that Coulomb explosion imaging could be used to produce so-called molecular movies has permeated the literature ever since the technique's emergence. The goal of observing

## 4.1.1 Coulomb explosion imaging

## 4.1.1.1 Molecular movies in momentum-space

## 4.1.1.2 Geometry reconstruction

Surprisingly, an attempt was made to arrive at an analytical solution for calculating geometries from measured momentum data. NAGAYA and BANDRAUK (2004) are able to derive so-called classical imaging formulas giving the position wavefunction squared for the Coulomb explosion of a diatomic molecule and a linear triatomic molecule (the cases of symmetric and asymmetric Coulomb explosion are treated).

LÉGARÉ et al. (2005a,b) was the first to use Coulomb explosion imaging and report molecular structures. To obtain structures, they assume the explosion proceeds under a purely Coulombic potential and use optimization methods to make guesses at the structure that most accurately reproduces the observed data consistent with minimizing a least-squares objective function. Unfortunately they provide very minimal information regarding their methods and there is a complete lack of discussion acknowledging the shortcomings of this method. Using 8 fs laser pulses they report on the structure of D<sub>2</sub>O and SO<sub>2</sub> (LÉGARÉ et al., 2005b). They also claim to have imaged vibrating D<sub>2</sub><sup>+</sup> and dissociating SO<sub>2</sub><sup>2+</sup> and SO<sub>2</sub><sup>3+</sup> however they provide no more than a couple of dissociation frames and infer the transient D<sub>2</sub><sup>+</sup> bond length from kinetic energy release ratios as a function of pump-probe time delay (LÉGARÉ et al., 2005a).

GAGNON et al. (2008) reported the reconstruction of dichloromethane (CH<sub>2</sub>Cl<sub>2</sub>) using a home-made stochastic-based simulated annealing algorithm that globally optimizes the molecular spatial configuration. They discuss uncertainties but are only able to obtain the structure in five cases.<sup>1</sup>

The best effort so far has probably been the one by KUNITSKI et al. (2015) in which they use a lookup table approach to image the Efimov state of the helium trimer.

*They seem to have worked hard to find an analytical solution but their unsaid conclusion seems to be that it is an intractable problem and their group went silent on this problem.*

*The main shortcomings being degenerate solutions and the fact that they employ convex optimization methods to a problem that is not convex. It is not clear if they even knew about these issues.*

*There is nothing wrong with writing your own code here but nonconvex optimization algorithms are tricky to get right and the reliance should be on professional code.*

<sup>1</sup> See BOCHAROVA et al. (2011) for a discussion of optimization algorithms.

#### 4.1.2 *Laser-induced electron diffraction*

#### 4.1.3 *Femtosecond X-ray scattering*

#### 4.1.4 *In the biological sciences*

Molecular movies are of course not only of interest in physics and chemistry as a means of probing fundamental processes, but also in the biological sciences where molecular structure play a crucial role in determining the function of biomolecules such as proteins. However, the molecules of interest there are much too large to be studied by any of the previous techniques. Thus molecular movies in the biological sciences tend to be annotated computer simulations amalgamated from multiple studies. That said, they are very impressive pieces of work.

A particularly impressive movie by CHEUNG and CRAMER (2012) showcases the process of RNA polymerase transcription and goes on for over six minutes.

### 4.2 PREVIOUS ATTEMPTS

#### 4.2.1 *Simplex algorithm*

BRICHTA, SEAMAN, and SANDERSON (2009) proposed the reconstruction of small triatomic molecules using a simplex algorithm. Unfortunately they only report on the reconstruction of molecular structures based on simulated data for carbon dioxide and formaldehyde. I could not use this algorithm to find the geometries of CO<sub>2</sub> or OCS from real data.

#### 4.2.2 *Lookup table*

In this approach, many Coulomb explosions are simulated for a wide variety of structures, and the resulting momentum vectors from each simulation are stored. Thus you have a mapping from molecular structures to momentum vectors. To determine the structure belonging to a certain set of observed momentum vectors, you simply read the table in reverse. This approach is simple to implement, very quick by design, and front-loads the computation which may be desirable for large data sets. However, of course, it has an exponential time and space complexity  $\mathcal{O}(e^{3N-6})$  where  $N$  is the number of atoms.

#### 4.2.3 *Convex optimization methods*

Much progress was made over the lookup table by treating the geometry reconstruction problem as a constrained nonlinear convex op-

timization such that MATLAB's `fmincon` function can be relied on. It relies on trust regions and uses an interior-point algorithm. This worked especially well in the case of triatomic molecules however four-atom systems proved incredibly difficult to tackle here. This was due to the exponential increase in the number of saddle points with dimensionality making the problem highly non-convex and unsuitable for `fmincon`.

#### 4.3 CURRENT ATTEMPT: MMMGRUBS

*Recall that triatomic molecules have three degrees of freedom resulting in a problem of dimension 3 while four-atom systems have six. That is,  $3N + 6$  for an  $N$ -atom system.*



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