

# Home Work #3

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## 1 Question 1

This homework used the below equation to simulate the position and velocity of the Hubble space telescope.

$$\begin{aligned}\ddot{x} - 2n\dot{y} - 3n^2x &= f_x \\ \ddot{y} + 2n\dot{x} &= f_y \\ \ddot{z} + n^2z &= f_z\end{aligned}$$

assumed that:

$$\begin{aligned}f_x &= 0 \\ f_y &= 0 \\ f_z &= 0\end{aligned}$$

where:

$$n = \sqrt{\frac{\mu}{r^3}}, \quad \mu = 398600.4418 \text{ km}^3 \text{ s}^{-2}, \quad r = r_{altitude} + r_{earth} = 590 + 6378 = 6968_{km}$$

and initial conditions:

$$r_{relative} = [0 \quad 0 \quad 0]^T, \quad v_{relative} = [-0.1 \quad -0.04 \quad -0.02]_{m/s}^T$$

Figure 1: position of the Hubble space telescope

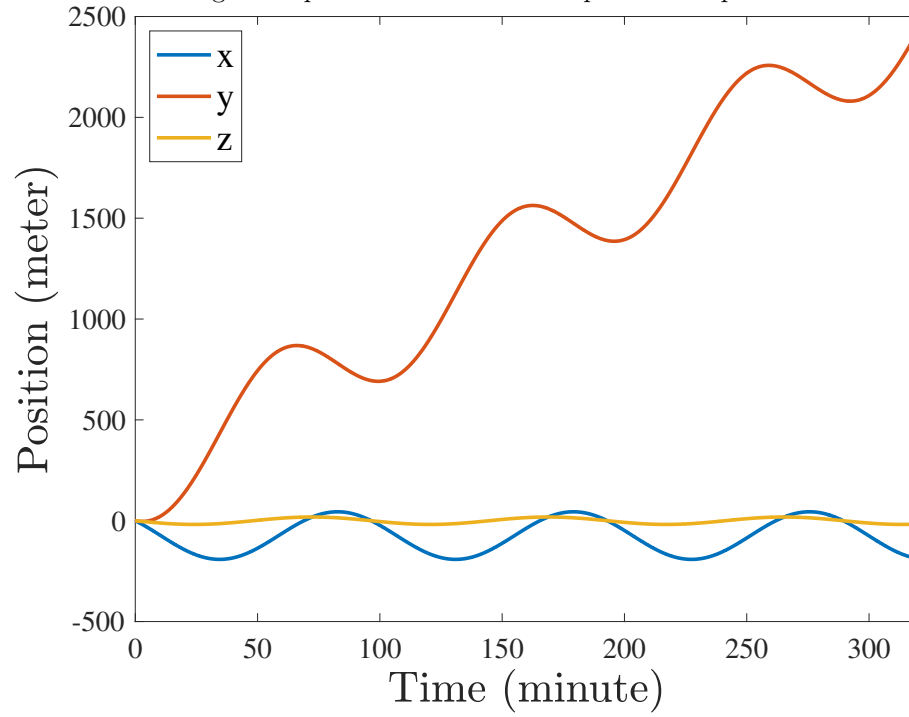
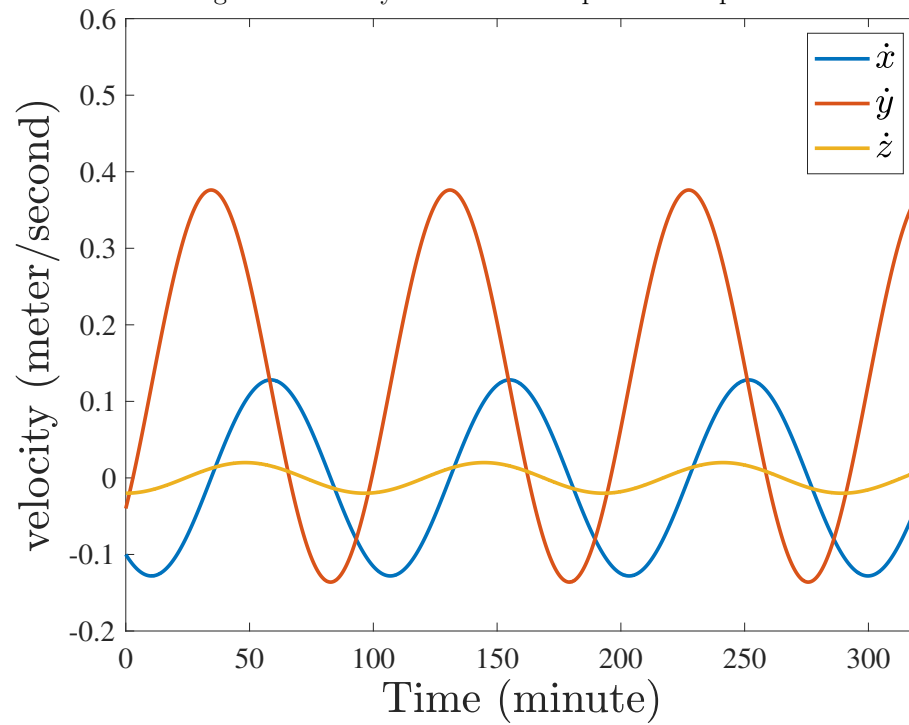


Figure 2: velocity of the Hubble space telescope



## 2 Question 2

Used below equations to find the orbital elements.

$$\mathbf{r} = [1600 \quad 5310 \quad 3800]_{km}^T, \quad \mathbf{v} = [-7.35 \quad 0.46 \quad 2.47]_{km/sec}^T$$

### 2.1 part a

$$\begin{aligned} \mathbf{h} &= \mathbf{r} \times \mathbf{v} \\ v_r &= \frac{\mathbf{r} \cdot \mathbf{v}}{r} \\ \mathbf{e} &= \frac{\mathbf{v} \times \mathbf{h} - \mu \frac{\mathbf{r}}{r}}{\mu} \\ a &= \frac{h^2}{\mu(1 - e^2)} \\ \mathbf{N} &= [0 \quad 0 \quad 1]^T \times \mathbf{h} \\ \theta &= \begin{cases} \arccos\left(\frac{\mathbf{e} \cdot \mathbf{r}}{er}\right), & v_r \geq 0 \\ 2\pi - \arccos\left(\frac{\mathbf{e} \cdot \mathbf{r}}{er}\right), & v_r < 0 \end{cases} \\ \Omega &= \begin{cases} \arccos\left(\frac{\mathbf{N}(1)}{N}\right), & \mathbf{N}(2) \geq 0 \\ 2\pi - \arccos\left(\frac{\mathbf{N}(1)}{N}\right), & \mathbf{N}(2) < 0 \end{cases} \\ \omega &= \begin{cases} \arccos\left(\frac{\mathbf{N} \cdot \mathbf{e}}{Ne}\right), & \mathbf{e}(3) \geq 0 \\ 2\pi - \arccos\left(\frac{\mathbf{N} \cdot \mathbf{e}}{Ne}\right), & \mathbf{e}(3) < 0 \end{cases} \\ i &= \arccos\left(\frac{\mathbf{h}(3)}{h}\right) \end{aligned}$$

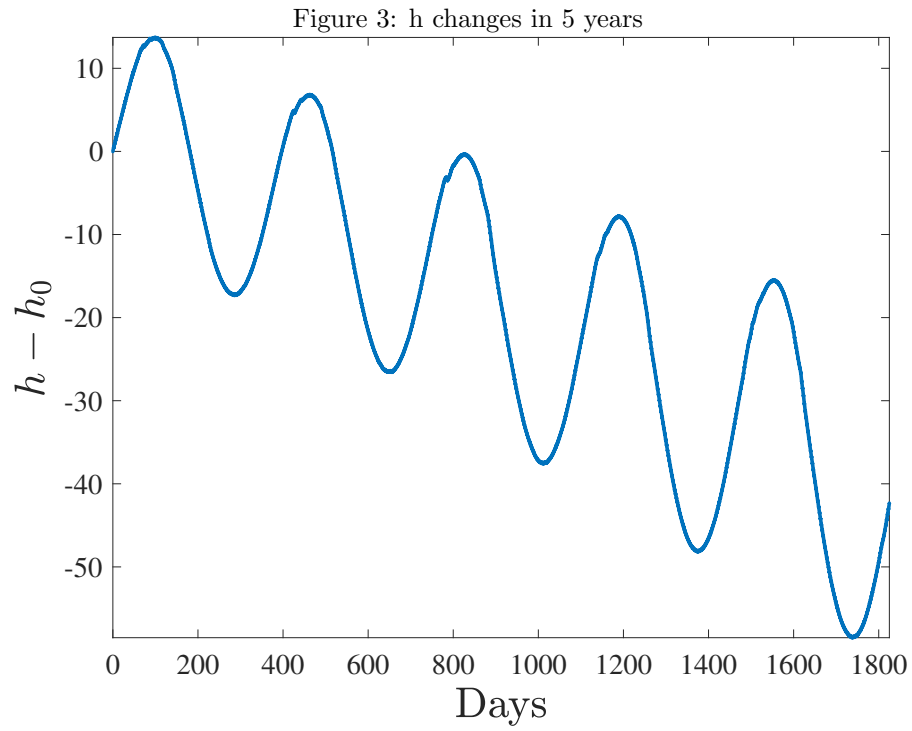
From the above equations, initial conditions will find. The below equation shows the force of solar radiation.

$$P_{SRP} = \nu \frac{S}{c} C_R \frac{A_s}{m}$$

$\nu$  calculates if the satellite is in the earth's shadow or not. Then used the below equations for rate changes.

$$\begin{aligned} \frac{dh}{dt} &= -p_{SR} r u_s \\ \frac{de}{dt} &= -p_{SR} \left( \frac{h}{\mu} \sin(\theta) u_r + \frac{1}{\mu h} ((h^2 + \mu r) \cos(\theta) \mu e r) u_s \right) \\ \frac{d\theta}{dt} &= \frac{h}{r^2} - \frac{p_{SR}}{eh} \left( \frac{h^2}{\mu} \cos(\theta) u_r - \left( r + \frac{h^2}{\mu} \right) \sin(\theta) u_s \right) \\ \frac{d\Omega}{dt} &= -p_{SR} \frac{r}{h \sin(i)} \sin(\omega + \theta) u_w \\ \frac{di}{dt} &= -p_{SR} \frac{r}{h} \cos(\omega + \theta) u_w \\ \frac{d\omega}{dt} &= -p_{SR} \left( \frac{1}{eh} \left( \frac{h^2}{\mu} \cos(\theta) u_r - \left( r + \frac{h^2}{\mu} \right) \sin(\theta) u_s \right) - \frac{r \sin(\omega - \theta)}{h \tan(i)} u_w \right) \end{aligned}$$

For this purpose, example 10.9 was used, the Gauss planetary equations for solar radiation pressure (Equations 10.106). The script file is Q2.m.



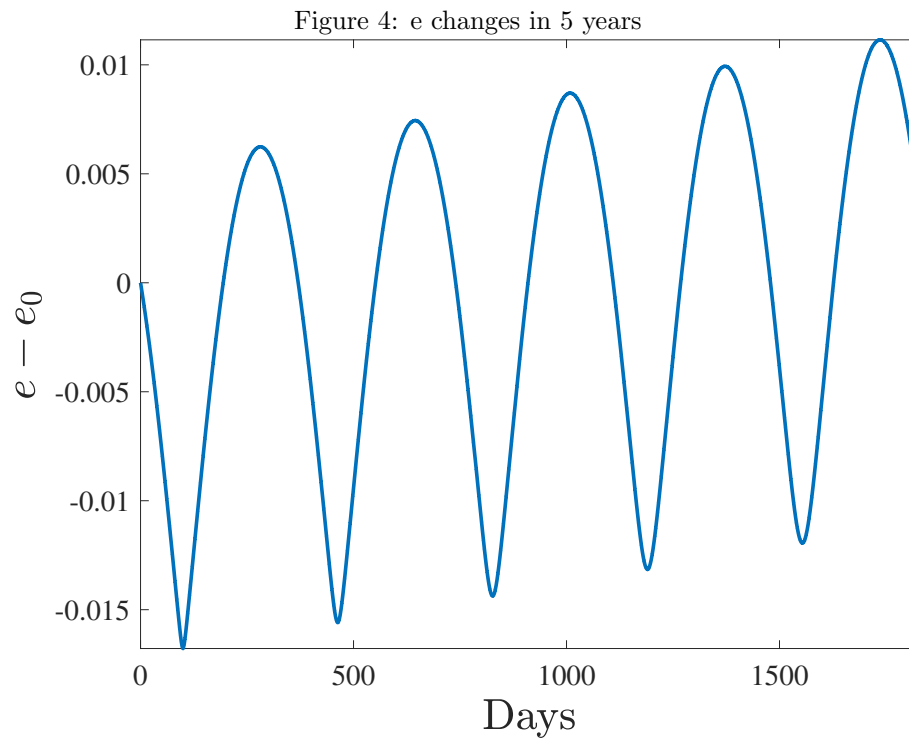


Figure 5:  $\theta$  changes in 5 years (the satellite has a short period of 5 years of changes)

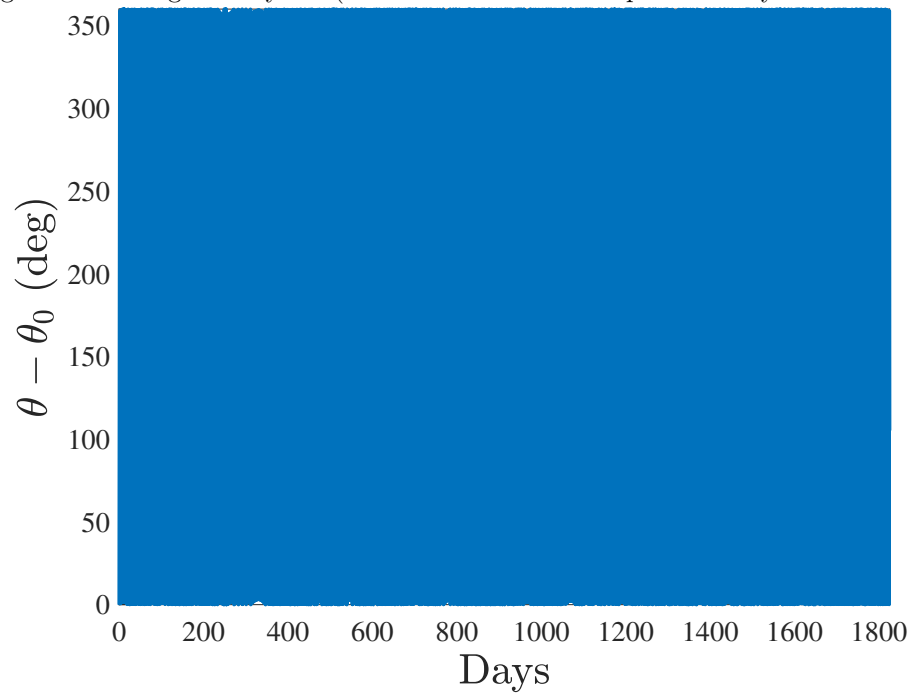


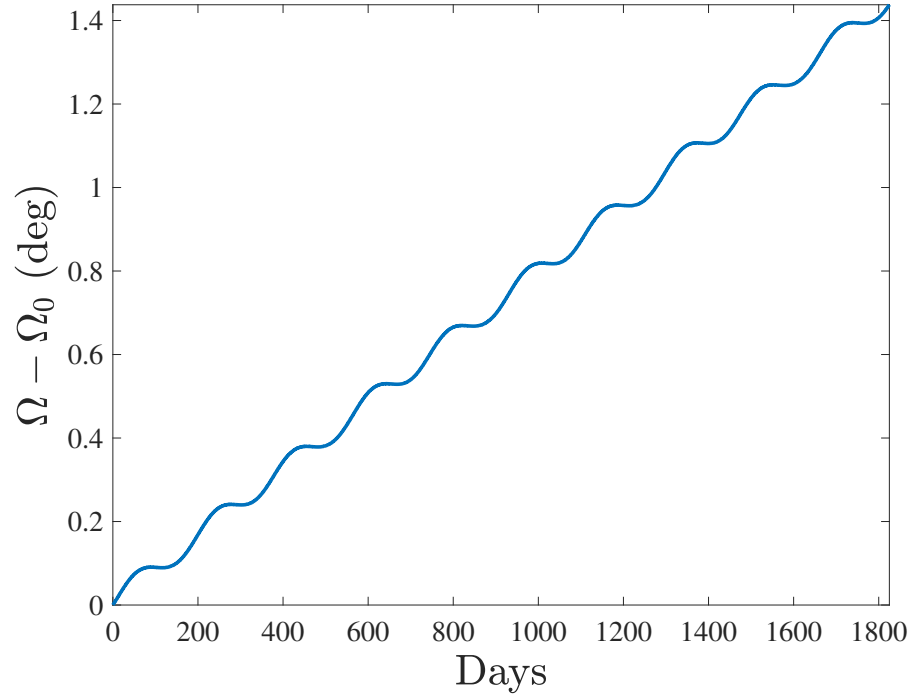
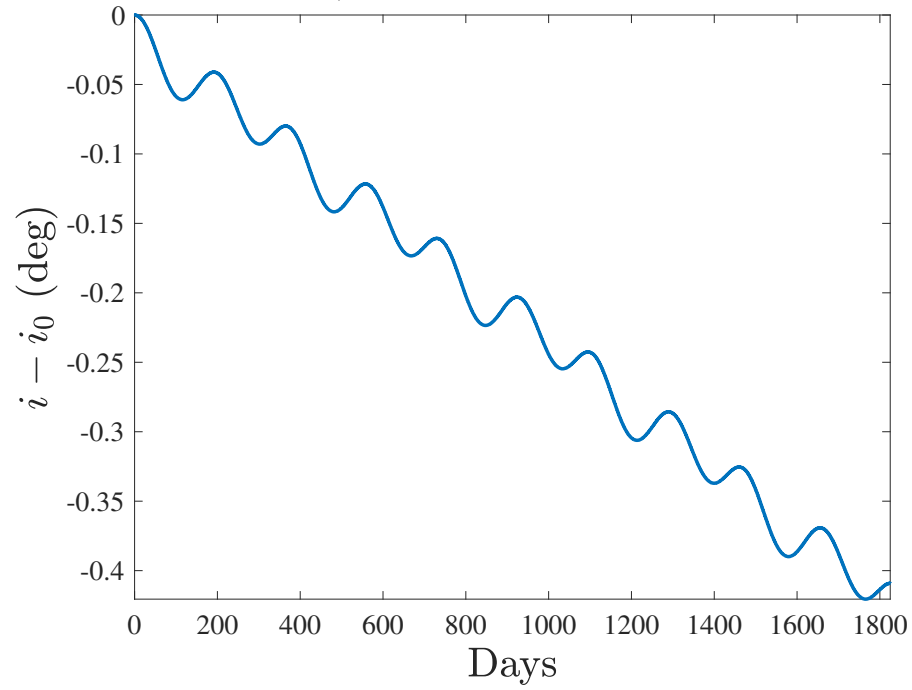
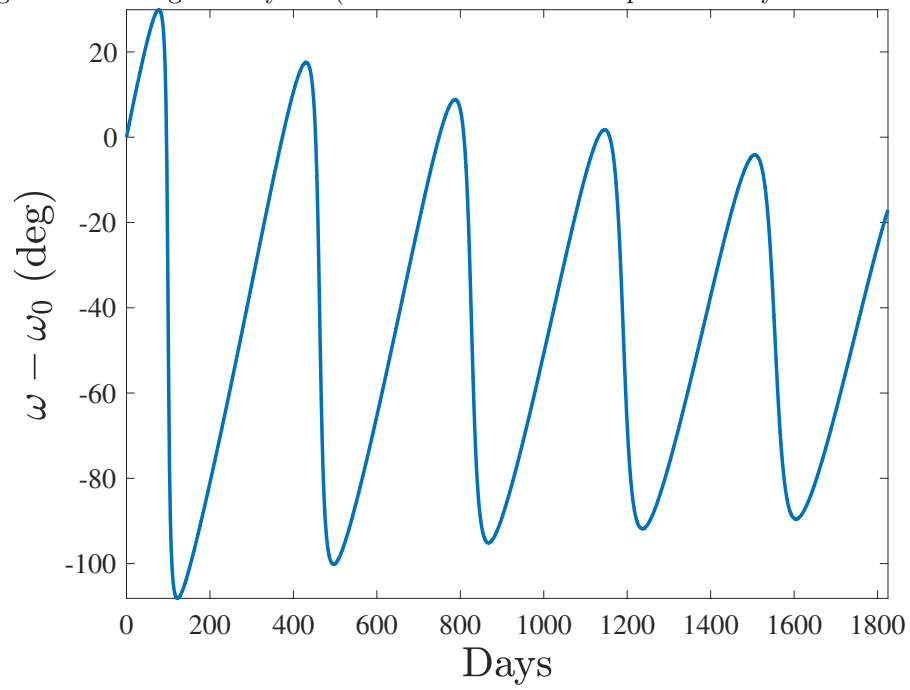
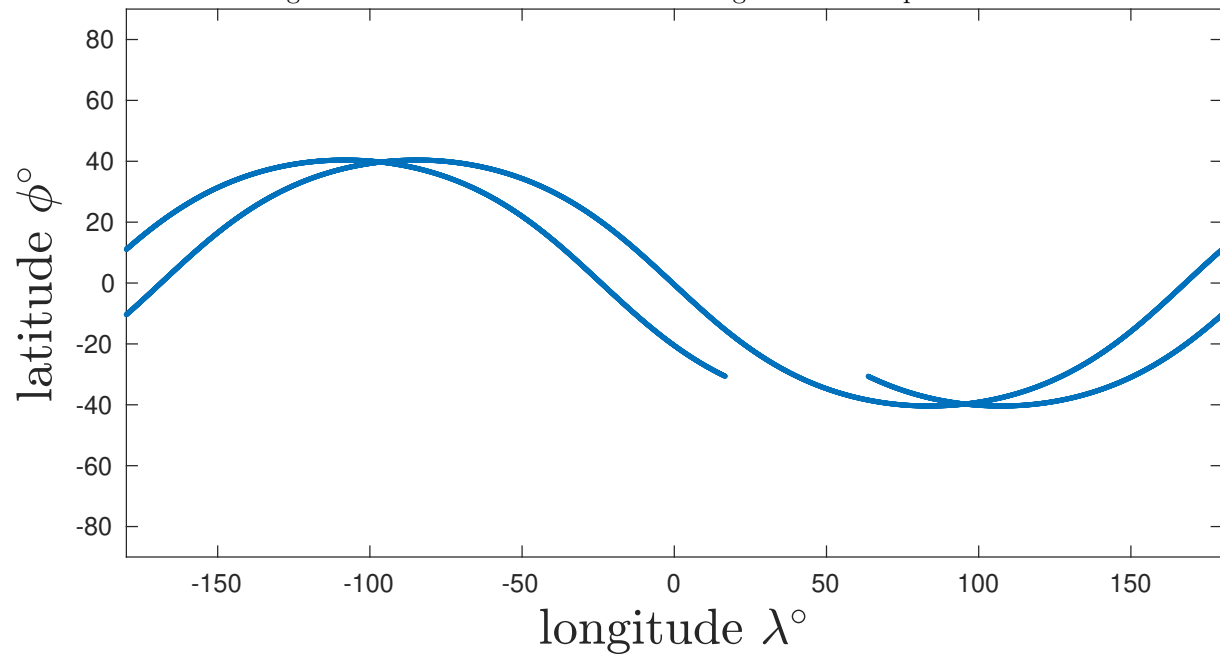
Figure 6:  $\Omega$  changes in 5 years (the satellite has a short period of 5 years of changes)Figure 7:  $i$  changes in 5 years (the satellite has a short period of 5 years of changes)

Figure 8:  $\omega$  changes in 5 years (the satellite has a short period of 5 years of changes)

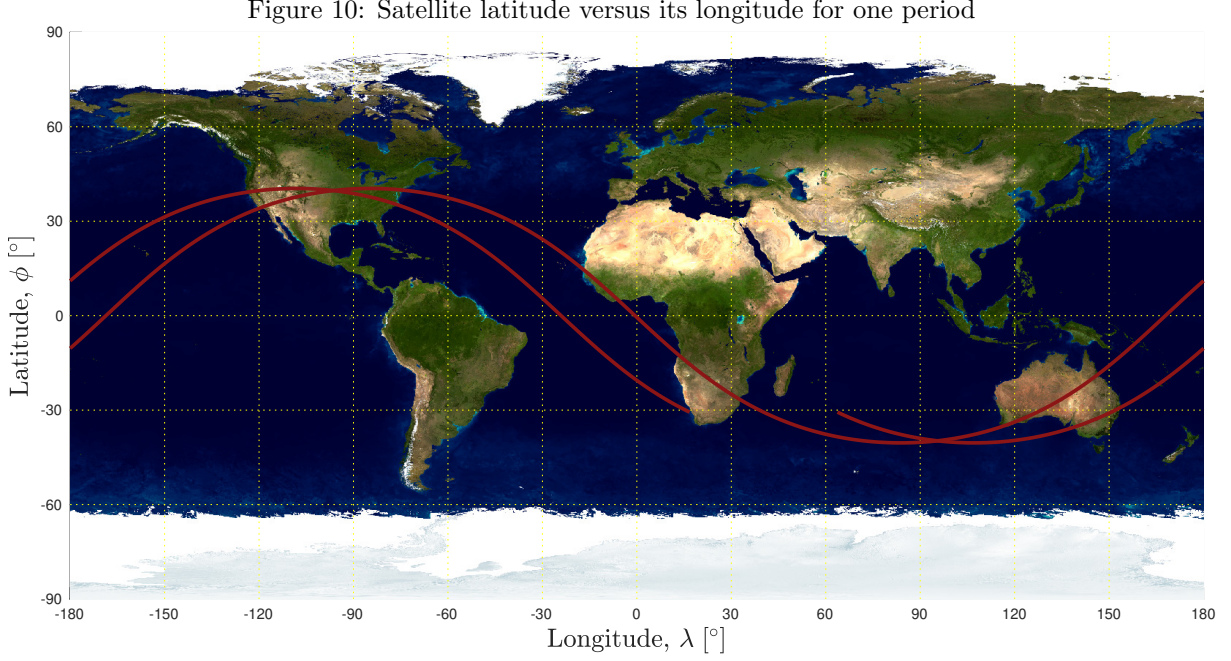
## 2.2 part b

Used orbital elements from above code and `sv.from.coe` function from Curtis book to get satellite position and used Q1 short project to plot ground track.

Figure 9: Satellite latitude versus its longitude for one period



Below the figure drawn provided by tamaskis, please click [here](#) to see the source code. Please use mentioned library to run code or skip part on earth fig.



### 3 Question 3

$$\Phi = -\frac{J_3 R^3 \mu \left( 3 \cos(\phi) - 5 \cos(\phi)^3 \right)}{2 r^4}$$

$$\frac{\partial \Phi}{\partial x} = \frac{xz}{r^3 \sin(\phi)}, \quad \frac{\partial \Phi}{\partial y} = \frac{yz}{r^3 \sin(\phi)}, \quad \frac{\partial \Phi}{\partial x} = \frac{\sin(\phi)}{r}$$

$$\frac{\partial \Phi}{\partial x} = \frac{\partial \Phi}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial \Phi}{\partial \phi} \frac{\partial \phi}{\partial x}$$

$$\frac{\partial \Phi}{\partial y} = \frac{\partial \Phi}{\partial r} \frac{\partial r}{\partial y} + \frac{\partial \Phi}{\partial \phi} \frac{\partial \phi}{\partial y}$$

$$\frac{\partial \Phi}{\partial z} = \frac{\partial \Phi}{\partial r} \frac{\partial r}{\partial z} + \frac{\partial \Phi}{\partial \phi} \frac{\partial \phi}{\partial z}$$

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\frac{\partial r}{\partial x} = \frac{x}{r}, \quad \frac{\partial r}{\partial y} = \frac{y}{r}, \quad \frac{\partial r}{\partial z} = \frac{z}{r}$$

$$\frac{\partial \Phi}{\partial x} = \frac{2 J R^3 \mu x \left( 3 \cos(\phi) - 5 \cos(\phi)^3 \right)}{r^6}$$

$$\frac{\partial \Phi}{\partial y} = \frac{2 J R^3 \mu y \left( 3 \cos(\phi) - 5 \cos(\phi)^3 \right)}{r^6}$$



$$\frac{\partial \Phi}{\partial z} = \frac{2 J R^3 \mu z \left( 3 \cos(\phi) - 5 \cos(\phi)^3 \right)}{r^6}$$

using the fact that  $\cos(\phi) = \frac{z}{r}$  leads to the following expressions for the gradient of perturbing potential  $\Phi$ :

$$\frac{\partial \Phi}{\partial x} = - \frac{2 J R^3 \mu x \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6}$$

$$\frac{\partial \Phi}{\partial y} = - \frac{2 J R^3 \mu y \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6}$$

$$\frac{\partial \Phi}{\partial z} = - \frac{2 J R^3 \mu z \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6}$$

$$[\mathbf{Q}]_{Xr} = \begin{bmatrix} -\sin(\Omega) \cos(i) \sin(u) + \cos(\Omega) \cos(u) & \cos(\Omega) \cos(i) \sin(u) + \sin(\Omega) \cos(u) & \sin(i) \sin(u) \\ -\sin(\Omega) \cos(i) \sin(u) - \cos(\Omega) \cos(u) & \cos(\Omega) \cos(i) \sin(u) - \sin(\Omega) \cos(u) & \sin(i) \sin(u) \\ \sin(\Omega) \sin(i) & -\cos(\Omega) \sin(i) & \cos(i) \end{bmatrix}$$

$$\begin{bmatrix} p_r \\ p_s \\ p_w \end{bmatrix} = [\mathbf{Q}]_{Xr} \begin{bmatrix} p_X \\ p_Y \\ p_Z \end{bmatrix}$$

$$\begin{bmatrix} p_r \\ p_s \\ p_w \end{bmatrix} = [\mathbf{Q}]_{Xr} \begin{bmatrix} - \frac{2 J R^3 \mu x \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6} \\ - \frac{2 J R^3 \mu y \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6} \\ - \frac{2 J R^3 \mu z \left( \frac{5 z^3}{r^3} - \frac{3 z}{r} \right)}{r^6} \end{bmatrix}$$

$$\frac{de}{dt} = \frac{(1-e^2)^{1/2}}{na} (\sin(\theta)p_r + (\cos(\theta) + \cos(E))p_s)$$

where:

$$n = \sqrt{\frac{\mu}{a^3}}$$

## 4 Question 4

$$\mathbf{r}(0) = [0.994 \quad 0 \quad 0], \quad \mathbf{v}(0) = [0 \quad -2.001585106 \quad 0]$$

$$m_{earth} = 5.974e24_{kg}, \quad m_{moon} = 7.348e22_{kg}, \quad r_{12} = 3.844e5_{km}$$

#### 4.1 part a

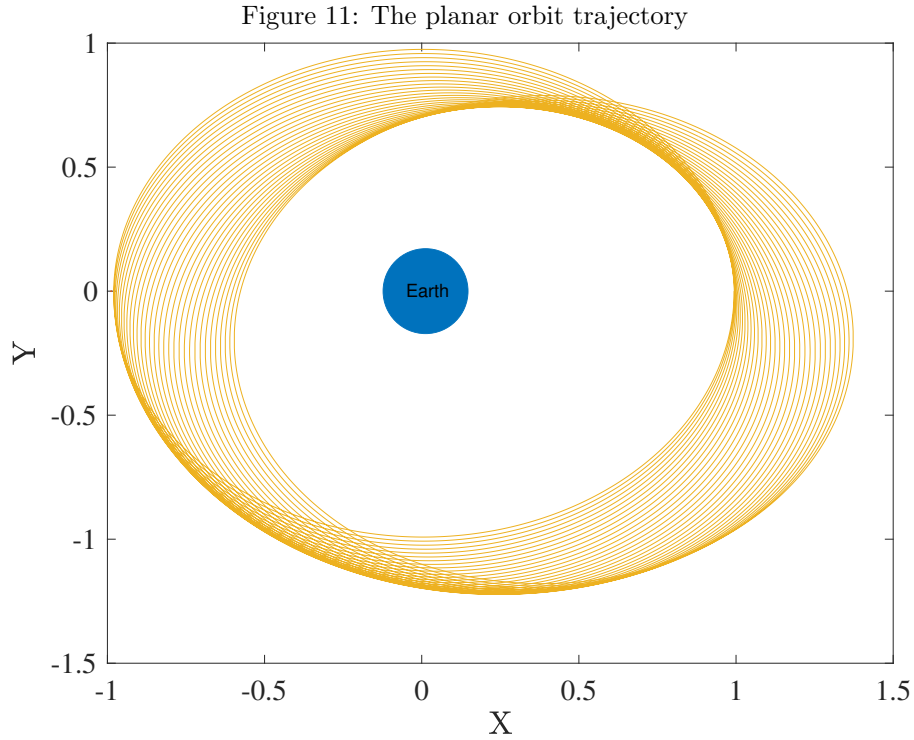
In Canonical units, jacobí constant is defined as:

$$C = \frac{1}{2}v^2 - \frac{1}{2}(x^2 + y^2) - \mu = 1.4970$$

#### 4.2 part b

Solve the equation of motion for the moon earth system in canonical units:

$$\begin{aligned}\ddot{x} &= 2\dot{y} + x - \frac{1-\mu}{r_1^3}(x-\mu) - \frac{\mu}{r_2^3}(x+1-\mu) \\ \ddot{y} &= -2\dot{x} + xy - \frac{1-\mu}{r_1^3}y - \frac{\mu}{r_2^3}y \\ \ddot{z} &= -\frac{1-\mu}{r_1^3}z - \frac{\mu}{r_2^3}z\end{aligned}$$



#### 4.3 part c

From Bong Wie space vehicle dynamics, we have:

$$\begin{aligned}U_{XX} &= \frac{\partial^2 U}{\partial X^2} \Big|_{X=X_0} = 1 - \left( (1-\mu) \left( \frac{1}{r_1^3} - 3 \frac{(X_0 - \mu)^2}{r_1^5} \right) + \mu \left( \frac{1}{r_2^3} - 3 \frac{(X_0 + 1 - \mu)^2}{r_2^5} \right) \right) \\ U_{YY} &= \frac{\partial^2 U}{\partial Y^2} \Big|_{Y=Y_0} = 1 - \left( (1-\mu) \left( \frac{1}{r_1^3} - 3 \frac{Y_0^2}{r_1^5} \right) + \mu \left( \frac{1}{r_2^3} - 3 \frac{Y_0^2}{r_2^5} \right) \right)\end{aligned}$$

$$\lambda^4 + (4 - U_{XX} - U_{YY})\lambda^2 + U_{XX}U_{YY} = 0$$

Above equation solved in Q4.m MATLAB file in part c section.

$$\lambda_{1,2} = 0.3460 \pm 0i$$

$$\lambda_{3,4} = 0 \pm 1.0803i$$

The in-plane motion has a divergent mode as well as an oscillatory mode with a nondimensional frequency  $\omega_{xy} = 1.0803$  and the period of the in-plane oscillatory mode is 25.25 days.

$$\omega_{xy} = 1.0803 \times \omega_{earth-moon} = 2.8794e - 6 \rightarrow \tau = \frac{2\pi}{\omega_{xy}} = 5.4553e5_{sec} = 25.2559_{day}$$

## 5 Bonus

General equation of mottion is:

$$\delta\ddot{x} - 3n^2\delta x - 2n\delta\dot{y} = 0$$

$$\delta\ddot{y} + 2n\delta\dot{x} = 0$$

$$\delta\ddot{z} + 2n\delta\dot{z} = 0$$

where:

$$n = \sqrt{\frac{\mu}{r^3}}$$

above differential equation can be solved in curtis book and final answer is:

$$\delta x = 4\delta x_0 + \frac{2}{n}\delta\dot{y}_0 + \frac{\delta\dot{x}_0}{n}\sin(nt) - \left(3\delta x_0 + \frac{2}{n}\delta\dot{y}_0\right)\cos(nt)$$

$$\delta y = \delta y_0 - \frac{2}{n}\delta\dot{x}_0 - 3(2n\delta x_0 + \delta\dot{y}_0)t + 2\left(3\delta x_0 + \frac{2}{n}\delta\dot{y}_0\right)\sin(nt) + \frac{2}{n}\delta\dot{x}_0\cos(nt)$$

$$\delta z = \frac{1}{n}\delta\dot{z}_0\sin(nt) + \delta z_0\cos(nt)$$

$$\delta\mathbf{r}(t) = \begin{bmatrix} \delta x(t) \\ \delta y(t) \\ \delta z(t) \end{bmatrix}, \quad \delta\mathbf{v}(t) = \begin{bmatrix} \delta u(t) \\ \delta v(t) \\ \delta w(t) \end{bmatrix}$$

$$\delta\mathbf{r}_0 = \begin{bmatrix} \delta x_0 \\ \delta y_0 \\ \delta z_0 \end{bmatrix}, \quad \delta\mathbf{v}_0 = \begin{bmatrix} \delta u_0 \\ \delta v_0 \\ \delta w_0 \end{bmatrix}$$

$$\delta\mathbf{r}(t) = [\Phi_{rr}(t)]\delta\mathbf{r}_0 + [\Phi_{rv}(t)]\delta\mathbf{v}_0$$

$$\delta\mathbf{v}(t) = [\Phi_{vr}(t)]\delta\mathbf{r}_0 + [\Phi_{vv}(t)]\delta\mathbf{v}_0$$

$$\Phi_{rr}(t) = \begin{bmatrix} 4 - 3\cos(nt) & 0 & 0 \\ 6(\sin(nt) - nt) & 1 & 0 \\ 0 & 0 & \cos(nt) \end{bmatrix}$$

$$\Phi_{rv}(t) = \begin{bmatrix} \frac{1}{n}\sin(nt) & \frac{2}{n}(1 - \cos(nt)) & 0 \\ \frac{2}{n}(\cos(nt) - 1) & \frac{1}{n}(4\sin(nt) - 3nt) & 0 \\ 0 & 0 & \frac{1}{n}\sin(nt) \end{bmatrix}$$

$$\Phi_{vr}(t) = \begin{bmatrix} 3n \sin(nt) & 2n \sin(nt) & 0 \\ 6n(\cos(nt) - 1) & 0 & 0 \\ 0 & 0 & -n \sin(nt) \end{bmatrix}$$

$$\Phi_{vv}(t) = \begin{bmatrix} \cos(nt) & 2 \sin(nt) & 0 \\ -2n \sin(nt) & 4n \cos(nt) - 3 & 0 \\ 0 & 0 & \cos(nt) \end{bmatrix}$$

The Clohessy-Wiltshire matrices, for  $t_f = 8h$  and  $n = 0.0011_{rad/sec}$ .

$$\Phi_{rr} = \begin{bmatrix} 1.0361 & 0 & 0 \\ -188.4918 & 1.0000 & 0 \\ 0 & 0 & 0.9880 \end{bmatrix}$$

$$\Phi_{rv} = \begin{bmatrix} 142 & 22 & 0 \\ -22 & -86970 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\Phi_{vr} = \begin{bmatrix} -0.0005 & -0.0003 & 0 \\ -0.0001 & 0 & 0 \\ 0 & 0 & 142.4195 \end{bmatrix}$$

$$\Phi_{vv} = \begin{bmatrix} 0.9880 & -0.3092 & 0 \\ 0.0003 & -2.9957 & 0 \\ 0 & 0 & 0.9880 \end{bmatrix}$$

At time  $t_f$ , Hubble arrives at spaceX, at the origin of the CW frame, which means  $\delta \mathbf{r}_f = \delta \mathbf{r}_f = \mathbf{0}$ . At  $t_f$  we find:

$$\mathbf{0} = \Phi_{rr}(t_f) \delta \mathbf{r}_0 + \Phi_{rv}(t_f) \delta \mathbf{v}_0^+ \rightarrow \delta \mathbf{v}_0^+ = -\Phi_{rr}(t_f)^{-1} \Phi_{rv}(t_f) \delta \mathbf{r}_0$$

$$\delta \mathbf{v}_0^+ = \begin{bmatrix} 0.9923 \\ -0.3086 \\ 104605.6 \end{bmatrix}_{m/sec}$$

$$\delta \mathbf{v}_f^- = \Phi_{vr}(t_f) \delta \mathbf{r}_0 + \Phi_{vv}(t_f) \delta \mathbf{v}_0^+ = \Phi_{vr}(t_f) \delta \mathbf{r}_0 + \Phi_{vv}(t_f) (-\Phi_{rr}(t_f)^{-1} \Phi_{rv}(t_f) \delta \mathbf{r}_0)$$

$$\delta \mathbf{v}_f^- = \begin{bmatrix} 0.9577 \\ 0.9136 \\ 105878.4 \end{bmatrix}_{m/sec}$$

So the change of velocity of Hubble is  $|\delta \mathbf{v}_0^- - \delta \mathbf{v}_0^+|$  and the change of velocity of SpaceX is  $|\delta \mathbf{v}_f^-|$ . The velocity change in this maneuver is too big.

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