

The changing notion of chimera states, a critical review

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Chimera states, states of coexistence of synchronous and asynchronous motion, have been a subject of extensive research since they were first given a name in 2004. Increased interest has led to their discovery in ever new settings, both theoretical and experimental. Less well-discussed is the fact that successive results have also broadened the notion of what actually constitutes a chimera state. In this article, we critically examine how the results for different model types and coupling schemes, as well as varying implicit interpretations of terms such as coexistence, synchrony and incoherence, have influenced the common understanding of what constitutes a chimera. We cover both theoretical and experimental systems, address various chimera-derived terms that have emerged over the years and finally reflect on the question of chimera states in real-world contexts.

Almost twenty years ago, Kuramoto and Battogtokh realized that a ring of coupled identical oscillators, when initialized in the right way, would split into two spatial domains¹: In one of these domains, the oscillators remained mutually synchronized, while in the other, they drifted with varying average frequencies. The coupling between the oscillators in their system was nonlocal, that is, of intermediate range between local (nearest-neighbor) and global (all-to-all symmetric) coupling. This was also the case when Abrams and Strogatz gave the phenomenon its name – chimera state – and expanded the mathematical understanding of how it arises². While they believed the chimera state to be peculiar to nonlocal coupling, Abrams and Strogatz never defined this as one of its definite properties, nor did they attempt to restrict the type of oscillator ensemble wherein it might be acceptable to occur. Thus, the research community was free to discover chimera states in other systems as well, ranging from various types of oscillators via maps to cellular automata; from local via various forms of nonlocal to global coupling; in both simulations and, from 2012 onward^{3,4}, in experiments. The general extent of this research is well known, and several reviews on chimera states have already been written^{5–8}. Less reflected upon is how the chimera community exploited the soft original chimera definition to associate ever new results with the illustrious term. Over time, the name itself implicitly changed and expanded its meaning, to the point where very different phenomena were all labeled chimera states, usually with little consideration of how strongly each of them is related to the others. Similarly, a lot of chimera-derived terminology was introduced, with little attempt, beyond a 2016 classification scheme⁹, to reflect on the extent of the relevance of many of the new terms. This may explain both why the field has been able to accelerate as quickly as it has, and why no unified theory of all chimera states has been established. Notably, while chimeras have now been realized in a wide range of experiments, the question of their potential existence outside laboratories has not yet been conclusively resolved.

I. INTRODUCTION

Towards the end of the 1980s, Kaneko discovered both clustering and chimera states in globally coupled logistic maps¹⁰. Altogether, he identified four different types of attractors, depending on the sizes of the clusters they contain (including “clusters” of size 1). These were (a) fully synchronized motion; as well as attractors with (b) a small number of clusters, much smaller than the system size N ; (c) a large number of clusters in the order of N , but at least one cluster N_1 comparable in size to N ; (d) only small clusters in the order of 1. The still unnamed chimera state was only one among several different states of the third type. See Fig. 1.

Not long after, clustering was found and studied in globally coupled phase oscillators by Golomb et al.¹¹ and Okuda¹², and in globally coupled Stuart-Landau oscillators by Hakim and Rappel¹³. Nakagawa and Kuramoto found a chimera state in the latter system in 1992, but did not treat it as more than a step from clustering to fully chaotic motion¹⁴. Not before ten years later, after studying the nonlocally coupled CGLE and a ring of nonlocally coupled phase oscillators, did Kuramoto and Battogtokh publish the idea that the coexistence of synchronized and non-synchronized motion might be an interesting phenomenon in its own right¹.

In 2004, Abrams and Strogatz² gave the presumptively new kind of symmetry breaking its name – *chimera state* – and in the years that followed, the number of reported chimeras has vastly increased, as summed up in several review papers^{5–8}. Chimera states have also been found in ever more different models, meaning that the general conditions under which one or more of them could be said to occur, have become increasingly varied as well. Over the course of this chimera research explosion, Schmidt and Krischer identified Kaneko’s earliest chimera in 2013¹⁵, while Sethia and Sen rediscovered Nakagawa and Kuramoto’s 1992 chimera only few months later¹⁶. Yet, what no one seems to have pointed out so far, is how the entirety of states considered to be chimeras in the first place has evolved into an ever more diverse collection of dynamical phenomena. This tacit liberalization of the term “chimera

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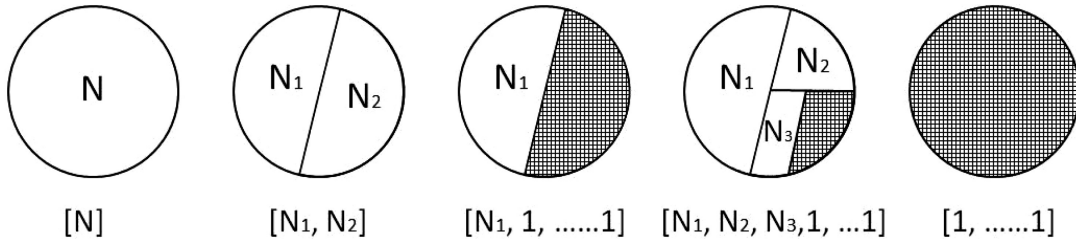


FIG. 1. Schematic examples of the different attractor types reported by Kaneko in 1989¹⁰. From left to right: $[N]$: Coherent (complete synchronization). $[N_1, N_2]$: Partition to two clusters as an example of attractor wherein the number of clusters does not grow with N . $[N_1, 1, 1, \dots, 1]$: Chimera state as an example of an attractor with $O(N)$ clusters, but at least one cluster comparable to N . $[N_1, N_2, N_3, 1, \dots, 1]$: A more complex attractor of the same fundamental type. $[1, 1, \dots, 1]$: Complete desynchronization and no clusters that grow in size with N . Reprinted from K. Kaneko, Chaos 25 (9), 097608, 2015, with the permission of AIP Publishing. doi:10.1063/1.4916925.

state” includes steps such as the extension from oscillators to maps, as well as ever new implicit interpretations of terms such as symmetric coupling, synchrony and incoherence, and is critically examined in this review article.

Also noticeable is how Kaneko’s original globally-coupled-maps chimera state was discovered and presented in the context of more or less related attractors right from the start. In contrast, Kuramoto and Battogtokh’s chimera under nonlocal coupling was carefully constructed¹ (for parameters where the fully synchronized solution is stable) and thus had little original context. Context was instead created around it whenever new results in a wide range of models were produced, a process that has yielded terms such as “alternating chimera”¹⁷, “multichimera”¹⁸, , to name but a few. Seldom, though, is any particular system found to exhibit more than a few of these derived phenomena. This issue will also be discussed below.

II. FROM NONLOCAL TO GLOBAL COUPLING AND VARIOUS COUPLING SCHEMES

When the term “chimera state” was coined by Abrams and Strogatz in 2004, they described it quite simply as “an array of identical oscillators split[ting] into two domains: one coherent and phase locked, the other incoherent and desynchronized”². In their article, they cited only two prior works discerning this kind of previously unnamed state: The first was a 2002 paper by Kuramoto and Battogtokh on its occurrence in the one-dimensional complex Ginzburg-Landau equation (CGLE) with nonlocal coupling. Here, the coexistence of synchronized and unsynchronized oscillators was also found to persist as the originally complex-valued oscillatory medium was reduced to a phase-oscillator approximation¹. The second was a 2004 paper on the two-dimensional equivalent, a spiral wave with an incoherent core, found both in the phase approximation of the nonlocally coupled, spatially two-dimensional CGLE and in a 2D array of FitzHugh-Nagumo oscillators¹⁹.

Kuramoto and Battogtokh considered their original chimera state to be “among the variety of patterns which are characteristic to nonlocally coupled oscillators”¹. In a 2006 paper, Abrams and Strogatz similarly draw the tentative conclusion that these dynamics are “peculiar to the intermediate case of nonlocal coupling”²⁰. The basis of their reasoning is two-fold:

Firstly, no other chimeras were known to them. Secondly, the coexistence of synchrony and incoherence in identical sine-coupled phase oscillators (the system to which Kuramoto and Battogtokh reduced their nonlocally coupled CGLE in 2002) does indeed become impossible when the oscillators are coupled globally²⁰. When Sethia and co-workers described chimera states in delay-coupled oscillators in 2008²¹, they were also under the impression that nonlocal coupling is indispensable. The same goes for Wolfrum et al. during their 2011 investigations of the Lyapunov spectra and stability of chimera states of various sizes^{22,23}.

The mentioned papers all concentrated their attention on the case of weak coupling, where the CGLE can be approximated by phase oscillators, if not outright restricting themselves to phase oscillators as the starting point of their investigations. So did a significant number of other papers published during the first decade of chimera-state research^{17,24–40}. An exception was Laing’s 2010 paper⁴¹, in which he analyses a chimera state in an ensemble of Stuart-Landau oscillators without resorting to the phase reduction, but here too, the coupling is relatively weak and the amplitudes of individual oscillators deviate only a few percent from their average value. Notably, all his non-synchronized oscillators are restricted to the same closed curve in the complex plane, which allows for effectively parametrizing their position by a phase alone, even though their amplitudes vary.

Only in 2013 did Sethia et al. show that a coexistence of coherent and incoherent dynamics can also occur in the nonlocally coupled CGLE in the case of strong coupling. Here, the amplitude in the incoherent part of the system varies strongly in both space and time, inspiring the name “amplitude-mediated chimeras” (AMC) and preventing any attempt at a phase reduction⁴². Less than half a year later, Schmidt et al. successfully took the next step and identified a chimera state in an ensemble of Stuart-Landau oscillators with nonlinear *purely global* (symmetric all-to-all) coupling¹⁵. Like the AMCs of Sethia et al., this state exhibits strong amplitude fluctuations in the incoherent oscillators. Moreover, it persists when adding diffusion and thus transitioning to a spatially extended 2D medium, where synchrony and incoherence form several clearly distinguishable intertwined islands. Shortly after, Sethia and Sen reported amplitude-mediated chimeras for Stuart-Landau oscillators with linear global coupling as well¹⁶. In 2015, Laing comple-

mented these efforts by producing chimera states in a ring of oscillators with only local (nearest-neighbor) coupling⁴³.

The chimera states reported before the mentioned paper by Schmidt et al. had all been found in nonlocally coupled systems in the broad sense of the coupling being neither next-neighbor/diffusional (“local coupling”) nor all-to-all symmetric (“global coupling”). However, the applied coupling schemes already encompassed a variety beyond Kuramoto and Battogtokh’s original 1D ring topology and exponentially decaying coupling kernel. Some of the alternative forms of nonlocal coupling were rather small variations, such as the use of a cosine kernel^{2,20} or a step function^{22,38} to limit the extent of the influence of each point in the system on the others. (As long as the extent of the ring is restricted to $-\pi < x \leq \pi$, the cosine kernel $G(x - x') \propto 1 + A \cos(x - x')$, and thus the coupling strength, also decreases monotonously with distance².) More radical was the idea of dividing the oscillators into two populations with symmetric all-to-all coupling within each group, as well as a weaker coupling between the groups²⁶. See Fig. 2. Besides Kuramoto and Battogtokh’s original system, this “simplest network of networks”⁴¹, is possibly the most influential theoretical model supporting chimera states, and similar models have been the subject of a large number of subsequent works^{3,17,28,29,31,32,35,36,39–41,44–49,50}.

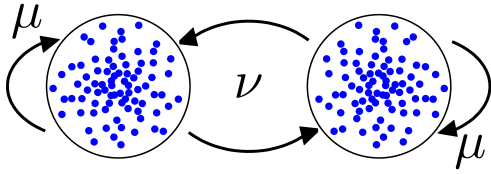


FIG. 2. Schematic representation of the two-groups model introduced to the field of chimera states by Abrams and Strogatz in 2008²⁶ and inspiring a large number of subsequent works. μ denotes the strength of the coupling within each group and ν that between the groups, usually with $\mu > \nu$. Reprinted figure with permission from M. Panaggio, D. M. Abrams, P. Ashwin & C. R. Laing, *Physical Review E* 93 (1), 012218, 2016. doi:10.1103/PhysRevE.93.012218. Copyright 2016 by the American Physical Society.

Among the chimera-supporting systems inspired by the two-groups model were networks of three^{35,36} and eight populations⁴⁰, respectively. It was possibly also the first model for which the question was asked, and answered in the affirmative, whether an observed chimera state will continue to exist if the originally identical oscillators are given heterogeneous frequencies^{29,30}. In a different variation from the original two-groups model, connections between oscillators are gradually removed in order to further test the robustness of the chimera, and it was found to be more sensitive to removal of intra-group than of inter-group links³⁹. Both the heterogeneous frequencies and the removal of connections were motivated by the aim to create a more realistic model, as neither really identical oscillatory units nor perfectly symmetric coupling schemes are likely to exist in nature^{29,30,39}.

Other variants of the two-groups model producing chimera states include one in which the individual links between the populations are randomly switched on and off at equally spaced

time intervals⁴⁶, one with different phase lags in the coupling within and between populations⁴⁸, and a delay-coupled version with different intra- and inter-population coupling delays^{31,32}. Schmidt et al. also invoke the two-groups model when explaining the stability of their globally coupled chimera. Here, the synchronized and incoherent oscillators effectively form two self-organized groups, and these groups exert different influences on the respective other group, thereby further reinforcing the chimera state once spontaneously formed¹⁵.

III. EXPERIMENTAL CHIMERAS AND A BROADER CHIMERA CONCEPT

One of the very first experimental realizations of a chimera state was also based on the two-groups model. It was a 2012 photochemical experiment by Tinsley et al., involving 40 photosensitive Belousov-Zhabotinsky (BZ) oscillator beads. Here, each of the beads emits light of a certain intensity, which is recorded with a CCD camera and projected selectively back on the beads by a spatial light modulator (SLM)³⁵¹. About a year later, the two-groups model also formed the basis for a purely mechanical chimera, without any computer-mediated coupling, found in metronomes placed on two swings connected by springs: The intra-group coupling is conveyed by vibrations of the respective swing, while inter-group coupling happens via the springs⁴⁵. Published by Hagerstrom et al. simultaneously with the BZ chimera paper, but inspired by a different model⁵², was an optical realization of an array of coupled maps⁴: Here, the different parts of the cross-section of a beam of circularly polarized light have different phases when they emerge from an SLM. As the beam passes through an optical setup, these phases are translated into intensities recorded by a camera, and these intensities in turn determine which phase shift the SLM is to apply to each part of the beam in the next iteration.

While the first laboratory chimeras had taken a full ten years since Kuramoto and Battogtokh’s 2002 chimera state⁵³, the next few years saw a much more rapid addition of experimental realizations, including mechanical models^{45,54,55}, networks of discrete electrochemical oscillators⁵⁶ and various electronic and optoelectronic systems^{57–61}. The first experimental chimera in a system with global coupling seems to have been observed in 2013 in an photoelectrochemical setup^{15,62}. Also of particular interest is the experimental chimera published by Totz et al. in 2017⁶³: Here, BZ beads of the type previously used by Tinsley et al.³ are coupled by means of the same kind of optical feedback and virtually arranged in a 40×40 grid. For suitable experimental parameters, this yields a spiral-wave chimera with an incoherent core – the qualitatively same kind of pattern as the first 2D numerical chimeras published by Shima and Kuramoto in 2004¹⁹.

Several of the experimental chimeras published from 2012 onward differ significantly from what a chimera state had originally been: Already one of the first experimental realizations had worked with chaotic maps⁴, and not with oscillators, as it said in Abrams and Strogatz’ 2004 definition, but this might just be taken as an extension of the phenomenon to a new do-

main. More remarkable in light of the original definition is the fact that this coupled-maps chimera is (partially) incoherent only in space, while the whole system is temporally periodic⁴. Something similar applies to the so-called “chimera states with quiescent and synchronous domains” found in the coupled electronic oscillators of Gambuzza et al. two years later⁵⁷: Here, the voltage of some constituent circuits is constant in time, while it is oscillating with the same frequency in all the others, but there is no desynchronized region. Arguably also softening the original concept was the mechanical chimera state reported by Wojewoda et al. in 2016, wherein two out of only three pendula are synchronized, while the third one, uncorrelated with the other two, is declared to constitute its own incoherent group⁵⁵. Similarly unprecedented was the optoelectronic chimera published by Larger et al. one year before, in which there are no physically distinct coupled units, but just a single semiconductor laser setup subject to time-delayed feedback⁵⁹. The result is a single time-varying signal, with features (among others) on the length-scale of the applied delay; only when this signal is chopped up into segments and these segments are stacked on top of each other to form the temporal evolution of a “virtual space” does the chimera state appear to the observer. In 2018, Brunner et al. successfully repeated the same procedure for a setup with two simultaneously applied delays of different magnitude, thereby creating a chimera in 2D virtual space⁶¹.

Actually, this broadening of the scientific community’s chimera concept had already begun in the theoretical systems: Already in 2008 did Omel’chenko et al. investigate a 1D array in which the force on each particular oscillator is not only proportional to its deviation from the common mean, but also dependent on its absolute location in the array²⁵. Such a coupling scheme is *not* symmetric in the sense that all oscillators are governed by the same equation of motion and would all feel the same force if they were fully synchronized. The resultant coexistence of synchrony (where the spatial modulation is strong) and incoherence (where the spatial modulation is weak) was nevertheless declared to be a chimera state. Something similar applies to Laing’s later gradual and random removal of individual links from the two-groups model³⁹ and to the randomly time-varying links which Buscarino et al. published in 2015⁴⁶.

Another chimera state, recognized in coupled maps by Omelchenko et al. in 2011, is notably periodic in time and incoherent only in space⁵² (thereby preceding the experimental chimera of Hagerstrom et al.⁴ in this regard). This breaks with the part of Abrams and Strogatz’ original chimera definition² that assigns the attribute of being phase locked to the coherent group only. By allowing for chimera states in coupled maps, Omelchenko et al. also laid the foundations for the later recognition of Kaneko’s much earlier globally-coupled-maps chimera¹⁰.

Similarly expanding the definition of chimera states were the “amplitude chimeras” of Zakharova et al.^{64,65}: Here, all oscillators oscillate in synchrony, with the “incoherent” ones oscillating around different points in the complex plane than the “coherent” ones, as well as having different radii of oscillation. See Fig. 3. Some of the cellular-automaton chimeras published by García-Morales in 2016 also have a well-known periodicity;

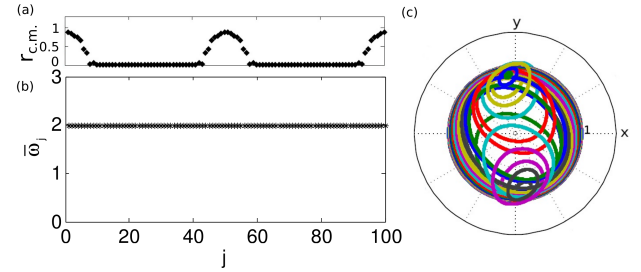


FIG. 3. Amplitude chimera in a ring of nonlocally coupled Stuart-Landau oscillators. (a) Averaged over a period, the position of all oscillators along the two “coherent” segments of the ring is centered on the origin, while that along the two “incoherent” ring segments form an arc-like shape, reminiscent of the distribution of average frequencies in the incoherent part of Kuramoto and Battogtokh’s 2002 chimera. (b) The average phase velocity $\bar{\omega}_j$ with which each oscillator orbits its own respective average position, is the same throughout. (c) Phase portrait of all oscillators in the complex plane. Reprinted from A. Zakharova, M. Kapeller & E. Schöll, *Journal of Physics: Conference Series* 727 (1), 012018, 2016. doi:10.1088/1742-6596/727/1/012018 under the terms of the Creative Commons Attribution 3.0 licence.

and while he found no periodicity for some of the others during the tested simulation time, it is of course fundamentally true that “because of the finiteness of the dynamics, the periodicity of any structure is bounded”⁶⁶, that is, because the states of the system are discrete, it is bound to repeat itself eventually. García-Morales was possibly also the first to recognize how the community’s chimera definition had broadened, mentioning how he would “regard chimera states as an experimental fact of nature rather than a feature of certain systems of differential equations or maps”⁶⁶.

IV. TYPES OF CHIMERA STATES AND CHIMERA-DERIVED CONCEPTS

The gradual expansion of the general concept of a chimera state was accompanied by the naming of an increasing number of derived phenomena, among them the aforementioned amplitude-mediated and amplitude chimera states. Discerned already in 2008 was the “breathing chimera”. Here, the phase coherence of the incoherent oscillators, quantified by the order parameter $r(t) = |\langle e^{i\theta_j(t)} \rangle_{\text{incoh.}}|$, where the sum is taken over the phases θ_j of all oscillators in the unsynchronized group, is either periodic²⁶ or quasiperiodic²⁸. This contrasts with what Abrams et al. call a “stable chimera”²⁶, such as the one discovered by Kuramoto and Battogtokh, where $r(t)$ is constant in time. A few later works have also used the term “breathing chimera” to denote a chimera in which the coherent and incoherent parts move through the system, while the global degree of clustering might remain constant throughout^{42,67}, possibly because this makes the *local* order parameter “breathe”.

Also coined in 2008 was the term “clustered chimera state” with several coherent regions phase-shifted relative to each other²¹. This was joined five years later by the concept of the “multichimera”, likewise containing several distinct co-

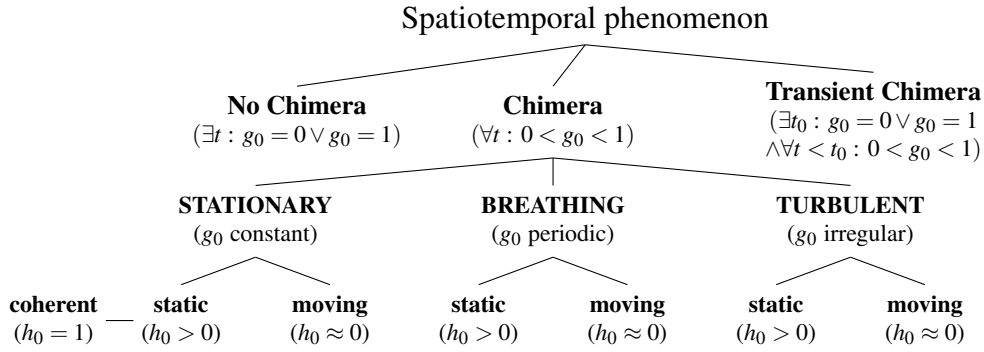


FIG. 4. Data-driven classification scheme of Kemeth et al.⁹. Any spatiotemporal phenomenon is assessed in the form of time-series vectors representing the different parts or units (e.g. oscillators, maps) of the system. For each time t , the measure $g_0 \in [0, 1]$ indicates how similar on average the value of each unit is to either its neighbors (in spatially ordered systems) or all other units. The measure $h_0 \in [0, 1]$ indicates the fraction of the units that are strongly temporally correlated over the evaluated interval. Phenomenologically, the synchronized part of the system remains fixed in “static chimeras”, while it moves in “moving chimeras.”

herent regions, but with no phase difference between them¹⁸. Either of these two phenomena have since come to be called both multicomponent⁶⁸, multicluster^{69,70}, multiheaded^{59,71} or multiple-headed chimera^{59,72}. In the case of equal coupling to a fixed number of nearest neighbors (that is, a step-function nonlocal coupling), the number of incoherent regions (referred to as “heads”⁷¹) increases if the coupling *range* is made shorter^{18,67,73}. If the coupling *strength* is made stronger, the number of heads may either increase¹⁸ or decrease⁷³, depending on the underlying type of oscillator. More complex coupling topologies can also produce various multiplicities of coherent and incoherent regions, but the determining factors are less obvious there^{74,75}. A 2D equivalent of the 1D multichimera is the multicore spiral chimera reported by Xie et al.⁷⁶. In the chimera found by Schmidt et al. in the CGLE with nonlinear global coupling in 2013¹⁵, the concept of a distinct number of chimera heads is less meaningful, as synchronized and incoherent regions, respectively, merge with time⁷⁷.

In 2015, Ashwin and Burylko proposed a rigorous chimera definition applicable to small ensembles. They did this by defining a “weak chimera” to be a state in which the average phase velocities of at least two oscillators converge in the limit of infinite time $T \rightarrow \infty$, while it remains different for at least one other oscillator⁷⁸. This definition was subsequently used to classify states in several later works^{47,66,79–83}. One may assume that it was inspired by the observation that the effective average frequencies of the incoherent oscillators differ from the average frequency of the synchronized cluster in both Kuramoto and Battogtokh’s 2002 chimera¹ as well as several later chimera states^{18,42,67,71,84,85}. However, not all identified classical chimera states (in the sense of some kind of coexistence of synchrony and incoherence) are actually weak chimeras as well. In particular they cannot be when the incoherent region drifts through the system with time and, as a consequence, all oscillators take turns being either coherent or incoherent⁷⁸. Not long after Ashwin and Burylko published their definition, Panaggio et al., for the purposes of their paper on two small oscillator populations, used it to define a chimera (without any qualifications) to mean a weak chimera in which

the frequency-synchronized oscillators have the same phase⁴⁷. Two years later, findings by Kemeth et al. implicitly challenged the potential use of this as a general definition: When scaled up, the two unsynchronized oscillators of a certain minimal (weak) chimera with a perfectly synchronized coherent part are namely not replaced by a greater number of incoherent oscillators⁸². Instead, the state becomes a three-cluster solution with one large and two small clusters. This contrasts with a different kind of minimal chimera that the authors also identify, wherein the dimensionality of the dynamics grows with the system size and which they thus coin an “extensive chimera state”⁸².

Also notable is the “alternating chimera”, in which two equivalent parts of a system take turns being synchronized and incoherent, respectively. This was first produced by external periodic forcing¹⁷ and later found to arise autonomously, in two pre-defined populations³⁹ as well as in a globally coupled oscillatory medium⁷⁷. Other chimera-inspired terms include the “globally clustered chimera”, denoting a system of several pre-defined populations that all split into both synchronized and incoherent oscillators³¹; “chimera death”, the coexistence of spatially coherent and incoherent oscillation death^{64,65}; and the poetically named “Bellerophon states” that occur when a certain chimera state is made unstable by parameter tuning⁸⁶. Additional chimeric phenomena are the “turbulent chimera”³³, the “intermittent chaotic chimera”⁵⁴ and the “blinking chimera”⁸⁷, as well as the “antichimera” and “dual chimera”⁸⁸. In contrast, Laing in a 2012 paper³⁹ reported a kind of imperfect chimera in which one half of the oscillators are more strongly clustered than the other, while none of the two groups is fully synchronized, but without giving this phenomenon an additional name.

When Kemeth et al. came up with a general classification scheme for chimera states in 2016, the wide variety of both existing chimeras and the systems in which they occur made the authors pick a data-driven approach⁹. Their scheme, reproduced in Fig. 4 uses some of the aforementioned labels, such as breathing^{26,28,42} and turbulent chimera³³, in addition to coining new terms, such as “moving chimera” and “static chimera”. In a moving chimera, most individual constituent units of the regarded system change from incoherent to syn-

chronized or vice versa within the regarded time interval, while in a static chimera, they do not⁹. Notably, the authors did not only apply their classification scheme to already declared chimera states, but to other dynamics as well. These include the localized turbulence in the CGLE with time-delayed linear global coupling, as reported by Battogtokh et al. already in 1997⁸⁹. It is classified as a “turbulent moving chimera”, which differs somewhat from the earlier conclusion of Schmidt et al., who, on finding localized turbulence in the CGLE with nonlinear global coupling, were more reluctant to call it a chimera state⁹⁰. Kemeth et al. also evaluate the gradual formation of incoherent patches on a uniformly oscillating background in Falcke and Engel’s 1994 work on a model of CO coverage on a platinum surface^{91–93}. The data-based classification scheme groups these dynamics in the same category of finite-lifetime states as the aforementioned amplitude chimera, a category Kemeth et al. suggest to call “transient chimera”. While not covered by the article on the classification scheme, what Yang et al. call “localized irregular clusters” in a 2000 paper on the Belousov-Zhabotinsky reaction with global feedback⁹⁴ also looks suspiciously like a chimera state.

V. DIFFERENT STANDARDS FOR NATURAL-WORLD CHIMERAS?

Any broader treatment of chimera states should reflect at least briefly on the possibility of chimeras outside of laboratories. Here, the phenomenon most readily invoked by the community is probably unihemispheric sleep, with Rattenborg, Amlaner and Lima’s extensive 2000 neuroscientific review paper⁹⁶ being cited in the introductions of many of the studies mentioned above. Very briefly explained, a unihemispherically sleeping animal sleeps with one half of its brain at a time. Aquatic mammals sleep this way, allowing them to surface to breathe, as do birds and at least some reptiles⁹⁷. When measured, the EEG activity in the sleeping brain hemisphere is high-amplitude, low-frequency, while that in the awake hemisphere is low-amplitude, high-frequency, implying that the individual neurons in the former are firing more strongly synchronized than those in the latter⁹⁶. First to notice the possible connection to early numerical chimeras were probably Abrams et al., whose 2008 chimera-supporting two-groups model is motivated by the question of what might be the simplest system of two oscillator populations to emulate this kind of brain behavior²⁶.

While the roles of the synchronized and incoherent group in the two-groups model are fixed once established, natural unihemispheric sleep tends to move from one half of the brain to the other several times over the course of an interval of sleeping. This was recognized by Ma, Wang and Liu in a 2010 paper, wherein they describe the first of the aforementioned alternating chimera states. However, in order to observe the switching of synchrony from one population to the other, they have to resort to an external periodic forcing, which they declare to “represent the varying environment”¹⁷. In 2015, Haugland et al. reported a fully self-organized alternating chimera without neither pre-defined groups nor any external force, claiming to

“tighten [...] the connection between chimera states and unihemispheric sleep”⁷⁷. As recently as 2019, chimera states have also been realized in two different two-layer networks more closely inspired by brain architecture^{98,99}, thereby completing the continuum of phenomena from actual unihemispheric sleep to the most ideal mathematical chimera.

Related to unihemispheric sleep and also mentioned as a motivation for chimera research is the “first-night effect” in humans, keeping one hemisphere more vigilant when sleeping in a novel environment^{98,100}. Other authors have likened chimera states to the regional highly synchronized brain activity during epileptic seizures or due to Parkinson’s disease¹⁶. Similarly, spiral-wave chimeras, with their incoherently fluctuating core¹⁹, have been compared to ventricular fibrillation, when rotating patterns of excitation occur on the heart, with possibly uncoordinated dynamics at their center⁵. But neither of these comparisons seem to have sparked in-depth deliberation like unihemispheric sleep.

In their 2013 article on chimera states in two pendulum populations, Martens et al. claim that their “model equations translate directly to recent theoretical studies of synchronization in power grids”⁴⁵, implying that chimera states might occur in power grids as well. Panaggio and Abrams also suggest that knowing the basins of stability of chimera states in power grids could be useful in avoiding them and thus maintaining the synchronous oscillation that the grid needs to function⁵. Several other chimera papers briefly refer to this possible connection^{7,58,64,66,75,101,102}, but they seldom elaborate on it. In fact, the article on power-grid modeling that is probably most often cited in chimera introductions, published by Motter et al. in 2013¹⁰³, contains just a single superficial reference to chimeras in its own introduction. Its main aim is to demonstrate a condition for when network synchrony is stable and it is little concerned with what kinds of unsynchronized states may exist. A few results are contributed by a 2014 paper by Pecora et al.¹⁰⁴, wherein they explain the onset of so-called isolated desynchronization by means of network topology and use two real power grids as models (among many others). However, since isolated desynchronization is caused by topology, it should only be relevant to some kinds of chimera states (in the widest sense) and less to those arising by spontaneous symmetry breaking.

Chimera states have also been linked^{3,15,42,69,77,105} to the turbulent-laminar patterns that may be observed in Taylor-Couette flow¹⁰⁶. Another article uses a social-agent model to suggest that “an analogue to a chimera state” could also exist in the behavior of interacting human populations¹⁰⁷.

At the end of his 2018 review article, Omel’chenko refers to all past attempts to identify a non-laboratory chimera as “rather speculative” and “requir[ing] more rigorous justification”⁸. With regards to most of the above examples, this indeed seems to be the case. As far as unihemispheric sleep is concerned, we could alternatively ask exactly what kind of justification is missing. Do actual brain-measurement data not show the coexistence of synchronized and desynchronized oscillation? Are these data not backed up by models modeled on the natural-world phenomenon, already declared to exhibit chimera states? Of course, the mechanism at work in the bird

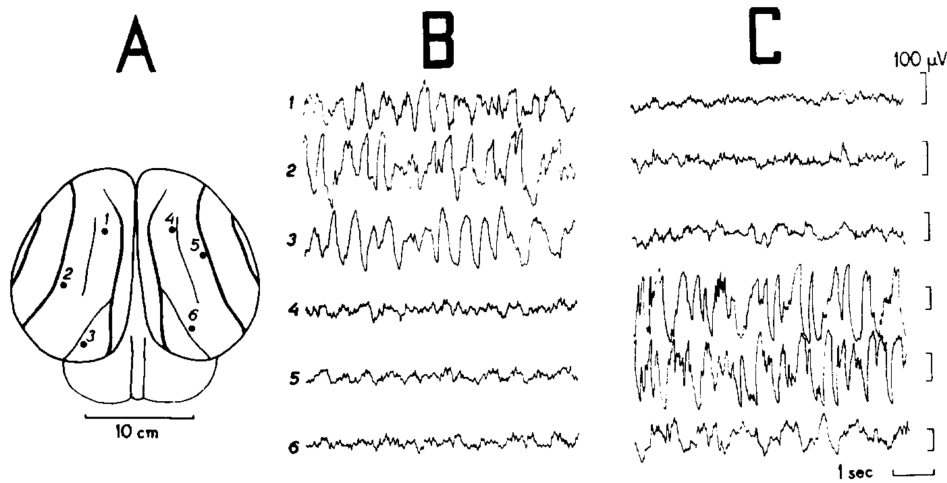


FIG. 5. EEG measurements of the bottlenose dolphin published by Mukhametov et al. in 1977⁹⁵ and later reprinted by Rattenborg et al.⁹⁶. (a) Location of the electrodes across the two brain hemispheres. (b,c) EEG activity measured by each electrode during two short intervals recorded one hour apart. The sleeping hemisphere (electrodes 1-3 in b and 4-6 in c) is characterized by high-amplitude low-frequency EEG activity, the awake hemisphere by low-amplitude high-frequency EEG activity. Reprinted figure with permission from L. M. Mukhametov, A. Y. Supin & I. G. Polyakova, *Brain Research* 134 (3), 581-584, 2016. doi:10.1016/0006-8993(77)90835-6. Copyright 1977 Published by Elsevier B.V.

or dolphin brain is not the same as that in all reported chimera states, but the latter also differ strongly among themselves. Could we thus be holding potential natural-world chimeras to a different standard than theoretical and experimental ones? And could this question possibly be better addressed, if the currently rather fluid and to a large extent implicit chimera definition were made more concrete?

VI. CONCLUDING REMARKS

Above, we have seen that chimera states were originally discovered in globally coupled logistic maps, earlier than often believed. Not until they had also been produced in nonlocally coupled oscillators, however, were they discerned as a special kind of state and given a name. Once they had a name, and once, a few years later, the versatile two-groups framework and the reference to unihemispheric sleep as a potential field of real-world relevance were introduced, a decade of expansive chimera research began. Their number and variety increased, as did the number of systems found to support them. Chimera states were found in a wide range of experimental settings as well. Various derived concepts emerged, though many of them remained mostly limited to their original context. Additional analogies to different real-world phenomena were also drawn up, though many of these have so far remained rather superficial. Notably, the research community has not yet arrived at a common conclusion that any natural-world phenomena actually are chimera states.

This might have something to do with the fact that chimeras are not a most definite physical phenomenon (like neutrons, the Hall effect or protein folding), with definite properties we just have to measure accurately enough to discover. The term is more abstract and might thus also be considered as a loose

collection of more or less related observations. New observations are then being given the same label as existing ones on rather discretionary grounds. This is probably both a blessing and a curse of the field, with unconstrained analogies enabling many a fruitful discovery, but at the same time counteracting the internal ordering of the entirety of results. In particular, as ever new results have broadened the scope of what is called a chimera state, there seems to have been rather limited reflection on how this has changed the object of the field itself. Future research could probably benefit from a more explicit consideration of this insight.

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¹Yoshiki Kuramoto and Dorjsuren Battogtokh. Coexistence of Coherence and Incoherence in Nonlocally Coupled Phase Oscillators. *Nonlinear Phenom. Complex Syst.*, 5(4):380–385, 2002.

²Daniel M. Abrams and Steven H. Strogatz. Chimera states for coupled oscillators. *Phys. Rev. Lett.*, 93(17):174102, oct 2004.

³Mark R. Tinsley, Simbarashe Nkomo, and Kenneth Showalter. Chimera and phase-cluster states in populations of coupled chemical oscillators. *Nat. Phys.*, 8(9):662–665, sep 2012.

⁴Aaron M. Hagerstrom, Thomas E. Murphy, Rajarshi Roy, Philipp Hövel, Iryna Omelchenko, and Eckehard Schöll. Experimental observation of chimeras in coupled-map lattices. *Nat. Phys.*, 8(9):658–661, sep 2012.

⁵Mark J. Panaggio and Daniel M. Abrams. Chimera states: Coexistence of coherence and incoherence in networks of coupled oscillators. *Nonlinearity*, 28(3):R67–R87, mar 2015.

⁶E. Schöll. Synchronization patterns and chimera states in complex networks: Interplay of topology and dynamics. *Eur. Phys. J. Spec. Top.*, 225(6-7):891–919, 2016.

- ⁷Bidesh K. Bera, Soumen Majhi, Dibakar Ghosh, and Matjaž Perc. Chimera states: Effects of different coupling topologies. *Epl*, 118(1):10001, apr 2017.
- ⁸O E Omel'chenko. The mathematics behind chimera states. *Nonlinearity*, 31(5):R121–R164, may 2018.
- ⁹Felix P. Kemeth, Sindre W. Haugland, Lennart Schmidt, Ioannis G. Kevrekidis, and Katharina Krischer. A classification scheme for chimera states. *Chaos*, 26(9):094815, sep 2016.
- ¹⁰Kunihiko Kaneko. Clustering, coding, switching, hierarchical ordering, and control in a network of chaotic elements. *Phys. D Nonlinear Phenom.*, 41(2):137–172, mar 1990.
- ¹¹D. Golomb, D. Hansel, B. Shraiman, and H. Sompolinsky. Clustering in globally coupled phase oscillators. *Phys. Rev. A*, 45(6):3516–3530, mar 1992.
- ¹²Koji Okuda. Variety and generality of clustering in globally coupled oscillators. *Phys. D Nonlinear Phenom.*, 63(3-4):424–436, 1993.
- ¹³Vincent Hakim and Wouter Jan Rappel. Dynamics of the globally coupled complex Ginzburg-Landau equation. *Phys. Rev. A*, 46(12):R7347–R7350, dec 1992.
- ¹⁴N. Nakagawa and Y. Kuramoto. Collective Chaos in a Population of Globally Coupled Oscillators. *Prog. Theor. Phys.*, 89(2):313–323, feb 1993.
- ¹⁵Lennart Schmidt, Konrad Schönleber, Katharina Krischer, and Vladimir García-Morales. Coexistence of synchrony and incoherence in oscillatory media under nonlinear global coupling. *Chaos*, 24(1):013102, mar 2014.
- ¹⁶Gautam C. Sethia and Abhijit Sen. Chimera states: The existence criteria revisited. *Phys. Rev. Lett.*, 112(14):144101, apr 2014.
- ¹⁷Rubao Ma, Jianxiong Wang, and Zonghua Liu. Robust features of chimera states and the implementation of alternating chimera states. *Epl*, 91(4):40006, aug 2010.
- ¹⁸Iryna Omelchenko, Oleh E. Omel'chenko, Philipp Hövel, and Eckehard Schöll. When nonlocal coupling between oscillators becomes stronger: Patched synchrony or multichimera states. *Phys. Rev. Lett.*, 110(22):224101, may 2013.
- ¹⁹Shin Ichiro Shima and Yoshiaki Kuramoto. Rotating spiral waves with phase-randomized core in nonlocally coupled oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 69(3 2):036213, mar 2004.
- ²⁰Daniel M. Abrams and Steven H. Strogatz. Chimera states in a ring of nonlocally coupled oscillators. *Int. J. Bifurc. Chaos*, 16(1):21–37, jan 2006.
- ²¹Gautam C. Sethia, Abhijit Sen, and Fatihcan M. Atay. Clustered chimera states in delay-coupled oscillator systems. *Phys. Rev. Lett.*, 100(14):144102, apr 2008.
- ²²M. Wolfrum, O. E. Omel'chenko, S. Yanchuk, and Y. L. Maistrenko. Spectral properties of chimera states. *Chaos*, 21(1):013112, mar 2011.
- ²³Matthias Wolfrum and Oleh E. Omel'chenko. Chimera states are chaotic transients. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 84(1):015201, jul 2011.
- ²⁴Yoji Kawamura. Chimera Ising walls in forced nonlocally coupled oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 75(5):056204, may 2007.
- ²⁵Oleh E. Omel'chenko, Yuri L. Maistrenko, and Peter A. Tass. Chimera states: The natural link between coherence and incoherence. *Phys. Rev. Lett.*, 100(4):044105, jan 2008.
- ²⁶Daniel M. Abrams, Rennie Mirollo, Steven H. Strogatz, and Daniel A. Wiley. Solvable model for chimera states of coupled oscillators. *Phys. Rev. Lett.*, 101(8):84103, aug 2008.
- ²⁷Edward Ott and Thomas M. Antonsen. Low dimensional behavior of large systems of globally coupled oscillators. *Chaos*, 18(3):037113, sep 2008.
- ²⁸Arkady Pikovsky and Michael Rosenblum. Partially integrable dynamics of hierarchical populations of coupled oscillators. *Phys. Rev. Lett.*, 101(26):1–4, 2008.
- ²⁹Carlo R. Laing. Chimera states in heterogeneous networks. *Chaos*, 19(1), 2009.
- ³⁰Carlo R. Laing. The dynamics of chimera states in heterogeneous Kuramoto networks. *Phys. D Nonlinear Phenom.*, 238(16):1569–1588, 2009.
- ³¹Jane H. Sheeba, V. K. Chandrasekar, and M. Lakshmanan. Globally clustered chimera states in delay-coupled populations. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 79(5):055203, may 2009.
- ³²Jane H. Sheeba, V. K. Chandrasekar, and M. Lakshmanan. Chimera and globally clustered chimera: Impact of time delay. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 81(4):046203, apr 2010.
- ³³Grigory Borduygov, Arkady Pikovsky, and Michael Rosenblum. Self-emerging and turbulent chimeras in oscillator chains. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 82(3):035205, sep 2010.
- ³⁴Erik A. Martens, Carlo R. Laing, and Steven H. Strogatz. Solvable model of spiral wave chimeras. *Phys. Rev. Lett.*, 104(4):044101, jan 2010.
- ³⁵Erik A. Martens. Bistable chimera attractors on a triangular network of oscillator populations. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 82(1):1–10, 2010.
- ³⁶Erik A. Martens. Chimeras in a network of three oscillator populations with varying network topology. *Chaos*, 20(4):043122, dec 2010.
- ³⁷Murray Shanahan. Metastable chimera states in community-structured oscillator networks. *Chaos*, 20(1):013108, mar 2010.
- ³⁸Oleh E. Omel'chenko, Matthias Wolfrum, and Yuri L. Maistrenko. Chimera states as chaotic spatiotemporal patterns. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 81(6):065201, jun 2010.
- ³⁹Carlo R. Laing, Karthikeyan Rajendran, and Ioannis G. Kevrekidis. Chimeras in random non-complete networks of phase oscillators. *Chaos*, 22(1):013132, mar 2012.
- ⁴⁰Mark Wildie and Murray Shanahan. Metastability and chimera states in modular delay and pulse-coupled oscillator networks. *Chaos*, 22(4):043131, dec 2012.
- ⁴¹Carlo R. Laing. Chimeras in networks of planar oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 81(6):6–9, 2010.
- ⁴²Gautam C. Sethia, Abhijit Sen, and George L. Johnston. Amplitude-mediated chimera states. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 88(4):042917, oct 2013.
- ⁴³Carlo R. Laing. Chimeras in networks with purely local coupling. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 92(5):4–9, 2015.
- ⁴⁴Carlo R. Laing. Disorder-induced dynamics in a pair of coupled heterogeneous phase oscillator networks. *Chaos*, 22(4):043104, dec 2012.
- ⁴⁵Erik Andreas Martens, Shashi Thutupalli, Antoine Fourrière, and Oskar Hallatschek. Chimera states in mechanical oscillator networks. *Proc. Natl. Acad. Sci. U. S. A.*, 110(26):10563–10567, jun 2013.
- ⁴⁶Arturo Buscarino, Mattia Frasca, Lucia Valentina Gambuzza, and Philipp Hövel. Chimera states in time-varying complex networks. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 91(2):1–7, 2015.
- ⁴⁷Mark J. Panaggio, Daniel M. Abrams, Peter Ashwin, and Carlo R. Laing. Chimera states in networks of phase oscillators: The case of two small populations. *Phys. Rev. E*, 93(1):012218, jan 2016.
- ⁴⁸Erik A. Martens, Christian Bick, and Mark J. Panaggio. Chimera states in two populations with heterogeneous phase-lag. *Chaos*, 26(9), 2016.
- ⁴⁹Erik A. Martens, Mark J. Panaggio, and Daniel M. Abrams. Basins of attraction for chimera states. *New J. Phys.*, 18(2):022002, feb 2016.
- ⁵⁰While the 2008 paper by Abrams et al. does not cite any prior source for the two-groups model, Laing points out in Ref.²⁹ that more general versions of it were actually implemented in earlier articles by Montbrió et al.¹⁰⁸ and Barreto et al.¹⁰⁹. However, these papers did not research chimera states.
- ⁵¹While their 2012 paper simply refers to “the projected image from a spatial light modulator”³, a later article by the same authors on a very similar experiment mentions that “[t]he experimental set-up consists of a modified video projector (SLM) with a 440–460 nm band pass filter”¹¹⁰. Thus it is not unlikely that the SLM used in the 2012 experiment was this custom-built apparatus as well.
- ⁵²Iryna Omelchenko, Yuri Maistrenko, Philipp Hövel, and Eckehard Schöll. Loss of coherence in dynamical networks: Spatial chaos and chimera states. *Phys. Rev. Lett.*, 106(23):234102, jun 2011.
- ⁵³Ashley G. Smart. Exotic chimera dynamics glimpsed in experiments. *Phys. Today*, 65(10):17–19, oct 2012.
- ⁵⁴Simona Olmi, Erik A. Martens, Shashi Thutupalli, and Alessandro Torcini. Intermittent chaotic chimeras for coupled rotators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 92(3):030901, sep 2015.
- ⁵⁵Jerzy Wojewoda, Krzysztof Czołczynski, Yuri Maistrenko, and Tomasz Kapitaniak. The smallest chimera state for coupled pendula. *Sci. Rep.*, 6(1):34329, dec 2016.
- ⁵⁶Mahesh Wickramasinghe and István Z. Kiss. Spatially organized dynamical states in chemical oscillator networks: Synchronization, dynamical differentiation, and chimera patterns. *PLoS One*, 8(11):e80586, nov 2013.
- ⁵⁷Lucia Valentina Gambuzza, Arturo Buscarino, Sergio Chessa, Luigi Fortuna, Riccardo Meucci, and Mattia Frasca. Experimental investigation

- of chimera states with quiescent and synchronous domains in coupled electronic oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 90(3):032905, sep 2014.
- ⁵⁸Fabian Böhm, Anna Zakharova, Eckehard Schöll, and Kathy Lüdge. Amplitude-phase coupling drives chimera states in globally coupled laser networks. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 91(4):040901, apr 2015.
- ⁵⁹Laurent Larger, Bogdan Penkovsky, and Yuri Maistrenko. Laser chimeras as a paradigm for multistable patterns in complex systems. *Nat. Commun.*, 6:7752, 2015.
- ⁶⁰Joseph D. Hart, Kanika Bansal, Thomas E. Murphy, and Rajarshi Roy. Experimental observation of chimera and cluster states in a minimal globally coupled network. *Chaos*, 26(9):094801, sep 2016.
- ⁶¹D. Brunner, B. Penkovsky, R. Levchenko, E. Schöll, L. Larger, and Y. Maistrenko. Two-dimensional spatiotemporal complexity in dual-delayed nonlinear feedback systems: Chimeras and dissipative solitons. *Chaos*, 28(10), 2018.
- ⁶²Konrad Schönleber, Carla Zensen, Andreas Heinrich, and Katharina Krischer. Pattern formation during the oscillatory photoelectrodisolution of n-type silicon: Turbulence, clusters and chimeras. *New J. Phys.*, 16(6):063024, jun 2014.
- ⁶³Jan Frederik Totz, Julian Rode, Mark R. Tinsley, Kenneth Showalter, and Harald Engel. Spiral wave chimera states in large populations of coupled chemical oscillators. *Nat. Phys.*, 14(3):282–285, 2018.
- ⁶⁴Anna Zakharova, Marie Kapeller, and Eckehard Schöll. Chimera death: Symmetry breaking in dynamical networks. *Phys. Rev. Lett.*, 112(15):154101, 2014.
- ⁶⁵Anna Zakharova, Marie Kapeller, and Eckehard Schöll. Amplitude chimeras and chimera death in dynamical networks. *J. Phys. Conf. Ser.*, 727(1):012018, jun 2016.
- ⁶⁶Vladimir García-Morales. Cellular automaton for chimera states. *Epl*, 114(1):18002, apr 2016.
- ⁶⁷Andrea Vüllings, Johanne Hizanidis, Iryna Omelchenko, and Philipp Hövel. Clustered chimera states in systems of type-I excitability. *New J. Phys.*, 16(12):123039, dec 2014.
- ⁶⁸Sangeeta Rani Ujjwal and Ramakrishna Ramaswamy. Chimeras with multiple coherent regions. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 88(3):1–6, 2013.
- ⁶⁹Jianbo Xie, Edgar Knobloch, and Hsien Ching Kao. Multiclustor and traveling chimera states in nonlocal phase-coupled oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 90(2):1–17, 2014.
- ⁷⁰Nan Yao, Zi Gang Huang, Celso Grebogi, and Ying Cheng Lai. Emergence of multiclustor chimera states. *Sci. Rep.*, 5:12988, 2015.
- ⁷¹Yuri L. Maistrenko, Anna Vasylenko, Oleksandr Sudakov, Roman Levchenko, and Volodymyr L. Maistrenko. Cascades of multiheaded chimera states for coupled phase oscillators. *Int. J. Bifurc. Chaos*, 24(8):1440014, aug 2014.
- ⁷²Alexander Schmidt, Theodoros Kasimatis, Johanne Hizanidis, Astero Provata, and Philipp Hövel. Chimera patterns in two-dimensional networks of coupled neurons. *Phys. Rev. E*, 95(3):032224, mar 2017.
- ⁷³Iryna Omelchenko, Anna Zakharova, Philipp Hövel, Julien Siebert, and Eckehard Schöll. Nonlinearity of local dynamics promotes multi-chimeras. *Chaos*, 25(8):083104, aug 2015.
- ⁷⁴Johanne Hizanidis, Evangelia Panagakou, Iryna Omelchenko, Eckehard Schöll, Philipp Hövel, and Astero Provata. Chimera states in population dynamics: Networks with fragmented and hierarchical connectivities. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 92(1):012915, jul 2015.
- ⁷⁵Stefan Ulonska, Iryna Omelchenko, Anna Zakharova, and Eckehard Schöll. Chimera states in networks of Van der Pol oscillators with hierarchical connectivities. *Chaos*, 26(9):094825, sep 2016.
- ⁷⁶Jianbo Xie, Edgar Knobloch, and Hsien Ching Kao. Twisted chimera states and multicore spiral chimera states on a two-dimensional torus. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 92(4), 2015.
- ⁷⁷Sindre W. Haugland, Lennart Schmidt, and Katharina Krischer. Self-organized alternating chimera states in oscillatory media. *Sci. Rep.*, 5:9883, apr 2015.
- ⁷⁸Peter Ashwin and Oleksandr Burylko. Weak chimeras in minimal networks of coupled phase oscillators. *Chaos*, 25(1):013106, jan 2015.
- ⁷⁹Yusuke Suda and Koji Okuda. Persistent chimera states in nonlocally coupled phase oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 92(6):060901, dec 2015.
- ⁸⁰Iryna Omelchenko, Oleh E. Omel’Chenko, Anna Zakharova, Matthias Wolfrum, and Eckehard Schöll. Tweezers for Chimeras in Small Networks. *Phys. Rev. Lett.*, 116(11):114101, mar 2016.
- ⁸¹Yuri Maistrenko, Serhiy Brezetsky, Patrycja Jaros, Roman Levchenko, and Tomasz Kapitaniak. Smallest chimera states. *Phys. Rev. E*, 95(1):010203, jan 2017.
- ⁸²Felix P. Kemeth, Sindre W. Haugland, and Katharina Krischer. Symmetries of Chimera States. *Phys. Rev. Lett.*, 120(21):214101, may 2018.
- ⁸³Jorge Luis Ocampo-Espindola, Christian Bick, and István Z. Kiss. Weak Chimeras in Modular Electrochemical Oscillator Networks. *Front. Appl. Math. Stat.*, 5(July):1–12, 2019.
- ⁸⁴Simbarashe Nkomo, Mark R. Tinsley, and Kenneth Showalter. Chimera states in populations of nonlocally coupled chemical oscillators. *Phys. Rev. Lett.*, 110(24):244102, jun 2013.
- ⁸⁵David P. Rosin, Damien Rontani, Nicholas D. Haynes, Eckehard Schöll, and Daniel J. Gauthier. Transient scaling and resurgence of chimera states in networks of Boolean phase oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 90(3):030902, 2014.
- ⁸⁶Tian Qiu, Stefano Boccaletti, Ivan Bonamassa, Yong Zou, Jie Zhou, Zonghua Liu, and Shuguang Guan. Synchronization and Bellerophon states in conformist and contrarian oscillators. *Sci. Rep.*, 6(1):36713, dec 2016.
- ⁸⁷Richard Janis Goldschmidt, Arkady Pikovsky, and Antonio Politi. Blinking chimeras in globally coupled rotators. *Chaos*, 29(7):071101, jul 2019.
- ⁸⁸Gabriela Petruccaro, Koichiro Uriu, and Luis G. Morelli. Mobility-induced persistent chimera states. *Phys. Rev. E*, 96(6):1–9, 2017.
- ⁸⁹D. Battogtokh, A. Preusser, and A. Mikhailov. Controlling turbulence in the complex Ginzburg-Landau equation II. Two-dimensional systems. *Phys. D Nonlinear Phenom.*, 106(3-4):327–362, aug 1997.
- ⁹⁰Lennart Schmidt and Katharina Krischer. Chimeras in globally coupled oscillatory systems: From ensembles of oscillators to spatially continuous media. *Chaos*, 25(6):64401, 2015.
- ⁹¹M. Falcke and H. Engel. Influence of global coupling through the gas phase on the dynamics of CO oxidation on Pt(110). *Phys. Rev. E*, 50(2):1353–1359, aug 1994.
- ⁹²M. Falcke and H. Engel. Pattern formation during the CO oxidation on Pt(110) surfaces under global coupling. *J. Chem. Phys.*, 101(7):6255–6263, 1994.
- ⁹³M Falcke. *Strukturbildung in Reaktions- Diffusionssystemen und globale Kopplung*. Wiss.-und-Technik-Verlag Gross, 1995.
- ⁹⁴Lingfa Yang, Milos Dolnik, Anatol M. Zhabotinsky, and Irving R. Epstein. Oscillatory clusters in a model of the photosensitive Belousov-Zhabotinsky reaction system with global feedback. *Phys. Rev. E - Stat. Physics, Plasmas, Fluids, Relat. Interdiscip. Top.*, 62(5):6414–6420, 2000.
- ⁹⁵L.M. Mukhametov, A.Y. Supin, and I.G. Polyakova. Interhemispheric asymmetry of the electroencephalographic sleep patterns in dolphins. *Brain Res.*, 134(3):581–584, oct 1977.
- ⁹⁶N. C. Rattenborg, C. J. Amlaner, and S. L. Lima. Behavioral, neurophysiological and evolutionary perspectives on unihemispheric sleep. *Neurosci. Biobehav. Rev.*, 24(8):817–842, dec 2000.
- ⁹⁷Christian G. Mathews, John A. Lesku, Steven L. Lima, and Charles J. Amlaner. Asynchronous eye closure as an anti-predator behavior in the western fence lizard (*Sceloporus occidentalis*). *Ethology*, 112(3):286–292, mar 2006.
- ⁹⁸Lukas Ramlow, Jakub Sawicki, Anna Zakharova, Jaroslav Hlinka, Jens Christian Claussen, and Eckehard Schöll. Partial synchronization in empirical brain networks as a model for unihemispheric sleep. *Epl*, 126(5):50007, 2019.
- ⁹⁹Ling Kang, Changhai Tian, Siyu Huo, and Zonghua Liu. A two-layered brain network model and its chimera state. *Sci. Rep.*, 9(1):1–12, 2019.
- ¹⁰⁰Masako Tamaki, Ji Won Bang, Takeo Watanabe, and Yuka Sasaki. Night Watch in One Brain Hemisphere during Sleep Associated with the First-Night Effect in Humans. *Curr. Biol.*, 26(9):1190–1194, 2016.
- ¹⁰¹Tanmoy Banerjee, Partha Sharathi Dutta, Anna Zakharova, and Eckehard Schöll. Chimera patterns induced by distance-dependent power-law coupling in ecological networks. *Phys. Rev. E*, 94(3):1–8, 2016.
- ¹⁰²Sarbendu Rakshit, Bidesh K. Bera, Matjaž Perc, and Dibakar Ghosh. Basin stability for chimera states. *Sci. Rep.*, 7(1):1–12, 2017.

- ¹⁰³Adilson E. Motter, Seth A. Myers, Marian Anghel, and Takashi Nishikawa. Spontaneous synchrony in power-grid networks. *Nat. Phys.*, 9(3):191–197, 2013.
- ¹⁰⁴Louis M. Pecora, Francesco Sorrentino, Aaron M. Hagerstrom, Thomas E. Murphy, and Rajarshi Roy. Cluster synchronization and isolated desynchronization in complex networks with symmetries. *Nat. Commun.*, 5(May):4079, 2014.
- ¹⁰⁵Christian Bick and Erik A. Martens. Controlling chimeras. *New J. Phys.*, 17(3):033030, mar 2015.
- ¹⁰⁶Dwight Barkley and Laurette S. Tuckerman. Computational study of turbulent laminar patterns in couette flow. *Phys. Rev. Lett.*, 94(1):14502, jan 2005.
- ¹⁰⁷J. C. González-Avella, M. G. Cosenza, and M. San Miguel. Localized coherence in two interacting populations of social agents. *Phys. A Stat. Mech. its Appl.*, 399:24–30, 2014.
- ¹⁰⁸Ernest Montbrió, Jürgen Kurths, and Bernd Blasius. Synchronization of two interacting populations of oscillators. *Phys. Rev. E - Stat. Physics, Plasmas, Fluids, Relat. Interdiscip. Top.*, 70(5):4, nov 2004.
- ¹⁰⁹Ernest Barreto, Brian Hunt, Edward Ott, and Paul So. Synchronization in networks of networks: The onset of coherent collective behavior in systems of interacting populations of heterogeneous oscillators. *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.*, 77(3):036107, mar 2008.
- ¹¹⁰Simbarashe Nkomo, Mark R. Tinsley, and Kenneth Showalter. Chimera and chimera-like states in populations of nonlocally coupled homogeneous and heterogeneous chemical oscillators. *Chaos*, 26(9), 2016.