

# Adaptive stepsize algorithms for Langevin dynamics

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## 1. Background

**Goal:** Sampling via the invariant distribution  $\rho_\infty(\mathbf{x}) \propto \exp(-\beta V(\mathbf{x}))$  of stochastic differential equation.

Using Langevin dynamics in the large friction limit, we get the overdamped stochastic differential equation:

$$dX = -\nabla V(X)dt + \sqrt{2\beta^{-1}}dW(t)$$

where  $W(t)$  is the Brownian motion and  $V(x)$  is a potential. With a numerical scheme, samples from the invariant distribution  $\rho_\infty(x) = C \exp(-\beta V(x))$  can be drawn integrating one path for a long time. Euler-Maruyama applied to the overdamped SDE yields:

$$X_{i+1} = X_i - \nabla V(X_i)h + \sqrt{2\beta^{-1}h}\Delta W_i,$$

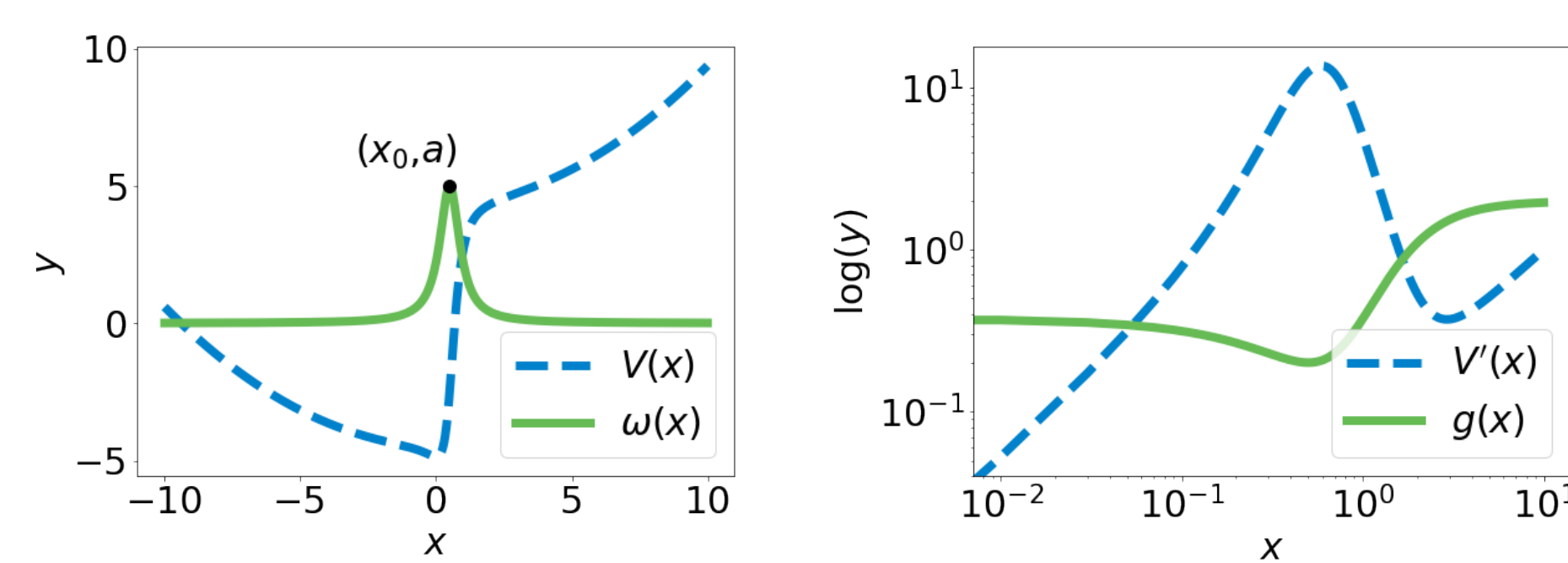
the time is discretised with  $h = 1/n_s$ ,  $n_s$  is the number of samples,  $X_i \approx X(t = ih)$  and  $\Delta W_i \sim \mathcal{N}(0, 1)$ .

## 2. Motivating example

Sampling from a modified harmonic potential with a sudden change of frequency:

$$V'(x) = (\omega(x)^2 + c)x$$

$$\omega(x) = \frac{b}{a + (x - x_0)^2}$$



**Figure 1:** On the left, the modified harmonic potential  $V(x)$  and the function  $\omega(x)$  with  $a = 5$ ,  $x_0 = 0.5$  and  $b = 1$ . A smaller  $b$  implies a steeper  $\omega(x)$  around  $x_0$ . On the right, the gradient  $V'(x)$  and an example of adaptive function  $g(x)$ .

To reduce the computational cost, the size of the steps could be smaller around singularities and bigger elsewhere driven by a function  $g(x)$ , for instance:

$$g(x) \propto \frac{1}{\|V'(x)\|},$$

with  $g(x) \neq 0$  and  $g(x) \in [\frac{\Delta t_{\min}}{\Delta t}, \frac{\Delta t_{\max}}{\Delta t}]$ .

Naive time re-scaling in the continuous dynamics:

$$dt \rightarrow g(x)d\tau$$

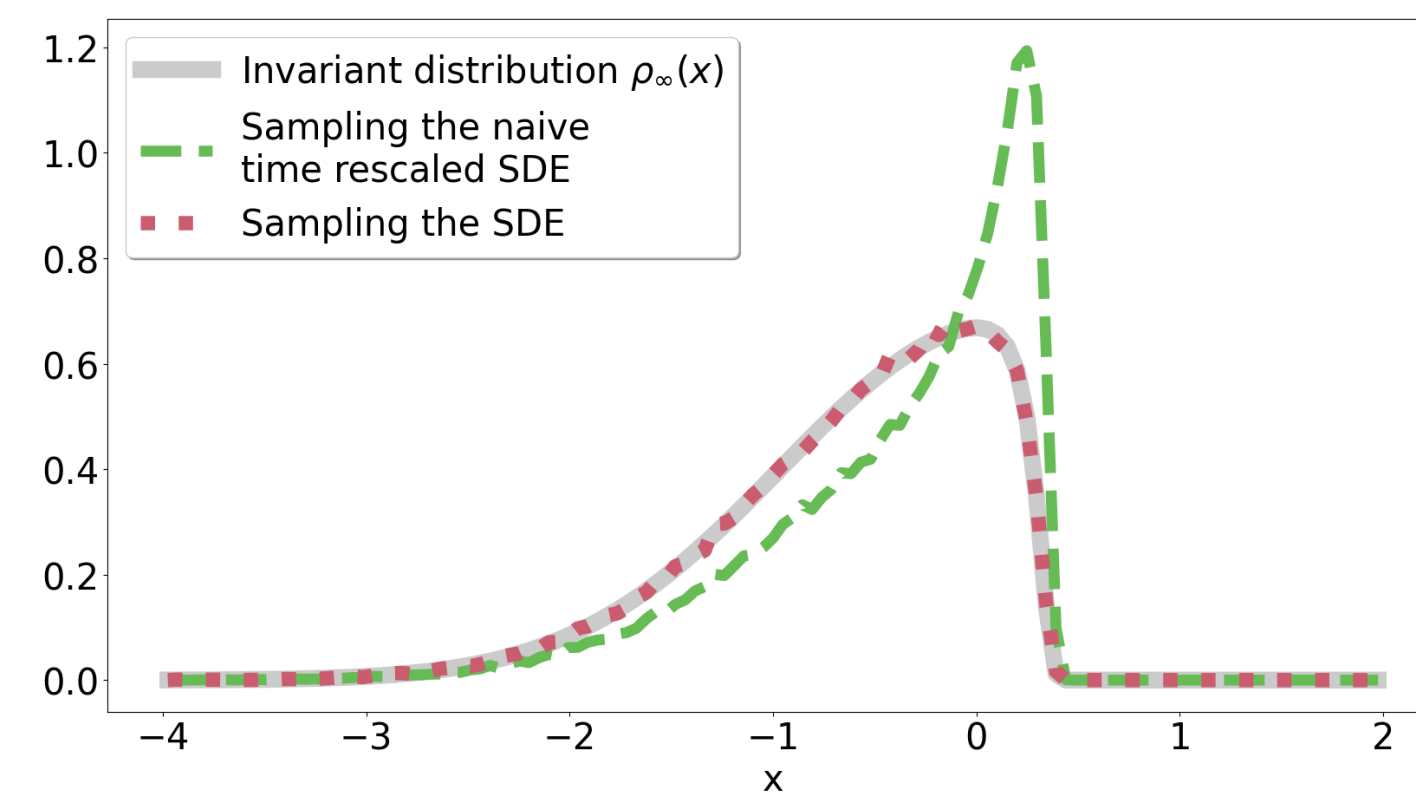
$$dW(t) \rightarrow \sqrt{g(x)}dW(\tau)$$

gives the naive time rescaled SDE:

$$d\tilde{X} = -\nabla V(\tilde{X})g(\tilde{X})d\tau + \sqrt{2\beta^{-1}g(\tilde{X})}dW(\tau).$$

Euler-Maruyama scheme applied to this SDE:

$$\tilde{X}_{i+1} = \tilde{X}_i - \nabla V(\tilde{X}_i)g(\tilde{X}_i)h + \sqrt{2\beta^{-1}g(\tilde{X}_i)h}\Delta W_i.$$



**Figure 2:** Histograms of the invariant distribution and samples from Euler-Maruyama scheme applied to the original SDE and to the naive time rescaled SDE after 50 000 steps,  $h = 0.001$ ,  $n_s = 10^5$ ,  $\frac{\Delta t_{\max}}{\Delta t} = 2$ ,  $\frac{\Delta t_{\min}}{\Delta t} = 0.001$ ,  $\tau = 0.1$  ( $c = 0.1$ ,  $b = 0.1$ ,  $a = 10$ ,  $x_0 = 0.5$ ).

In fact, using the backward Kolmogorov equation, one can show that  $\rho_\infty(x)$  is **not the invariant distribution** for the naive time rescaled SDE.

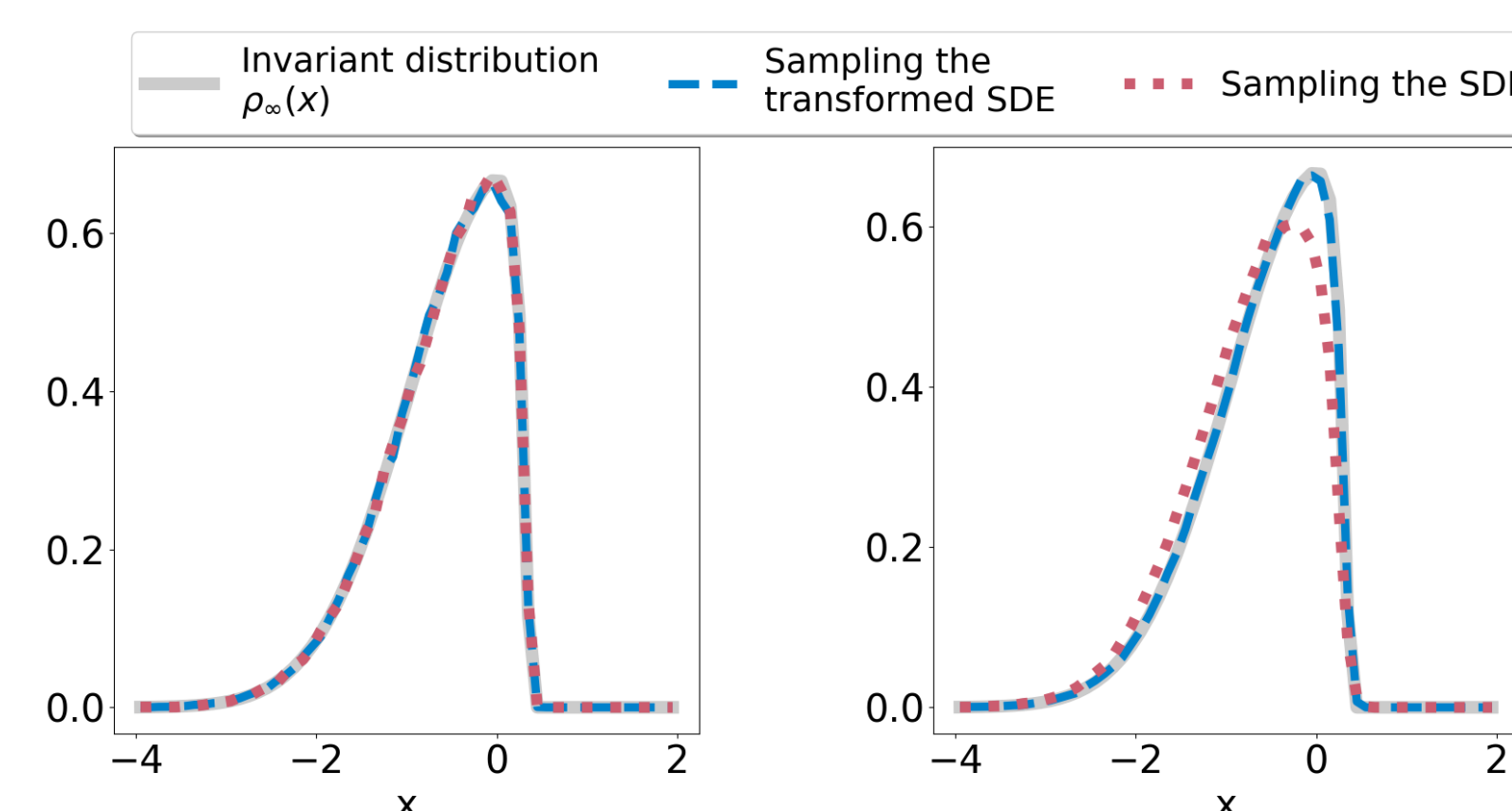
## 3. An IP-transformed SDE

The bias in the invariant distribution arises from the factor  $\sqrt{g(x)}$  in the diffusion term which gives rise to an increasing variance of displacement. The **invariant preserving (IP)-transformed SDE** leading to the appropriate distribution is:

$$dY = -\nabla V(Y)g(Y)ds + \beta^{-1}\nabla g(Y)ds + \sqrt{2\beta^{-1}g(Y)}dW_s.$$

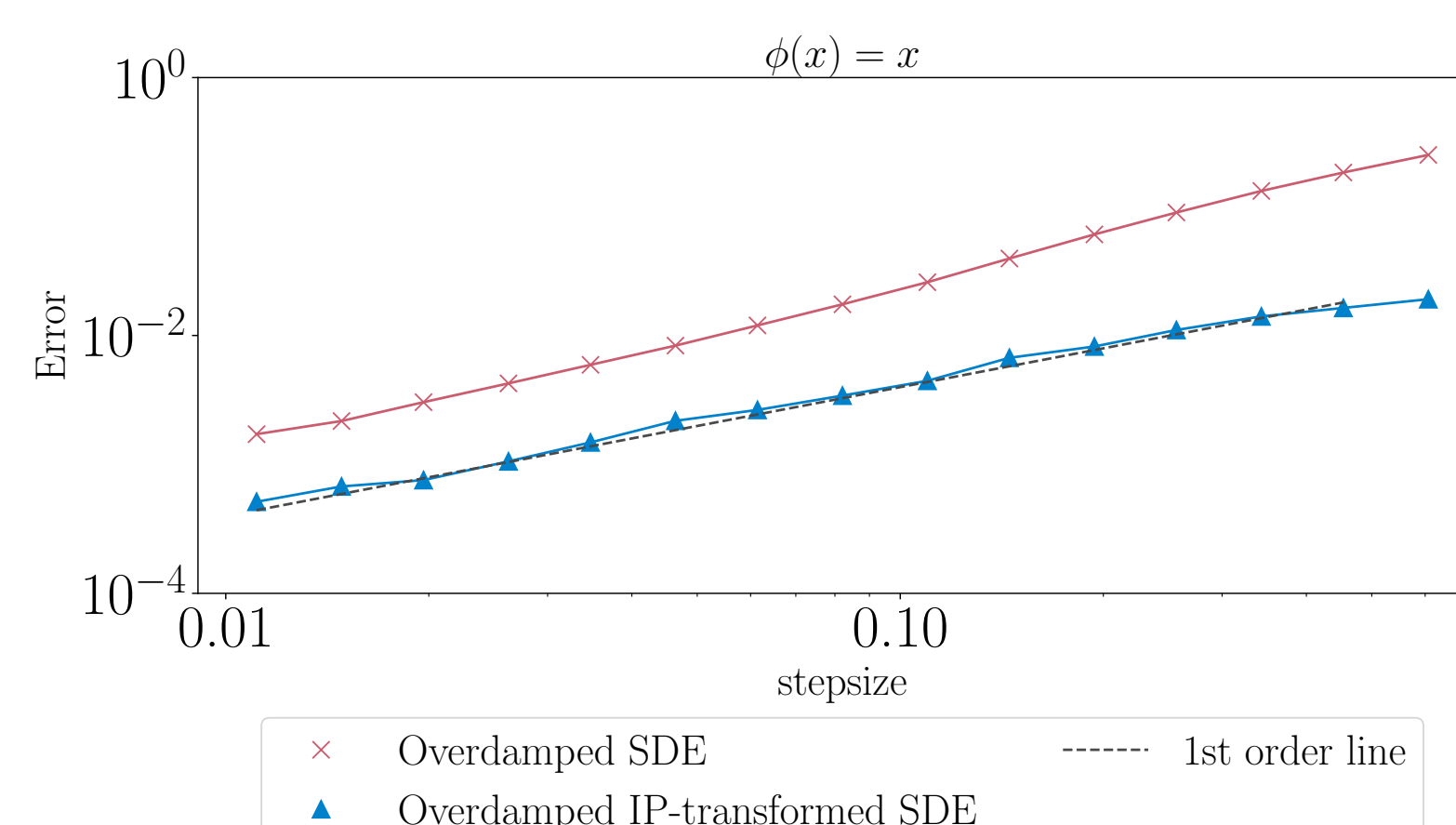
We can verify using the backward Kolmogorov equation that  $\rho_\infty(x)$  is **the invariant distribution** of this direct time-rescaled transformed SDE. Euler-Maruyama on this SDE:

$$Y_{i+1} = Y_i - \nabla V(Y_i)g(Y_i)h + \beta^{-1}\nabla g(Y_i)h + \sqrt{2\beta^{-1}g(Y_i)h}\Delta W_i.$$



**Figure 3:** On the left, histograms with a small steps  $h = 0.001$ , both SDEs converge. On the right histograms with a larger stepsize  $h = 0.05$ , only the transformed SDE converges. Simulations use similar parameters to figure 2.

For a large stepsize, the **IP-transformed SDE converges** toward the invariant distribution while the **overdamped SDE fails** to converge. In Figure 5, the rescaling factor is **1.18** but the error is significantly smaller for the IP-transformed SDE.



**Figure 4:** The order of convergence of the Euler-Maruyama method is recovered for larger stepsize for both dynamics, but the error is significantly smaller with the IP-transformed SDE.

## 4. IP-transformed underdamped

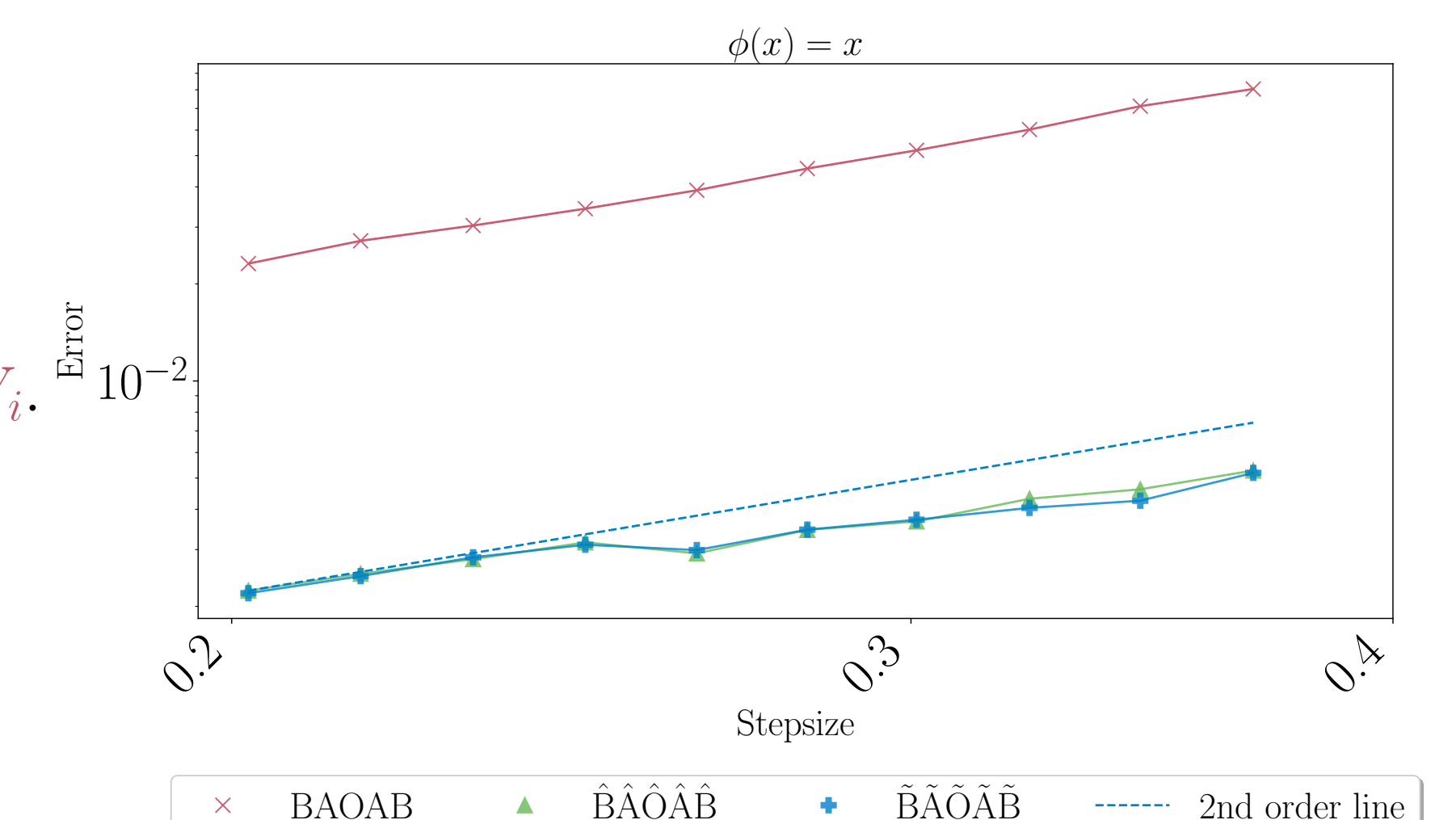
A separable Hamiltonian dynamics,  $H(q, p) = E_{\text{kin}}(p) + V(q)$ , with position  $q \in \mathbb{R}^d$  and momentum  $p \in \mathbb{R}^d$ , with the friction  $\gamma$  has the corresponding **IP-transformed dynamic**:

$$\begin{cases} dq = pg(q)ds, \\ dp = -\nabla V(q)g(q)ds + \beta^{-1}\nabla g(q)ds - \gamma pg(q)ds + \sqrt{2\gamma\beta^{-1}g(q)}dW(s). \end{cases}$$

with invariant distribution is  $\rho(q, p) \propto \exp(-\beta H(q, p))$ . We use numerical timestepping integration methods based on splitting schemes, which break the equations into separate parts to be solved independently

$$\begin{bmatrix} dq \\ dp \end{bmatrix} = \underbrace{\begin{bmatrix} g(q)p \\ 0 \end{bmatrix}}_A dt + \underbrace{\begin{bmatrix} 0 \\ -\nabla V(q)g(q) \end{bmatrix}}_B dt + \underbrace{\begin{bmatrix} 0 \\ \beta^{-1}\nabla g(q)dt - \gamma pg(q)dt + \sqrt{2\gamma\beta^{-1}g(q)}dW(t) \end{bmatrix}}_O.$$

The IP-transformed system of SDEs requires the computation of extra terms such as  $\Delta g(x)\beta^{-1}$  and renders the step A implicit. We design novel schemes based on modified building blocks such as  $\hat{B}\hat{A}\hat{O}\hat{A}\hat{B}$  and  $\tilde{B}\tilde{A}\tilde{O}\tilde{A}\tilde{B}$ .



**Figure 5:** The schemes  $\hat{B}\hat{A}\hat{O}\hat{A}\hat{B}$  and  $\tilde{B}\tilde{A}\tilde{O}\tilde{A}\tilde{B}$  have much higher efficiency and display a quadratic decay in the error, similarly to the original schemes BAOAB.

## 5. Conclusions

- Invariant-preserving transformed dynamics are based on an embedded adaptive time stepping based on a monitor function.
- We provide several novel splitting schemes to integrate the invariant-preserving transformed underdamped system.
- Higher efficiency is shown for the numerical integration of the IP-transformed dynamics.

- Preprint with detailed analytical considerations

