

Ab Initio Molecular Dynamics Simulations of Phosphate Hydrolysis Using Neural Network Potentials

Albert MAKHMUDOV

Supervisor: Prof. J. Harvey
KU Leuven

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List of abbreviations

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Chapter 1

Introduction

- 1.1 Role of phosphates in biological systems**
- 1.2 Enzymes involved in phosphate hydrolysis**
- 1.3 Reaction mechanism**
- 1.4 Research goals**

Chapter 2

Theory

2.1 A brief introduction to statistical mechanics

2.1.1 Classical forcefields and molecular dynamics

2.1.2 The canonical ensemble and free energy calculations

2.1.3 Enhanced sampling techniques

Metadynamics and its well-tempered flavour

Kinetics from metadynamics

2.2 Density functional theory

2.2.1 The Kohn-Sham approach

2.2.2 Generalised gradient approximation and PBE functional

2.2.3 *Ab initio* molecular dynamics and GPW method

2.3 Extended tight binding

2.4 Neural network potentials

2.4.1 Deep neural networks

Multilayer perceptron

Graph neural networks

Message passing neural networks

2.4.2 Invariance and equivariance

2.4.3 Robust Barrinello neural network potentials

Chapter 3

Computational Details

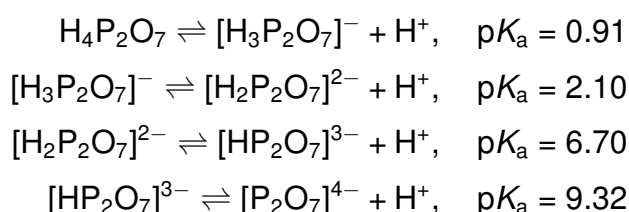
This chapter provides the details of the computational methods used in this work. The first section describes the generation of the training dataset, including the preparation of the system, initial equilibration using the molecular mechanics, exploration of the configuration space at the xTB level, further data labeling, and iterative training of the neural network potential. The second section discusses the production runs at different temperatures using the fitted neural network potential. The third section describes the workflow of validating the transition states obtained from the simulations following the partial Hessian formalism. Finally, the fourth section presents the data analysis and visualisation techniques employed to interpret the results.

3.1 Training dataset generation

3.1.1 System preparation

The systems were prepared using the CHARMM-GUI webserver's functionality [1]. In particular, the Multicomponent Assembler interface [2] was utilised.

As a first step, the singly protonated and deprotonated forms of the methyl diphosphate were parametrised in CGenFF [3], i.e. CHARMM General Forcefield. These states of the methyl diphosphate were chosen based on the fact that pyrophosphoric (diphosphoric) acid has the following dissociation constants [4]:



Thus, at the physiological pH of 7.4 this acid exists as an equilibrium between the doubly and singly protonated forms. As an assumption, the methyl group can be considered as a proton, therefore we considered the methyl diphosphate molecule to exist as a mixture of the singly (MeHDP) and deprotonated (MeDP) forms at the physiological pH.

After successfully parametrising the molecules, the system was solvated in a cubic box of TIP3 water molecules together with the sodium counterions Na^+ to neutralise the charge. The final system composition can be seen in Table 3.1.

3.1.2 Initial equilibration using the classical forcefields

The equilibration of the system was performed following the standard protocol generated by the CHARMM-GUI webserver [1]. The system was first energy minimised using the steepest descent algorithm for 5000 steps.

Subsequently the system was equilibrated in the NVT (constant number of particles, volume, and temperature) ensemble for 5 ns. During the minimisation and NVT equilibration, the heavy atoms of the solute were restrained using a harmonic potential with a force constant of $400 \text{ kJ mol}^{-1} \text{ nm}^{-2}$.

As a last step, the system was equilibrated in the NPT (constant number of particles, pressure, and temperature) ensemble for 45 ns. Throughout the whole protocol, the temperature was set to 300 K and the pressure was set to 1 bar. To ensure the constant temperature and pressure, the system was coupled to a ν -rescale thermostat [5] with a coupling constant of 1 ps and an isotropic c -rescale barostat [6] with a coupling constant of 5 ps. During the NPT run, the cut-off for the non-bonded interactions was set to 0.6 nm and the long-range electrostatics were treated using the Particle Mesh Ewald (PME) method.

All simulations were conducted in GROMACS 2021.4 [7] using the CHARMM36m forcefield [8] and the leap-frog integration method with a time step of 1 fs. All hydrogen involving bonds were constrained using the LINCS algorithm. The final dimensions of the box for all further calculations were obtained after the NPT run and are shown in Table 3.1. The last frame of the NPT runs was used as a starting point in all further calculations unless stated otherwise.

Table 3.1: System composition and simulation box details. ¹The final dimensions were obtained after the NPT run using the CHARMM36m forcefield.

System	Equilibrated box dimensions ¹ (Å)	No. of water molecules	No. of Na^+
MeDP	$15.877 \times 15.877 \times 15.877$	119	3
MeHDP	$15.901 \times 15.901 \times 15.901$	124	2

3.1.3 Collective variables

In order to properly sample the reaction space, we used two types of collective variables (CVs) to bias the system, e.g. distances and coordination numbers (CNs). The coordination number is defined as the smooth coordination number function:

$$\sum_{i \in A} \sum_{j \in B} CN_{ij} = \frac{1 - \left(\frac{r_{ij} - d_0}{r_0} \right)^n}{1 - \left(\frac{r_{ij} - d_0}{r_0} \right)^m} \quad (3.1)$$

where r_{ij} is the distance between the atoms i and j from the groups A and B , d_0 is a shift in the distance where CN starts to decay, r_0 is a scaling parameter or a characteristic decay length that defines how fast the CN drops off with distance, and n and m are the positive integers that control the steepness of the function. Normally, $m > n$ and the purpose of these parameters is to control how quickly the neighbours stop contributing as the distance increases. Therefore, the CN_{ij} ranges from ≈ 1 when $r_{ij} \leq r_0$ and ≈ 0 when $r_{ij} \geq r_0$. The CVs used to bias the systems in this work are shown in Figure 3.1. The corresponding parameters to describe each collective variable are as follows:

- The distance between the β -phosphorus and the oxygen atom connecting it to the rest of the molecule (CV₁, d(O_{remaining} - P_{leaving}),
- The coordination number of all oxygens around the β -phosphorus atom (CV₂, CN(P_{leaving} - O_{all}): $d_0 = 0$, $r_0 = 2.1$ Å, $n = 8$, $m = 16$,
- The coordination number of all hydrogens excluding the methyl ones around the oxygen atoms bound to the β -phosphorus atom (CV₃, CN(O_{leaving} - H_{all}): $d_0 = 0$, $r_0 = 1.4$ Å, $n = 6$, $m = 12$.

Besides the above mentioned CVs that were biased during the simulations, we also used the following collective variables to track the number of H₃O⁺ and OH⁻ in the solution:

- The number of H₃O⁺ ($n_{\text{H}_3\text{O}^+}$): TODO,
- The number of OH⁻ (n_{OH^-}): TODO.

To prevent the system from exploring the unphysical regions of the potential energy surface, we additionally applied quadratic (harmonic-like) wall potentials to softly restrain the accessible regions of some of the degrees of freedom.

The aforementioned wall potentials have the following form:

$$\text{for upper walls: } \sum_i k_i \left(\frac{CV_i - a_i + o_i}{s_i} \right)^{e_i} \quad (3.2)$$

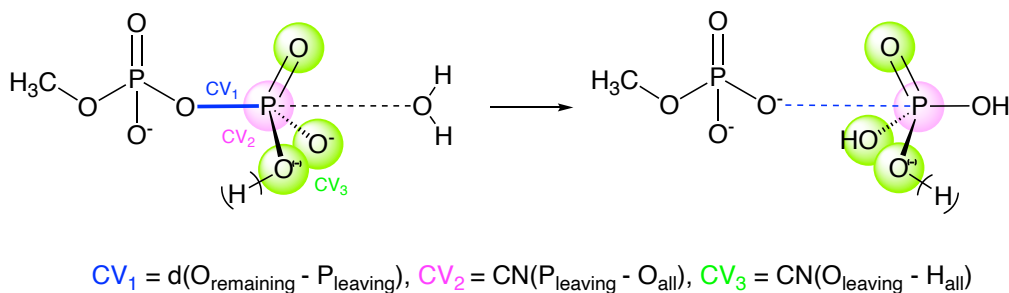


Figure 3.1: The definition of the collective variables (CVs) used in this work. CN stands for the coordination number.

$$\text{for lower walls: } \sum_i k_i \left| \frac{CV_i - a_i - o_i}{s_i} \right|^{e_i} \quad (3.3)$$

where CV_i is the value of the collective variable, k_i is the force constant (wall strength) that controls how strong the wall pushes back, a_i is the position of the wall, o_i is the offset to shift the wall slightly, s_i is the rescaling factor, and e_i is the exponent that determines the sharpness of the wall. When $e_i = 2$ the wall is harmonic-like.

Thus this potential is equal to zero while the CV is within the boundaries but once the CV tries to go beyond the boundaries, the potential increases and penalises the system.

The wall potentials were applied to the following collective variables:

- CV_1 , $d(O_{\text{remaining}} - P_{\text{leaving}})$: $k = 500 \text{ kcal mol}^{-1} \text{ \AA}^{-2}$, $a = 5.0 \text{ \AA}$ (upper wall), $o = 0 \text{ \AA}$, $s = 1 \text{ \AA}$, $e = 2$,
- TODO (CONSIDER TRANSFORMING INTO A TABLE)

All CV related computations were carried out using either the CP2K 2023.1 in-built tools [9] or PLUMED 2.9.3 [10]. It's important to note that the number of the CVs and the restraints used in the simulations varies depending on the part of the overall workflow. In the next sections we will clearly mention the respective collective variables and wall potentials used in each part of the workflow.

3.1.4 GFN1-xTB based exploration of the configuration space

To generate the first set of the configurations for the training dataset, the system was subjected to molecular dynamics simulations using the semi-empirical GFN1-xTB [11] level of theory. GFN1-xTB gives a good first approximation of the potential energy surface and is computationally efficient hence making it suitable for relatively long MD simulations of big systems.

Each system was first equilibrated for 5 ps in the NVT ensemble at 300 K to relax the structures at the GFN1-xTB level of theory with a D3 dispersion correction [12]. Afterwards, we performed 50 ps long well-tempered metadynamics (WTMD) [13] simulations in the NVT ensemble as well. In the latter simulation, the biasing potential was applied to force the system to explore the configuration space outside of the reactants basin. The biasing potential was added along two collective variables (CV): the distance between the β -phosphorus and the oxygen atom connecting it to the rest of the molecule (CV₁) and the coordination number of all oxygens around the β -phosphorus atom (CV₂).

All calculations were performed using the CP2K 2023.1 package [9]. The temperature was controlled by means of the ν -rescale thermostat [5] with a time constant of 50 fs for the equilibration and 100 fs for the WTMD run. The self-consistent field (SCF) convergence was set to 10^{-5} a.u. The biasing potential was applied every 25 fs with a gaussian hill height of 2 kcal mol⁻¹ and a width of 0.07 for each CV. The bias factor was set to 30. Lastly, the time step for the integration was set to 0.5 fs. Throughout all simulations, the periodic boundary conditions (PBC) were applied in all directions.

3.1.5 Data labeling

All data points were labelled by means of running the single-point calculations in order to obtain the energy and force values. The single-point calculations were performed using the Perdew–Burke–Ernzerhof (PBE) exchange-correlation functional [14] with the D3 dispersion correction and the Becke-Johnson damping function [12, 15]. In all calculations, the Goedecker-Teter-Hutter (GTH) pseudopotentials [16, 17] were used to represent the core electrons together with the triple- ζ valence basis set with two polarisation functions (TZV2P).

The single-point calculations were carried out using the Gaussian Plane Wave (GPW) method implemented in QUICKSTEP module [18] of the CP2K 2023.1 package [9]. The SCF convergence was set to 10^{-6} a.u. A plane-wave cutoff of 800 Ry for the total density and a cutoff of 60 Ry for the Kohn-Sham orbitals were utilised.

The above mentioned cutoffs were determined based on the convergence test per-

formed on one of the configurations as described in [19]. The error in total energy less than 10^{-8} a.u. for the convergence test has been considered as sufficient. The convergence test was performed using the following parameters: the cutoff for the total density was varied from 400 to 1500 Ry and the cutoff for the Kohn-Sham orbitals was varied from 10 to 200 Ry. The results of the convergence test are shown in Table A.1.

3.1.6 Iterative training of the neural network potential

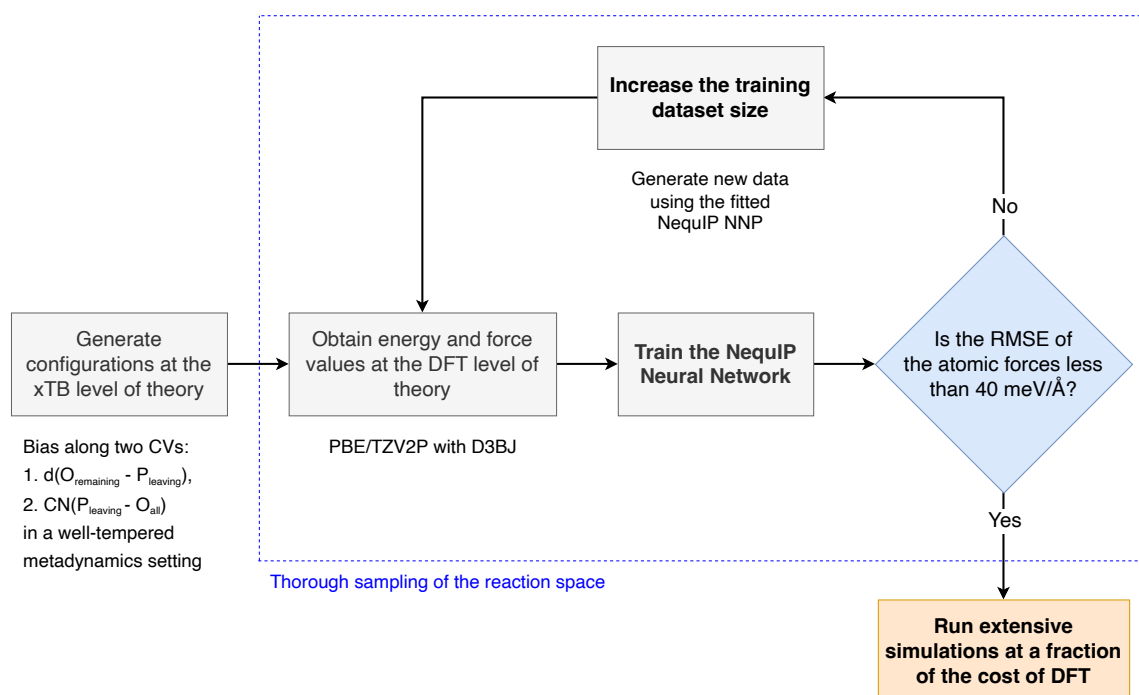


Figure 3.2: Iterative training of the NequIP neural network potential.

First round

Second round

Third round

3.2 Production runs at different temperatures

3.3 Validation of the transition states

3.4 Lifetime of the transition states

3.5 Data analysis and visualisation

Chapter 4

Results and Discussion

Chapter 5

Conclusions

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Appendix A

Supplementary information

Table A.1: The plane-wave cutoff convergence test for DFT calculations. The calculation of ΔE involves subtracting the previous energy, e.g. $\Delta E(450 \text{ Ry}) = E(450 \text{ Ry}) - E(400 \text{ Ry})$. When the cutoff ≥ 800 and the rel cutoff ≥ 60 the error in total energy reduces to ca. 10^{-8} a.u. Only part of the results is shown for the sake of clarity.

Cutoff (Ry)	Rel cutoff (Ry)	Total energy (a.u.)	ΔE (a.u.)
400	60	-2352.6355962810	—
450	60	-2352.6262868887	9.31×10^{-3}
500	60	-2352.6262867349	1.54×10^{-7}
550	60	-2352.6254866602	8.00×10^{-4}
600	60	-2352.6243443853	1.14×10^{-3}
650	60	-2352.6242425582	1.02×10^{-4}
700	60	-2352.6224669798	1.78×10^{-3}
750	60	-2352.6209571227	1.51×10^{-3}
800	60	-2352.6212901605	-3.33×10^{-4}
850	60	-2352.6212901727	-1.22×10^{-8}
900	60	-2352.6212901873	-1.46×10^{-8}
950	60	-2352.6213082173	-1.80×10^{-5}
1000	60	-2352.6208957304	4.12×10^{-4}
10	800	-2354.4562984779	—
20	800	-2352.6775968461	1.78
30	800	-2352.6281701514	4.94×10^{-2}
40	800	-2352.6213637375	6.81×10^{-3}
50	800	-2352.6212892865	7.45×10^{-5}
60	800	-2352.6212901605	-8.74×10^{-7}
70	800	-2352.6212901729	-1.24×10^{-8}
80	800	-2352.6212901739	-1.00×10^{-9}
90	800	-2352.6212901739	0.00
100	800	-2352.6212901739	0.00

Quantum Chemistry and Physical Chemistry

Celestijnenlaan 200F bus 2404

3001 LEUVEN, BELGIË

tel. + 32 16 37 21 98

jeremy.harvey@kuleuven.be

www.kuleuven.be

