




UNIVERSITY OF LUXEMBOURG  
Physics and Materials Science  
Research Unit (PHYMS)

# 12 – Schrödinger Equation

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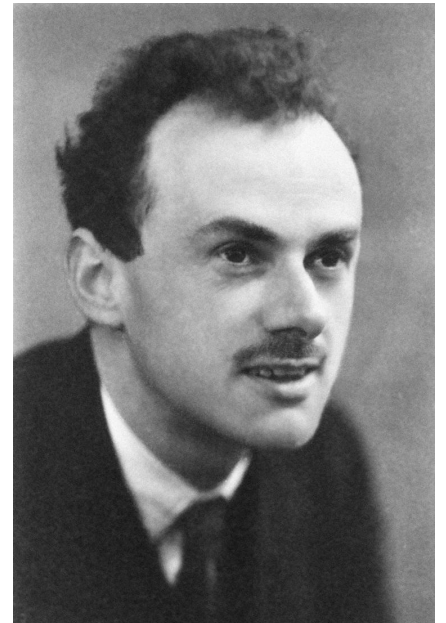



 (1933)

## Time-dependent Schrödinger equation

$$i\hbar \frac{\partial |\psi\rangle}{\partial t} = \hat{H} |\psi\rangle$$

$$|\psi\rangle = |\tilde{\psi}\rangle \cdot e^{-\frac{iEt}{\hbar}}$$



 (1933)

## Time-independent (stationary) Schrödinger equation

$$\hat{H} |\psi\rangle = E |\psi\rangle$$

*“The underlying physical laws necessary for the mathematical theory of a large part of physics and the whole of chemistry are thus completely known, and the difficulty is only that the exact application of these laws leads to equations much too complicated to be soluble”.*

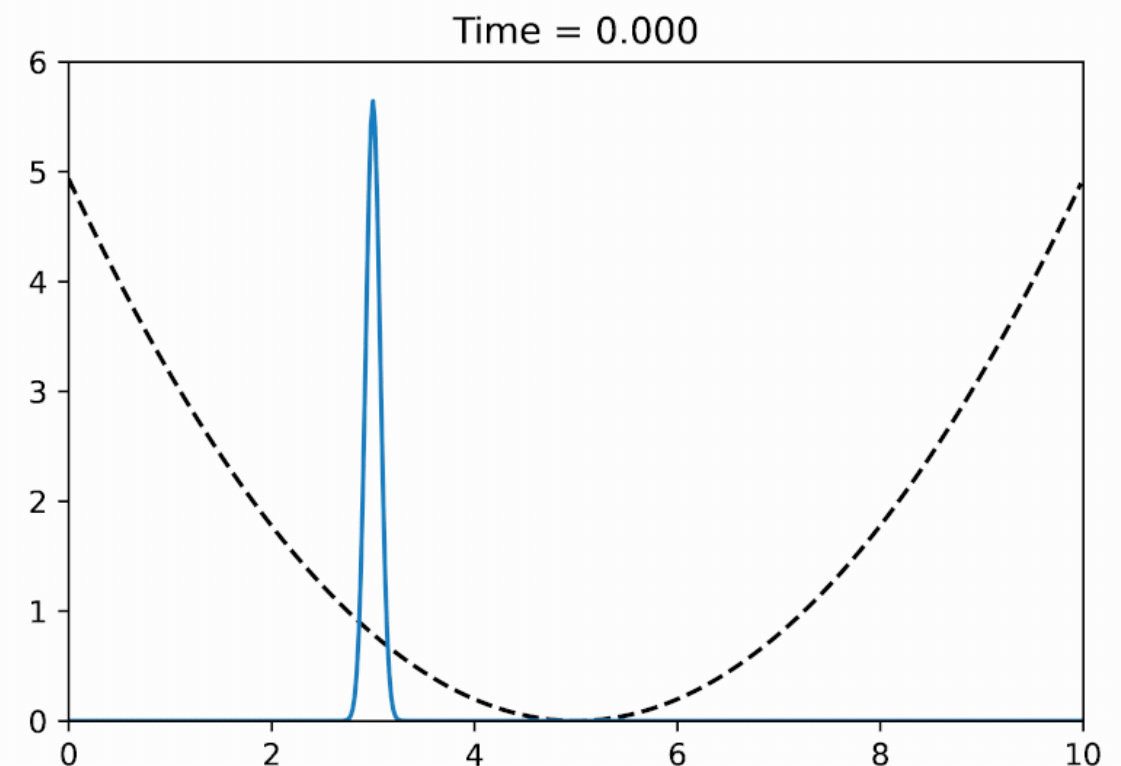
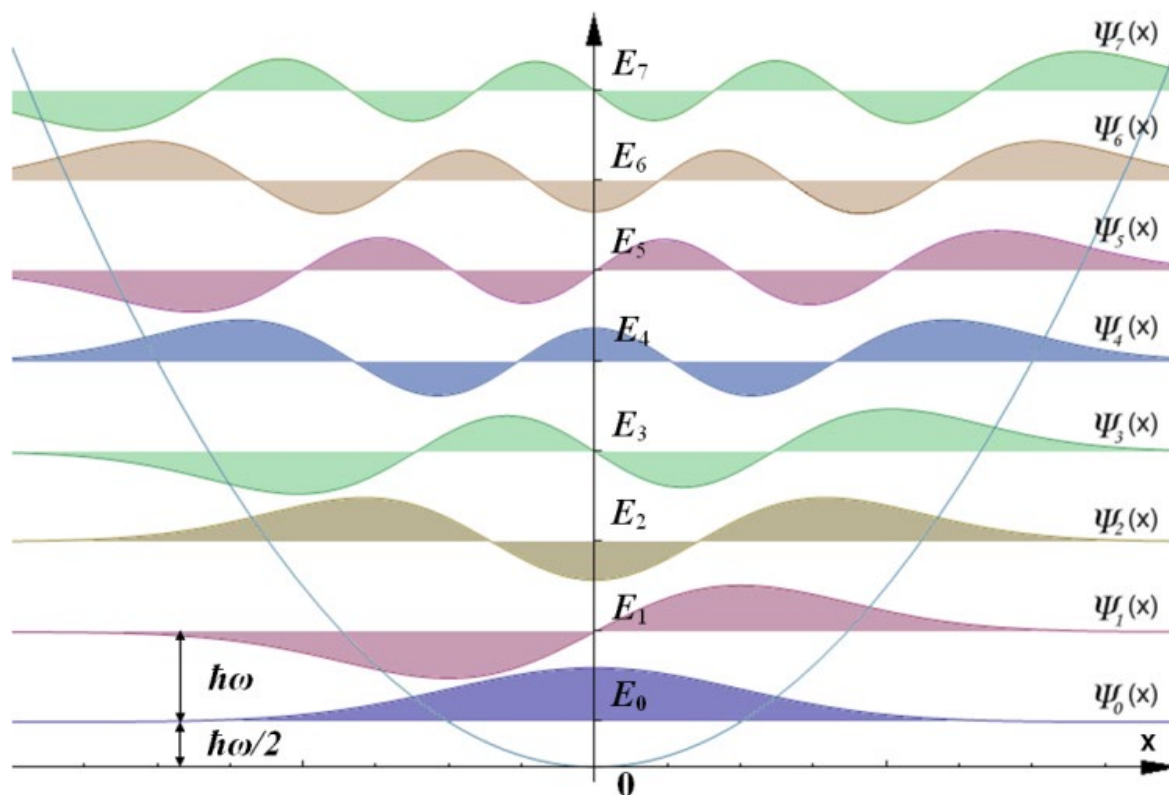
P.A.M. Dirac, 1929

# Harmonic Oscillator

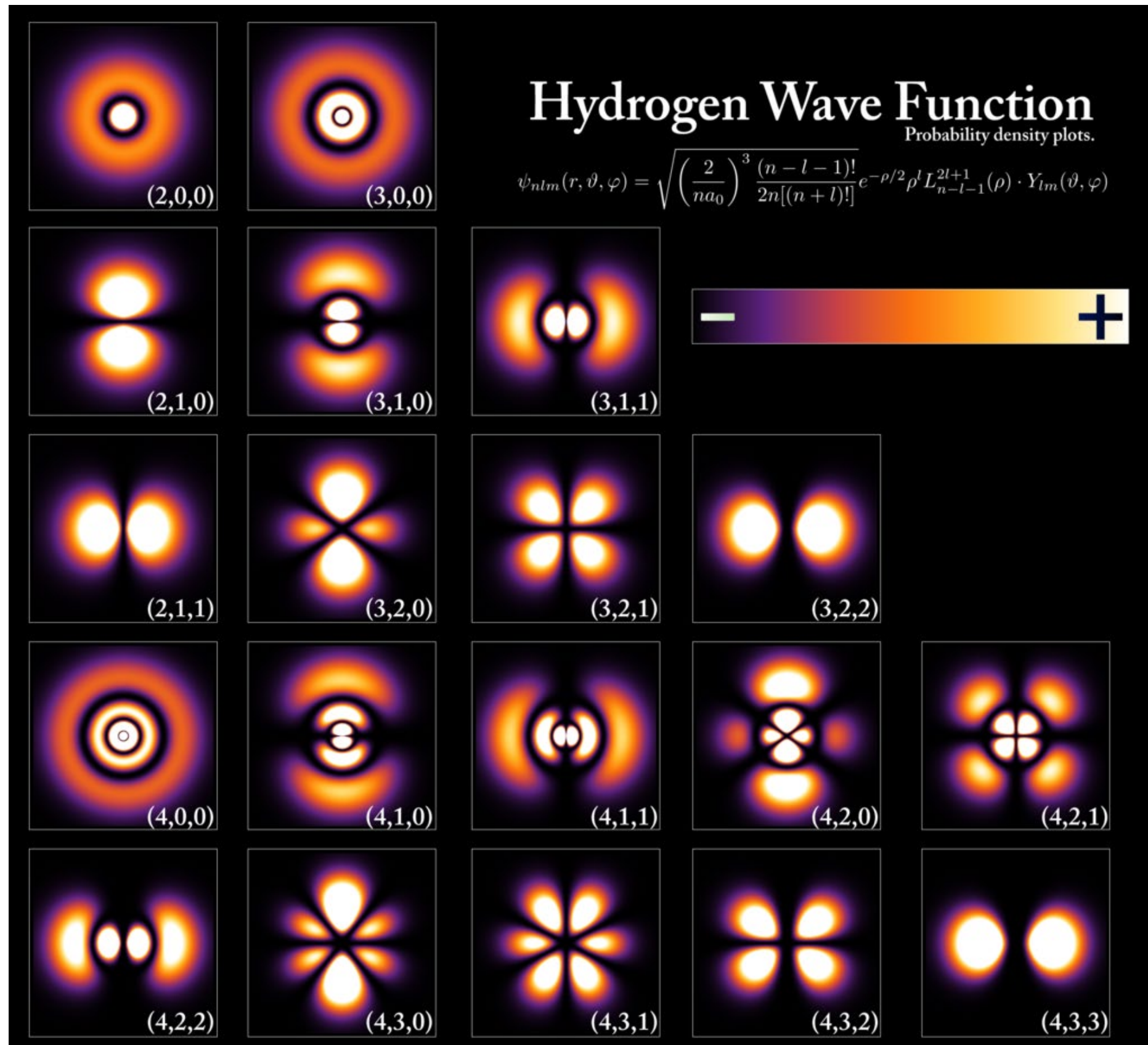
$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 x^2 = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2}m\omega^2 x^2$$

$$E_n = \hbar\omega \left( n + \frac{1}{2} \right)$$

$$\psi_n = \frac{1}{\sqrt{2^n n!}} \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\frac{m\omega x^2}{2\hbar}} H_n \left( \sqrt{\frac{m\omega}{\hbar}} x \right)$$

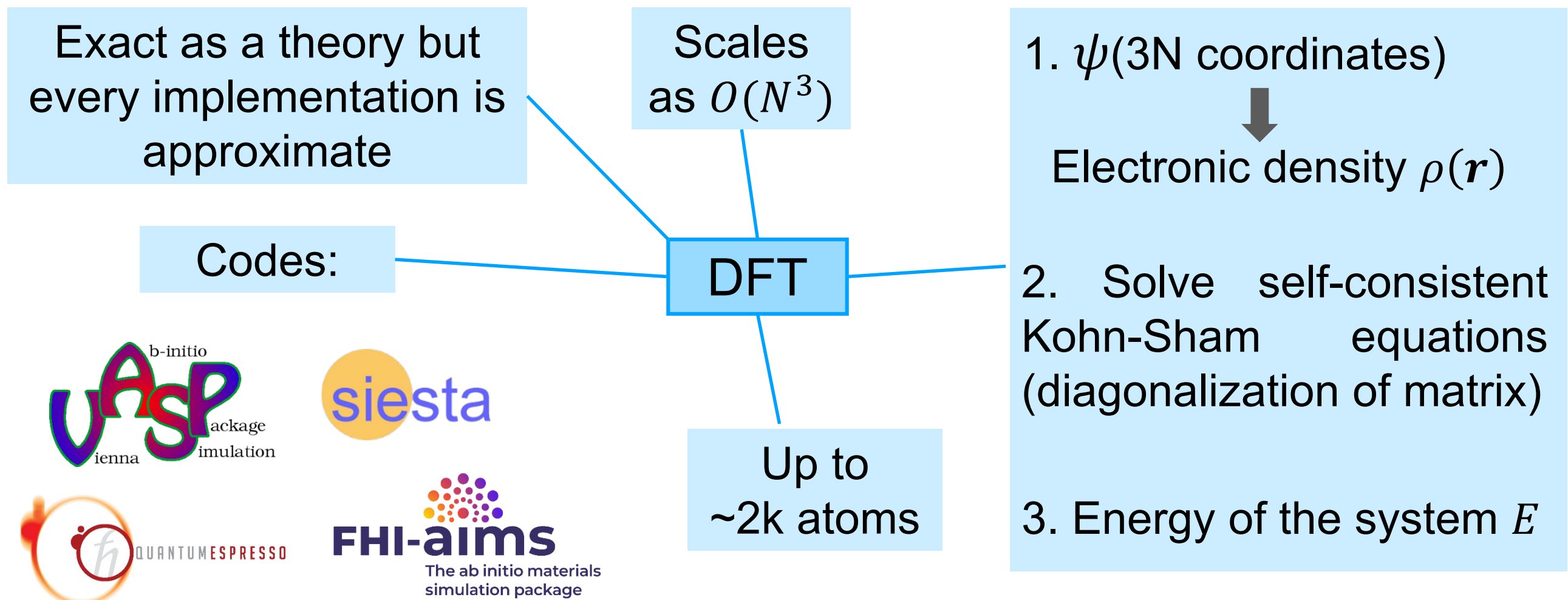


# Hydrogen Atom





- Coupled cluster method CCSD(T) – “golden standard” in the modern quantum chemistry; scales as  $O(N^7)$ ; limited to ~20-30 atoms
- Quantum Monte Carlo – similar size limitations, very long calculations
- Density-Functional Theory (DFT) – the main “workhorse” in modern simulations of materials and molecules (~40k papers per year)



One-dimensional time-independent Schrödinger equation:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$$

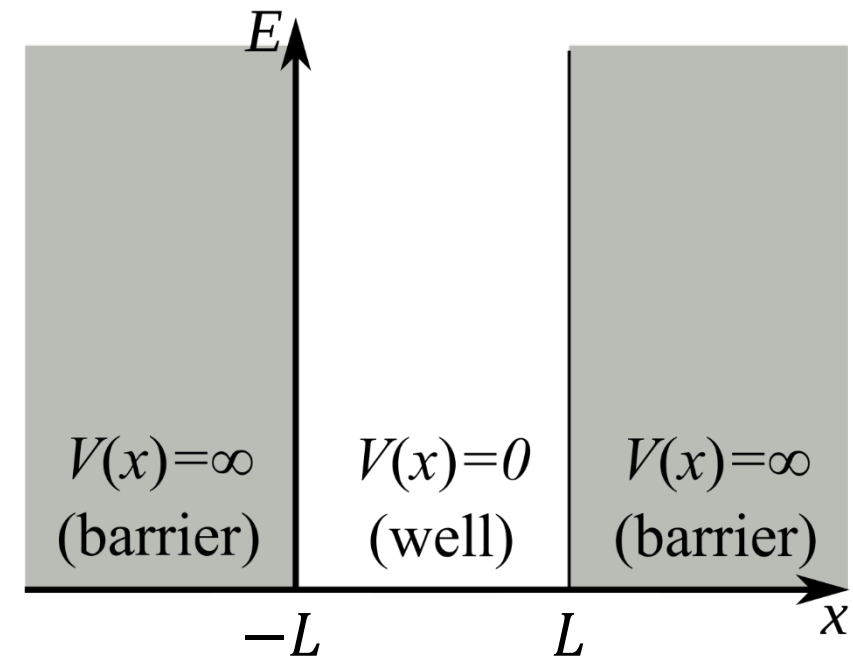
- If  $\psi(x)$  is a solution, then  $C \cdot \psi(x)$  is also a solution for any  $C \in \mathbb{C}$
- The probability density function is given by  $P(x) = \psi^*(x)\psi(x) = |\psi(x)|^2 \rightarrow$  should be normalized to give  $\int_{-\infty}^{+\infty} |\psi(x)|^2 dx = 1$
- If the potential  $V(x)$  is symmetric (i.e. has even parity), then the solution have either even or odd parities
- If we order solutions by their energies  $E_n$ , the  $n^{\text{th}}$  solution  $\psi_n(x)$  has  $(n - 1)$  nodes (i.e. points where  $\psi_n(x) = 0$ )

# Particle in a Box

The simplest possible potential:

$$V(x) = \begin{cases} +\infty, & |x| > L \\ 0, & -L \leq x \leq L \end{cases}$$

The wave function has to vanish  
at  $x = 0$  and  $x = L$



The even solutions:

$$\psi_k(x) = A \cdot \cos(\pi k x / 2L);$$
$$k = 2n - 1; \quad n \in \mathbb{N}$$

The odd solutions:

$$\psi_k(x) = A \cdot \sin(\pi k x / 2L);$$
$$k = 2n; \quad n \in \mathbb{N}$$

The energy levels:

$$E_k = \frac{\hbar^2 \pi^2 k^2}{8mL^2}$$

$$\psi''(x) + \frac{2mE}{\hbar^2} \psi(x) = 0$$

Looks very similar to the equation for harmonic oscillator

but...

We do not know  $E$  in advance. It should be determined consistently with  $\psi(x)$  to make the wave function satisfy the boundary conditions:

$$\psi(-L) = \psi(L) = 0$$

Mathematically this is a two-point **boundary value problem**. One possible method to solve such problems is the **shooting method**.



Reduce the boundary value problem to the **initial value problem**:

$$\begin{cases} \frac{d}{dx} \psi(x) = \psi'(x) \\ \frac{d}{dx} \psi'(x) = -\frac{2mE}{\hbar^2} \psi(x) \end{cases}$$

Initial conditions:

$$\begin{cases} \psi(0) = \psi_0 \\ \psi'(0) = \psi'_0 \end{cases}$$

$\psi(x)$  are either even or odd, so we can restrict ourselves to the interval  $[0, L]$  to reduce the amount of computations. Afterwards the solution can be continued on  $[-L, L]$ .

“Even” conditions:

$$\begin{cases} \psi(0) = 1 \\ \psi'(0) = 0 \end{cases}$$

“Odd” conditions:

$$\begin{cases} \psi(0) = 0 \\ \psi'(0) = 1 \end{cases}$$

# The Shooting Method

The Euler-Cromer integration scheme:

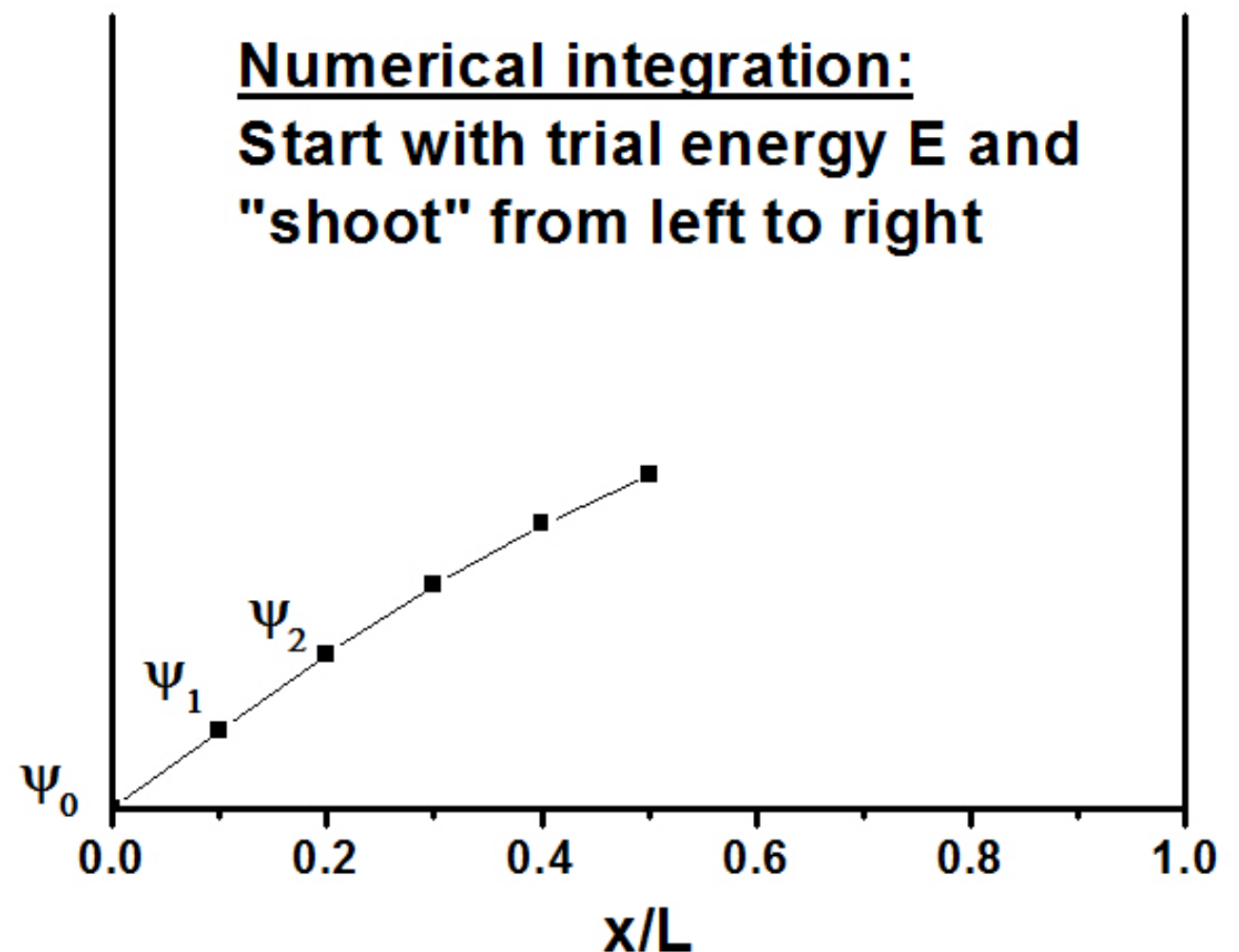
$$\begin{cases} \psi'(x + \Delta x) = \psi'(x) - \frac{2mE}{\hbar^2} \psi(x) \Delta x \\ \psi(x + \Delta x) = \psi(x) + \psi'(x + \Delta x) \Delta x \end{cases}$$

But we still do not know  $E$ ...

**Solution:** fix some trial energy  $E$  and just try

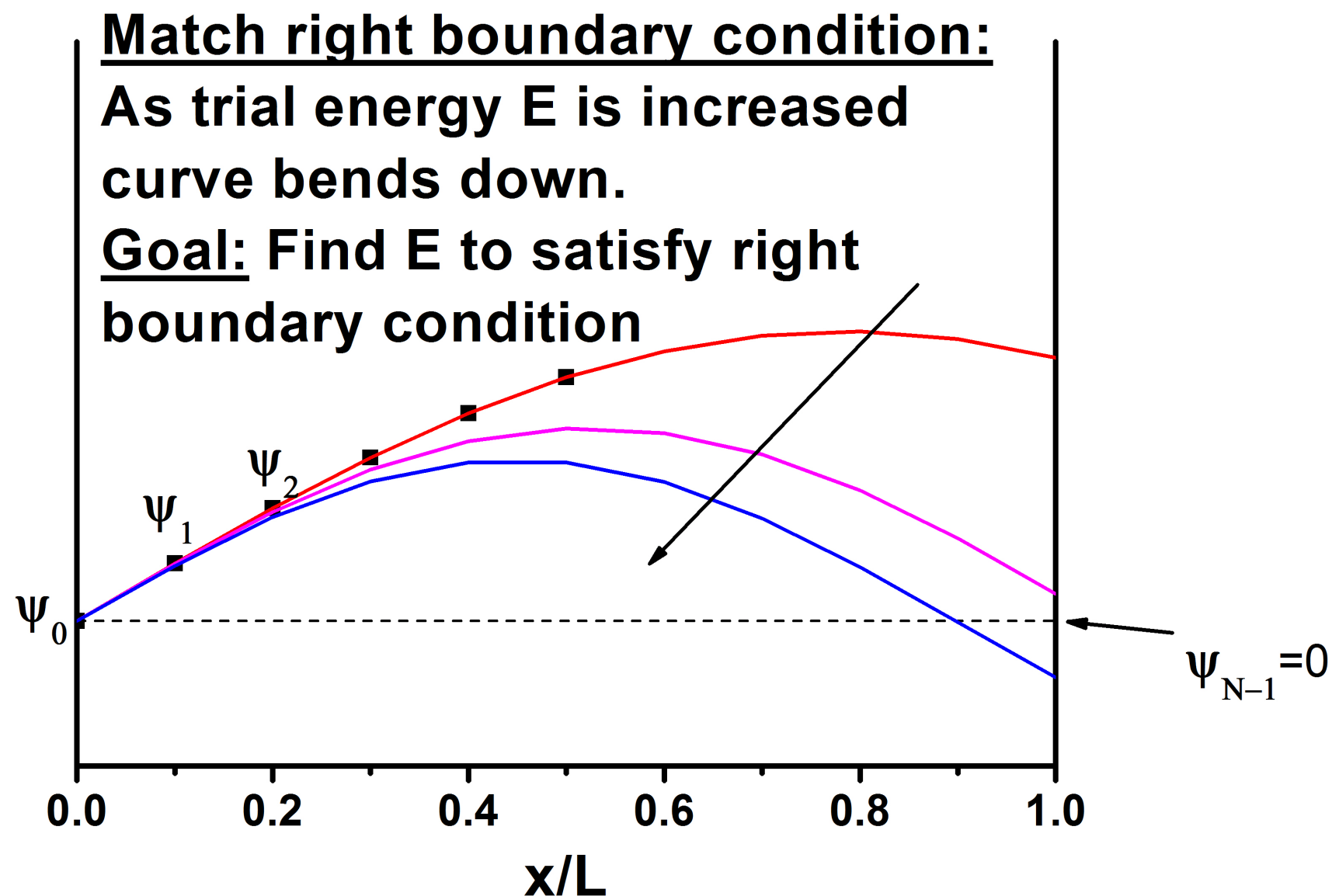
Basically, we simply solve two coupled ODEs using our integration scheme.

**Our goal:** to have  $\psi(L) = 0$

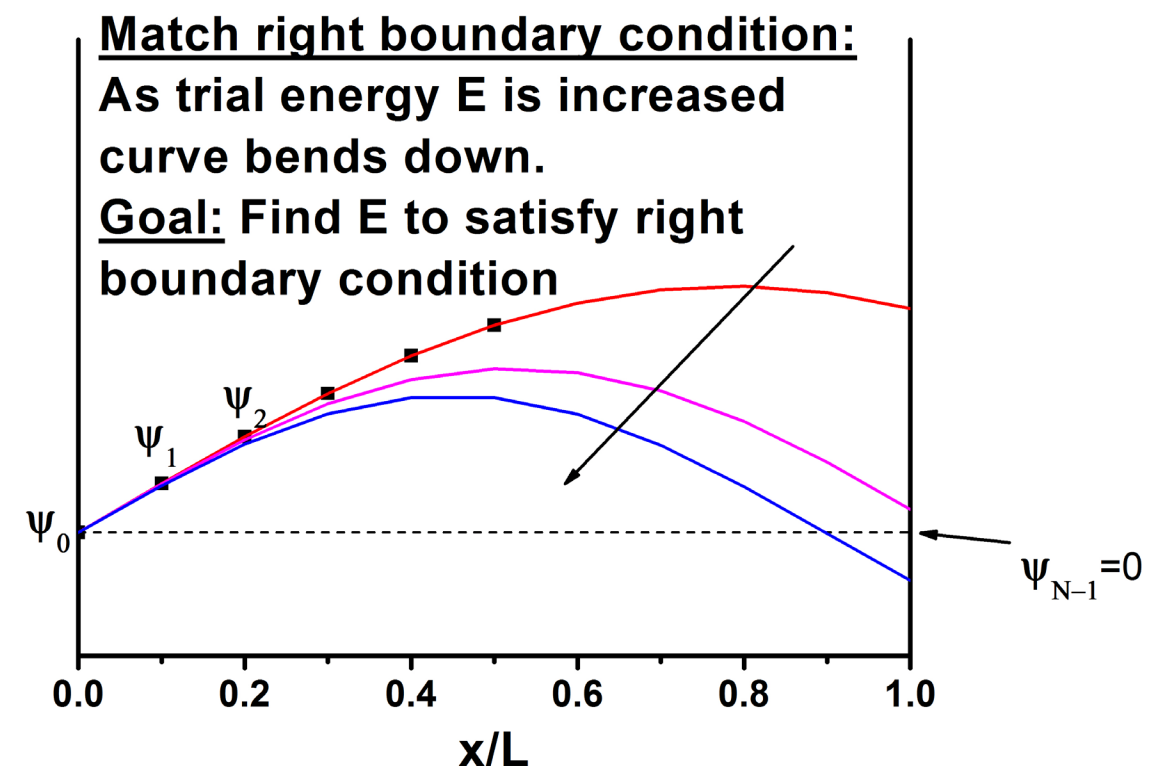


# The Shooting Method

Now shoot to the right boundary. Unless your trial energy matches an eigenvalue  $E_k$ ,  $\psi(x)$  will not hit 0 at the right boundary. Need to increment trial energy step by step. When overshooting, need to decrease it with a smaller decrement.

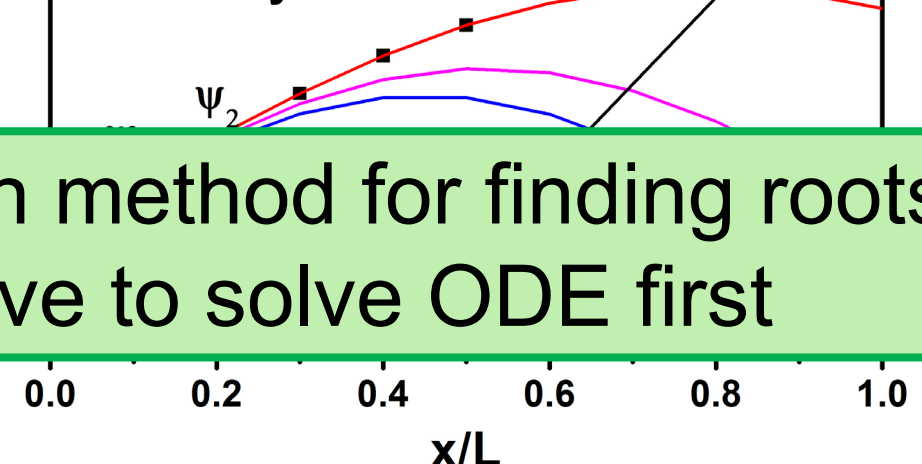


1. Set the trial energy  $E$ , choose the initial conditions (even or odd)
2. Make a first “shoot”: evaluate  $\psi(x)$  on  $[0, L]$  using Euler-Cromer
3. Store the value of obtained  $\psi(L)$
4. Define the energy increment  $dE$  and tolerance  $\varepsilon \ll dE$
5. Update  $E = E + dE$
6. “Shoot” again with a new  $E$ , obtain a new  $\tilde{\psi}(L)$
7. If  $\tilde{\psi}(L) \cdot \psi(L) < 0$ , make  $dE = -0.1 \cdot dE$
8. Update  $\psi(L) = \tilde{\psi}(L)$
9. Repeat 5) – 8) while  $dE > \varepsilon$



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**Match right boundary condition:**  
As trial energy  $E$  is increased  
curve bends down.  
**Goal:** Find  $E$  to satisfy right  
boundary condition



The idea is very similar to the bisection method for finding roots, but to evaluate our function we just have to solve ODE first