

Propagation and standing wave in transmission lines

Alejandro García Lampérez

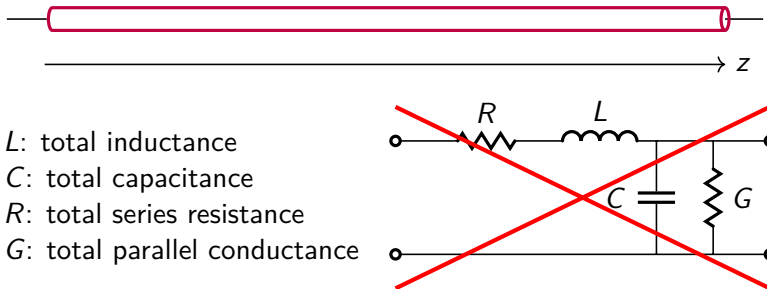
Dpto. de Teoría de la Señal y Comunicaciones
Universidad Carlos III de Madrid

aglamper@ing.uc3m.es

February 14, 2022

Transmission line

Coaxial cable, twin-lead line, ...



L : total inductance

C : total capacitance

R : total series resistance

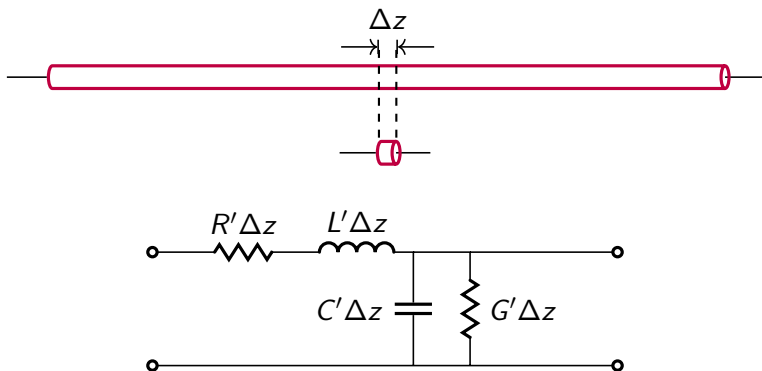
G : total parallel conductance

Primary parameters

- L' : **inductance** per unit length H/m
- C' : **capacitance** per unit length F/m
- R' : series **resistance** per unit length Ω /m
- G' : parallel **conductance** per unit length S/m

Transmission line

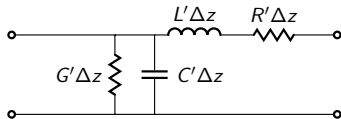
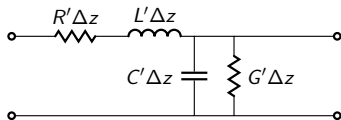
Model of a short section



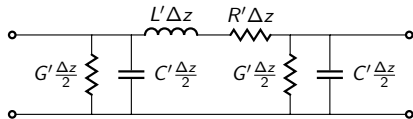
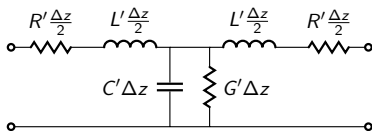
- $L'\Delta z$: total inductance
- $C'\Delta z$: total capacitance
- $R'\Delta z$: total series resistance
- $G'\Delta z$: total parallel conductance

Transmission line

Alternative circuit models for a short section



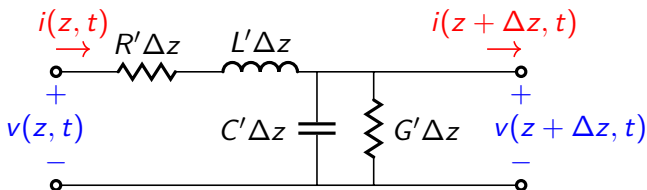
- $L'\Delta z$: total inductance
- $R'\Delta z$: total series resistance



- $C'\Delta z$: total capacitance
- $G'\Delta z$: total parallel conductance

Circuit analysis of the short section model

Kirchoff's laws



$$v(z, t) - v(z + \Delta z, t) = R'\Delta z i(z, t) + L'\Delta z \frac{\partial i(z, t)}{\partial t}$$

$$i(z, t) - i(z + \Delta z, t) = G'\Delta z v(z + \Delta z, t) + C'\Delta z \frac{\partial v(z + \Delta z, t)}{\partial t}$$

Telegrapher's equations

Circuit differential equations

$$\begin{cases} v(z, t) - v(z + \Delta z, t) = R' \Delta z i(z, t) & + L' \Delta z \frac{\partial i(z, t)}{\partial t} \\ i(z, t) - i(z + \Delta z, t) = G' \Delta z v(z + \Delta z, t) & + C' \Delta z \frac{\partial v(z + \Delta z, t)}{\partial t} \end{cases}$$

Divide by Δz and change signs

$$\begin{cases} \frac{v(z + \Delta z, t) - v(z, t)}{\Delta z} = -R' i(z, t) - L' \frac{\partial i(z, t)}{\partial t} \\ \frac{i(z + \Delta z, t) - i(z, t)}{\Delta z} = -G' v(z, t) - C' \frac{\partial v(z + \Delta z, t)}{\partial t} \end{cases}$$

Limit at $\Delta z \rightarrow 0$

$$\begin{cases} \frac{\partial v(z, t)}{\partial z} = -R' i(z, t) - L' \frac{\partial i(z, t)}{\partial t} \\ \frac{\partial i(z, t)}{\partial z} = -G' v(z, t) - C' \frac{\partial v(z, t)}{\partial t} \end{cases}$$

Telegrapher's equations

Time domain t

$$\begin{cases} \frac{\partial v(z, t)}{\partial z} = -R' i(z, t) - L' \frac{\partial i(z, t)}{\partial t} \\ \frac{\partial i(z, t)}{\partial z} = -G' v(z, t) - C' \frac{\partial v(z, t)}{\partial t} \end{cases}$$

Frequency domain ω

$$\begin{cases} \frac{dV(z)}{dz} = -(R' + j\omega L') I(z) \\ \frac{dI(z)}{dz} = -(G' + j\omega C') V(z) \end{cases}$$

How to solve: differentiate one equation, substitute the other

$$\frac{d^2 V(z)}{dz^2} = -(R' + j\omega L') \frac{dI(z)}{dz} = \overbrace{(R' + j\omega L')(G' + j\omega C')}^{\gamma^2} V(z)$$

Helmholtz wave equations

$$\frac{d^2 V(z)}{dz^2} - \gamma^2 V(z) = 0$$

$$\frac{d^2 I(z)}{dz^2} - \gamma^2 I(z) = 0$$

Propagation

Helmholtz equations

$$\frac{d^2 V(z)}{dz^2} - \gamma^2 V(z) = 0$$

$$\frac{d^2 I(z)}{dz^2} - \gamma^2 I(z) = 0$$

General solutions

$$V(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z}$$

$$I(z) = I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z}$$

Propagation constant

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')} = \alpha + j\beta$$

- α : attenuation constant Np/m
- β : phase constant rad/m

Characteristic impedance

$$\begin{aligned}V(z) &= V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z} \\I(z) &= I_0^+ e^{-\gamma z} + I_0^- e^{\gamma z}\end{aligned}$$

But the solutions are not independent: $\frac{dV(z)}{dz} = -(R' + j\omega L')I(z)$

Definition: characteristic impedance

$$\frac{V_0^+}{I_0^+} = -\frac{V_0^-}{I_0^-} \triangleq Z_0 = \sqrt{\frac{R' + j\omega L'}{G' + j\omega C'}}$$

$$\left. \begin{aligned}\gamma &= \sqrt{(R' + j\omega L')(G' + j\omega C')} \\Z_0 &= \sqrt{\frac{R' + j\omega L'}{G' + j\omega C'}}\end{aligned}\right\} \text{Secondary parameters}$$

Physical interpretation of the solution

Voltage/current waves

$$V(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z}$$
$$I(z) = \frac{V_0^+}{Z_0} e^{-\gamma z} - \frac{V_0^-}{Z_0} e^{\gamma z}$$

- Let be $V_0^+ = |V_0^+| e^{j\phi^+}$ $V_0^- = |V_0^-| e^{j\phi^-}$ $\gamma = \alpha + j\beta$
- Time domain: $v(t, z) = \text{Re}[V(z)e^{j\omega t}]$

Time-domain solution

$$v(z, t) = \underbrace{|V_0^+| e^{-\alpha z} \cos(\omega t - \beta z + \phi^+)}_{\text{Incident wave (towards } +z)} + \underbrace{|V_0^-| e^{\alpha z} \cos(\omega t + \beta z + \phi^-)}_{\text{Reflected wave (towards } -z)}$$

<https://www.desmos.com/calculator/y3jnr0xqwy>

Lossless (non-dissipative) line

Particular case

$$R = G = 0 \quad \Rightarrow \quad \begin{cases} \gamma = j\beta = j\omega\sqrt{LC} & (\alpha = 0) \\ Z_0 = \sqrt{\frac{L}{C}} \end{cases}$$

Voltage and current waves

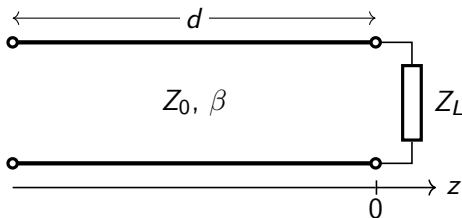
$$\begin{aligned} V(z) &= V_0^+ e^{-j\beta z} + V_0^- e^{j\beta z} \\ I(z) &= \frac{V_0^+}{Z_0} e^{-j\beta z} - \frac{V_0^-}{Z_0} e^{j\beta z} \end{aligned}$$

Time-domain solution

$$v(z, t) = \underbrace{|V_0^+| \cos(\omega t - \beta z + \phi^+)}_{\text{Incident wave (towards } +z)} + \underbrace{|V_0^-| \cos(\omega t + \beta z + \phi^-)}_{\text{Reflected wave (towards } -z)}$$

Terminated line

Boundary condition: impedance



$$V(z) = V_0^+ e^{-\gamma z} + V_0^- e^{\gamma z}$$
$$I(z) = \frac{V_0^+}{Z_0} e^{-\gamma z} - \frac{V_0^-}{Z_0} e^{\gamma z}$$

For $z = 0$:

$$\left. \begin{aligned} V(0) &= V_0^+ + V_0^- \\ I(0) &= \frac{V_0^+ - V_0^-}{Z_0} \end{aligned} \right\} \Rightarrow Z_L = \frac{V(0)}{I(0)} = Z_0 \frac{V_0^+ + V_0^-}{V_0^+ - V_0^-}$$

Definition: reflection coefficient

$$\Gamma \triangleq \frac{V_0^-}{V_0^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

Reflection coefficient

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$$

$$Z_L = Z_0 \frac{1 + \Gamma}{1 - \Gamma}$$

Voltage and current waves

$$V(z) = V_0^+ \left(e^{-j\beta z} + \Gamma e^{j\beta z} \right)$$

$$I(z) = \frac{V_0^+}{Z_0} \left(e^{-j\beta z} - \Gamma e^{j\beta z} \right)$$

Particular case: **matched load**

- $V_0^- = 0$
- $\Gamma = 0$
- $Z_L = Z_0$

Average power

$$\begin{aligned}P(z) &= \frac{1}{2} \operatorname{Re}[V(z)I(z)^*] \\&= \frac{1}{2} \frac{|V_0^+|^2}{Z_0} \operatorname{Re}[(e^{-j\beta z} + \Gamma e^{j\beta z})(e^{j\beta z} - \Gamma^* e^{-j\beta z})] \\&= \frac{1}{2} \frac{|V_0^+|^2}{Z_0} \operatorname{Re}[1 - |\Gamma|^2 + \Gamma e^{2j\beta z} - \Gamma^* e^{-2j\beta z}] \\&= \frac{1}{2} \frac{|V_0^+|^2}{Z_0} (1 - |\Gamma|^2)\end{aligned}$$

- Independent of z

- Two terms: $P^+ = \frac{1}{2} \frac{|V_0^+|^2}{Z_0}$ $P^+ = \frac{1}{2} \frac{|V_0^-|^2}{Z_0} = \frac{1}{2} \frac{|V_0^+|^2}{Z_0} |\Gamma|^2$

- $\Gamma = 0$: Maximum power transfer $|\Gamma| = 1$: no power transfer

Standing wave

Voltage amplitude

Let be $\Gamma = |\Gamma| e^{j\theta}$

$$\begin{aligned}|V(z)| &= |V_0^+| \left| e^{-j\beta z} + \Gamma e^{j\beta z} \right| = |V_0^+| \left| 1 + \Gamma e^{2j\beta z} \right| \\ &= |V_0^+| \left| 1 + |\Gamma| e^{j(2\beta z + \theta)} \right|\end{aligned}$$

Maxima and minima condition: $|\Gamma| e^{j(2\beta z + \theta)} \in \mathbb{R}$

- $V_{\max} = |V_0^+| |1 + |\Gamma|| \quad 2\beta z + \theta = 2n\pi \text{ rad}$
- $V_{\min} = |V_0^+| |1 - |\Gamma|| \quad 2\beta z + \theta = (2n + 1)\pi \text{ rad}$
- Distance: $z_{\max} - z_{\min} = \frac{\pi}{2\beta} = \frac{\lambda}{4}$
- Distance: $z_{\max} - z'_{\max} = \frac{\pi}{\beta} = \frac{\lambda}{2}$

<https://www.desmos.com/calculator/2pwe46lry0>