

COULOMB WAVE FUNCTIONS FOR ALL REAL η AND ρ *

A.R. BARNETT, D.H. FENG**, J.W. STEED† and L.J.B. GOLDFARB
 Physics Department, Schuster Laboratory, The University, Manchester M13 9PL, UK

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PROGRAM SUMMARY

Title of program: RCWFN

Catalogue number: ABPC

Program obtainable from: CPC Program Library, Queen's University of Belfast, N. Ireland (see application form in this issue).

<i>Computer:</i>	<i>Installation:</i>
CDC 7600	University of Manchester Regional Computing Centre
IBM 370/165	Daresbury Nuclear Physics Laboratory, Warrington, Lancs.
IBM 360/195	Rutherford High Energy Laboratory, Didcot, Berks.

Operating system: UMRCC Scope V2.0

Programming language used: ASA Fortran

High-speed core required: < 1 Kwords

No. of bits in a word: 60

Overlay structure: None

No. of magnetic tapes required: None

Other peripherals used: Card reader, line printer

No. of cards in the combined program and test deck: 282

Card punching code: EBCDIC

Keywords: Nuclear, atomic, Schrodinger equation, reactions, scattering, heavy ion, wavefunction, Coulomb, potential, hypergeometric function, continued fraction, Runge-Kutta.

Nature of the physical problem

The subroutine RCWFN calculates the regular and irregular Coulomb wavefunctions, $F_l(\eta, \rho)$ and $G_l(\eta, \rho)$ and their radial derivatives for real, positive energy for all (η, ρ) , $(-1000 \leq \eta \leq 1000, \rho > 0)$ rapidly and with high accuracy. It is well suited

for calculations of Coulomb matrix elements for heavy-ion scattering [1] ($1 < \eta < 200, \eta < \rho < 1000$) and for all programs which require matching to asymptotic Coulomb wavefunctions in both nuclear and atomic physics.

Method of solution

An entirely new technique, developed by Steed [2], is used, which differs greatly from standard methods [3]. The method is based on the continued-fraction expansion of the quantities F_l/F_l and $(G_l + F_l)/(G_l + iF_l)$ and it is applicable for all values of $\rho \geq \rho_{TP} = \eta + [\eta^2 + l(l+1)]^{1/2}$ without restriction. A somewhat modified method is used for values of $\rho < \rho_{TP}$. The results are given for all l -values between MINL and MAXL inclusive ($MAXL \geq MINL \geq 0$) and for any required ρ -value. In general results are obtained without integrating the differential equation. The subroutine is fast and compact ‡; the results can be obtained to any desired accuracy in the range 10^{-6} – 10^{-12} .

Restriction on the complexity of the problem

Subroutine RCWFN has been programmed for real values of η and ρ : the value of η is unrestricted (positive, negative or zero) while the value of ρ cannot be zero. There is a loss of accuracy in cases when $\rho < 0.2\rho_{TP}$.

Typical running time

For the $l = 0$ functions the averaging running time (CDC 7600 CP time) is 0.005 sec for ACCUR = 10^{-7} and STEP = 100.0 (i.e. practical values), and there is no significant increase with l -value.

References

- [1] A.R. Barnett, D.H. Feng and L.J.B. Goldfarb, Phys. Lett. 48B (1974) 290; and Computer Phys. Commun., to be published.
- [2] J.W. Steed, PhD thesis, University of Manchester (1967) unpublished.
- [3] C.E. Froberg, Rev. Mod. Phys. 27 (1955) 399; C. Bardin et al., Computer Phys. Commun. 3 (1972) 73.

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** Present address: Center for Nuclear Studies, University of Texas, Texas, USA.

† Now with Software Sciences Limited, Park Green, Macclesfield, UK.

‡ When the \$-separator format on the CDC card input is used the entire subroutine occupies 36 cards.

LONG WRITE-UP

1. Introduction

The Coulomb functions $F_l(\eta, \rho)$ and $G_l(\eta, \rho)$, together with their derivatives with respect to ρ , are basic to the solution of charged-particle scattering problems for both positive and negative values of η . An example taken from nuclear physics is the calculation of Coulomb excitation, in which the radial matrix elements

$$M_{l'l'}^{-\lambda-1} = (k^\lambda/kk') \times \int_0^\infty f_l(\eta, \rho) \rho^{-\lambda-1} f_{l'}(\eta', \rho') d\rho \quad (1)$$

are integrals over the regular solution $f_l = F_l$ if there is no nuclear distortion. Nuclear effects give rise to nuclear phase shifts δ_l different from zero for the lower partial waves and in this case

$$f_l = F_l + \frac{1}{2}i[1 - \exp(2i\delta_l)] (G_l + iF_l)$$

in the region outside the range of nuclear interactions. For problems of this type we require the Coulomb functions in the region

$$\rho \geq \rho_{TP} = \eta + [\eta^2 + l(l+1)]^{1/2}$$

since, otherwise, the integrand in (1) is vanishingly small. The value ρ_{TP} is the turning point of the functions which separates the monotonic region $\rho < \rho_{TP}$ from the oscillatory one $\rho > \rho_{TP}$.

The subroutine RCWFN calculates the Coulomb functions for the region $0.1\rho_{TP} < \rho < \infty$ and the method it uses represents a substantial break with tradition [1-4]. It relies on a continued-fraction representation and does not solve the differential equation (save in the small region $0.1\rho_{TP} < \rho < \rho_{TP}$). The subroutine is thus exceptionally fast, it evaluates F_l, F'_l, G_l and G'_l at a given ρ -value independently of any other, it can evaluate a single l -value if desired (i.e. there is no need for recursion from $l=0$) and the results can be obtained to any desired degree of accuracy. The second of these properties makes the evaluation of (1) by gaussian quadrature methods very attractive [5,6]. The methods of continued fractions are discussed in section 3 below.

2. General properties of the Coulomb functions

The Coulomb functions are solutions of the differential equation

$$\frac{d^2 U_l}{d\rho^2} + \left(1 - \frac{2\eta}{\rho} - \frac{l(l+1)}{\rho^2}\right) U_l = 0, \quad (2)$$

where the parameters ρ and η are defined as

$$\rho = kr = (2\mu E)^{1/2} r/\hbar$$

and

$$\eta = Z_1 Z_2 e^2 / \hbar v = \mu Z_1 Z_2 e^2 / \hbar^2 k.$$

The two charges are $Z_1 e$ and $Z_2 e$, μ is the reduced mass of the system in amu, k is the wave number and E the centre-of-mass energy. The numerical expressions obtained from the 1969 Atomic Constants [7] are

$$k \text{ (fm}^{-1}\text{)} = 0.218\,741\,0 [\mu(\text{amu}) \cdot E(\text{MeV})]^{1/2}$$

and

$$\eta = (1.439\,906 \times 0.218\,741\,0/2.0) Z_1 Z_2 \times [\mu(\text{amu})/E(\text{MeV})]^{1/2}.$$

The regular and irregular solutions of eq. (2), $F_l(\eta, \rho)$ and $G_l(\eta, \rho)$, have boundary conditions

$$\begin{aligned} F_l(\eta, \rho) &\xrightarrow{\rho \rightarrow \infty} \sin \theta_l, & F_l(\eta, 0) &= 0, \\ G_l(\eta, \rho) &\xrightarrow{\rho \rightarrow \infty} \cos \theta_l, & G_l(\eta, 0) &= \infty, \end{aligned} \quad (3)$$

where

$$\theta_l = \rho - \eta \ln 2\rho - \frac{1}{2}l\pi + \sigma_l.$$

The Coulomb phase shift σ_l is given by

$$\sigma_l = \arg \Gamma(l+1+i\eta).$$

The solutions for successive l -values are connected by the following recurrence relations, with $U_l = F_l$ or G_l ,

$$R_l U_{l-1} = U'_l + S_l U_l, \quad (4)$$

$$U'_{l-1} = S_l U_{l-1} - R_l U_l, \quad (5)$$

$$R_{l+1}U_{l+1} = T_l U_l - R_l U_{l-1}, \quad (6)$$

where the coefficients are given for $l \geq 1$ by

$$\begin{aligned} R_l &= (l^2 + \eta^2)^{1/2}/l, & S_l &= l/\rho + \eta/l, \\ T_l &= \frac{(2l+1)[l(l+1)/\rho + \eta]}{l(l+1)}. \end{aligned} \quad (7)$$

The primes denote differentiation with respect to ρ . Further necessary relations are the wronskian

$$W = F_l' G_l - F_l G_l' = 1, \quad \text{and} \quad (8)$$

$$W_l = F_l G_{l+1} - F_{l+1} G_l = 1/R_{l+1}. \quad (9)$$

3. Continued fractions and the Coulomb functions

A continued fraction is defined as

$$h = b_0 + \frac{a_1}{b_1 + \frac{a_2}{b_2 + \dots}},$$

which we shall write as

$$h = b_0 + \frac{a_1}{b_1 +} \frac{a_2}{b_2 +} \dots \frac{a_n}{b_n + \dots}. \quad (10)$$

We state two necessary theorems, whose proof can be found in ref. [8].

Theorem 1. The continued fraction h is said to converge if the n th convergent h_n (the result of terminating the fraction at b_n) is such that the sequence h_n as $n \rightarrow \infty$ tends to a limit h . The limit h is then the value of the continued fraction.

Theorem 2. Let two sequences, A_n and B_n satisfy the two-term recurrence relations (expressed in matrix form)

$$\begin{pmatrix} A_n \\ B_n \end{pmatrix} = \begin{pmatrix} A_{n-1} & A_{n-2} \\ B_{n-1} & B_{n-2} \end{pmatrix} \begin{pmatrix} b_n \\ a_n \end{pmatrix}, \quad (11)$$

with the initial set of values

$$\begin{pmatrix} A_0 & A_{-1} \\ B_0 & B_{-1} \end{pmatrix} = \begin{pmatrix} b_0 & 1 \\ 1 & 0 \end{pmatrix}.$$

Then the n th convergent of the continued fraction

$$h_n = A_n/B_n.$$

The equivalence transformation allows us to rewrite (10) as

$$h = b_0 + \frac{c_1 a_1}{c_1 b_1 +} \frac{c_1 c_2 a_2}{c_2 b_2 +} \frac{c_2 c_3 a_3}{c_3 b_3 +} \dots \frac{c_{n-1} c_n a_n}{c_n b_n + \dots} \quad (12)$$

for all $c_p \neq 0$ and it will be used in manipulating the results into their final form.

3.1. Steed's procedure for evaluating Coulomb functions

The mathematical methods and the algorithms described in this program were developed as original work by Steed [9] which is unpublished. His program has been translated into Fortran IV and extended to any arbitrary value of MINL.

We derive in sections 3.2 and 3.3 two continued fractions involving the Coulomb functions, the one being

$$f_m \equiv F_m'/F_m, \quad (13)$$

while the other is

$$p_m + iq_m \equiv (G_m' + iF_m')/(G_m + iF_m). \quad (14)$$

The use of eqs. (13) and (14) enables us to solve for F_m' , G_m' and G_m in terms of F_m , thus (dropping the subscript)

$$F' = fF, \quad (15)$$

$$G = (F' - pF)/q = (f - p)F/q, \quad (16)$$

$$G' = pG - qF = p(f - p)F/q - qF, \quad (17)$$

and then the value of F itself, i.e. the scaling factor for eqs. (15) – (17) can be found by substitution in the wronskian equation (8)

$$F_m = \pm 1/(F'G - FG')^{1/2}. \quad (18)$$

The sign of F_m is obtained during the evaluation of f_m (section 3.2).

The convergence of the continued-fraction expansion of eq. (13) is sufficiently rapid and is valid for all values of ρ , except for special cases with negative η (section 3.2). On the other hand eq. (14) can only be used for $\rho \geq \rho_{TP}$ since it very rapidly tends to G_m'/G_m for smaller

values of ρ . For example, at $\rho = 0.95 \rho_{TP}$ and $m = 0$, the real and imaginary parts of eq. (14), p and q , differ by a factor of 10^{-12} , and a loss of significance results. Therefore in the region $\rho < \rho_{TP}$ the differential equation (2) for G_m and G'_m is integrated inward to the desired ρ value from the starting values of $G_m(\eta, \rho_{TP})$ and $G'_m(\eta, \rho_{TP})$ and then combined with $f_m(\eta, \rho)$ as above. The method of integration is the standard 4th-order Runge–Kutta algorithm used simultaneously for G_m and G'_m [10].

In subroutine RCWFN the Coulomb functions are calculated at a required ρ -values for a range of l -values, from $l = \text{MAXL}$ to $l = \text{MINL}$, by using the recurrence relations eqs. (4)–(5) as is described in section 4.

3.2. A continued fraction for F'_l/F_l

Consider the recurrence relation (6) and let $V_n = U_{k+1}$ and also $k = n + l$ with l fixed at the desired value. Then

$$V_n = V_{n-1}b_n + V_{n-2}a_n, \quad (19)$$

where

$$b_n = T_k/R_{k+1} = \frac{(2k+1)[(k+1)/\rho + \eta/k]}{[(k+1)^2 + \eta^2]^{1/2}}$$

and

$$a_n = -R_k/R_{k+1} = -\frac{(k^2 + \eta^2)^{1/2}(k+1)}{[(k+1)^2 + \eta^2]^{1/2}k}$$

in the notation of eq. (7). Since V_n is a Coulomb function the sequence (19) is satisfied by the two independent solutions F_n and G_n . The most general sequences V_n are thus

$$A_n = \alpha F_n + \beta G_n \equiv \alpha F_{n+1+l} + \beta G_{n+1+l}$$

and

$$B_n = \gamma F_n + \delta G_n \equiv \gamma F_{n+1+l} + \delta G_{n+1+l}$$

from which we can find those coefficients which yield initial values of the sequences A_n, B_n in conformity with Theorem 2, namely $A_0 = b_0, A_{-1} = 1, B_0 = 1, B_{-1} = 0$. These coefficients are

$$\begin{aligned} \alpha &= (b_0 G_l - G_{l+1})R_{l+1}, & \beta &= (b_0 F_l - F_{l+1})R_{l+1}, \\ \gamma &= -G_l R_{l+1}, & \delta &= F_l R_{l+1}. \end{aligned} \quad (20)$$

The conditions of Theorem 2 are met and so the limit as $n \rightarrow \infty$ of the continued fraction defined by a_n, b_n is $\lim_{n \rightarrow \infty} A_n/B_n$ which, since $F_{n+1+l} \rightarrow 0$ and $G_{n+1+l} \rightarrow \infty$ is just the ratio β/δ . From eqs. (19), (20), this limit is both $b_0 + (a_1/b_1 + [a_2/b_2 + \dots])$ and $b_0 - F_{l+1}/F_l$ which means that

$$\frac{F_{l+1}}{F_l} = \frac{-a_1}{b_1} + \frac{a_2}{b_2 + \dots} \quad (21)$$

with a_n, b_n defined as above. This continued fraction for the ratio F_{l+1}/F_l is recast into one for F'_l/F_l with the help of (5), by eliminating F_{l+1} , and we reach the penultimate result

$$\begin{aligned} \frac{F'_l}{F_l} &= S_{l+1} + R_{l+1} \left(\frac{a_1}{b_1} + \frac{a_2}{b_2 + \dots} \right) \\ &+ \frac{H_{l+1}}{K_{l+1}} + \frac{H_{l+2}}{K_{l+2}} + \frac{H_{l+3}}{K_{l+3} + \dots}, \end{aligned} \quad (22)$$

with

$$H_p = -\rho^2(p^2 - 1)(p^2 + \eta^2), \quad p = l+2, l+3, \dots,$$

and

$$K_p = (2p+1)[p(p+1) + \eta\rho], \quad p = l+1, l+2, \dots,$$

with

$$H_{l+1} = -\rho[(l+1)^2 + \eta^2](l+2)/(l+1).$$

To obtain the sign of F_l we observe that as $n \rightarrow \infty$ the value of $B_n \rightarrow \delta G_{n+1+l} = F_l R_{l+1} G_{n+1+l}$ and $G_{n+1+l} \rightarrow \infty$ while remaining positive. The sign of F_l is hence the sign of the n th denominator when the continued fraction has converged.

3.2.1. A minor difficulty with eq. (22)

For the scattering of like charges, the value of η is positive and eq. (22) holds for all $\rho > 0$; however when η is negative, a peculiar difficulty can arise at the first step of the algorithm (section 3.4) in which $D_1 =$

$1.0/K_{p+1}$. A zero of K_{p+1} arises when

$$\eta\rho = -p(p+1) = -(L+1)(L+2)$$

and here $L = \text{MAXL}$ and, if $\rho < \rho_{\text{TP}}$, also $L = \text{MINL}$. Elaborate methods are available to circumvent this difficulty but in our subroutine we choose to avoid this value of ρ by setting

$$\rho = \rho + 10^{-6} \text{ (cards 37 - 39).}$$

The accuracy of the results is reduced for values of ρ which make $F_l(\eta, \rho) = 0.0$ within the value of ACCUR. For the evaluation of integrals over products of F_l values such a reduced accuracy is not a problem.

3.3. A continued fraction for $(G_l' + iF_l')/(G_l + iF_l)$

The boundary conditions at infinity, eq. (3), suggest the substitution

$$G_l + iF_l = y_l \exp(i\theta_l) \quad (23)$$

which leads to the following differential equation for $y_l(\eta, \rho)$:

$$\begin{aligned} d^2 y_l / d\rho^2 + 2i(1 - \eta/\rho) dy_l / d\rho \\ + (i\eta - l)(i\eta + l + 1) y_l = 0. \end{aligned}$$

After the substitution $Z = (2i\rho)^{-1}$ this equation can be recognised as that for the hypergeometric function ${}_2F_0(a, b; Z)$:

$$Z^2 {}_2F_0'' + [(a+b+1)Z - 1] {}_2F_0' + ab {}_2F_0 = 0 \quad (24)$$

with $a = i\eta - l$ and $b = i\eta + l + 1$. Thus the solution of eq. (23) is

$$\begin{aligned} y_l(\eta, \rho) &= {}_2F_0(i\eta - l, i\eta + l + 1; (2i\rho)^{-1}) \\ &\equiv 1 + \frac{ab}{1!(2i\rho)} + \frac{a(a+1)b(b+1)}{2!(2i\rho)^2} + \dots \quad (25) \end{aligned}$$

The series expansion is the definition of the (specialised) hypergeometric function, and the result (25) is a known asymptotic form for $G_l + iF_l$ when ρ is large, quoted for example as eq. 12.16 of ref. [2].

The series (25) however is convergent only in an asymptotic sense, viz. only for very small Z (i.e. large ρ) and whose summation is terminated after the term with the smallest absolute magnitude. Nevertheless the ratio of two such hypergeometric functions has a very strong convergence property, namely that it can be ex-

panded as a continued fraction which is convergent throughout the complex Z -plane exterior to the cut along the real axis from $Z = 1$ to $Z = \infty$, i.e. for all ρ . We can now proceed to derive such a ratio.

Since ${}_2F_0'(a, b; Z) = dF/dZ = ab {}_2F_0(a+1, b+1; Z)$ we find that the complex ratio

$$\begin{aligned} \frac{G_l' + iF_l'}{G_l + iF_l} &= \frac{i(\rho - \eta)}{\rho} \\ &+ i \frac{ab}{2\rho^2} \frac{{}_2F_0(a+1, b+1; (2i\rho)^{-1})}{{}_2F_0(a, b; (2i\rho)^{-1})}. \quad (26) \end{aligned}$$

We derive a continued fraction for the ratio of F 's directly from eq. (24) using the above differentiation property: eq. (24) reads on dividing by ${}_2F_0'$ and rearranging,

$$\begin{aligned} \frac{{}_2F_0(a, b; Z)}{{}_2F_0(a+1, b+1; Z)} &= 1 - (a+b+1)Z \\ &- Z^2(a+1)(b+1) \frac{{}_2F_0(a+2, b+2; Z)}{{}_2F_0(a+1, b+1; Z)}. \end{aligned}$$

Inverting this equation, incrementing a and b by 1 and using the equation again, and repeating, gives us immediately the required continued fraction:

$$\begin{aligned} \frac{{}_2F_0(a+1, b+1; Z)}{{}_2F_0(a, b; Z)} &= \frac{1}{1 - (a+b+1)Z -} \\ &\frac{(a+1)(b+1)Z^2}{1 - (a+b+3)Z -} \frac{(a+2)(b+2)Z^2}{1 - (a+b+5)Z -} \dots \quad (27) \end{aligned}$$

Now $Z = (2i\rho)^{-1}$ and we can make an equivalence transformation with $c_p = 2\rho$ to reach the final result:

$$\begin{aligned} p_l + iq_l &= \frac{G_l' + iF_l'}{G_l + iF_l} \\ &= \frac{1}{\rho} \left((\rho - \eta)i + \frac{i(i\eta - l)(i\eta + l + 1)}{2(\rho - \eta + i) +} \right. \\ &\quad \left. \frac{(i\eta - l + 1)(i\eta + l + 2)}{2(\rho - \eta + 2i) + \dots} \right) \quad (28) \end{aligned}$$

The continued-fraction (28) is programmed [see eqs. (32)]:

$$\begin{aligned} a_1 &= -\eta^2 - l(l+1) + i\eta, & a_{k+1} - a_k &= 2k + 2i\eta, \\ b_1 &= 2(\rho - \eta) + 2i, & b_{k+1} - b_k &= 2i, \end{aligned}$$

and

$$(\Delta p)_1 + i(\Delta q)_1 = ia_1/b_1.$$

We note in passing that the result for $(G'_l - iF'_l)/(G_l - iF_l)$ is exactly $p_l - iq_l$; its convergence takes the same number of iterations.

We stress that (28) can only be evaluated in practice for

$$\rho \geq \rho_{TP} = \eta + [\eta^2 + l(l+1)]^{1/2}$$

for otherwise it tends too rapidly towards G'_l/G_l . A further point is of academic interest. If we take one of the Gauss relations [8]

$$aZ \frac{F(a+1, b+1; Z)}{F(a, b; Z)} = \frac{F(a, b+1; Z)}{F(a, b; Z)} - 1$$

we can re-express eq. (26) in terms of the ratio on the right-hand side, which itself can be developed as the *continued-fraction of Gauss*

$$\frac{F(a, b+1; Z)}{F(a, b; Z)} = \frac{1}{1-} \frac{aZ}{1-} \frac{(b+1)Z}{1-} \frac{(a+1)Z}{1-} \frac{(b+2)Z}{1-} \dots$$

This version of eq. (26) converges only half as rapidly as does eq. (28).

3.4. Steed's algorithm for evaluating continued fractions

Several methods are available to compute the continued fraction

$$h_n \equiv \frac{A_n}{B_n} = b_0 + \frac{a_1}{b_1 +} \dots \frac{a_n}{b_n}$$

and hence to find its limit. One is to calculate the sequences A_n, B_n by their two-term recurrence relations, eq. (11), and to divide the results. A second is to use the algorithm

$$A_1/B_1 = b_0 + a_1/b_1, \quad \Delta_1 = a_1,$$

$$B_0 = 1, \quad B_1 = b_1.$$

$$\Delta_k = -a_k \Delta_{k-1}, \quad B_k = b_k B_{k-1} + a_k B_{k-2}.$$

$$\frac{A_k}{B_k} = \frac{A_{k-1}}{B_{k-1}} + \frac{\Delta_k}{B_k B_{k-1}} \text{ for } k = 1, 2, 3, \dots$$

A significant amount of labour can be bypassed by the use of an algorithm due to Steed [9] which we now establish. If we define the quantity D_n to be the ratio of denominators B_{n-1}/B_n then the relation

$$B_n = B_{n-1} b_n + B_{n-2} a_n$$

is equivalent to the relation

$$D_n = 1/(b_n + a_n D_{n-1}). \quad (29)$$

The identical expression results from A_{n-1}/A_n but the two sequences are different because of their different starting values, eq. (11). Now

$$\begin{aligned} h_n &\equiv A_n/B_n \\ &= \frac{A_{n-1} b_n + A_{n-2} a_n}{B_{n-1} b_n + B_{n-2} a_n} \\ &= \frac{(A_{n-1}/B_{n-1}) b_n + (A_{n-2} B_{n-2}/B_{n-2} B_{n-1}) a_n}{b_n + (B_{n-2}/B_{n-1}) a_n} \\ &= (h_{n-1} b_n + h_{n-2} D_{n-1}) D_n. \end{aligned} \quad (30)$$

Furthermore, from eq. (29), $b_n D_n + a_n D_n D_{n-1} = 1$ which then yields for the difference between h_n and h_{n-1} the recurrence relation

$$\begin{aligned} \Delta h_n &\equiv h_n - h_{n-1} \\ &= (b_n D_n - 1)(h_{n-1} - h_{n-2}) \\ &= (b_n D_n - 1) \Delta h_{n-1}. \end{aligned} \quad (31)$$

The expression (31) forms the basis of *Steed's algorithm* as follows

$$\begin{aligned} h_0 &= b_0, & D_1 &= 1/b_1, \\ \Delta h_1 &= a_1 D_1, & D_n &= 1/(D_{n-1} a_n + b_n), \\ \Delta h_n &= (b_n D_n - 1) \Delta h_{n-1}, & h_n &= h_{n-1} + \Delta h_n, \end{aligned} \quad (32)$$

which can be iterated until $\Delta h_n/h_n$ is sufficiently small. This method is used in subroutine RCWFN.

4. Notes on the program

The calling sequence for the subroutine is CALL RCWFN (RHO, ETA, MINL, MAXL, FC, FCP, GC, GCP, ACCUR, STEP) where the arguments have the following type and meaning:

RHO	real, $\rho > 0$. Accuracy reduced if ρ such that $\rho < 0.2 \rho_{TP}$.
ETA	real, unlimited, RCWFN has been tested for $-1000 \leq \eta \leq 1000$.
MINL	integer
MAXL	integer
Coulomb functions are calculated for $\text{MINL} \leq l \leq \text{MAXL}$ with $\text{MINL} \geq 0$	
FC, FCP	real arrays of dimension at least $\text{MAXL} - \text{MINL} + 1$ for F_l, F'_l, G_l, G'_l .
GC, GCP	Normally the dimension will be $\geq \text{MAXL} + 1$.
ACCUR	real, accuracy parameter of the continued fractions. If $\text{ACCUR} < 10^{-15}$ or $\text{ACCUR} > 10^{-6}$ it is reset to 10^{-6} within RCWFN. A suitable value on the CDC 7600 is 10^{-10} .
STEP	real, number of steps in Runge-Kutta integration below the turning point. A suitable value is 100.0 and RCWFN chooses this as a minimum value, which it increases to 1000.0 when necessary.

The subroutine sets up the variables:

R	the desired RHO value
TURN	$\eta + (\eta^2 + \text{MINL}(\text{MINL} + 1))^{1/2}$, the turning point for MINL
KTR	control variables = 1 if $R \geq \text{TURN}$ and -1 if $R < \text{TURN}$
KTRP	= -1 if a second evaluation of F'_l/F_l at $R = \text{TURN}$ is required,

and then begins the evaluation of $F'_{\text{MAXL}}/F_{\text{MAXL}}$. This is done in terms of two ratios

$$x_l = F_l/F_L \quad \text{and} \quad x'_l = F'_l/F'_L$$

for $L = \text{MAXL}$. Starting values for these ratios are $x_L = 1$ and $x'_L = f_L$ [which is evaluated by the continued fraction eq. (22)]. Both x_l and x'_l satisfy the downward recurrence relations, eq. (4) and eq. (5), which are used to calculate (and store) values of x_l and x'_l until the value $l = m = \text{MINL}$ is reached. This procedure gives us all the values of F_l and F'_l between $l = \text{MAXL}$, MINL to within a common multiplier and this is determined from the wronskian at $l = m = \text{MINL}$. The ratio x'_m/x_m gives f_m , and the continued fraction eq. (28) for $(G'_m + iF'_m)/(G_m + iF_m)$ is used to evaluate $p_m + iq_m$, and hence the complete solution for the four Coulomb functions at $l = \text{MINL}$. The values of G_l and G'_l are then obtained by upward recurrence to $l = \text{MAXL}$, and the values of

F_l, F'_l result by renormalising the stored values of x_l, x'_l by the quantity $(x'_m G_m - x_m G'_m)^{-1/2}$.

The modification of this method of solution if $\rho < \text{TURN}$ arises after the x_l, x'_l have been obtained. The subroutine sets $\rho = \text{TURN}$ and $l = \text{MINL} = m$ and proceeds to evaluate F_m, F'_m, G_m and G'_m . These are functions at the turning point, and now the differential equations for $G_m(\rho)$ and $G'_m(\rho)$ are integrated inwards, simultaneously, until the required ρ -value is reached. The rest of the calculation proceeds as before. The 4th-order Runge-Kutta technique [10] is employed for the integration. We found this method to be satisfactory over 6–10 orders of magnitude ($G, G' \sim 10^6 - 10^{10}$) in that 6D accuracy was maintained in the final results with 1000 steps in the integration. In extreme cases of say 50–100 orders of magnitude only 1D or 2D survived. This loss of accuracy is not merely a function of the step length but also of the truncation errors in the 4th-order algorithm. Consequently this method of extrapolation towards small ρ eventually cannot compete with approximations designed for small ρ [1, 2, 4]. The use of the logarithmic variables defined by Froberg (eqs. 6.1, 6.2) seems a more efficient procedure but in practice it offered no advantages.

The subroutine has been extensively tested on the CDC 7600 at the University of Manchester Regional Computing Centre, on the IBM 370/165 at the Daresbury Laboratory, and on the IBM 360/195 at the Rutherford High Energy Laboratory. For these latter two machines, and others of smaller word length, if accuracies exceeding 10^{-6} are desired it is important to run the routine in double word length, including the arrays FC, FCP, GC, GCP. Suitable declaration statements for IBM machines are:

```

IMPLICIT REAL*8 (A-H, O-Z)
REAL*8 K, K1, K2, K3, K4, M1, M2, M3, M4
FLOAT(I) = DFLOAT(I)
EXP(X) = DEXP(X)
ALOG(X) = DLOG(X)
ABS(X) = DABS(X)
IFIX(X) = X
SQRT(X) = DSQRT(X)

```

Ensure that card 156 which defines R3 reads

```
R3 = 1.0/3.0D0
```

(as it does in our listing).

5. Test deck and accuracy checks

The Coulomb wavefunction test program CFNCHK is a short routine to call RCWFN for specified arguments and to compute two test quantities. These are the wronskian

$$W = F_l' G_l - F_l G_l' = 1 \quad (8)$$

and the relation

$$W_l = F_l G_{l+1} - F_{l+1} G_l = 1/R_{l+1}. \quad (9)$$

Two input cards are required for each case:

CARD 1

ETA, N, (L(I), I = 1, 10) FORMAT (F10.3, 11I5)

CARD 2

RHO(I), I = 1, N FORMAT (7F10.3)

where $N \leq 7$ is the number of ρ -values to be found on CARD 2, and $N = 0$ calls EXIT. If N is made negative a heading on a new page is produced, and $(-N)$ ρ -values are sought. Repetitions of CARD 2 can be used for 14 or 21 ρ -values. The array L contains up to 10 l -values which will be calculated for each ρ -value. The program outputs the input data, the values of F_l, F_l', G_l, G_l' and the quantities

$$\text{TEST 1} = W - 1 \quad \text{and} \quad \text{TEST 2} = W_l - 1/R_{l+1}$$

and both are expected to be zero.

The value of TEST 1 is not an independent test since the Wronskian relation is used in the subroutine. It can thus be regarded as a check on the arithmetic and it is usually less than 2×10^{-14} with single-precision arithmetic on the CDC 7600. The quantity TEST 2, on the other hand, is a test of the results since it is evaluated after two separate calls to RCWFN. The first has MINL = 0, MAXL = l while the second has MINL = 1, MAXL = $l + 1$ and this means that evaluations of the continued fractions for different l -values are called for. Since the definition of TEST 2 embodies all the properties of the Coulomb functions which determine their solution it does not follow that if TEST 2 is 10^{-12} , for example, that the wavefunctions themselves have this accuracy, or indeed even a related accuracy. Nevertheless our experience does suggest that its value gives a reasonable guide to the errors. We display in the test run output the test cases of Bardin et al. [4] for which $\rho > 0.2\rho_{\text{TP}}$ and some others. The results, given to 12D, are ob-

tained with ACCUR = 10^{-14} and STEP = 999.0; the column labelled W-1 contains TEST 1 and that headed TEST contains TEST 2, as described above. The agreement in most cases is excellent having regard to the values of TEST and of the approximate errors discussed in [4], but a systematic discrepancy is evident in those cases where the asymptotic Riccati method (no. 8 of [4]) is used for the lowest values of ρ . The values of F_l, G_l agree to 9D-12D while those for F_l', G_l' in ref. [4] are seriously different from our results. An example is for $\eta = -50.0, \rho = 5.0$ where F_0, G_0 agree to 12D whereas F_0' differs at 7D and G_0' at 5D; the expected accuracy claimed by Bardin et al. is about 10D. This behaviour is illustrated in table 1 and it indicates presumably that the boundary line for the use of the asymptotic Riccati method was drawn somewhat too low in ρ -value. Since the two methods of solution agree for F_l, G_l , a numerical differentiation can be used to obtain the derivative independently, and this agrees with the values provided by RCWFN. It is clearly an extremely difficult task to assess the errors of all such high-accuracy calculations (except in special cases as above) and beyond about 8D it is often unclear which program is checking which.

In practice, however, such precision is irrelevant. The same test deck calculated with ACCUR = 10^{-7} and STEP = 100.0 produced results which are in agreement with those of Bardin et al. to the 6th digit (except, of course, for the cases in table 1), and very similar results were obtained on the IBM machines. For normal nuclear-reaction calculations such precision is adequate. Further comments on the range and limitations of RCWFN will be found in the last section.

6. Limitations of the method

Arising out of the considerations of accuracy in the previous section, it is natural to investigate the limits of the present approach. We will restrict ourselves to a few observations derived from extensive use and development of the subroutine. A mathematical study will appear later. For all values of $\rho \geq \rho_{\text{TP}}$ the algorithms are only limited by the word-size used and by the round-off errors inherent in calculating the continued fractions until convergence. A convenient feature of the algorithms is that relatively few additional iterations are necessary to obtain results to a considerably higher

Table 1
Comparison of continued fraction method with some asymptotic Riccati results. The first row is from the test run output; the second (with the same power of 10) from ref. [4].

η	ρ	$F_0(\eta, \rho)$	$F'_0(\eta, \rho)$	$G_0(\eta, \rho)$	$G'_0(\eta, \rho)$	TEST
-500	0.61	-1.56747 426448 E-1 -1.56747 426448	-5.16358 838507 E-1 -5.16398 330790	-1.11662 577817 E-2 -1.11662 577815	6.34290 607911 E+0 6.34290 326579	1E-14 4E-13
-200	0.155	9.26335 589426 E-2 9.26335 591744	-5.20338 827163 E+0 -5.20248 460204	-1.05498 227730 E-1 -1.05498 227893	-4.86920 468457 E+0 -4.87023 387696	2E-15 4E-9
-50	5	1.52236 975714 E-1 1.52236 975714	2.03091 041166 E+0 2.03091 614165	4.41680 690236 E-1 4.41680 690236	-6.76485 374767 E-1 -6.76468 750507	1E-14 9E-13
-50	50	5.81412 475321 E-1 5.81412 475320	8.49247 350399 E-1 8.49247 565739	4.89200 660675 E-1 4.89200 660676	-1.00539 231599 E+0 -1.00539 213480	-4E-13 -4E-14
-10	3.1	1.92806 827635 E-1 1.92806 828167	-1.55175 356287 E+0 -1.55165 77983 ^a	-5.74163 450399 E-1 -5.74163 451255	-5.65539 205794 E-1 -5.65824 394424	-8E-15 3E-9
34	116.4	-1.22893 716648 E+0 -1.22893 716643	1.33208 933613 E-1 1.33217 506068	2.00811 761092 E-1 2.00811 761083	7.91944 540370 E-1 7.91943 139574	-3E-12 -8E-11
50	153.3	-1.29566 554688 E+0 -1.29566 554683	-7.28583 466960 E-2 -7.28495 928556	-1.30262 077773 E-1 -1.30262 077755	7.64479 168841 E-1 7.64480 048891	-2E-12 -8E-11
100	1000	-1.65581 311974 E-1 -1.65581 311968	-9.34063 273097 E-1 -9.34063 273043	-1.04432 596295 E+0 -1.04432 596295	1.48165 711328 E-1 1.48165 711669	-1E-13 -4E-14
200	476	4.20046 440799 E-1 4.20046 440740	-6.10636 723020 E-1 -6.10639 246490	-1.52505 808348 E+0 -1.52505 808337	-1.63656 974114 E-1 -1.63647 811979	6E-13 -2E-10
500	1095	1.46929 027367 E+0 1.46929 027350	3.23944 539005 E-1 3.23935 776975	1.11156 684610 E+0 1.11156 684581	-4.35525 914746 E-1 -4.35532 543385	-2E-12 -4E-10

Table 2
A sample calculation to 18D of Coulomb functions with $\rho = 2\eta$ and $L = 0$ [11]. The parameter ACCUR was set at 10^{-21} in the double precision version of RCWFN.

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - RCWFN									
ETA	RHO	L	F	FP	G	GP			
2.5	5.0	0	348512502280354298	725140521049396111	1.148208486965474131	4802921226691679932			
5.0	1.0	0	516601500314181575	592924553360720776	1.197486970798485546	561323515538948353			
7.5	1.5	0	6055419890014514342	5232903399603027	1.237532738992948680	580796804239306086			
1.00	2.0	0	661781613832681298	4617574857099492083	1.275778784768276589	582728813097184743			
1.25	2.5	0	700411103046531477	453547031343177808	1.310604085222607841	579059078471091182			
1.50	3.0	0	730129112311153282	433000424831788472	1.34290560734846743	573580220114165108			
1.75	3.5	0	754460672412615498	416997407026594634	1.371121194810441860	567618241047685957			
2.00	4.0	0	775197191073771120	404007816494934739	1.3974834267930118884	561670987933261970			
2.25	4.5	0	793351758259822650	393147580377034250	1.421741252870013592	555927266629963202			
2.50	5.0	0	809551965068304961	383844008546991053	1.444202721094029777	55045577490027672			
2.75	5.5	0	824215134429597640	375789975554407382	1.465120421448198788	54527310668750535			
3.00	6.0	0	837634054540468182	3686699598817686	1.48470030304953169	54037391489506708			
3.25	6.5	0	85022806838043657	362319457586325567	1.503111000288887515	535742845976427728			
3.50	7.0	0	861542986950039350	356601224972436893	1.520491699186036555	53136183241547779			
3.75	7.5	0	872319866950039350	3514106326223010	1.536582728808690755	527210876264827421			
4.00	8.0	0	882452654489153772	3466624232351198408	1.552608151535213074	52372454238079663			
4.25	8.5	0	892021351143997427	342303533864408391	1.56752388948435348	519529081432638804			
4.50	9.0	0	901091483333817700	33827058589988515	1.581776045432882446	515965008142906423			
4.75	9.5	0	909717444651540790	334525096436581679	1.59542537450763258	512565935144917499			
5.00	10.0	0	917944918945897700	331032101931063326	1.608524555509838509	509318942457823266			
5.25	10.5	0	925812672729550718	327762477354522733	1.621119598954193636	506212380834634570			
5.50	11.0	0	933535905198317896	324891602217860823	1.63325073869094811	503235751484668172			
5.75	11.5	0	940597281428529682	321798915631097252	1.64400846487446741	500379586793187456			
6.00	12.0	0	94756773446372403	319066320910719973	1.656260078075885120	49763538641548594			
6.25	12.5	0	954287095934420434	31647854490719144	1.66719968651929902	49495276216347290			
6.50	13.0	0	960774597898490051	314022260088792310	1.6779267734565351	492452394697346902			
6.75	13.5	0	967047275962114927	311685840084069795	1.688066777506918472	490000334123593851			
7.00	14.0	0	973120296666177565	30945908762891084	1.69804073799436540	487816561485292902			
7.25	14.5	0	979007225442517074	299888471097679328	1.707233096439878743	485346039550612073			
7.50	15.0	0	984720247788432251	305299612001220344	1.71716056829029703	483133708102033583			
7.75	15.5	0	990270353176653246	303351841401840432	1.726338564312094101	48099184893439711			
8.00	16.0	0	995667489034255634	301483333419602548	1.735281048937105334	478916561485292902			
8.25	16.5	0	1.00092069048678669	299888471097679328	1.74400846487446741	476903971773195223			
8.50	17.0	0	1.006038190332746411	297962083677327411	1.75250973800118190	474950899311727625			
8.75	17.5	0	1.011027512778279866	296299597978959900	1.760818572397683218	473053578115389050			
9.00	18.0	0	1.01589553744867882	294596799957335341	1.768937363188798799	471209680940693404			
9.25	18.5	0	1.02048650007829034	293149929716171825	1.776875372801135853	469416296533851127			
9.50	19.0	0	1.02529263899764838	291855523220975984	1.784641196785169267	467670909436361035			
9.75	19.5	0	1.029832910700812731	290210426589853516	1.792242778515140158	465971182031829076			
10.00	20.0	0	1.034274453020203063	288811752905226998	1.7996875629697511865	464314938552149343			
11.00	22.0	0	1.051141102578809651	283631381712959554	1.828030344913430501	458086194401178475			
12.00	24.0	0	1.066745825490252472	279015514828006492	1.854343697745818603	454212880675066787			
13.00	26.0	0	1.081280026096515860	274861048424339390	1.878921290712321990	447208597641140071			
14.00	28.0	0	1.094893532080730428	271090091115407894	1.901996020321938888	44240227263446124			
15.00	30.0	0	1.107706056074175111	267542635756946072	1.923755848780808422	437951244059142786			
16.00	32.0	0	1.1198146567822311	264470773751189782	1.944354775729600062	433799440113401204			
17.00	34.0	0	1.131299508700884521	261537079960348111	1.963320632728151733	429914385544036685			
18.00	36.0	0	1.142227261840811469	258810565539925006	1.98260754520918599	426265723268157505			
19.00	38.0	0	1.15265426753700929	256265866101158525	2.000366183406406425	422827903563157789			
20.00	40.0	0	1.162628485736277329	253881940579569691	2.017414842373797009	41957929738988896			
21.00	42.0	0	1.172191118653625588	25164107969413620	2.03377397174558210	416501127910900059			
22.00	44.0	0	1.18137789682654270	24932823643018539	2.04950196475792945	413794683862956136			
23.00	46.0	0	1.190219752864032480	24732059269956542	2.064650039335835626	410794683862956136			
24.00	48.0	0	1.198744151244637430	245637118813482713	2.079263159118716626	408140292347680451			
25.00	50.0	0	1.206975123287461823	24383824337522006	2.093381125553251018	405603739682524577			

order of accuracy. As an example we give in table 2 a short list of $l = 0$ wavefunctions to 18D for selected values of $\rho = 2\eta$ which can be compared with the calculations of Abramowitz and Rabinowitz [11] in 1954. The agreement is complete to within 1 unit in the seventh decimal place, as was indeed claimed by them. A double-precision version of the subroutine was used for table 2 and each derivative was checked by numerical differentiation. The value of the parameter ACCUR when $\rho \geq \rho_{TP}$ approximately determines the final accuracy of the results; it was taken as 10^{-21} .

For large values of ρ the limitation on the efficiency of RCWFN is the calculation of $f_l = F'_l/F_l$ which requires about $\rho + 50$ iterations. Thus a practical upper limit to ρ is perhaps 20 000. For small ρ , say 0.001, the limit is in the calculation of

$$p_l + iq_l = (G'_l + iF'_l)/(G_l + iF_l)$$

which may require 50 000 iterations. Clearly these are inefficient methods if many results are required. The calculation in RCWFN will cease if more than 20 000 iterations for f_l , or 23 000 for $p_l + iq_l$, are required or if $|G'_l|$ approaches the computer word size during the integration, and zeros are returned for the functions.

In calculations involving high partial waves, $l \sim 500$ –1000, errors may propagate by successive iterations from the $l = 0$ wavefunctions. We have run some checks on this point, examples of which appear at the end of the test run output. For an η -value of 50.0 and a ρ -value of 600.0, which exceeds the turning point for $l = 500$ ($\rho_{TP} = 552.99$), we evaluate the Coulomb functions for MINL = 0, 100, 200, 300, 400, 500 and MAXL = 500 and hence show that rounding errors are of minor importance. Similar conclusions are evident for $\eta = 50.0$, $\rho = 1200.0$, MINL = 0(200) 1000 and MAXL = 1000. It is to be noted that in our method we do not calculate the $l = 0$ (or, in general, $l = \text{MINL}$) wavefunctions directly but combine the calculation of e.g. F'_{1000}/F_{1000} , after downward recurrence to F'_0/F_0 with the value of $(G'_0 + iF'_0)/(G_0 + iF_0)$. The method thus differs from the classical one [1–4].

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Note added in proof

A complete review of recent methods of computing Coulomb wavefunctions is that of Kölbig [12], which has only just come to our attention.

TEST RUN OUTPUT

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - RCMFN

ETA	RHO	L	F	PP	G	GP	M-1	TEST
50.0	120.000	0	2.00259933341E-01	6.32274340306E-01	1.55059921847E+00	-9.7A53412082E+02	-3.E-15	-9.E-14
50.0	80.000	0	1.20366249591E-03	1.9770431959F-04	A.32243179902E+02	-4.0227246887E+02	-3.E-15	-9.E-11
9.0	50.000	0	9.35708567796E-01	-4.92126155892E-01	-6.11601989096E-01	-7.46937949161E-01	-4.E-15	5.E-12
5.0	20.000	0	-2.29347373002E-01	8.26679430870F-01	1.16571667134E+00	1.482323529679E-01	-1.E-14	-3.E-13
4.0	10.000	0	1.39920928925E+00	7.39287733242E-02	3.56558327881E-01	-6.95850221937E-01	-1.E-14	-1.E-13
-1000.0	1.000	0	8.68222518991E-02	5.12063399274E+00	1.139339784380E-01	-4.30243486522E+00	-4.E-15	9.E-15
-1000.0	10.000	0	-1.1532232817E-01	3.38892408671E+00	2.39239735828E-01	1.64091977791E+00	0.	-4.E-14
-1000.0	100.000	0	2.62318102583E-01	-1.77068805220F-00	3.8653227941E-01	-1.20301232363F+00	-3.E-15	-3.E-13
-500.0	1000.000	0	5.84941362837E-01	8.40086562103F-01	4.84967926715E-01	-1.01306729907E+00	-1.E-14	8.E-15
-500.0	.590	0	9.83996748953E-02	-5.01886992919E+00	-1.20880038910E-01	3.99878108841E+00	-7.E-15	1.E-14
-500.0	.610	0	-1.5674742644E-01	5.16358838507E-01	1.11662578781E-02	6.34290607911E+00	-7.E-15	1.E-14
-500.0	100.000	0	5.06441621224E-01	7.04935480214E-01	2.12199193856E-01	-1.07919306735E+00	-7.E-15	1.E-14
-500.0	1000.000	0	8.33639868856E-01	1.55997627391F-01	1.10233300886E-01	-6.91861702588E+00	-1.E-14	9.E-11
-200.0	.145	0	1.30429892607E-01	-2.15329440026E+00	-4.53284624318E-02	-6.91861702588E+00	-1.E-14	1.E-11
-200.0	.155	0	9.2633589426E-02	5.20338927163E+00	1.05498227730E-01	-8.8692046645E+00	-9.E-15	2.E-15
-200.0	100.000	0	3.86835282992E-01	1.22055184960F+00	5.45502478229E-01	-4.63857260489E-01	-5.E-15	-1.E-13
-200.0	1000.000	0	3.56540148972E-01	1.00264538165E+00	A.47368822143E-01	-4.2183449477E-01	-5.E-15	-3.E-12
-50.0	5.000	0	1.52336975714E-01	2.03091041166F+00	4.1680690236E-01	8.76485374767E-01	-1.E-14	1.E-14
-50.0	5.000	1	1.92207798328E-01	-1.93857353492E+00	-4.26299008702E-01	1.19942377798E+00	-5.E-15	-2.E-14
-50.0	5.000	2	2.67261843675E-01	1.76660377667E+00	3.64620081792E-01	1.62987457269E+00	4.E-15	1.E-14
-50.0	5.000	3	-3.62923967642E-01	1.36860370000F+00	2.9846378931E-01	1.62987457269E+00	2.E-14	1.E-14
-50.0	5.000	5	-4.69406488262E-01	2.75219646325E-01	6.65491186260E-02	2.0913310737E+00	4.E-15	1.E-14
-50.0	5.000	10	-3.68114360219E-01	1.33848751032E+00	3.31586324611E-01	1.51088862814E+00	1.E-14	1.E-14
-50.0	5.000	20	5.72508610339E-01	-8.8809281094F-01	3.31586324611E-01	1.51088862814E+00	1.E-14	1.E-14
-50.0	5.000	30	1.17551067129E-01	4.92521055485E-02	1.05772803735E-03	-4.07522179361E+03	-5.E-14	-3.E-14
-50.0	5.000	50	4.20266913301E-23	3.83028028856E-22	1.32186136031E+21	-1.17470506213E+22	4.E-14	1.E-14
-50.0	5.000	100	1.86693584349F-91	3.67291554441E-90	1.36847205283E+89	-2.66410746916E+90	1.E-14	-4.E-14
-50.0	50.000	0	6.81412475321E-01	8.49247350399E-01	4.89200660678E-01	1.00339231599E+00	0.	-4.E-13
-50.0	50.000	1	-5.98045260245E-01	-8.13864797554E-01	4.68803382789E-01	1.03413147948E+00	0.	-4.E-13
-50.0	50.000	2	6.29138669000E-01	7.40243490434E-01	4.26347094843E-01	1.05848414109E+00	4.E-14	3.E-13
-50.0	50.000	3	-6.70084506913E-01	-6.23348321659F-01	3.58896812233E-01	1.15848414109E+00	1.E-13	-6.E-13
-50.0	50.000	5	7.48260826432E-01	-2.34353261214E-01	1.36481036469E-01	1.29296330882E+00	6.E-13	-3.E-13
-50.0	50.000	10	2.77129451057E-01	-1.22068610371E+00	-7.10519441453E-01	1.29296330882E+00	1.E-13	-6.E-13
-50.0	50.000	20	7.58876294799E-01	-2.52368602573E-01	-1.51281009646E-01	1.27246037372E+01	9.E-14	6.E-12
-50.0	50.000	30	-5.86392849287E-01	2.34026371225E-04	2.19814428489E+03	1.17874820130E+00	-4.E-14	7.E-13
-50.0	50.000	50	-8.37683933712E-01	1.3309920321F-01	9.45277610412E-02	1.17874820130E+00	-3.E-13	-2.E-13
-50.0	50.000	100	2.23304593664E-04	1.3309920321F-01	9.45277610412E-02	1.17874820130E+00	2.E-13	-2.E-13
-10.0	2.900	0	4.58985850692E-01	-1.03688988335E-01	-3.81736647052E-01	-1.31650857017E+00	3.E-14	3.E-14
-10.0	3.100	0	1.92806827635E-01	-1.56175356287E+00	-5.81736647052E-01	-1.31650857017E+00	1.E-14	3.E-14
-10.0	100.000	0	9.45619474305E-01	1.56175356287E+00	-5.81736647052E-01	-1.31650857017E+00	1.E-14	3.E-14
-6.0	6.900	0	3.85247881197E-01	1.50103536499E-01	1.36665935802E-01	-1.053814073796E+00	7.E-15	2.E-13
-6.0	7.100	0	1.57079204316E-01	9.26590385003E-01	9.17459327551E-01	-3.89076788674E-01	-9.E-15	2.E-12
-6.0	17.400	0	3.66765336014E-01	1.26309945681E+00	7.61568011342E-01	-2.42321740803E-01	-3.E-15	2.E-12
-6.0	100.000	0	4.9928090574E-01	1.11101020211E+00	6.72881591082E-01	-6.35323383974E-01	-1.E-14	0.
-6.0	17.600	0	1.96558034259E-01	-1.0059660393E-01	-7.75838032295E-01	-6.35323383974E-01	-4.E-15	2.E-14
-6.0	100.000	0	6.83030680099E-01	1.10869089826E+00	8.5506713933E-01	-2.59791041854E-01	3.E-15	2.E-14
-6.0	5.000	0	-2.31316696922E-01	7.32156892237E-01	6.91651921576E-01	-2.26664871390E-01	-1.E-14	9.E-13
-4.0	5.000	0	4.07842723006E-01	1.09881233363E+00	-9.6981753563E-01	2.32697882573E-01	-4.E-15	2.E-14
-4.0	5.000	1	-6.42286329143E-01	-7.58912993040E-01	6.74327035383E-01	-1.36110427280E-01	-2.E-14	4.E-15
-4.0	5.000	5	2.83637834139E-01	-1.04544995379F-01	-6.7204210149E-01	1.00494559453E+00	-2.E-14	1.E-13
-4.0	5.000	10	7.21985521560E-02	1.10084333023E-01	-8.70189296614E-01	-1.8202528554E-01	-2.E-14	3.E-14
-4.0	5.000	20	3.72990728261E-09	1.4783504452E-08	3.55670502218E+07	-1.30042852498E+08	-5.E-15	-6.E-14
-4.0	5.000	30	4.60279973470E-19	2.75619476973E-18	3.55670502218E+07	-1.30042852498E+08	-5.E-15	-6.E-14
-4.0	5.000	50	4.4063085555E-43	4.43813906512E-42	1.13810237178E+41	-1.12314953439E+42	-4.E-14	-3.E-14
-4.0	5.000	100	1.12267597835E-116	2.26058262062E-115	2.222990834565E+114	-4.43131600080E+115	-7.E-14	-6.E-14

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - RCMFN

ETA	RM0	L	F	FP	G	GP	N=1	TEST
-4.0	50,000	0	-2.77250198046E-01	-9.94100934435E-01	-9.22332107813E-01	2.97966188586E-01	-3.E-15	6.E-14
-4.0	50,000	1	5.08732230779E-01	8.16224222333E-01	8.18834394187E-01	-5.47187525860E-01	2.E-15	-5.E-14
-4.0	50,000	2	-6.01949969000E-01	-5.09226263557E-01	4.7273239043E-01	9.03657029357E-01	-4.E-15	-5.E-14
-4.0	50,000	3	9.4744263766E-01	-1.9391341872E-01	-1.80993170988E-01	-1.01842503613E+00	0.E-16	-1.E-13
-4.0	50,000	4	-4.0046636605E-01	9.21749731522E-01	8.6004867449E-01	4.70973589630E-01	8.E-16	3.E-12
-4.0	50,000	10	4.21084204243E-01	9.26703490800E-01	8.7709381067E-01	-4.44553277267E-01	8.E-16	3.E-12
-4.0	50,000	20	7.67805341241E-01	-6.4050934806E-01	-6.4378734761E-01	7.65326028717E-01	7.E-15	7.E-13
-4.0	50,000	30	1.01812491622E+00	-2.62178848644E-01	-2.99525470598E-01	9.05066502455E-01	7.E-15	3.E-14
-4.0	50,000	50	1.59533986031E+00	-8.64299712691E-02	5.15310112968E-03	6.26548507323E-01	-1.E-14	-2.E-15
-4.0	50,000	100	3.14191623347E-19	5.37489508855E-19	9.37788817810E+17	-1.57855264866E+18	-1.E-14	-2.E-15
-3.3	7,900	0	4.84136409096E-01	-8.78940011038E-01	-8.48600513729E-01	-5.14293320446E-01	-7.E-15	4.E-14
-3.3	7,900	1	3.09016485748E-01	9.65595782501E-01	9.32538631329E-01	3.22133721542E-01	-7.E-15	4.E-14
-3.3	100,000	0	9.38314793063E-01	3.42473282788E-01	3.4143656803E-01	-9.41120298466E-01	-1.E-14	3.E-13
-3.3	100,000	0	2.71583185525E-01	-9.62547844082E-01	9.62599189051E-01	2.71664503970E-01	-1.E-14	3.E-13
0.0	1,000	0	9.9999833333E-04	9.99998000000E-01	9.99998000000E-01	9.9999833333E-04	-1.E-10	1.E-10
0.0	1,000	0	9.99983333417E-03	9.99950000417E-01	9.99950000417E-01	-9.99983333417E-03	-1.E-15	9.E-11
0.0	1,000	0	9.98334166468E-02	9.95004165278E-01	9.95004165278E-01	-9.98334166468E-02	-1.E-14	9.E-15
0.0	1,000	0	8.41470984808E-01	5.40302305868E-01	5.40302305868E-01	-8.41470984808E-01	-1.E-15	2.E-15
0.0	10,000	0	-6.44021110889E-01	8.39071529076E-01	8.39071529076E-01	6.44021110889E-01	-1.E-14	3.E-14
0.0	100,000	0	5.06365641110E-01	8.62318872287E-01	8.62318872287E-01	5.06365641110E-01	-1.E-14	3.E-14
0.0	100,000	0	8.26879540536E-01	5.62379076286E-01	5.62379076286E-01	-8.26879540536E-01	0.E-14	5.E-15
1.0	2,000	0	6.61781613833E-01	4.81537455710E-01	4.81537455710E-01	-6.61781613833E-01	-2.E-14	1.E-12
1.0	7,900	0	1.05570358181E+00	1.8470582985E-01	1.2079638427E+00	9.12103548935E-01	-1.E-15	1.E-13
1.0	8,100	0	-1.00317220511E+00	3.39504302422E-01	3.79306166870E-01	8.6846884820E-01	-1.E-14	8.E-12
1.0	9,100	0	-3.46718446657E-01	8.91613091252E-01	1.00525985795E+00	2.9908237584E-01	9.E-15	2.E-14
1.0	9,200	0	2.36320123198E-01	9.15167922867E-01	1.03120585991E+00	2.10463267175E-01	-1.E-14	8.E-14
1.0	100,000	0	1.57407746175E-01	9.82676202096E-01	9.92660472446E-01	-1.55876438714E-01	1.E-15	8.E-14
1.0	100,000	0	6.01756274039E-01	7.98507100149E-01	7.99306854825E-01	6.01153816434E-01	7.E-15	3.E-12
8.0	1,050	0	9.88270778102E-09	4.00516771647E+08	1.33332774448E+07	-4.71591379833E+07	0.E-16	3.E-08
8.0	2,010	0	2.38960314934E-07	6.66487750754E-07	7.90791806745E+05	-1.97928543055E+06	-5.E-15	9.E-09
8.0	16,000	0	9.95667489034E-01	3.0148335420E-01	1.73528104894E+00	-4.78914561495E-01	-1.E-15	7.E-12
8.0	20,800	0	1.00861490379E+00	-5.31608367267E-01	-1.0129081251E+00	-4.57588870202E-01	-1.E-14	3.E-13
8.0	20,900	0	9.34304912831E-01	-5.54441105976E-01	1.06747192361E+00	-4.33503057118E-01	-1.E-14	3.E-13
8.0	100,000	0	1.02637831266E+00	-1.77302473353E-01	-1.93984775563E-01	9.40789578233E-01	4.E-15	5.E-13
8.0	100,000	0	-1.25742178869E-01	9.86134028826E-01	9.96135632648E-01	1.24728136388E-01	-1.E-14	2.E-13
10.0	5,000	1	1.72074342808E-06	3.09759950467E-06	1.67637564193E+05	-2.79370763893E+05	-2.E-15	7.E-10
10.0	5,000	2	1.43822713899E-06	2.62336071014E-06	1.97940193028E+05	-3.34252891322E+05	-2.E-14	3.E-10
10.0	5,000	3	1.00863819538E-06	1.86650199163E-06	2.75176377328E+05	-4.76651734756E+05	-1.E-14	7.E-10
10.0	5,000	5	5.9780090985E-07	1.15881372654E-06	4.47978845520E+05	-8.0440821794E+05	0.E-15	6.E-10
10.0	5,000	10	3.1826036427E-07	2.80453363740E-07	1.84919797798E+06	-3.65176286902E+06	2.E-14	2.E-10
10.0	5,000	20	3.26278064114E-10	9.21968559846E-10	5.59559146826E+08	-1.47467684346E+09	3.E-14	3.E-13
10.0	5,000	30	4.04283871422E-18	1.83778759676E-17	2.77870711499E+16	-1.21035896966E+17	2.E-14	9.E-10
10.0	5,000	50	3.18942998567E-28	2.05356149068E-27	2.47223679219E+26	-1.343597040281E+27	1.E-14	9.E-10
10.0	5,000	80	2.13937218584E-52	2.21349851593E-51	5.28070929210E+50	-2.31454433640E+51	3.E-14	9.E-10
10.0	5,000	100	4.16264130910E-126	8.43944020525E-125	5.95385803310E+123	-1.19523119358+125	-2.E-14	2.E-14
10.0	50,000	0	2.37518766416E-02	8.79909263353E-01	1.13888213569E+00	-2.21787413752E-02	-2.E-15	2.E-14
10.0	50,000	1	-6.38729759079E-02	8.78710662060E-01	1.13471232589E+00	4.6656784785E-02	-2.E-14	3.E-13
10.0	50,000	2	-2.35462997144E-01	8.61043057098E-01	1.11262436375E+00	1.78293397961E-01	3.E-15	3.E-13
10.0	50,000	3	-4.77011395975E-01	7.99223383111E-01	1.03365112652E+00	3.44523238412E-01	8.E-15	2.E-13
10.0	50,000	8	-1.00346497670E+00	4.21525638376E-01	9.48907634434E-01	7.67475808062E-01	-4.E-15	8.E-13
10.0	50,000	10	4.16672475541E-01	-8.07592601896E-01	-1.08034481290E+00	-3.0604772789E-01	2.E-14	1.E-13
10.0	50,000	20	4.28928799671E-01	-7.64035879316E-01	-1.1350952567E+00	-2.72407317071E-01	5.E-14	4.E-13
10.0	50,000	30	-3.52635436429E-01	-6.64436751932E-01	-1.39731750175E+00	2.0260027088E-01	-5.E-14	1.E-12
10.0	50,000	50	3.76729101128E-03	3.6965549530E-03	1.34226818062E+02	-8.27249102369E+01	1.E-14	-1.E-14
10.0	50,000	100	7.69173185651E-26	1.43607614052E-25	3.50484330733E+24	-6.45729516259E+24	-0.E-14	-9.E-14

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - RCHFN

ETA	RHO	L	F	FP	G	GP	M=1	TEST
26.0	24.300	0	2.97292502747E+08	3.23118014790E+08	1.57544174746E+07	-1.65139172688E+07	4.E-15	-4.E-10
26.0	24.500	0	3.69173028846E+08	3.98202039842E+08	1.27653537642E+07	-1.32968701118E+07	0.	-4.E-10
26.0	50.700	0	9.07720503067E+01	2.25471559616E+01	2.64080995466E+00	-4.43893246791E-01	2.E-15	5.E-12
26.0	50.900	0	9.53263074069E+01	2.29864424951E+01	2.56058413409E+00	-4.31563169569E-01	0.	1.E-11
26.0	52.000	0	1.214933404198E+00	2.42125624592E+01	2.10703938758E+00	-7.03175600979E-01	-2.E-14	-2.E-11
26.0	100.000	0	6.33002441304E+01	7.0575980379F+01	1.02108209693E+00	-7.41360804438E-01	-7.E-15	-4.E-13
26.0	100.000	0	1.00121809251E+00	1.52767131867E+01	1.56915115126E+01	-9.74841081281E-01	-7.E-15	-8.E-12
34.0	40.800	0	5.85629750847E+07	4.59552376883E+07	1.02663527256E+06	-3.3785057753E+05	-2.E-15	-1.E-10
34.0	40.500	0	7.54363068649E+07	6.32953581764E+07	8.04606710722E+08	-3.50810769855E+05	-7.E-15	-2.E-10
34.0	67.900	0	1.24766281927E+00	2.30713705592E+01	2.24165682521E+00	-3.86978789245E-01	-5.E-15	-3.E-12
34.0	68.000	0	1.27074034088E+00	2.30806101527E+01	2.20298989948E+00	-7.86814748778E-01	-2.E-15	-2.E-11
34.0	68.100	0	1.29381781050E+00	2.30710261777E+01	2.16428307650E+00	-7.86974679008E-01	-1.E-14	-2.E-12
34.0	100.000	0	1.23024654108E+00	2.78457249564E+01	5.03670986464E-01	-6.98865617145E-01	-4.E-15	2.E-12
34.0	116.200	0	1.24830758487E+00	3.03100736723E+02	4.11948933130E+02	8.02011806516E-01	4.E-15	2.E-12
34.0	116.400	0	1.22893716648E+00	1.33208933613E+01	2.00811761020E+01	7.91944540370E-01	4.E-15	-3.E-12
34.0	100.000	0	5.26610693682E+01	8.40807046862E+01	8.70930332585E+01	5.08374862631E-01	-7.E-15	-4.E-12
50.0	67.800	0	3.4872031775E+07	4.46639955061E+07	1.13796294833E+06	-7.74550232622E+05	0.	2.E-13
50.0	67.800	0	7.30561336690E+07	5.11675235419E+07	9.93452150841E+05	-6.73010070079E+05	0.	2.E-12
50.0	99.900	0	1.33379109294E+00	2.15575177151E+01	2.38521990283E+00	-3.64229300808E-01	2.E-15	-9.E-12
50.0	100.000	0	1.3535307794E+00	2.15642270714E+01	2.34680489911E+00	-3.64110567292E-01	-9.E-15	-3.E-11
50.0	100.100	0	1.37691502927E+00	2.15573829979E+01	2.31238995999E+00	-3.64226716433E-01	-1.E-14	-8.E-12
50.0	153.100	0	1.27213570977E+00	-1.62131094239E-01	-2.1900180075E-01	7.50192053754E-01	-7.E-15	7.E-14
50.0	153.300	0	1.29566554688E+00	-7.28563466960E-02	-1.3026207773E-01	7.64759168841E-01	-7.E-15	-2.E-12
50.0	100.000	0	9.02460923791E+01	4.64400827350E+01	4.89547545830E+01	-8.56163232319E-01	-1.E-14	1.E-11
100.0	100.000	0	1.65581311974E+01	-9.34063273097E-01	-1.04432596255E+00	1.48165711328E-01	-7.E-15	-1.E-13
100.0	100.000	1	1.56234523525E+01	-9.353505066214E+01	-1.04576577634E+00	1.398056036618E-01	6.E-15	-4.E-13
100.0	100.000	1	1.37506264479E+01	-9.37700128836E-01	-1.0489406123E+00	1.23054420858E-01	-1.E-13	4.E-14
100.0	100.000	3	1.0933844943E+01	-9.40661432623E-01	-1.0517068765E+00	9.78603939247E-02	-4.E-14	2.E-12
100.0	100.000	5	2.44816942225E+02	-9.45477732181E-01	-1.05709767609E+00	2.19687751248E-02	-1.E-13	3.E-11
100.0	100.000	10	3.46330582628E+01	-8.9356690370E-01	-9.99082523040E-01	-3.0963701194E-01	6.E-14	3.E-12
100.0	100.000	20	1.04777556029E+00	1.27943911132E+01	1.43156773333E+01	-9.36921989821E-01	6.E-14	-9.E-12
100.0	100.000	30	7.51833277569E+01	6.58051104114E+01	7.33929454337E+01	6.72022224704E-01	3.E-14	-1.E-11
100.0	100.000	50	1.02728883766E+01	2.26855983258E+01	2.53963904196E+01	9.1735264491E-01	-4.E-14	1.E-11
100.0	100.000	100	9.36270607960E+01	4.43161964617E+01	-4.98554253144E+01	-8.32088192329E-01	-8.E-14	1.E-11
200.0	343.000	0	3.617559822E+07	9.74581592350E+08	5.19490430979E+06	-2.09045401279E+06	-2.E-15	2.E-10
200.0	344.000	0	3.5608658724E+07	1.45476847461E+07	3.48128926652E+06	-1.38604968934E+06	-5.E-15	2.E-10
200.0	399.000	0	1.5385091324E+00	1.67860189723E+01	3.25118563457E+00	-2.95191274837E-01	-7.E-15	-1.E-10
200.0	400.000	0	1.70803104755E+00	1.69857322090E+01	2.95864746041E+00	-2.91242987695E-01	-5.E-15	-1.E-10
200.0	401.000	0	1.87714230558E+00	1.67584821532E+01	2.66623392522E+00	-2.94692307136E-01	-1.E-14	8.E-12
200.0	475.000	0	1.681975837E+01	-4.97941021547E-01	-1.24599573540E+00	-3.86650279840E-01	-2.E-15	-4.E-12
200.0	476.000	0	4.20046440799E+01	-6.10636723026E-01	-1.52505808348E+00	-1.63656974114E-01	-2.E-15	6.E-13
200.0	100.000	0	1.12173949248E+00	-1.39872626369E-01	-1.80816071552E+01	8.68926152391E-01	-4.E-15	-6.E-11
500.0	918.000	0	9.6777017096E+08	2.92443068045E+08	1.2878744622E+07	-5.10790996398E+06	4.E-15	3.E-10
500.0	920.000	0	1.76434005999E+07	5.26126911617E+08	9.61309194998E+06	-2.80121373882E+06	2.E-15	3.E-10
500.0	999.000	0	1.84475784730E+00	1.44518679335E+01	3.5978284737E+00	-2.52384160504E-01	4.E-15	-1.E-09
500.0	1000.000	0	1.98990418398E+00	1.45466034536E+01	3.44670129710E+00	-2.50575899128E-01	4.E-15	-1.E-09
500.0	1001.000	0	2.13502662979E+00	1.44423361189E+01	3.1955212237E+00	-2.52214527754E-01	-1.E-14	-6.E-10
500.0	1094.000	0	1.0869119464E+00	4.35124353444E+01	1.493133335619E+00	-3.22200648502E-01	-2.E-14	3.E-11
500.0	1095.000	0	1.46929027367E+00	3.23944839003E+01	1.11156684610E+00	-4.35525914746E-01	-1.E-14	-2.E-12
500.0	1000.000	0	5.31450210236E+01	-8.3357932907E-01	-8.78437947200E+01	5.04180378468E-01	2.E-15	1.E-11
1000.0	1920.000	0	2.9493306403E+05	4.66213784110E+06	1.08963466775E+05	-2.18728710472E+04	2.E-15	-8.E-09
1000.0	1999.000	0	2.10432530931E+00	1.28927279390E+01	4.09255913389E+00	-2.24469801810E-01	-5.E-15	-2.E-10
1000.0	2000.000	0	2.23360882198E+00	1.28464300508E+01	3.68878233956E+00	-2.23464997529E-01	-2.E-14	-1.E-08
1000.0	2001.000	0	2.36288164524E+00	1.28884534129E+01	3.64498433934E+00	-2.24394604184E-01	-7.E-15	-7.E-10
1000.0	1000.000	0	7.55288842589E+01	6.61854934048E+01	7.39981588961E+01	-6.755555503365E-01	-1.E-14	-3.E-11

TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PROGRAM - RCMFN

ETA	RHO	L	F	FP	G	GP	M=1	TEST
50.0	600.000	500	1.73798850811E+01	6.04848225640E-01	1.63257196022E+00	-7.21624255451E-02	-2.1E-13	6.E-13
50.0	600.000	500	1.73798850811E+01	6.04848225639E-01	1.63257196022E+00	-7.21624255448E-02	-2.1E-13	7.E-13
50.0	600.000	500	1.73798850811E+01	6.04848225639E-01	1.63257196022E+00	-7.21624255448E-02	-9.E-14	7.E-13
50.0	600.000	500	1.73798850811E+01	6.04848225639E-01	1.63257196022E+00	-7.21624255449E-02	-7.E-14	7.E-13
50.0	600.000	500	1.73798850810E+01	6.04848225638E-01	1.63257196022E+00	-7.21624255449E-02	-7.E-14	7.E-13
50.0	600.000	500	1.73798850810E+01	6.04848225638E-01	1.63257196022E+00	-7.21624255450E-02	3.E-15	8.E-13
50.0	1200.000	1000	-1.35051870473E+00	2.59965760159E-01	5.48342608435E-01	6.34903977241E-01	-3.E-13	-4.E-12
50.0	1200.000	1000	-1.35051870473E+00	2.59965760156E-01	5.48342608437E-01	6.34903977242E-01	-2.E-13	-3.E-12
50.0	1200.000	1000	-1.35051870473E+00	2.59965760156E-01	5.48342608438E-01	6.34903977242E-01	-3.E-13	-3.E-12
50.0	1200.000	1000	-1.35051870473E+00	2.59965760159E-01	5.48342608439E-01	6.34903977243E-01	-3.E-13	-3.E-12
50.0	1200.000	1000	-1.35051870472E+00	2.59965760155E-01	5.48342608439E-01	6.34903977243E-01	-2.E-13	-3.E-12
50.0	1200.000	1000	-1.35051870472E+00	2.59965760155E-01	5.48342608440E-01	6.34903977244E-01	-7.E-15	-2.E-12

PROGRAM LISTING

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ABPCRCWFN. COULOMB WAVE FUNCTIONS FOR ALL REAL ETA AND RHO.
1  BARNETT, A.R., FENG, D.H., STEED, J.W., GOLDFARB, L.J.B.
JOB ARBFF, :MBCNP, CP76(T36,P2000)
CONVERT, INPUT, FILE, IBM360.
REWIND, FILE.
FTN, I=FILE, R=1, LX.
LDSET, MAP=B/ZZZZMP, PRESET=NGINF.
LGO, FILE, PL=1500.
####S
PROGRAM CFNCHK(INPUT, OUTPUT, TAPE5=INPUT, TAPE6=OUTPUT)
C *** MAIN PROGRAM TO TEST COULOMB WAVEFUNCTIONS MANCHESTER MAY 1974
DIMENSION FC(1200), FCP(1200), GC(1200), GCP(1200), RHO(21), LLL(10)
WRITE(6,1)
100 READ(5,2) ETA, N, (LLL(I), I=1,10), M
IF(N.EQ.0) CALL EXIT
IF(N.LT.0) WRITE(6,1)
IF(N.LT.0) N = -N
READ(5,3) (RHO(I), I=1, N)
ACCUR = 1.0E-14
STEP = 999.0
DO 10 I = 1, N
DO 10 J = 1, 10
IF(LLL(J).EQ.0.AND.J.GT.1) GO TO 10
L = LLL(J)
CALL RCWFN(RHO(I), ETA, M, L, FC, FCP, GC, GCP, ACCUR, STEP)
F = FC(L+1)
FP = FCP(L+1)
G = GC(L+1)
GP = GCP(L+1)
RLP1 = FLOAT(L+1)/SQRT(ETA*ETA + FLOAT((L+1)**2))
C *** MINL = M + 1 MAXL = L + 1 IN THIS CALL TO CHECK RESULTS ***
CALL RCWFN(RHO(I), ETA, M+1, L+1, FC, FCP, GC, GCP, ACCUR, STEP)
TEST1 = FP*G - F*GP - 1.00
TEST2 = F*GC(L+2) - G*FC(L+2) - RLP1
WRITE(6,4) ETA, RHO(I), L, F, FP, G, GP, TEST1, TEST2
10 CONTINUE
GO TO 100
1 FORMAT(1H1,30X,61H TEST OF THE MANCHESTER COULOMB WAVEFUNCTION PRO
1GRAM - RCWFN //,4X,4H ETA,6X,4H RHO,5X,2H L,8X,2H F,17X,2HFP,
117X,2H G,19X,2HGP,10X,4H W-1,5X,4HTEST,/)
2 FORMAT(F10.3,12I5)
3 FORMAT(7F10.3)
4 FORMAT(F10.1,F10.3,I5,4(1X,1PE18.11),2(2X,E7.0))
END
SUBROUTINE RCWFN(RHO, ETA, MINL, MAXL, FC, FCP, GC, GCP, ACCUR, STEP)
REAL K, K1, K2, K3, K4, M1, M2, M3, M4
DIMENSION FC(1), FCP(1), GC(1), GCP(1)
C *** COULOMB WAVEFUNCTIONS CALCULATED AT R = RHO BY THE
C *** CONTINUED-FRACTION METHOD OF STEED MINL, MAXL ARE ACTUAL L-VALUES
C *** SEE BARNETT FENG STEED AND GOLDFARB COMPUTER PHYSICS COMMUN 1974
PACE = STEP
ACC = ACCUR
IF(PACE.LT.100.0) PACE = 100.0
IF(ACC.LT.1.0E-15.OR.ACC.GT.1.0E-6) ACC = 1.0E-6
R = RHO
KTR = 1
LMAX = MAXL
LMIN1 = MINL + 1
XLL1 = FLOAT(MINL*LMIN1)
ETA2 = ETA*ETA
TURN = ETA + SQRT(ETA2 + XLL1)
IF(R.LT.TURN.AND.ABS(ETA).GE.1.0E-6) KTR = -1
KTRP = KTR
GO TO 2

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1      R      = TURN
      TF      = F
      TFP     = FP
      LMAX    = MINL
      KTRP    = 1
2      ETAR   = ETA*R
      RHO2    = R*R
      PL      = FLOAT(LMAX + 1)
      PMX     = PL + 0.5
C *** CONTINUED FRACTION FOR FP(MAXL)/F(MAXL)  XL IS F  XLPRIME IS FP **
      FP      = ETA/PL + PL/R
      DK      = ETAR*2.0
      DEL     = 0.0
      D       = 0.0
      F       = 1.0
      K       = (PL*PL - PL + ETAR)*(2.0*PL - 1.0)
      IF(PL*PL+PL+ETAR.NE.0.0) GO TO 3
      R       = R + 1.0E-6
      GO TO 2
3      H      = (PL*PL + ETA2)*(1.0 - PL*PL)*RHO2
      K       = K + DK + PL*PL*6.0
      D       = 1.0/(D*K + K)
      DEL     = DEL*(D*K - 1.0)
      IF(PL.LT.PMX) DEL = -R*(PL*PL + ETA2)*(PL + 1.0)*D/PL
      PL      = PL + 1.0
      FP      = FP + DEL
      IF(D.LT.0.0) F = -F
      IF(PL.GT.20000.) GO TO 11
      IF(ABS(DEL/FP).GE.ACC) GO TO 3
      FP      = F*FP
      IF(LMAX.EQ.MINL) GO TO 5
      FC(LMAX+1) = F
      FCP(LMAX+1) = FP
C *** DOWNWARD RECURSION TO MINL FOR F AND FP, ARRAYS GC,GCP ARE STORAGE
      L       = LMAX
      DO 4 LP = LMIN1,LMAX
      PL      = FLOAT(L)
      GC(L+1) = ETA/PL + PL/R
      GCP(L+1) = SQRT(ETA2 + PL*PL)/PL
      FC(L)   = (GC(L+1)*FC(L+1) + FCP(L+1))/GCP(L+1)
      FCP(L)  = GC(L+1)*FC(L) - GCP(L+1)*FC(L+1)
4      L      = L - 1
      F       = FC(LMIN1)
      FP      = FCP(LMIN1)
5      IF(KTRP.EQ.-1) GO TO 1
C *** REPEAT FOR R = TURN IF RHO LT TURN
C *** NOW OBTAIN P + I.Q FOR MINL FROM CONTINUED FRACTION (32)
C *** REAL ARITHMETIC TO FACILITATE CONVERSION TO IBM USING REAL*8
      P       = 0.0
      Q       = R - ETA
      PL      = 0.0
      AR      = -(ETA2 + XLL1)
      AI      = ETA
      BR      = 2.0*Q
      BI      = 2.0
      WI      = 2.0*ETA
      DR      = BR/(BR*BR + BI*BI)
      DI      = -BI/(BR*BR + BI*BI)
      DP      = -(AR*DI + AI*DR)
      DQ      = (AR*DR - AI*DI)
6      P      = P + DP
      Q       = Q + DQ
      PL      = PL + 2.0
      AR      = AR + PL
      AI      = AI + WI
      BI      = BI + 2.0
      D       = AR*DR - AI*DI + BR
      DI      = AI*DR + AR*DI + BI

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T = 1.0/(D*D + DI*DI)
DR = T*D
DI = -T*DI
H = BR*DR - BI*DI - 1.0
K = BI*DR + BR*DI
T = DP*H - DQ*K
DQ = DP*K + DQ*H
DP = T
IF(PL.GT.46000.) GO TO 11
IF(ABS(DP)+ABS(DQ).GE.(ABS(P)+ABS(Q))*ACC) GO TO 6
P = P/R
Q = Q/R
C *** SOLVE FOR FP,G,GP AND NORMALISE F AT L=MINL
G = (FP - P*F)/Q
GP = P*G - Q*F
W = 1.0/SQRT(FP*G - F*GP)
G = W*G
GP = W*GP
IF(KTR.EQ.1) GO TO 8
F = TF
FP = TFP
LMAX = MAXL
C *** RUNGE-KUTTA INTEGRATION OF G(MINL) AND GP(MINL) INWARDS FROM TURN
C *** SEE FOX AND MAYERS 1968 PG 202
IF(RHO.LT.0.2*TURN) PACE = 999.0
R3 = 1.0/3.0DO
H = (RHO - TURN)/(PACE + 1.0)
H2 = 0.5*H
I2 = IFIX(PACE + 0.001)
ETAH = ETA*H
H2LL = H2*XLL1
S = (ETAH + H2LL/R )/R - H2
RH2 = R + H2
T = (ETAH + H2LL/RH2)/RH2 - H2
K1 = H2*GP
M1 = S*G
K2 = H2*(GP + M1)
M2 = T*(G + K1)
K3 = H*(GP + M2)
M3 = T*(G + K2)
M3 = M3 + M3
K4 = H2*(GP + M3)
RH = R + H
S = (ETAH + H2LL/RH )/RH - H2
M4 = S*(G + K3)
G = G + (K1 + K2 + K2 + K3 + K4)*R3
GP = GP + (M1 + M2 + M2 + M3 + M4)*R3
R = RH
I2 = I2 - 1
IF(ABS(GP).GT.1.0E300) GO TO 11
IF(I2.GE.0) GO TO 7
W = 1.0/(FP*G - F*GP)
C *** UPWARD RECURSION FROM GC(MINL) AND GCP(MINL),STORED VALUES ARE R,S
C *** RENORMALISE FC,FCP FOR EACH L-VALUE
8 - GC(LMIN1) = G
- GCP(LMIN1) = GP
IF(LMAX.EQ.MINL) GO TO 10
DO 9 L = LMIN1,LMAX
- T = GC(L+1)
- GC(L+1) = (GC(L)*GC(L+1) - GCP(L))/GCP(L+1)
- GCP(L+1) = GC(L)*GCP(L+1) - GC(L+1)*T
- FC(L+1) = W*FC(L+1)
- FCP(L+1) = W*FCP(L+1)
9 - FC(LMIN1) = FC(LMIN1)*W
- FCP(LMIN1) = FCP(LMIN1)*W
RETURN
10 - FC(LMIN1) = W*F
- FCP(LMIN1) = W*FP
RETURN
11 W = 0.0
G = 0.0
GP = 0.0
GO TO 8
END

```

FCP → FCP

FC → F

GCP → GP

GC → G

[illegible]