## laplace/body\_ax - an excerpt from the book C.Pozrikidis 'A Practical Guide to Boundary-Element Methods'

with additional notes and commentaries composed by Alexander Samoilov

## Introduction

- This is a code that computes potential flow past an axisymmetric, compact (singly connected) or toroidal (doubly connected) body with arbitrary geometry, as illustrated in Figure
- original Fortran77 code is here CFDLAB code for laplace/body\_ax

## **Mathematical Formulation**

The velocity u is decomposed into three parts as

$$u = u^{\infty} + v + u^{D} \tag{1}$$

where:

- $u^{\infty}$  is the far-field component prevailing far from the body, expressing uniform (streaming) flow along the x axis of symmetry.
- v is the velocity due to a line vortex ring with specified strength situated in the interior of the body, generating circulation around the toroidal body. In the case of flow past a compact body, this component is inconsequential, and serves only to modify the disturbance velocity.
- $u^D$  is a disturbance velocity expressed by the gradient of the single-valued harmonic potential  $\Phi^D$ ,

$$u^D = \nabla \Phi^D \tag{2}$$

Requiring the no-penetration boundary condition  $u \cdot n = 0$  around the contour of the body in a meridional plane of constant angle  $\varphi$ , where n is the unit vector normal to the body, we derive a boundary condition for the normal derivative of the disturbance potential,

$$n \cdot \nabla \phi^D \equiv \frac{\partial \phi^D}{\partial n} = -\left(u^\infty + v\right) \cdot n \tag{3}$$

Using the standard boundary-integral formulation, we find that the disturbance potential satisfies the integral equation of the second kind

$$\phi^{D}(x_{0}) = -2 \int_{C} G(x, x_{0})[n(x) \cdot \nabla \phi^{D}(x)] dl(x) + 2 \int_{C}^{PV} \phi^{D}(x)[n(x) \cdot \nabla G(x, x_{0})] dl(x)$$
(4)

where  $G(x, x_0)$  is the free-space Green's function of Laplace's equation in an axisymmetric domain, and the point  $x_0$  lies on the contour of the body C. Inserting the boundary condition (3) into (4) and rearranging, we obtain

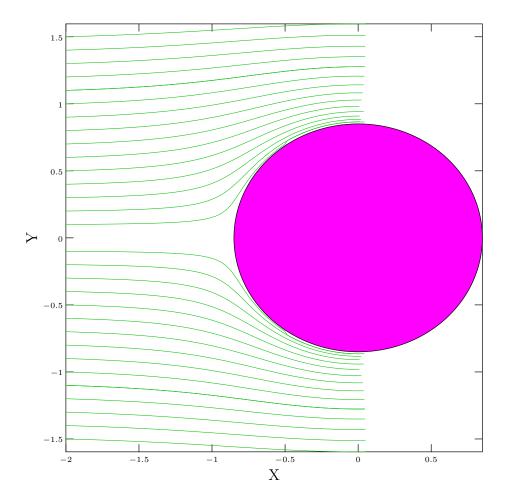


Figure 1: Streamlines of flow past a sphere

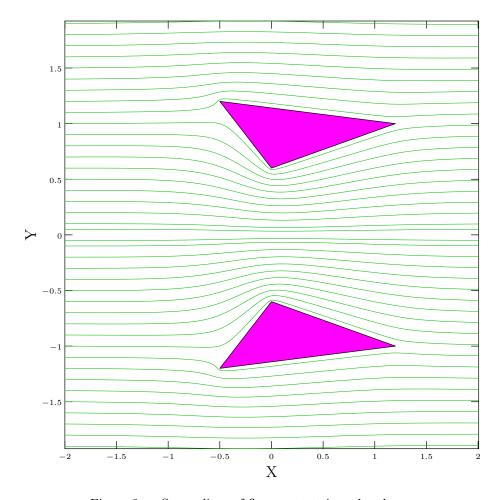


Figure 2: Streamlines of flow past a triangular thorus

$$\phi^{D}(x_{0}) - 2 \int_{C}^{PV} \phi^{D}(x) [n(x) \cdot \nabla G(x, x_{0})] dl(x) = 2 \int_{C} (u^{\infty} + v) \cdot G(x, x_{0}) n(x) dl(x)$$
(5)

## **Numerical Method**

The contour of the body in a meridional plane is discretized into a collection of N boundary elements denoted by  $E_i, i = 1, \ldots, N$ . The boundary elements can be straight segments or circular arcs. The disturbance potential and its normal derivative are approximated with constant functions the ith element denoted respectively by  $\varphi_i^D$  and

$$\left(\frac{\partial \phi^D}{\partial n}\right)_i = (u^\infty + v^{(i)}) \cdot n^{(i)} \tag{6}$$

where  $n^{(i)}$  and  $v^{(i)}$  are the normal vector and velocity induced by the line vortex ring evaluated at the mid-point of the *i*-th element. Subject to these approximations, the integral equation (5) assumes the discretized form

$$\phi_{j}^{D} - 2\sum_{i=1}^{N} \phi_{i}^{D} \int_{E_{i}}^{PV} n(x) \nabla G\left(x, x_{j}^{M}\right) dl(x) = 2\sum_{i=1}^{N} \left(u_{\infty} + v^{(i)}\right) \cdot n^{(i)} \int_{E_{i}} G\left(x, x_{j}^{M}\right) dl(x)$$
(7)