

① We know $F = \frac{mv^2}{R}$ (centripetal force)

$$\Rightarrow q \times B = \frac{mv^2}{R} \Rightarrow p = qBR \quad (\text{in S.I units})$$

$$p \text{ kg m/s} = p' \frac{\text{GeV}}{c} \Rightarrow p \left(\frac{\text{J}}{\text{m/s}} \right) = p' \left(\frac{\text{GeV}}{c} \right) \left(\frac{1.6 \times 10^{-19} \text{ J}}{10^{-9} \text{ GeV}} \right) \left(\frac{1}{3 \times 10^8 \text{ m/s}} \right)$$

$$\Rightarrow p = p' \frac{1.6 \times 10^{-10}}{3 \times 10^8}$$

$$1) q(C) = q'(e) \left(\frac{1.6 \times 10^{-19} \text{ C}}{e} \right) \Rightarrow q = q' (1.6 \times 10^{-19})$$

$$\therefore p' \left(\frac{1.6 \times 10^{-18}}{3} \right) = q' (1.6 \times 10^{-19}) BR \Rightarrow p' = 0.3 q' BR$$

where p' in $\frac{\text{GeV}}{c}$ & q' in e (can be assumed = 1) ~~8.0~~

$$2) \text{ ~~Boxed~~ } [(511c^2)^2 + 1]^2 = (511c^2)^2 + p^2 c^2$$

$$\Rightarrow (511c^2)^2 + 1 + 2(511c^2) = (511c^2)^2 + p^2 c^2 \Rightarrow p^2 = \left(1022 + \frac{1}{c^2} \right) \left(\frac{\text{KeV}}{c} \right)^2$$

$$\Rightarrow p = \sqrt{1022 + \frac{1}{c^2}} \frac{\text{KeV}}{c}$$

$$2) @ 1 \text{ KeV } e^- \Rightarrow (511 + 1)^2 = (511)^2 + p^2 c^2$$

$$\Rightarrow 1023 \frac{\text{KeV}^2}{c^2} = p^2 \Rightarrow p = 31.98 \frac{\text{KeV}}{c} = 31.98 \times 10^{-6} \frac{\text{GeV}}{c}$$

$$\therefore R = \frac{31.98 \times 10^{-6}}{0.3 \times 1 \times 1} = 10.66 \times 10^{-5} \text{ m} \Rightarrow 0.1066 \text{ mm}$$

$$(511 + 100)^2 = (511)^2 + p^2 c^2 \Rightarrow p = \frac{334.96 \text{ GeV}}{c} \approx 334.96 \times 10^{-6} \frac{\text{GeV}}{c}$$

$$\Rightarrow R = \frac{334.96 \times 10^{-6}}{0.3 \times 1} = 1116.54 \times 10^{-6} \text{ m} \approx 1.11 \text{ mm}$$

Yes, modern HEP detectors / trackers can detect these electrons due to their high radius of curvature.

③ Using Synchrotron Radiation formula $\rightarrow P_{\gamma} = \frac{q^2 c}{6\pi\epsilon_0} \frac{\beta^4 \gamma^4}{R^2}$
where R = radius of curvature of particle.

$$\text{We know } E = \gamma m_0 c^2 \Rightarrow (m_0 c^2 + KE) = \gamma m_0 c^2 \Rightarrow \gamma = 1 + \frac{KE}{m_0 c^2}$$

$$\text{For a proton at LHC} \rightarrow KE = 6.8 \text{ TeV}, m_0 c^2 = 0.938 \text{ GeV}$$

$$R = \frac{27 \times 10^3}{2\pi} \text{ m}$$

$$\therefore \gamma = 1 + \frac{6800}{0.938} \approx 7250 \quad \therefore \beta \approx 1$$

$$\therefore P_{\gamma} = \frac{\frac{2}{3} (1.6 \times 10^{-19})^2 \times (3 \times 10^8)^2 (7250)^4 \times 9 \times 10^9}{\left(\frac{27 \times 10^3}{2\pi}\right)^2} = 6.89 \times 10^{-12} \text{ W}$$

$$\text{for } e^- \text{ at LEP} \rightarrow KE = 209 \text{ GeV}, m_0 c^2 = 0.511 \text{ GeV}$$

$$R = \frac{27 \times 10^3}{2\pi} \text{ m}$$

$$\therefore \gamma = 1 + \frac{209}{0.511} \approx 410 \Rightarrow \beta = 0.99999103$$

$$\therefore P_{\gamma} = \frac{2}{3} \times 9 \times 10^9 \times \frac{(1.6 \times 10^{-19})^2 \times (3 \times 10^8)^2 \times (410)^4}{\left(\frac{27000}{2\pi}\right)^2} = 7.04 \times 10^{-17} \text{ W}$$

$$\therefore \frac{P_{p^+}}{P_e} = 97869.32$$

For e^- at Super KEKB $\rightarrow KE = 7 \text{ GeV}$, $m_e c^2 = 0.511 \text{ GeV}$

$$P = \frac{3016}{2\pi} \text{ m}$$

$$\therefore \gamma = 1 + \frac{7}{0.511} \Rightarrow \beta = 0.997 \therefore P_{\gamma} \approx 1.15 \times 10^{-22} \text{ W}$$

$$= 1.147$$

$$\therefore \frac{P_{\text{LEP}}}{P_{\text{SKEKB}}} = 6.12173 \times 10^5$$

A) $\therefore -\frac{dE}{dx} = \frac{E}{X_0} \Rightarrow \frac{dE}{E} = \frac{dx}{X_0} \Rightarrow \ln E = -\frac{x}{X_0} \Rightarrow E = E_0 e^{-\frac{x}{X_0}}$

$$\Rightarrow E(X_0) = \frac{E_0}{e} \therefore P_{\text{loss}} = E_0 \left(1 - \frac{1}{e}\right) = 0.63 E_0$$

Radiation Length of Be = 35.28 cm.

$$\therefore E(1 \text{ mm}) = 0.5 e^{-\frac{1}{35.28}} \text{ GeV} \Rightarrow E_{\text{loss}} = 0.5 \left(1 - e^{-\frac{1}{35.28}}\right)$$

$$= 0.0001415 \text{ MeV}$$

5) scattering angle $\theta_{ms} \approx \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{\frac{z}{X_0}} \left[1 + 0.038 \log_{10} \left(\frac{z^2}{X_0 \beta^2} \right) \right]$

for 1 GeV pion (π^-) $\rightarrow z=1$, $X_0 = 8.897 \text{ cm}$
through Aluminium.

$$E^2 = m^2 c^4 + p^2 c^2 \Rightarrow (139.57 + 1000)^2 = (139.57)^2 + p^2 c^2$$

$$\Rightarrow (1279.14)(1000) = p^2 c^2 \Rightarrow p \approx \cancel{1131} 1131 \frac{\text{MeV}}{c}$$

$$\gamma = 1 + \frac{KE}{mc^2} = 1 + \frac{1000}{139.57} \approx 8.164 = \frac{1}{\sqrt{1-\beta^2}} \Rightarrow \beta = \sqrt{1 - \frac{1}{(8.164)^2}} \approx 0.9925$$

$$\therefore \theta_{ms} = \frac{13.6 \text{ MeV}}{(0.9925)(1131 \text{ MeV})} \sqrt{\frac{1}{8.897}} \left[1 + 0.038 \log_{10} \left(\frac{1}{8.897 (0.9925)^2} \right) \right]$$

$$= \cancel{0.232} 3.72 \text{ mrad}$$

6) In rest frame $\rightarrow \tau_0 = 8.954 \times 10^{-11} \text{ s}$, $m_0 = 497.611 \frac{\text{MeV}}{c^2}$

$$\text{In lab frame} \rightarrow \tau_L = \gamma \tau_0 = \left(1 + \frac{10000}{497.611} \right) \tau_0 = 21 \times 8.954 \times 10^{-11} \text{ s} \\ = 1.88089 \times 10^{-9} \text{ s} \\ = 1.88 \text{ ns}$$

$$\therefore \text{Flight Distance} = v \tau_L = \beta c \tau_L$$

$$= \sqrt{1 - \frac{1}{(21)^2}} (1.88 \text{ ns}) c \approx 0.498 \text{ m}$$

7) given $\frac{\sigma_E}{E} = \frac{2.03\%}{(E/\text{GeV})^{1/4}} \quad \textcircled{1} 1.04\%$

For 5 GeV $\gamma \rightarrow \frac{\sigma_E}{E} = \frac{2.03\%}{5^{1/4}} \quad \textcircled{1} 1.04\% \quad = 1.54\% \quad \textcircled{4} 1.04\%$

$= \sqrt{(1.54)^2 + (1.04)^2} \approx 2.08\% \quad \therefore \frac{\sigma_E}{E} = 2.08\%$

8) From PDG, we find $|\frac{dE}{dx}|_{\min} = 3.876 \text{ MeV/cm}$ for Si

$\therefore n_{eh} = \frac{3.876 \times 10^6 \times 10^2 \times 100 \times 10^{-6}}{3.6}$

$\approx 1.0766 \times 10^4 = 10766 \text{ e-hole pairs}$

9) For Photoelectric effect $\sigma_{PE} \propto \frac{Z^5}{E^3}$

Compton Scattering $\sigma_{CS} \propto \frac{Z}{E}$

Pair Production $\sigma_{PP} \propto E, Z \uparrow$

10) For C, $\sigma_{PE} = \sigma_{CS} \text{ at } \sim 100 \text{ KeV}$

$\therefore K = \alpha \frac{Z^5}{E^3} = \alpha \frac{6^5}{(100)^3} \quad \& \quad K = \beta \frac{Z}{E} = \beta \frac{6}{(100)}$

\therefore For $E = 1000 \text{ KeV}$ in Al (13)

$\frac{\sigma_{PE}}{\sigma_{CS}} = \frac{\alpha \left(\frac{13^5}{1000^3}\right)}{\beta \left(\frac{13}{1000}\right)} = K \frac{\left(\frac{13}{6}\right)^5 \left(\frac{100}{1000}\right)^3}{\left(\frac{13}{1000}\right) \left(\frac{100}{6}\right)} = \left(\frac{13}{6}\right)^4 \left(\frac{100}{1000}\right)^2$

\therefore Compton scattering dominates

b) As before, for $H (Z=2)$ at $E = 100 \text{ keV}$.

$$\frac{\sigma_{PE}}{\sigma_{CS}} = \frac{10^5 \cdot \frac{2^5}{100^3} \cdot 100^3}{10^5 \cdot \frac{2^5}{100^3} \cdot 100^3} \cdot \left(\frac{2}{6}\right)^4 = \frac{1}{81}$$

\therefore Compton scattering is dominant

c) For Fe ($Z=26$) at $E = 100 \text{ keV}$

$$\frac{\sigma_{PE}}{\sigma_{CS}} = \left(\frac{26}{6}\right)^4 \gg 1 \quad \therefore \text{Photo electric effect is more dominant}$$

d) From the figure ~~mentions~~ ^{34.15 in referred document,} in $C (Z=6)$ at 10 MeV , Compton scattering is slightly more dominant than pair production.

e) From the same figure, we note that at 10 MeV in $Pb (Z=82)$, pair production is slightly more dominant than Compton scattering.