

Quantum Discord and Decoherence of inflationary perturbations

Asia-Pacific Workshop on Gravitation and Cosmology 2022

20th March 2022

Amaury Micheli^{1,2}

arXiv:2112.05037 AM, Jérôme Martin¹, Vincent Vennin^{1,3}

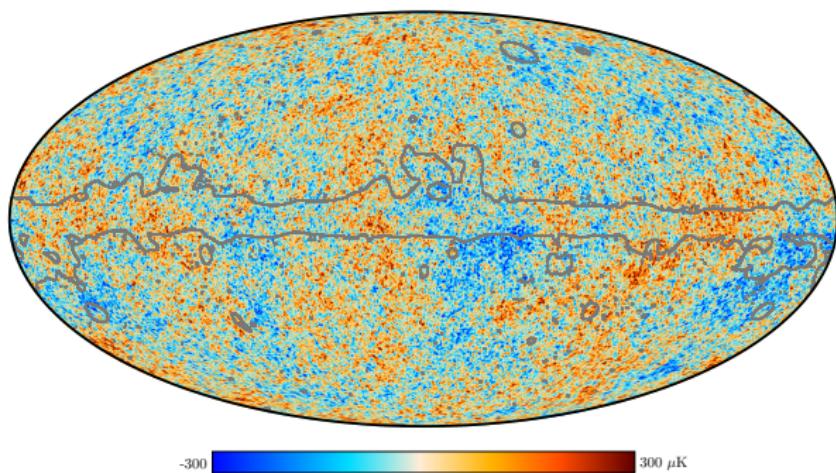
Institut d'Astrophysique de Paris¹

Irene-Joliot Curie Laboratory, Orsay² APC, Paris³

INTRODUCTION : QUANTUM FEATURES IN THE EARLY UNIVERSE ?

CONTEXT I, INHOMOGENEITIES IN THE EARLY UNIVERSE

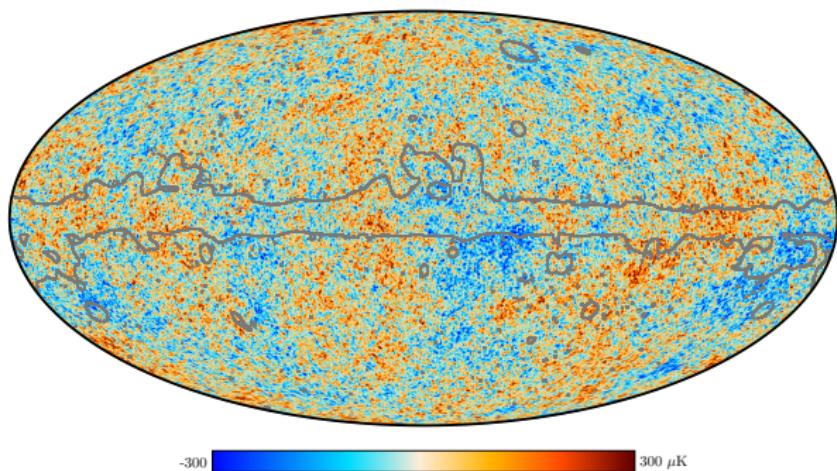
- CMB¹ : isotropic temperature $T \sim 3K$ + **small anisotropies**
 $\Delta T/T \sim 10^{-4}$



1. [Planck-Collaboration et al., 2020b]

CONTEXT I, INHOMOGENEITIES IN THE EARLY UNIVERSE

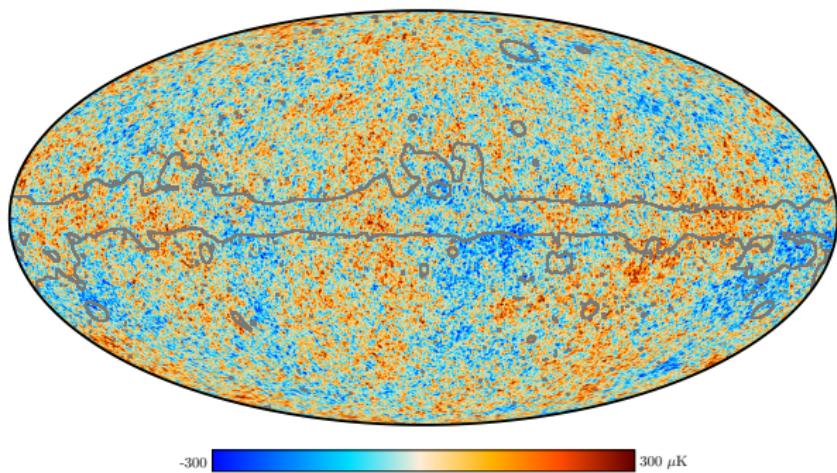
- CMB¹ : isotropic temperature $T \sim 3K$ + **small anisotropies**
 $\Delta T/T \sim 10^{-4}$
- Early Universe : homogeneous + **small inhomogeneities**



1. [Planck-Collaboration et al., 2020b]

CONTEXT I, INHOMOGENEITIES IN THE EARLY UNIVERSE

- CMB¹ : isotropic temperature $T \sim 3K$ + **small anisotropies**
 $\Delta T/T \sim 10^{-4}$
- Early Universe : homogeneous + **small inhomogeneities**
Origin of inhomogeneities ?



1. [Planck-Collaboration et al., 2020b]

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition $\sim 80s^2$: Inhomogeneities come from minimal (quantum) vacuum fluctuations at the beginning of inflation stretched to cosmological scales by expansion!

2. [Mukhanov and Chibisov, 1981]

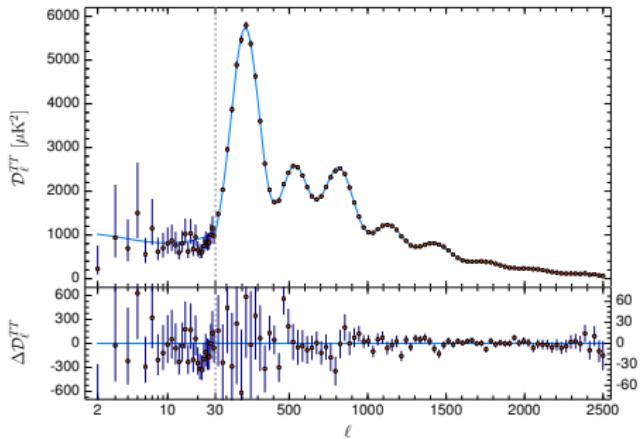
CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition $\sim 80s^2$: Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !

2. [Mukhanov and Chibisov, 1981]

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition ~80s : Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !
- Indirect proof : very good agreement with observational data.²



2. [Planck-Collaboration et al., 2020a]

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition ~80s : Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !
- Indirect proof : very good agreement with observational data.

Quantum then, classical now : how and when the transition happened ?

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition ~80s : Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !
- Indirect proof : very good agreement with observational data.

Quantum then, classical now : how and when the transition happened ?

- Need tools to **measure the quantumness of a state** :
Quantum Discord

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition ~80s : Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !
- Indirect proof : very good agreement with observational data.

Quantum then, classical now : how and when the transition happened ?

- Need tools to **measure the quantumness of a state** : **Quantum Discord**
- Need mechanism for a **quantum-to-classical transition** : **Decoherence**

CONTEXT II, INHOMOGENEITIES IN THE EARLY UNIVERSE

- Proposition ~80s : Inhomogeneities come from minimal (quantum) **vacuum fluctuations** at the beginning of inflation stretched to **cosmological scales** by expansion !
- Indirect proof : very good agreement with observational data.

Quantum then, classical now : how and when the transition happened ?

- Need tools to **measure the quantumness of a state** :
Quantum Discord
- Need mechanism for a **quantum-to-classical transition** :
Decoherence

CHARACTERIZING QUANTUMNESS OF INFLATIONNARY PERTURBATIONS

QUANTUMNESS OF A STATE

- Paradigm : **Quantumness of a state** for a system \mathcal{S} =
Quantumness of correlations of subsystems $\mathcal{S} = \mathcal{S}_1 \cup \mathcal{S}_2$
for this state.

QUANTUMNESS OF A STATE

- Paradigm : **Quantumness of a state** for a system \mathcal{S} = **Quantumness of correlations** of subsystems $\mathcal{S} = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Formalised in the definition of **Quantum Discord** $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2)$.

QUANTUMNESS OF A STATE

- Paradigm : **Quantumness of a state** for a system \mathcal{S} = **Quantumness of correlations** of subsystems $\mathcal{S} = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Formalised in the definition of **Quantum Discord** $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2)$.

$$\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) \equiv \mathcal{I}(\mathcal{S}_1, \mathcal{S}_2) - \max_{\{\hat{\Pi}_j^{\mathcal{S}_2}\}} \mathcal{J} \left(\mathcal{S}_1, \mathcal{S}_2, \{\hat{\Pi}_j^{\mathcal{S}_2}\} \right)$$

with \mathcal{I}, \mathcal{J} two measures of **mutual information** between $\mathcal{S}_{1/2}$.

QUANTUMNESS OF A STATE

- Paradigm : **Quantumness of a state** for a system \mathcal{S} = **Quantumness of correlations** of subsystems $\mathcal{S} = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Formalised in the definition of **Quantum Discord** $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2)$.

$$\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) \equiv \mathcal{I}(\mathcal{S}_1, \mathcal{S}_2) - \max_{\{\hat{\Pi}_j^{\mathcal{S}_2}\}} \mathcal{J} \left(\mathcal{S}_1, \mathcal{S}_2, \{\hat{\Pi}_j^{\mathcal{S}_2}\} \right)$$

with \mathcal{I}, \mathcal{J} two measures of **mutual information** between $\mathcal{S}_{1/2}$.

If \mathcal{S}_i described by classical probabilities $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) = 0$.
Quantum state $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) \geq 0$.

QUANTUMNESS OF A STATE

- Paradigm : **Quantumness of a state** for a system \mathcal{S} = **Quantumness of correlations** of subsystems $\mathcal{S} = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Formalised in the definition of **Quantum Discord** $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2)$.

$$\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) \equiv \mathcal{I}(\mathcal{S}_1, \mathcal{S}_2) - \max_{\{\hat{\Pi}_j^{\mathcal{S}_2}\}} \mathcal{J} \left(\mathcal{S}_1, \mathcal{S}_2, \{\hat{\Pi}_j^{\mathcal{S}_2}\} \right)$$

with \mathcal{I}, \mathcal{J} two measures of **mutual information** between $\mathcal{S}_{1/2}$.

If \mathcal{S}_i described by classical probabilities $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) = 0$.

Quantum state $\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) \geq 0$.

$\mathcal{D}(\mathcal{S}_1, \mathcal{S}_2) > 0$ = quantum state

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

→ Subsystems ?

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At quadratic order inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At **quadratic order** inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.
- **Gaussian state**

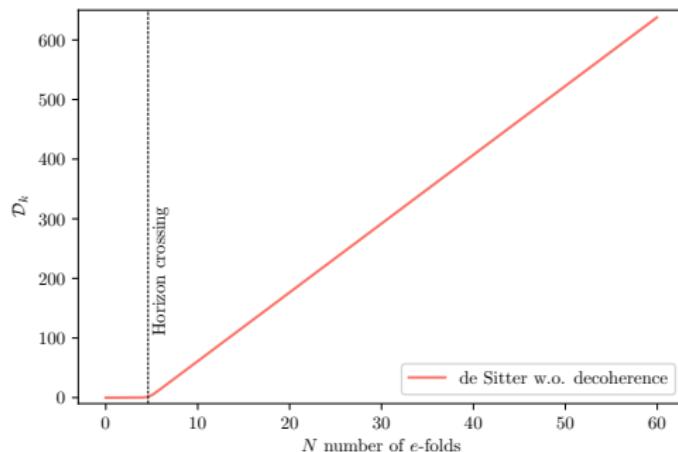
DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At **quadratic order** inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.
- **Gaussian state** → Discord can be computed.²

2. [Adesso and Datta, 2010]

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At **quadratic order** inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.
- **Gaussian state** → Discord can be computed. **It increases to very large values!**²



2. [Martin and Vennin, 2016]

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At **quadratic order** inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.
- **Gaussian state** → Discord can be computed. **It increases to very large values !**

Take-home message 1

Without decoherence Quantum Discord is strongly amplified by inflation and final state is very quantum in this sense.²

2. [Martin and Vennin, 2016]

DISCORD OF PERTURBATIONS W.O. DECOHERENCE

- Subsystems ? At **quadratic order** inflation creates perturbations in independent $\pm \mathbf{k}$ pairs.
- **Gaussian state** → Discord can be computed. **It increases to very large values !**

Take-home message 1

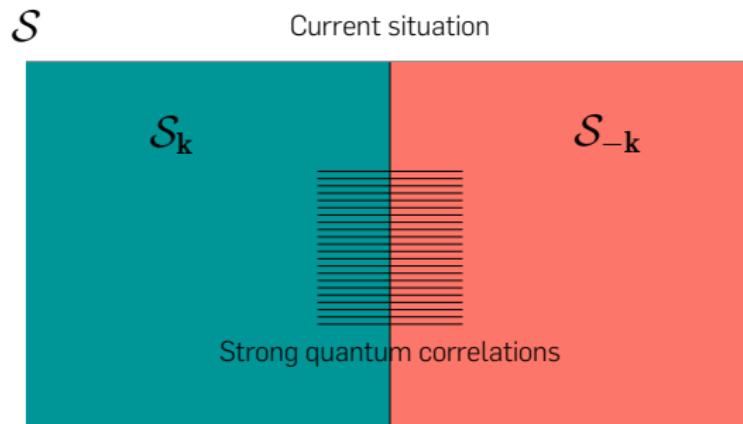
Without decoherence Quantum Discord is strongly amplified by inflation and final state is very quantum in this sense.²

Can this result be due to oversimplified models ?

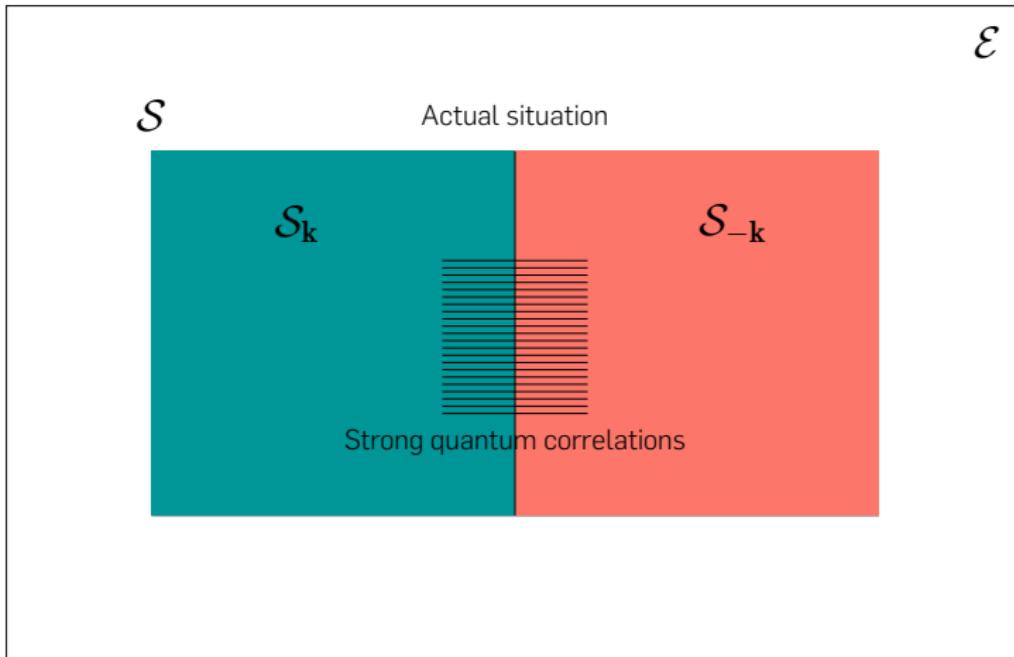
2. [Martin and Vennin, 2016]

DECOHERENCE AND LOSS OF QUANTUMNESS

NON-LINEARITIES, INTERACTIONS : DECOHERENCE

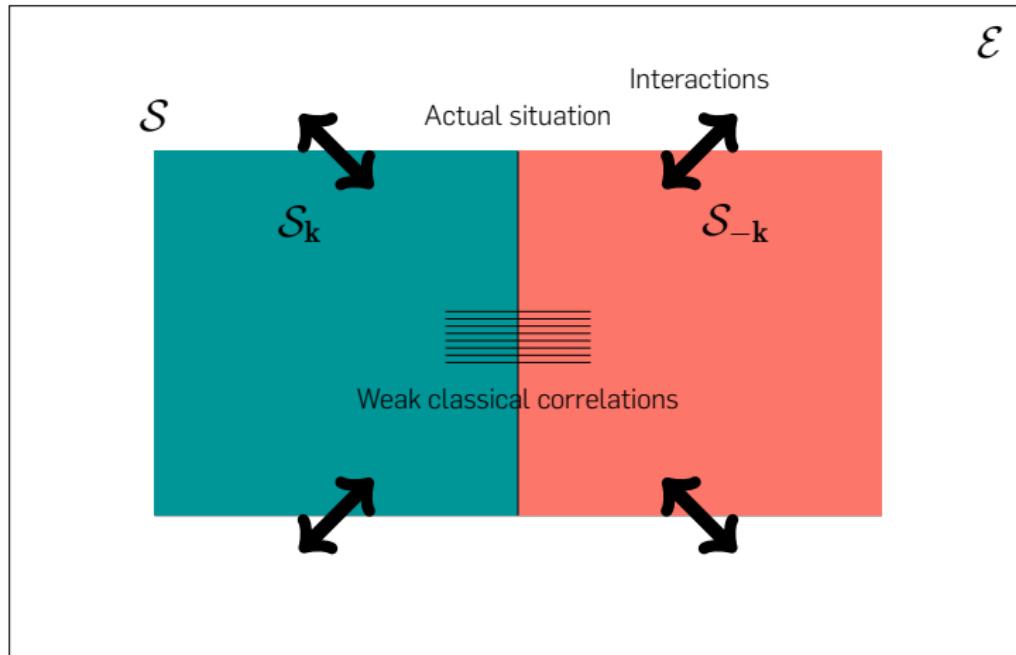


NON-LINEARITIES, INTERACTIONS : DECOHERENCE



In fact \mathcal{S} has an environment \mathcal{E} (e.g. $\mathcal{S}_{\pm k'}$ with $k' \neq k$) or other fields.

NON-LINEARITIES, INTERACTIONS : DECOHERENCE



Interactions $\mathcal{S} / \mathcal{E}$ destroy correlations S_k / S_{-k} : decoherence.

DECOHERENCE MODEL FOR MUKHANOV-SASAKI \hat{v}

- Perturbations \mathcal{S} described by the Mukhanov-Sasaki variable \hat{v}

DECOHERENCE MODEL FOR MUKHANOV-SASAKI \hat{v}

- Perturbations \mathcal{S} described by the Mukhanov-Sasaki variable \hat{v}
- Interaction with the environment \mathcal{E} **Linear** to preserve Gaussianity, independence of $\pm \mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm \mathbf{k}}$

$$\hat{H}_{\text{int}} = \lambda \int d^3 \mathbf{x} \sqrt{-g} \frac{\hat{v}}{a} \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

DECOHERENCE MODEL FOR MUKHANOV-SASAKI \hat{v}

- Perturbations \mathcal{S} described by the Mukhanov-Sasaki variable \hat{v}
- Interaction with the environment \mathcal{E} **Linear** to preserve Gaussianity, independence of $\pm \mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm \mathbf{k}}$

$$\hat{H}_{\text{int}} = \lambda \int d^3 \mathbf{x} \sqrt{-g} \frac{\hat{v}}{a} \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

- Leave \mathcal{E} unspecified to a certain degree, can we get evolution of $\hat{\rho}_{\mathcal{S}}$ **only**?

DECOHERENCE MODEL FOR MUKHANOV-SASAKI \hat{v}

- Perturbations \mathcal{S} described by the Mukhanov-Sasaki variable \hat{v}
- Interaction with the environment \mathcal{E} **Linear** to preserve Gaussianity, independence of $\pm \mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm \mathbf{k}}$

$$\hat{H}_{\text{int}} = \lambda \int d^3 \mathbf{x} \sqrt{-g} \frac{\hat{v}}{a} \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

- Leave \mathcal{E} unspecified to a certain degree, can we get evolution of $\hat{\rho}_{\mathcal{S}}$ **only**? **Lindblad equation**

Assumptions :

- Perturbation $\lambda \ll 1$.
- \mathcal{E} stationary and not perturbed by \mathcal{S} .
- Consider evolution of \mathcal{S} for $\eta \gg \eta_C$ auto-correlation time of \mathcal{E} .

DECOHERENCE MODEL FOR MUKHANOV-SASAKI \hat{v}

- Perturbations \mathcal{S} described by the Mukhanov-Sasaki variable \hat{v}
- Interaction with the environment \mathcal{E} **Linear** to preserve Gaussianity, independence of $\pm \mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm \mathbf{k}}$

$$\hat{H}_{\text{int}} = \lambda \int d^3 \mathbf{x} \sqrt{-g} \frac{\hat{v}}{a} \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

- Leave \mathcal{E} unspecified to a certain degree, can we get evolution of $\hat{\rho}_{\mathcal{S}}$ **only**? **Lindblad equation**

Assumptions :

- Perturbation $\lambda \ll 1$.
- \mathcal{E} stationnary and not perturbed by \mathcal{S} .
- Consider evolution of \mathcal{S} for $\eta \gg \eta_C$ auto-correlation time of \mathcal{E} .

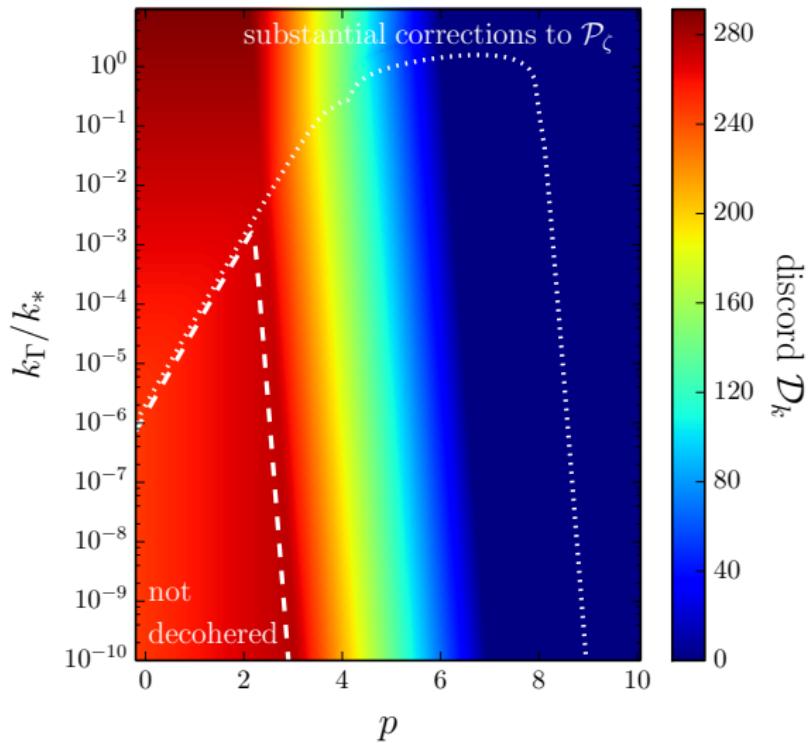
- Further computations³ decoherence term $\propto k_{\Gamma}^2 a^{p-3}$.

3. [Martin and Vennin, 2016]

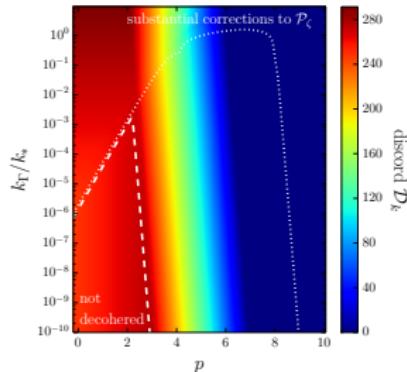
COMPETITION OF ENTANGLEMENT AND DECOHERENCE

Is Quantum Discord spoiled by decoherence?

COMPETITION OF ENTANGLEMENT AND DECOHERENCE



COMPETITION OF ENTANGLEMENT AND DECOHERENCE



Take-home message 2⁴

Decoherence does not always destroy Quantum Discord. Its fate is the result of a competition with generation of quantum correlations by inflation.

4. [arXiv:2112.05037 Martin et al., 2021]

FUTURE DIRECTIONS

FUTURE DIRECTIONS

- Compare the effect of decoherence on different criteria (Bell Inequalities, non-separability etc.).

FUTURE DIRECTIONS

- Compare the effect of decoherence on different criteria (Bell Inequalities, non-separability etc.).
- Use a more realistic interaction for decoherence, for instance non-linearities of pure gravity and see whether quantum discord is destroyed or not.

Thank you for your attention !

-  Adesso, G. and Datta, A. (2010).
Quantum versus classical correlations in Gaussian states.
Physical Review Letters, 105(3):030501.
-  arXiv:2112.05037 Martin, J., Micheli, A., and Vennin, V. (2021).
Discord and Decoherence.
arXiv:2112.05037 [astro-ph, physics:hep-th, physics:quant-ph].
-  Martin, J. and Vennin, V. (2016).
Quantum Discord of Cosmic Inflation : Can we Show that CMB
Anisotropies are of Quantum-Mechanical Origin ?
Physical Review D, 93(2):023505.
-  Mukhanov, V. F. and Chibisov, G. V. (1981).
Quantum fluctuations and a nonsingular universe.
ZhETF Pisma Redaktsiiu, 33:549–553.

 Planck-Collaboration, Aghanim, N., Akrami, Y., Ashdown, M., Aumont, J., Baccigalupi, C., Ballardini, M., Banday, A. J., Barreiro, R. B., Bartolo, N., Basak, S., Battye, R., Benabed, K., Bernard, J.-P., Bersanelli, M., Bielewicz, P., Bock, J. J., Bond, J. R., Borrill, J., Bouchet, F. R., Boulanger, F., Bucher, M., Burigana, C., Butler, R. C., Calabrese, E., Cardoso, J.-F., Carron, J., Challinor, A., Chiang, H. C., Chluba, J., Colombo, L. P. L., Combet, C., Contreras, D., Crill, B. P., Cuttaia, F., de Bernardis, P., de Zotti, G., Delabrouille, J., Delouis, J.-M., Di Valentino, E., Diego, J. M., Doré, O., Douspis, M., Ducout, A., Dupac, X., Dusini, S., Efstathiou, G., Elsner, F., Enßlin, T. A., Eriksen, H. K., Fantaye, Y., Farhang, M., Fergusson, J., Fernandez-Cobos, R., Finelli, F., Forastieri, F., Frailis, M., Fraisse, A. A., Franceschi, E., Frolov, A., Galeotta, S., Galli, S., Ganga, K., Génova-Santos, R. T., Gerbino, M., Ghosh, T., González-Nuevo, J., Górski, K. M., Gratton, S., Gruppuso, A., Gudmundsson, J. E., Hamann, J., Handley, W., Hansen, F. K., Herranz, D., Hildebrandt, S. R.,

Hivon, E., Huang, Z., Jaffe, A. H., Jones, W. C., Karakci, A., Keihänen, E., Keskitalo, R., Kiiveri, K., Kim, J., Kisner, T. S., Knox, L., Krachmalnicoff, N., Kunz, M., Kurki-Suonio, H., Lagache, G., Lamarre, J.-M., Lasenby, A., Lattanzi, M., Lawrence, C. R., Jeune, M. L., Lemos, P., Lesgourgues, J., Levrier, F., Lewis, A., Liguori, M., Lilje, P. B., Lilley, M., Lindholm, V., López-Caniego, M., Lubin, P. M., Ma, Y.-Z., Macías-Pérez, J. F., Maggio, G., Maino, D., Mandolesi, N., Mangilli, A., Marcos-Caballero, A., Maris, M., Martin, P. G., Martinelli, M., Martínez-González, E., Matarrese, S., Mauri, N., McEwen, J. D., Meinhold, P. R., Melchiorri, A., Mennella, A., Migliaccio, M., Millea, M., Mitra, S., Miville-Deschénes, M.-A., Molinari, D., Montier, L., Morgante, G., Moss, A., Natoli, P., Nørgaard-Nielsen, H. U., Pagano, L., Paoletti, D., Partridge, B., Patanchon, G., Peiris, H. V., Perrotta, F., Pettorino, V., Piacentini, F., Polastri, L., Polenta, G., Puget, J.-L., Rachen, J. P., Reinecke, M., Remazeilles, M., Renzi, A., Rocha, G., Rosset, C., Roudier, G., Rubiño-Martín, J. A., Ruiz-Granados, B., Salvati, L., Sandri, M., Savelainen, M., Scott,

D., Shellard, E. P. S., Sirignano, C., Sirri, G., Spencer, L. D., Sunyaev, R., Suur-Uski, A.-S., Tauber, J. A., Tavagnacco, D., Tenti, M., Toffolatti, L., Tomasi, M., Trombetti, T., Valenziano, L., Valiviita, J., Van Tent, B., Vibert, L., Vielva, P., Villa, F., Vittorio, N., Wandelt, B. D., Wehus, I. K., White, M., White, S. D. M., Zacchei, A., and Zonca, A. (2020a).

Planck 2018 results. VI. Cosmological parameters.

Astronomy & Astrophysics, 641:A6.



Planck-Collaboration, Akrami, Y., Ashdown, M., Aumont, J., Baccigalupi, C., Ballardini, M., Banday, A. J., Barreiro, R. B., Bartolo, N., Basak, S., Benabed, K., Bersanelli, M., Bielewicz, P., Bond, J. R., Borrill, J., Bouchet, F. R., Boulanger, F., Bucher, M., Burigana, C., Calabrese, E., Cardoso, J.-F., Carron, J., Casaponsa, B., Challinor, A., Colombo, L. P. L., Combet, C., Crill, B. P., Cuttaia, F., de Bernardis, P., de Rosa, A., de Zotti, G., Delabrouille, J., Delouis, J.-M., Valentino, E. D., Dickinson, C., Diego, J. M., Donzelli, S., Doré, O., Ducout, A., Dupac, X.,

Efstathiou, G., Elsner, F., Enßlin, T. A., Eriksen, H. K., Falgarone, E., Fernandez-Cobos, R., Finelli, F., Forastieri, F., Frailis, M., Fraisse, A. A., Franceschi, E., Frolov, A., Galeotta, S., Galli, S., Ganga, K., Génova-Santos, R. T., Gerbino, M., Ghosh, T., González-Nuevo, J., Górski, K. M., Gratton, S., Gruppuso, A., Gudmundsson, J. E., Handley, W., Hansen, F. K., Helou, G., Herranz, D., Hildebrandt, S. R., Huang, Z., Jaffe, A. H., Karakci, A., Keihänen, E., Keskitalo, R., Kiiveri, K., Kim, J., Kisner, T. S., Krachmalnicoff, N., Kunz, M., Kurki-Suonio, H., Lagache, G., Lamarre, J.-M., Lasenby, A., Lattanzi, M., Lawrence, C. R., Jeune, M. L., Levrier, F., Liguori, M., Lilje, P. B., Lindholm, V., López-Caniego, M., Lubin, P. M., Ma, Y.-Z., Macías-Pérez, J. F., Maggio, G., Maino, D., Mandolesi, N., Mangilli, A., Marcos-Caballero, A., Maris, M., Martin, P. G., Martínez-González, E., Matarrese, S., Mauri, N., McEwen, J. D., Meinhold, P. R., Melchiorri, A., Mennella, A., Migliaccio, M., Miville-Deschénes, M.-A., Molinari, D., Moneti, A., Montier, L., Morgante, G., Natoli, P., Oppizzi, F., Pagano, L., Paoletti, D.,

Partridge, B., Peel, M., Pettorino, V., Piacentini, F., Polenta, G.,
Puget, J.-L., Rachen, J. P., Reinecke, M., Remazeilles, M., Renzi,
A., Rocha, G., Roudier, G., Rubiño-Martín, J. A., Ruiz-Granados,
B., Salvati, L., Sandri, M., Savelainen, M., Scott, D., Seljebotn,
D. S., Sirignano, C., Spencer, L. D., Suur-Uski, A.-S., Tauber,
J. A., Tavagnacco, D., Tenti, M., Thommesen, H., Toffolatti, L.,
Tomasi, M., Trombetti, T., Valiviita, J., Tent, B. V., Vielva, P., Villa,
F., Vittorio, N., Wandelt, B. D., Wehus, I. K., Zacchei, A., and
Zonca, A. (2020b).

Planck 2018 results - IV. Diffuse component separation.
Astronomy & Astrophysics, 641:A4.

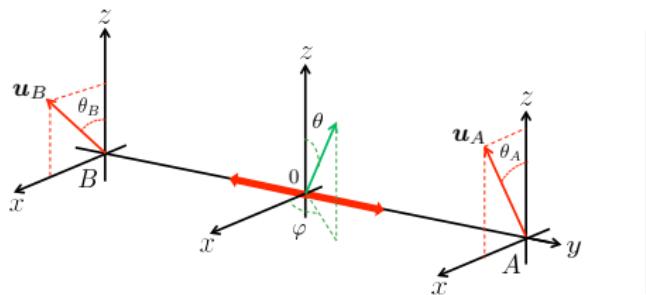
EXTRA

ANOTHER CRITERION : BELL INEQUALITIES

- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.

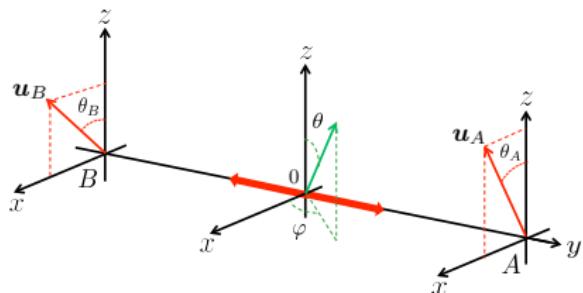
ANOTHER CRITERION : BELL INEQUALITIES

- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Ex: Bell inequalities for two 2-valued spins $\pm 1/2$



ANOTHER CRITERION : BELL INEQUALITIES

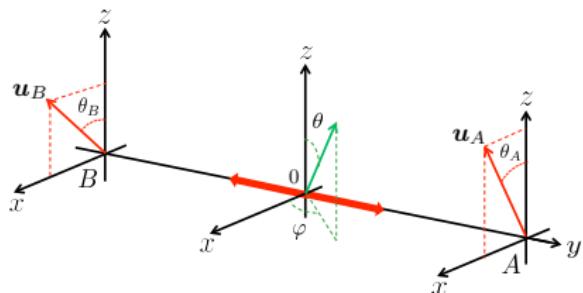
- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Ex: Bell inequalities for two 2-valued spins $\pm 1/2$



→ Smart combination of measurements \mathcal{O} .

ANOTHER CRITERION : BELL INEQUALITIES

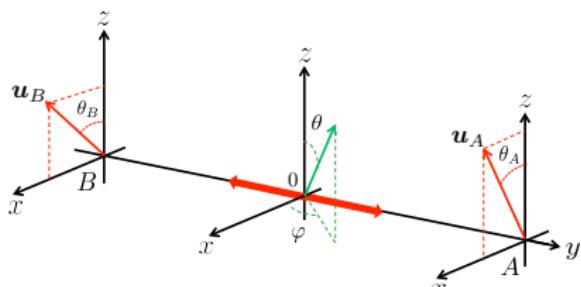
- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Ex: Bell inequalities for two 2-valued spins $\pm 1/2$



- Smart combination of measurements \mathcal{O} .
- Classical local probability for A and B : $\langle \mathcal{O} \rangle \leq 2$

ANOTHER CRITERION : BELL INEQUALITIES

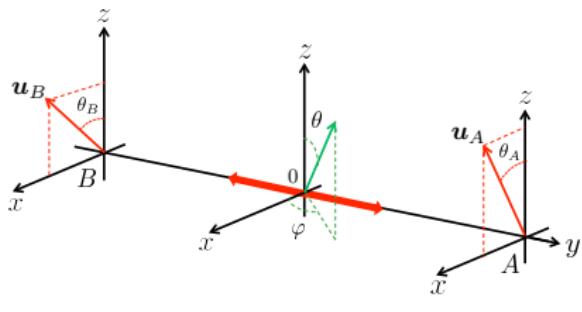
- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Ex: Bell inequalities for two 2-valued spins $\pm 1/2$



- Smart combination of measurements \mathcal{O} .
- Classical local probability for A and B : $\langle \mathcal{O} \rangle \leq 2$
- **Quantum state can reach** $\langle \hat{\mathcal{O}} \rangle = 2\sqrt{2}$

ANOTHER CRITERION : BELL INEQUALITIES

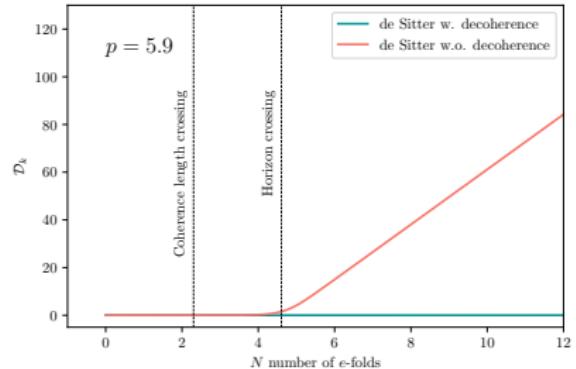
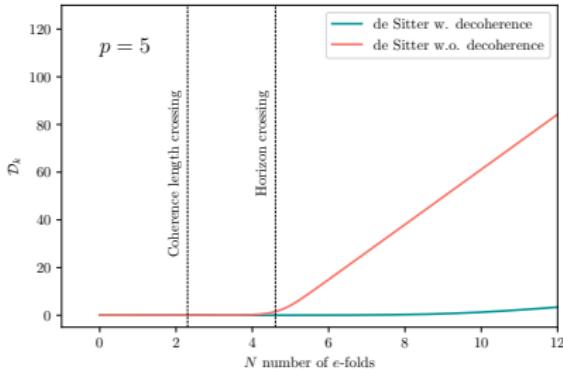
- Quantumness of a state for a system \mathcal{S} = Quantumness of correlations of subsystems $S = \mathcal{S}_1 \cup \mathcal{S}_2$ for this state.
- Ex: Bell inequalities for two 2-valued spins $\pm 1/2$



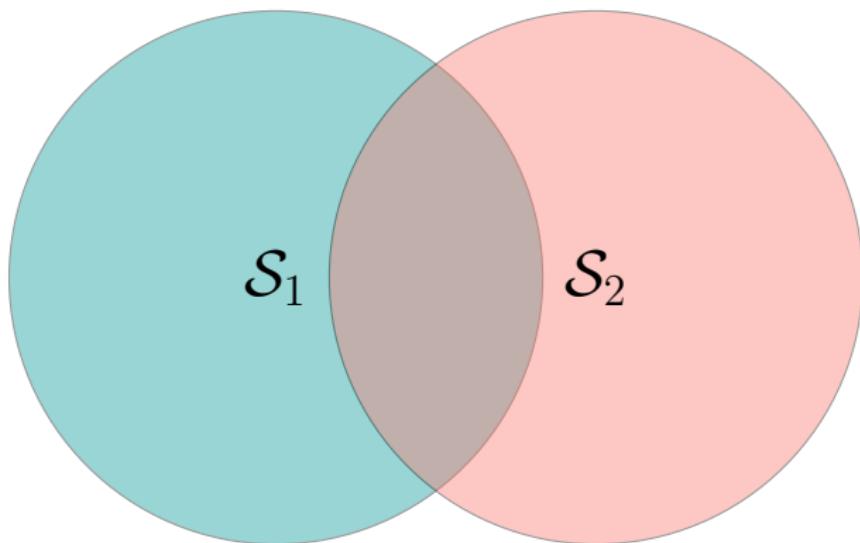
- Smart combination of measurements \mathcal{O} .
- Classical local probability for A and B : $\langle \mathcal{O} \rangle \leq 2$
- **Quantum state can reach $\langle \hat{\mathcal{O}} \rangle = 2\sqrt{2}$**

If measure $\langle \hat{\mathcal{O}} \rangle > 2$, correlations stronger than classical ones
→ quantum state.

GROWTH OF \mathcal{D}_K AND VALUES p



MUTUAL INFORMATION



$$\mathcal{I}(\mathcal{S}_1, \mathcal{S}_2) = H(\mathcal{S}_1) + H(\mathcal{S}_2) - H(\mathcal{S})$$

DECOHERED INFLATIONNARY FLUCTUATIONS

- Environment for \mathcal{S} ? Modeled by Lindblad equation + linear interaction with strength

$$k_{\Gamma}^2 \left(\frac{a}{a_*} \right)^{p-3} H \left(1 - \frac{k\ell_E}{a} \right) . \quad (1)$$

DECOHERENCE, LINDBLAD EQUATION

→ Interaction of \mathcal{S} with the environment \mathcal{E} ?

DECOHERENCE, LINDBLAD EQUATION

→ Interaction of \mathcal{S} with the environment \mathcal{E} ?

$$\hat{H}_{\text{int}} = g \int d^3x \hat{O}_{\mathcal{S}}(\mathbf{x}) \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

DECOHERENCE, LINDBLAD EQUATION

→ Interaction of \mathcal{S} with the environment \mathcal{E} ?

$$\hat{H}_{\text{int}} = g \int d^3x \hat{O}_{\mathcal{S}}(\mathbf{x}) \otimes \hat{O}_{\mathcal{E}}(\mathbf{x}) .$$

→ Evolution of \mathcal{S} only?

DECOHERENCE, LINDBLAD EQUATION

→ Interaction of \mathcal{S} with the environment \mathcal{E} ?

$$\hat{H}_{\text{int}} = g \int d^3x \hat{O}_{\mathcal{S}}(x) \otimes \hat{O}_{\mathcal{E}}(x) .$$

→ Evolution of \mathcal{S} only? **Lindblad equation**

$$\frac{d\hat{\rho}_{\mathcal{S}}}{d\eta} = -g^2 \eta_C \int d^3x d^3y C_{\mathcal{E}}(x, y) \left[\hat{O}_{\mathcal{S}}(x), \left[\hat{O}_{\mathcal{S}}(y), \hat{\rho}_{\mathcal{S}} \right] \right] .$$

DECOHERENCE, LINDBLAD EQUATION

→ Interaction of \mathcal{S} with the environment \mathcal{E} ?

$$\hat{H}_{\text{int}} = g \int d^3x \hat{O}_{\mathcal{S}}(x) \otimes \hat{O}_{\mathcal{E}}(x).$$

→ Evolution of \mathcal{S} only? **Lindblad equation**

$$\frac{d\hat{\rho}_{\mathcal{S}}}{d\eta} = -g^2 \eta_C \int d^3x d^3y C_{\mathcal{E}}(x, y) \left[\hat{O}_{\mathcal{S}}(x), \left[\hat{O}_{\mathcal{S}}(y), \hat{\rho}_{\mathcal{S}} \right] \right].$$

Assumptions:

- Perturbation $g \ll 1$.
- \mathcal{E} stationnary and not perturbed by \mathcal{S} .
- Consider evolution of \mathcal{S} for $\eta \gg \eta_C$ auto-correlation time of \mathcal{E} .

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

→ **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve
Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

- **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.
- \mathcal{E} only correlated over a physical length $\ell_{\mathcal{E}}$ and stat.
homogeneous : $C_{\mathcal{E}}(\mathbf{x}, \mathbf{y}) = \bar{C}_{\mathcal{E}}\Theta(|\mathbf{x} - \mathbf{y}|a/\ell_{\mathcal{E}})$

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

- **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.
- \mathcal{E} only correlated over a physical length $\ell_{\mathcal{E}}$ and stat.
homogeneous : $C_{\mathcal{E}}(\mathbf{x}, \mathbf{y}) = \bar{C}_{\mathcal{E}}\Theta(|\mathbf{x} - \mathbf{y}|a/\ell_{\mathcal{E}})$

$$\frac{d\hat{\rho}_{\pm\mathbf{k}}}{d\eta} = -g^2\eta_C \sqrt{\frac{2}{\pi}} \frac{(2\pi)^{3/2}\ell_{\mathcal{E}}^3}{a^3} \bar{C}_{\mathcal{E}} a^6 \int d^3\mathbf{k} \Theta\left(\frac{k\ell_{\mathcal{E}}}{a}\right) [\hat{v}_{-\mathbf{k}}, [\hat{v}_{\mathbf{k}}, \hat{\rho}_{\pm\mathbf{k}}]]$$

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

- **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.
- \mathcal{E} only correlated over a physical length $\ell_{\mathcal{E}}$ and stat.
homogeneous : $C_{\mathcal{E}}(\mathbf{x}, \mathbf{y}) = \bar{C}_{\mathcal{E}}\Theta(|\mathbf{x} - \mathbf{y}|a/\ell_{\mathcal{E}})$

$$\frac{d\hat{\rho}_{\pm\mathbf{k}}}{d\eta} = - \underbrace{g^2 \eta_C \sqrt{\frac{2}{\pi}} \frac{(2\pi)^{3/2} \ell_{\mathcal{E}}^3}{a^3} \bar{C}_{\mathcal{E}} a^6}_{\equiv k_{\Gamma}^2 \left(\frac{a}{a_*}\right)^{5-3}} \int d^3\mathbf{k} \Theta\left(\frac{k\ell_{\mathcal{E}}}{a}\right) [\hat{v}_{-\mathbf{k}}, [\hat{v}_{\mathbf{k}}, \hat{\rho}_{\pm\mathbf{k}}]]$$

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

- **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.
- \mathcal{E} only correlated over a physical length $\ell_{\mathcal{E}}$ and stat.
homogeneous : $C_{\mathcal{E}}(\mathbf{x}, \mathbf{y}) = \bar{C}_{\mathcal{E}}\Theta(|\mathbf{x} - \mathbf{y}|a/\ell_{\mathcal{E}})$

$$\frac{d\hat{\rho}_{\pm\mathbf{k}}}{d\eta} = -g^2\eta_C \underbrace{\sqrt{\frac{2}{\pi}} \frac{(2\pi)^{3/2}\ell_{\mathcal{E}}^3}{a^3} \bar{C}_{\mathcal{E}} a^6 \int d^3\mathbf{k} \Theta\left(\frac{k\ell_{\mathcal{E}}}{a}\right) [\hat{v}_{-\mathbf{k}}, [\hat{v}_{\mathbf{k}}, \hat{\rho}_{\pm\mathbf{k}}]]}_{\equiv k_{\Gamma}^2 \left(\frac{a}{a_*}\right)^{p-3} \text{ if } g(\eta), \bar{C}_{\mathcal{E}}(\eta)}$$

LINBLAD FOR INFLATIONNARY PERTURBATIONS

Details of interaction when \mathcal{S} is Mukhanov-Sasaki variable \hat{v} ?

- **Linear** $\hat{O}_{\mathcal{S}} = \sqrt{-\det(g_{\mu\nu})}\hat{\phi} = a^4\hat{v}/a$ to preserve Gaussianity, independence of $\pm\mathbf{k}$ pairs $\hat{\rho}_{\mathcal{S}} = \bigotimes_{\mathbf{k} \in \mathbb{R}^3, +} \hat{\rho}_{\pm\mathbf{k}}$.
- \mathcal{E} only correlated over a physical length $\ell_{\mathcal{E}}$ and stat. homogeneous : $C_{\mathcal{E}}(\mathbf{x}, \mathbf{y}) = \bar{C}_{\mathcal{E}}\Theta(|\mathbf{x} - \mathbf{y}|a/\ell_{\mathcal{E}})$

$$\frac{d\hat{\rho}_{\pm\mathbf{k}}}{d\eta} = -g^2\eta_C \underbrace{\sqrt{\frac{2}{\pi}} \frac{(2\pi)^{3/2}\ell_{\mathcal{E}}^3}{a^3} \bar{C}_{\mathcal{E}} a^6}_{\equiv k_{\Gamma}^2 \left(\frac{a}{a_*}\right)^{p-3} \text{ if } g(\eta), \bar{C}_{\mathcal{E}}(\eta)} \int d^3\mathbf{k} \Theta\left(\frac{k\ell_{\mathcal{E}}}{a}\right) [\hat{v}_{-\mathbf{k}}, [\hat{v}_{\mathbf{k}}, \hat{\rho}_{\pm\mathbf{k}}]]$$

- Free-parameters : k_{Γ} and p .

5. [Martin and Vennin, 2016]