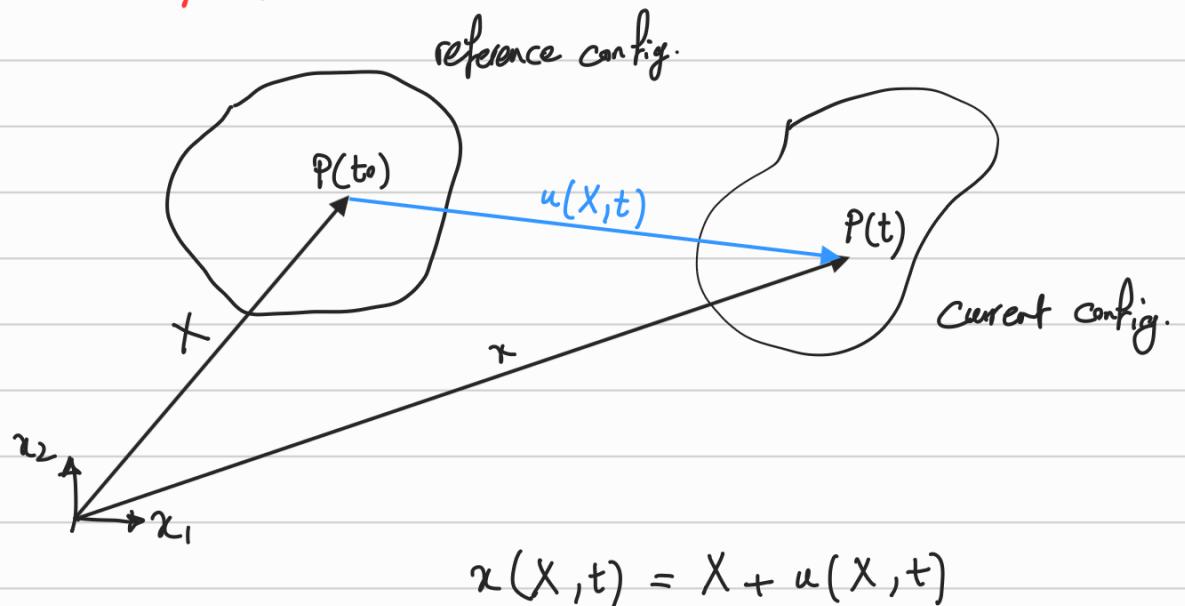


Infinitesimal deformation:



$$\text{material coordinate: } u(x, t) = x(x, t) - x$$

$$\text{special coordinate: } u(x, t) = x - X(x, t)$$

Kinematic Equation for Rigid Body Motion:

- Rigid Body Translation: $x = X + c(t)$ (★)

$$u = x - X = X + c(t) - X = c(t)$$

- Rigid Body Rotation: $x = R(t) X$ center of rotation is at

origin

$$x - b = R(t) X - b \quad (\star\star)$$

Example 3.6.1: Rigid Body Rotation does not change the distance between any pair of material points.

first point: $x' - b = R(t) X' - b \quad (1)$ } pair

Second point: $x^2 - b = R(t) X^2 - b \quad (2)$

$$(2) - (1) : x^2 - x^1 = R(t)(x^2 - x^1)$$

$$\Delta x = R(t) \Delta X$$

let us assume the length of $\Delta x = l$ and $\Delta X = L$

$$\begin{aligned} l^2 &= \Delta x \cdot \Delta x = (R(t) \Delta X) \cdot (R(t) \Delta X) = (\Delta X R^T(t)) \cdot (R(t) \Delta X) \\ &= \Delta X \cdot I \Delta X = \Delta X \cdot \Delta X = L^2 \quad \therefore l^2 = L^2 \\ &\quad l = L \end{aligned}$$

- General Rigid Body Motion

Combining (**) and (***)

$$x = R(t)(X - b) + c(t)$$

Example 3.6-2

$$x = R(t)(X - b) + c(t) \xrightarrow{\partial / \partial t} v = \dot{R}(t)(X - b) + \dot{c}(t) \quad (4)$$

$$x - c(t) = R(t)(X - b) \xrightarrow{R(t)^T} R^T(t)(x - c(t)) = (X - b) \quad (3)$$

substitute (3) in (4) : $v = \dot{R}(t) R^T(t) [x - c(t)] + \dot{c}(t)$

$$\dot{R}(t) R^T(t) = \frac{\partial (R(t) R^T(t))}{\partial t} = \dot{R}(t) R^T(t) + R(t) \dot{R}^T(t) = 0$$

$\underbrace{R(t) R^T(t)}_I = I$

$$\Rightarrow \dot{R}(t) R^T(t) = -R(t) \dot{R}^T(t) = -(\dot{R}(t) R^T(t))^T$$

$\underbrace{\dot{R}(t)}_{T(t)} \quad \underbrace{R^T(t)}_{T(t)}$

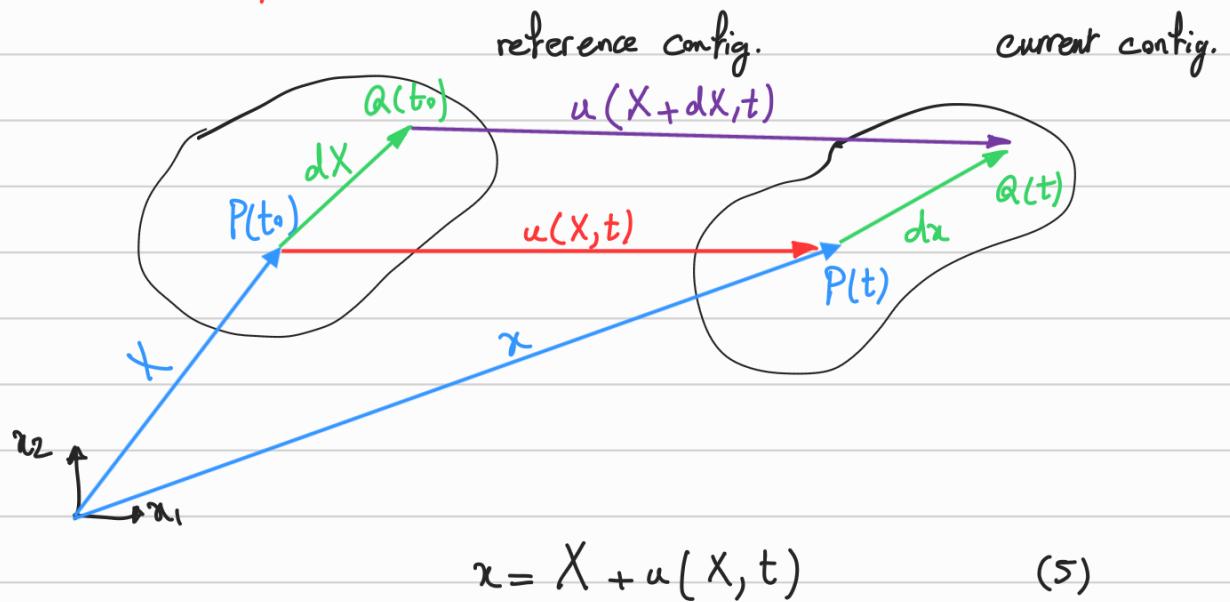
$$T(t) = -(\dot{R}(t) R^T(t))^T$$

$T(t) = \dot{R}(t) R^T(t)$ is antisymmetric tensor

Now that we know $R(t) R(t)^T$ is an antisymmetric tensor, we can write it in form of a dual vector, let's assume $R(t) R(t)^T \equiv \omega$

$$v = \omega \times \underbrace{[x - c(t)]}_{r} + \dot{c}(t) = \omega \times r + \dot{c}(t)$$

Infinitesimal Deformation:



$$x + dx = X + dX + u(X + dx, t) \quad (6)$$

$$(6) - (5) : \left\{ \begin{array}{l} dx = dX + u(X + dx, t) - u(X, t) \\ \nabla u = \frac{u(X + dx, t) - u(X, t)}{dx} \end{array} \right.$$

$$\Rightarrow dx = dX + (\nabla u) dX \quad \nabla u : \text{displacement gradient}$$

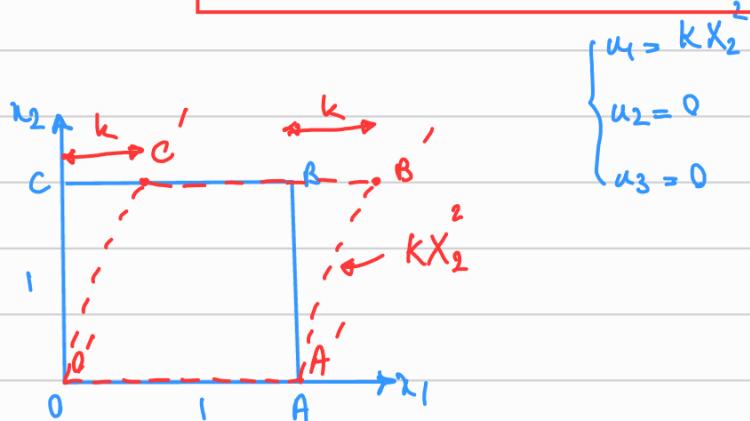
$$\nabla u = \left[\begin{array}{ccc} \frac{\partial u_1}{\partial X_1} & \frac{\partial u_1}{\partial X_2} & \frac{\partial u_1}{\partial X_3} \\ \frac{\partial u_2}{\partial X_1} & \frac{\partial u_2}{\partial X_2} & \frac{\partial u_2}{\partial X_3} \\ \frac{\partial u_3}{\partial X_1} & \frac{\partial u_3}{\partial X_2} & \frac{\partial u_3}{\partial X_3} \end{array} \right] \quad \text{in Cartesian Coordinate System}$$

indicial notation :

$$dx_i = dx_i + \frac{\partial x_i}{\partial x_j} dx_j$$

Example 3.7.1

Please watch the corrected solution of this video!



(a) Every coordinate on OA has an $x_2 = 0 \therefore OA$ remains the same

$$OA = O'A'$$

$$\text{, " , " , CB , " , } x_2 = 1 \therefore C'B' = \begin{cases} k \\ 0 \\ 0 \end{cases}$$

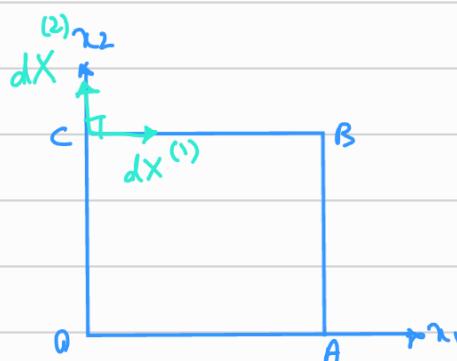
" " " CB " a different x_2 value where $0 \leq x_2 \leq 1$

$$\therefore A'B' = \begin{cases} k x_2^2 \\ 0 \\ 0 \end{cases}$$

" " " OC " a different x_2 value where $0 \leq x_2 \leq 1$

$$\therefore O'C' = \begin{cases} k x_2^2 \\ 0 \\ 0 \end{cases}$$

(b) $dx^{(1)} = dx_1 \hat{e}_1$ & $dx^{(2)} = dx_2 \hat{e}_2$ on point C



$$\nabla u = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix} = \begin{bmatrix} 0 & 2kx_2 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

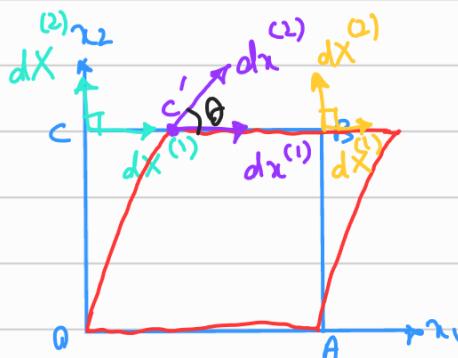
$$dx^{(1)} = dX^{(1)} + (\nabla u) dX^{(4)} = \begin{Bmatrix} dX_1 \\ 0 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 & 2k & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{Bmatrix} \begin{Bmatrix} dX_1 \\ 0 \\ 0 \end{Bmatrix}$$

$$x_2=1$$

$$= \begin{Bmatrix} dX_1 \\ 0 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} dX_1 \\ 0 \\ 0 \end{Bmatrix}$$

$$dx^{(2)} = dX^{(2)} + (\nabla u) dX^{(2)} = \begin{Bmatrix} 0 \\ dX_2 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 & 2k & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{Bmatrix} \begin{Bmatrix} 0 \\ dX_2 \\ 0 \end{Bmatrix}_{x_2=1} \\ = \begin{Bmatrix} 0 \\ dX_2 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 2k dX_2 \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} 2k dX_2 \\ dX_2 \\ 0 \end{Bmatrix}$$

$$dx^{(2)} = \begin{Bmatrix} 2k \\ 1 \\ 0 \end{Bmatrix} |dX_2| \quad dX_2 = \begin{Bmatrix} 0 \\ 1 \\ 0 \end{Bmatrix} |dX_2|$$



$$|dx^{(1)}| = |dX_1|$$

$$|dx^{(2)}| = \sqrt{4k^2 + 1} |dx_2|$$

$$dX^{(1)} = dX_1 \hat{e}_1 \quad \& \quad dX^{(2)} = dX_2 \hat{e}_2 \quad \text{on point } B$$

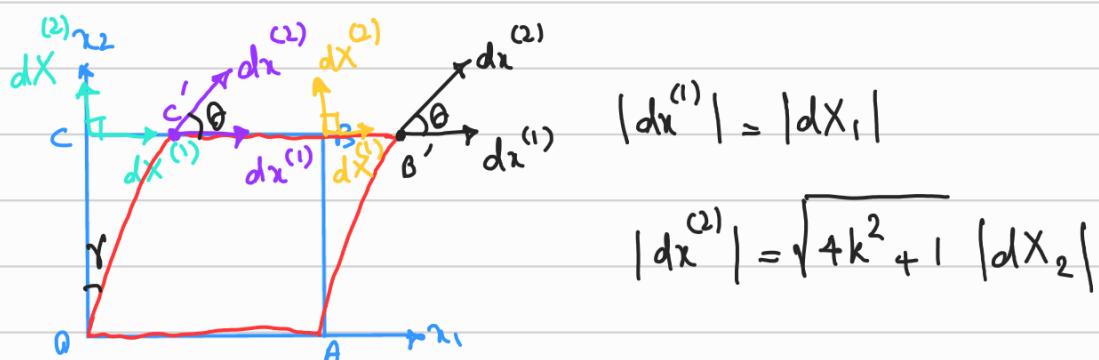
$$dx^{(1)} = dX^{(1)} + (\nabla u) dX^{(2)} = \begin{Bmatrix} dx_1 \\ 0 \\ 0 \end{Bmatrix} + \begin{bmatrix} 0 & 2k & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} dx_1 \\ 0 \\ 0 \end{Bmatrix} \Big|_{x_2=1}$$

$$= \begin{Bmatrix} dX_1 \\ 0 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 0 \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} dx_1 \\ 0 \\ 0 \end{Bmatrix} \quad \therefore \quad \overset{(1)}{dx} = \overset{(1)}{dX}$$

$$dx^{(2)} = dX^{(2)} + (\nabla u) dX^{(1)} = \begin{Bmatrix} 0 \\ dX_2 \\ 0 \end{Bmatrix} + \begin{bmatrix} 0 & 2k & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{Bmatrix} 0 \\ dx_2 \\ 0 \end{Bmatrix} \Big|_{x_2=1}$$

$$= \begin{Bmatrix} 0 \\ dX_2 \\ 0 \end{Bmatrix} + \begin{Bmatrix} 2k dx_2 \\ 0 \\ 0 \end{Bmatrix} = \begin{Bmatrix} 2k \\ 1 \\ 0 \end{Bmatrix} \begin{Bmatrix} dx_2 \\ 0 \\ 0 \end{Bmatrix}$$

$$\overset{(2)}{dx} = \begin{Bmatrix} 2k \\ 1 \\ 0 \end{Bmatrix} \begin{Bmatrix} |dX_2| \\ |dX_2| \\ |dX_2| \end{Bmatrix}, \quad dX^{(2)} = \begin{Bmatrix} 0 \\ 1 \\ 0 \end{Bmatrix} \begin{Bmatrix} |dX_2| \\ |dX_2| \\ |dX_2| \end{Bmatrix}$$



Let's determine θ :

$$\cos \theta = \frac{dX^{(1)} \cdot dX^{(2)}}{|dX^{(1)}| |dX^{(2)}|} = -\frac{\begin{Bmatrix} 1 \\ 0 \\ 0 \end{Bmatrix} (dX_1) \cdot \begin{Bmatrix} 2k \\ 1 \\ 0 \end{Bmatrix} (dX_2)}{|dX_1| \cdot \sqrt{4k^2 + 1} |dX_2|} = -\frac{2k}{\sqrt{4k^2 + 1}}$$

(c) Assuming $k \ll 1$ and $k \neq 0$

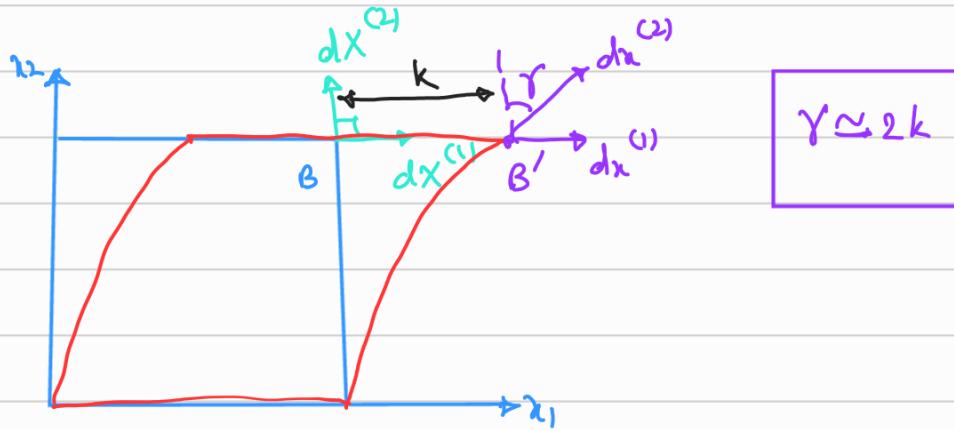
$$\cos \theta = \frac{2k}{\sqrt{4k^2 + 1}}$$

Using Taylor's expansion $\sqrt{1+x} \simeq 1 + \frac{x}{2} + \dots$

$$= \frac{2k}{1 + \frac{2k^2}{0} + \dots} = 2k \rightarrow \cos \theta \simeq 2k$$

$$\theta = \frac{\pi}{2} - \gamma \Rightarrow \cos \theta = \cos \left(\frac{\pi}{2} - \gamma \right) = \sin \gamma \simeq 2k \Rightarrow \gamma = 2k$$

Small deformation $\Rightarrow \sin \gamma \simeq \gamma$



Deformation Gradient Tensor:

$$dx = dX + (\nabla u) dX = I dX + (\nabla u) dX = (I + \nabla u) dX$$

$$dx = F dX \quad \text{where} \quad F = I + \nabla u$$

\uparrow deformation gradient

let us assume that the length of $|dx| = ds$ and $|dX| = dS$

$$dx \cdot dx = (F dX) \cdot (F dX) = (dX F^T) \cdot (F dX) = dX \cdot C dX$$

where $C = F^T F$ (right Cauchy-Green deformation tensor)

$$C = F^T F = (I + \nabla u)^T (I + \nabla u) = (I + \nabla u^T) (I + \nabla u)$$

$$C = I + \nabla u + (\nabla u)^T + (\nabla u)^T (\nabla u)$$

$$C = I + 2E^*$$

where $E^* \equiv \frac{1}{2} [\nabla u + (\nabla u)^T + (\nabla u)^T (\nabla u)]$ Lagrange strain tensor

To get the infinitesimal strain tensor, we neglect the higher-order terms

of displacement gradient in E^*

Infinitesimal strain tensor :

$$E = \frac{1}{2} [\nabla u + (\nabla u)^T]$$

the symmetric part
of ∇u

indicial notation: $E_{ij} = \frac{1}{2} \left[\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right] = \frac{1}{2} [u_{i,j} + u_{j,i}]$

C for the infinitesimal strain :

$$C \simeq I + 2E$$

∴ if we have two elements, dx^1 and dx^2

$$dx^1 \cdot dx^2 = dx^1 \cdot C dx^2 = dx^1 \cdot (I + 2E) dx^2$$

$$dx^1 \cdot dx^2 = dx^1 \cdot dx^2 + 2 dx^1 \cdot E dx^2$$

⇒

Cartesian Coordinates

$$E_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad E_{11} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_1} + \frac{\partial u_1}{\partial x_1} \right) = \frac{\partial u_1}{\partial x_1}$$

$$E_{12} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right)$$

Geometrical Representation / Interpretation of [E]:

Diagonal elements:

$$dX^1 = dX^2 = dX = dS \hat{n} \quad , \quad d\hat{x}^1 = d\hat{x}^2 = dS \hat{m}$$

$$\rightarrow d\hat{x} \cdot d\hat{x} = dX \cdot dX + 2dX \cdot E dX$$

$$ds^2 = dS^2 + 2 dS (n \cdot E n) \Rightarrow ds^2 - dS^2 = 2dS^2 (n \cdot E n)$$

$$\rightarrow \frac{(ds - dS)(ds + dS)}{2dS^2} = n \cdot E n \xrightarrow[\text{strains}]{\text{infinitesimal}} ds \approx dS$$

$$\frac{(ds - dS) 2dS}{2dS^2} = \boxed{\frac{ds - dS}{dS} = \hat{n} \cdot \hat{E} \hat{n}}$$

E_{ii} is the unit elongation for an element in the \hat{x}_i -direction.

$$E_{22} \quad " \quad x_2 \quad "$$

$$E_{33} \quad " \quad x_3 \quad "$$

Off-diagonal elements

$$dX^1 = dS_1 \hat{m} \quad \text{where } \hat{m} \cdot \hat{n} = 0$$

$$dX^2 = dS_2 \hat{n}$$

$$\rightarrow d\hat{x}^1 \cdot d\hat{x}^2 = dX^1 \cdot dX^2 + 2dX^1 \cdot E dX^2$$

$$ds_1 ds_2 \cos\theta = dS_1 dS_2 \hat{m} \cdot \hat{n} + 2 dS_1 dS_2 \hat{m} \cdot \hat{E} \hat{n}$$

$$ds_1 ds_2 \cos\theta = 2 dS_1 dS_2 \hat{m} \cdot \hat{E} \hat{n}$$

Due to the assumption of infinitesimal deformation $\frac{ds_1}{ds_1} = \frac{ds_2}{ds_2} \approx 1$

And since dX^1 and dX^2 are orthogonal $\therefore \theta = \frac{\pi}{2} - \gamma$

$$\therefore \cos \theta = \cos\left(\frac{\pi}{2} - \gamma\right) = \sin \gamma \equiv \gamma$$

We finally get $Y = 2^{\hat{m}} \cdot E^{\hat{n}}$

$2E_{12}$ is the decrease in angle between two elements originally in the γ_1 and γ_2 directions.

Example 3.8.2

$$\begin{cases} u_1 = k(2x_1 + x_2^2), \\ u_2 = k(x_1^2 - x_2^2) \\ u_3 = 0 \end{cases}, \quad k = 10^{-4}$$

$$(a) \quad dx^i = dX_i e_i^{\wedge}$$

$$E = \frac{1}{2} \left[\nabla u + (\nabla u)^T \right]$$

$$dx^2 = dx, e^2$$

$$\nabla u = \begin{bmatrix} 2k & 2kx_2 & 0 \\ 2kx_1 & -2kx_2 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\nabla u = \begin{bmatrix} 2k & -2k & 0 \\ 2k & 2k & 0 \\ 0 & 0 & 0 \end{bmatrix} k$$

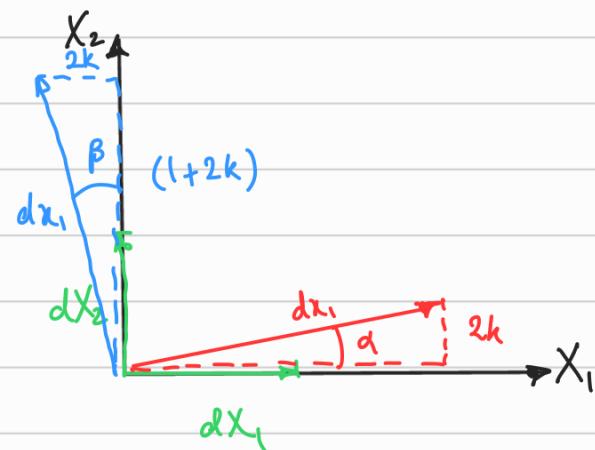
$$E = \frac{1}{2} \left\{ \left[\begin{array}{ccc} 2k & -2k & 0 \\ 2k & 2k & 0 \\ 0 & 0 & 0 \end{array} \right] + \left[\begin{array}{ccc} 2k & 2k & 0 \\ -2k & 2k & 0 \\ 0 & 0 & 0 \end{array} \right] \right\} = \left[\begin{array}{ccc} 2k & 0 & 0 \\ 0 & 2k & 0 \\ 0 & 0 & 0 \end{array} \right]$$

$$(b) \quad dx^1 = dX^1 + \nabla u \cdot dX^1 = \left\{ \begin{array}{c} dX_1 \\ 0 \\ 0 \end{array} \right\} + \left\{ \begin{array}{c} 2k & -2k & 0 \\ 2k & 2k & 0 \\ 0 & 0 & 0 \end{array} \right\} \left\{ \begin{array}{c} dX_1 \\ 0 \\ 0 \end{array} \right\}$$

$$= \left\{ \begin{array}{c} dX_1 + 2k dX_1 \\ 2k dX_1 \\ 0 \end{array} \right\} = \left\{ \begin{array}{c} (1+2k) dX_1 \\ 2k dX_1 \\ 0 \end{array} \right\}$$

$$dx^2 = dX^2 + \nabla u \cdot dX^2 = \left\{ \begin{array}{c} 0 \\ dX_2 \\ 0 \end{array} \right\} + \left\{ \begin{array}{c} 2k & -2k & 0 \\ 2k & 2k & 0 \\ 0 & 0 & 0 \end{array} \right\} \left\{ \begin{array}{c} dX_2 \\ 0 \\ 0 \end{array} \right\}$$

$$= \left\{ \begin{array}{c} -2k dX_2 \\ dX_2 + 2k dX_2 \\ 0 \end{array} \right\} = \left\{ \begin{array}{c} -2k dX_2 \\ (1+2k) dX_2 \\ 0 \end{array} \right\}$$

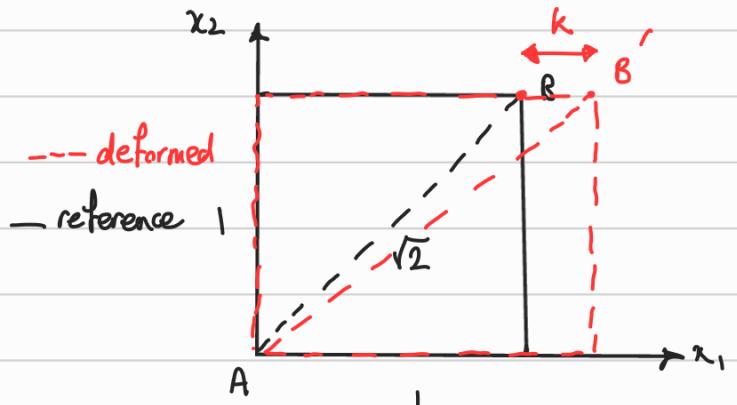


$$\tan \alpha \equiv \alpha = \frac{2k dX_1}{(1+2k) dX_1} = 2k \left(\frac{1}{1+2k} \right) = 2k (1 + \dots) \approx 2k$$

$$\beta = \frac{(2k) dX_2}{(1+2k) dX_2} \approx 2k$$

Example 3.8.3

$$\begin{cases} u_1 = kx_1, & k = 10^{-4} \\ u_2 = 0 \\ u_3 = 0 \end{cases}$$



$$(a) \nabla u = \begin{bmatrix} k & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$E = \frac{1}{2} (\nabla u + (\nabla u)^T) = \begin{bmatrix} k & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

the unit elongation in x_1 -direction

We have to find the unit vector in \vec{AB} direction, \hat{n}_{AB}

$$n \cdot E n_{AB} = \left\{ \frac{\sqrt{2}}{2} \quad \frac{\sqrt{2}}{2} \quad 0 \right\} \cdot \begin{bmatrix} k & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \left\{ \begin{array}{l} \frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} \\ 0 \end{array} \right\}$$

$$\left\{ \frac{\sqrt{2}}{2} \quad \frac{\sqrt{2}}{2} \quad 0 \right\} \cdot \left\{ \begin{array}{l} \frac{k\sqrt{2}}{2} \\ 0 \\ 0 \end{array} \right\} = \frac{k}{2} \quad \text{the unit elongation in } \hat{n}_{AB} \text{-direction}$$

elongation in x_2 -direction:

$$\left\{ 0 \quad 1 \quad 0 \right\} \cdot \begin{bmatrix} k & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \left\{ \begin{array}{l} 0 \\ 1 \\ 0 \end{array} \right\} = \left\{ 0 \quad 1 \quad 0 \right\} \cdot \left\{ \begin{array}{l} 0 \\ 0 \\ 0 \end{array} \right\} = 0$$

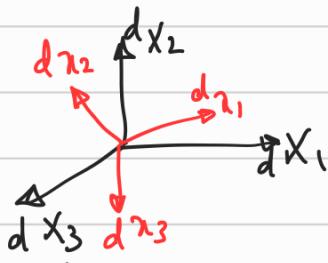
Principal strain:

$$[E] = \begin{bmatrix} E_1 & 0 & 0 \\ 0 & E_2 & 0 \\ 0 & 0 & E_3 \end{bmatrix}$$

symmetric tensor

Dilatation:

$$dV_1 = dS_1 dS_2 dS_3$$



$$dV_2 = (dS_1 + dS_1 E_1) (dS_2 + dS_2 E_2) (dS_3 + dS_3 E_3)$$

$$= dS_1 dS_2 dS_3 (1 + E_1) (1 + E_2) (1 + E_3)$$

$$(\text{dilatation}) e = \frac{dV_2 - dV_1}{dV_1} = \frac{dS_1 dS_2 dS_3 (1 + E_1) (1 + E_2) (1 + E_3) - dS_1 dS_2 dS_3}{dS_1 dS_2 dS_3}$$

$$= (1 + E_1) (1 + E_2) (1 + E_3) - 1$$

$$= E_1 + E_2 + E_3 + \text{higher-order terms}$$

Since we are dealing with infinitesimal deformation, we can neglect the higher-order terms.

$$e = E_1 + E_2 + E_3 = E_{ii} = \text{div}(u) = \frac{\partial u_i}{\partial x_i}$$

Infinitesimal Rotation Tensor (Antisymmetric part of ∇u):

$$du = dX + (\nabla u) dX = dX + (E + \Omega) dX$$

symmetric part Antisymmetric part of ∇u
of ∇u

$$\Omega = \frac{1}{2} [\nabla u - (\nabla u)^T] \quad \therefore \quad \Omega dX = t^A x dX$$

where t^A is the dual vector of Ω and is

$$t^A = -(\Omega_{23} \hat{e}_1^1 + \Omega_{31} \hat{e}_2^1 + \Omega_{12} \hat{e}_3^1)$$

$$\nabla u = E + \Omega$$

elongation

change in angle

* Zero dilatation or $\text{div}(u) = 0$ does not indicate zero deformation.