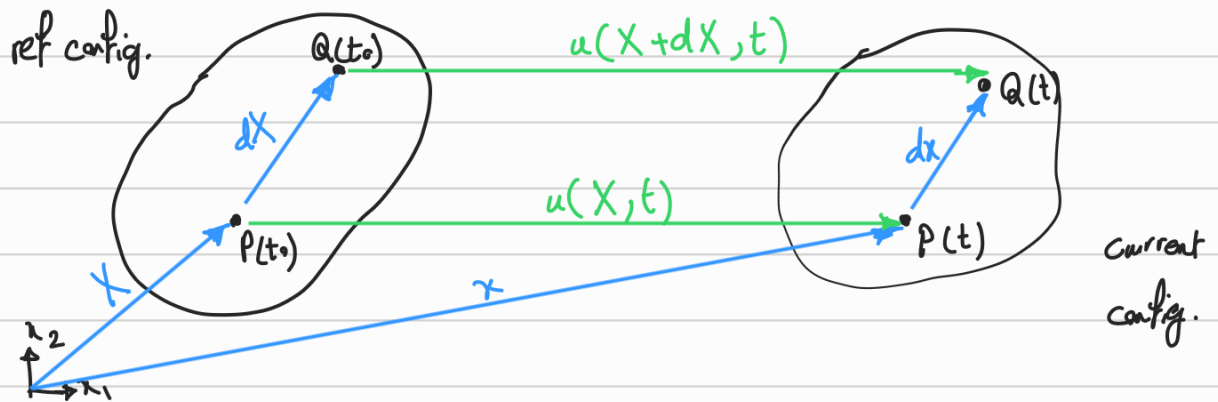


Deformation gradient  $\left\{ \begin{array}{l} \text{local rigid body motion} \\ \text{pure stretch} \end{array} \right.$

Deformation Gradient:



$$dx = x(X + dX, t) - x(X, t) = (\nabla x) dX, \quad \boxed{dx = F dX}$$

$$F = \nabla x$$

$$\nabla x = \begin{bmatrix} \frac{\partial x_1}{\partial X_1} & \frac{\partial x_1}{\partial X_2} & \frac{\partial x_1}{\partial X_3} \\ \frac{\partial x_2}{\partial X_1} & \frac{\partial x_2}{\partial X_2} & \frac{\partial x_2}{\partial X_3} \\ \frac{\partial x_3}{\partial X_1} & \frac{\partial x_3}{\partial X_2} & \frac{\partial x_3}{\partial X_3} \end{bmatrix}$$

$$x = X + u(X, t)$$

$$x + dx = X + dX + u(X + dX, t)$$

$$\left. \begin{array}{l} x = X + u(X, t) \\ x + dx = X + dX + u(X + dX, t) \end{array} \right\} \begin{array}{l} dx = dX + u(X + dX, t) - u(X, t) \\ dx = dX + \nabla u dX \end{array}$$

$$\boxed{dx = (I + \nabla u) dX}$$

1. There's this transformation of each node from reference to current configuration and vice-versa. Therefore,  $F^{-1}$  exists.

2. When the reference configuration is deformed, it deforms such that the initial unit vectors elongate and rotate. If we have  $e_1 \cdot (e_2 \times e_3) = e_2 \cdot (e_3 \times e_1) = e_3 \cdot (e_1 \times e_2)$  then it remains the same in the current configuration; therefore, for the transformation  $F$

$$Fe_1 \cdot (Fe_2 \times Fe_3) = Fe_2 \cdot (Fe_3 \times Fe_1) = Fe_3 \cdot (Fe_1 \times Fe_2) \\ = \det(F) > 0$$

If  $\det(F) < 0$  then it shows that the particle does not exist in the current configuration. If  $\det(F) < 0$  then we cannot simply move back from current to the reference because the particle does not exist.

(Assumption: We use right-handed basis)

Local Rigid Body Motion:

Neither length nor angle of element are changed.

$\det(F) > 0$  and  $F^{-1}$  exists.

$$FF^T = I \quad \text{and} \quad \det(F) = 1$$

Pure Stretch:

Pure Stretch is a deformation tensor that is positive definite.

$$\forall \vec{a} : \vec{a} \cdot U \vec{a} \geq 0 \quad \text{where } U = \text{pure stretch}$$

It can be shown that if a tensor is positive definite and symmetric, all of its eigenvalues are positive.

If we rewrite  $U$  in the form of its eigenvalues,

$$U = \begin{bmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{bmatrix}$$

$$dx = U dX \Rightarrow \begin{Bmatrix} dx^1 \\ dx^2 \\ dx^3 \end{Bmatrix} = \begin{bmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{bmatrix} \begin{Bmatrix} dX^1 \\ dX^2 \\ dX^3 \end{Bmatrix}$$

$$\begin{Bmatrix} dx^1 \\ dx^2 \\ dx^3 \end{Bmatrix} = \begin{Bmatrix} \lambda_1 dX^1 \\ \lambda_2 dX^2 \\ \lambda_3 dX^3 \end{Bmatrix}$$

$$\lambda_1 = \frac{|dx^1|}{|dX^1|}, \quad \lambda_2 = \frac{|dx^2|}{|dX^2|}, \quad \lambda_3 = \frac{|dx^3|}{|dX^3|}$$