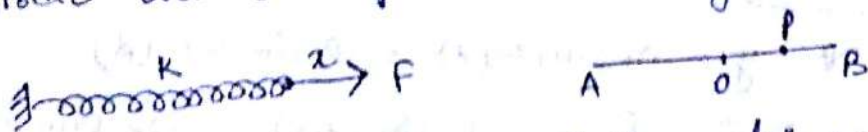


SHM Motion: Translation, rotation, vibration/oscillation
 periodic motion $f(t) = f(t+T)$ e.g. $\sin \frac{2\pi t}{T}$, $\cos \frac{2\pi t}{T}$

if periodic over same path \rightarrow oscillatory motion

Elasticity & Inertia



SHM is a special type of periodic motion where restoring force on particle is proportional to displacement & directed to the mean position.

oscillation between point A & B, mean position O. at time t , particle is at P & displacement is x . F = restoring force

$$F \propto -x \quad \text{or} \quad F = -kx \quad \text{or} \quad ma = -kx$$

"Small oscillation approximation"

$$\therefore a = -\frac{k}{m}x = -\omega^2 x$$

Characteristics

- (1) linear motion \rightarrow to-n-fro in straight line.
- (2) $F \propto -x$.

Linear harmonic motion \leftrightarrow angular harmonic motion.
 (pendulum) (torsional pendulum)
 $F \propto -x$ $\tau \propto -\theta$

Complete oscillation: one point to same point. (time period)

amplitude: maximum displacement on both sides.

frequency: no. of oscillations in 1 second.

phase: displacement, velocity, acceleration & direction of motion. After 1 oscillation, phase is same.

$t=0$, initial phase.

Relation between SHM & uniform circular motion.

$$OA = x, OB = y$$

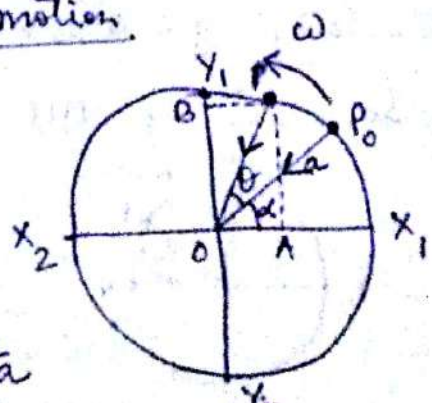
$$\theta = \omega t$$

$$s = a\theta$$

$$= OP \cos(\theta + \alpha) = a \cos(\theta + \alpha)$$

$$= a \cos(\omega t + \alpha)$$

$$\text{Speed } v = \omega a, \text{ Centripetal acc } f_r = \frac{v^2}{a} = \omega^2 a$$



Acceleration of A is component of f_r along X_1OX_2 .

$$f_A = -f_r \cos(\omega t + \alpha) = -\omega^2 a \cos(\omega t + \alpha) = -\omega^2 x$$

$$\therefore f_A \propto -x.$$

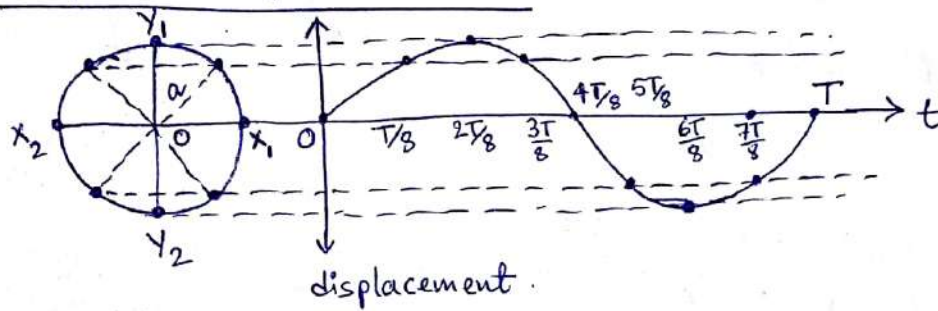
Similarly, $OB = y = OP \sin(\theta + \alpha) = a \sin(\omega t + \alpha)$

Acceleration of B is $f_B = -f_r \sin(\theta + \alpha) = -\omega^2 a \sin(\omega t + \alpha) = -\omega^2 y$

$$\therefore f_B \propto -y.$$

\therefore SHM is defined as the projection of uniform circular motion along diameter of circle.

Graphical representation



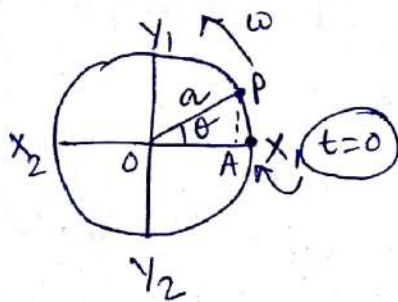
Time period = T.

$$y = a \sin \frac{2\pi}{T} t$$

(SHM along y-axis)

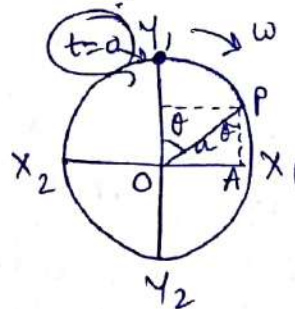
Displacement

In SHM, displacement at time t is the distance of the particle from the mean position.



$$OA = OP \cos \theta$$

$$x = a \cos \omega t$$

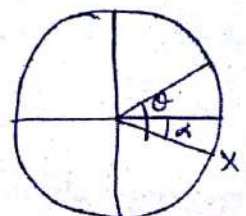
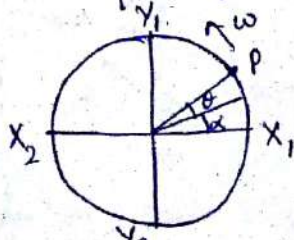


$$OA = OP \cos (\frac{\pi}{2} - \theta)$$

$$x = a \sin \theta = a \sin \omega t$$

Similarly, $y = a \cos \omega t$ & $y = a \sin \omega t$.

So, eqⁿ of SHM can be derived from any instant t .



$$x = a \cos(\theta + \alpha) = a \cos(\omega t + \alpha)$$

Similarly, $x = a \sin(\theta + \alpha) = a \sin(\omega t + \alpha)$.

If initial position is x_1 (2nd pic) then $x = a \cos(\omega t - \alpha)$
 $\therefore x = a \sin(\omega t - \alpha)$

Velocity & acceleration

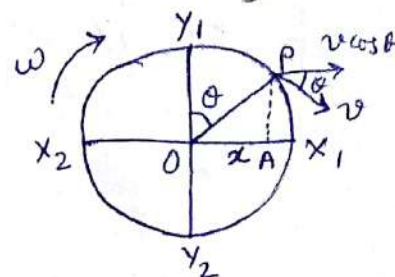
velocity of SHM is component of the particle's velocity along x-axis at time t.

$$V = a\omega, \quad V \text{ parallel to } OA = v \cos \theta$$

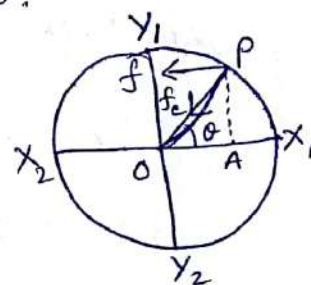
$$= a\omega \cos \theta = a\omega \sqrt{1 - \frac{x^2}{a^2}}$$

$$\therefore \boxed{v = \omega \sqrt{a^2 - x^2}}$$

v_{\max} is at $x=0$, $v_{\max} = a\omega$. $\& \ x=a$, $v_{\min} = 0$.



$$x = a \sin \theta$$



Same with acceleration \Rightarrow SHM is the projection along x-axis is component of acceleration along x-axis. $f_c = -\omega^2 a$ & component around $x_1 x_2$ is $-\omega^2 a \cos \theta = -\omega^2 a \cos \omega t = -\omega^2 x$.

$$\therefore f = -\omega^2 x$$

$$f_{\max} = -\omega^2 a \text{ when } x = \pm a, \quad f_{\min} = \pm \omega^2 a$$

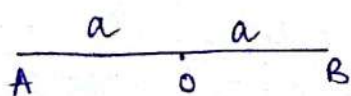
$$f_{\min} = 0 \text{ when } x = 0$$

Calculus: $x = a \sin \omega t$, $v = \dot{x} = a\omega \cos \omega t = a\omega \sqrt{1 - \frac{x^2}{a^2}}$
 $= \omega \sqrt{a^2 - x^2}$

$$f = \ddot{x} = -a\omega^2 \sin \omega t = -\omega^2 x$$

$$\omega^2 = f/x \text{ (neglect)}$$

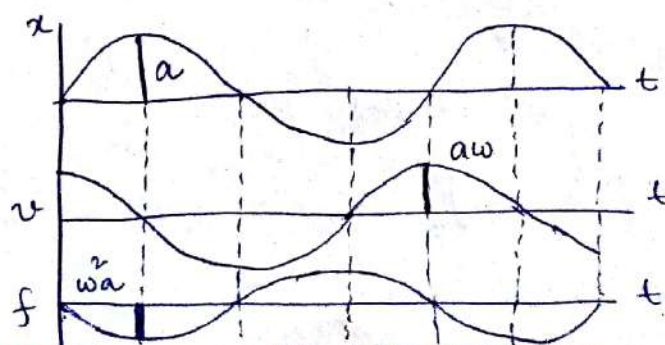
Time period $T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{a}{f}}$



$$x = a \sin \omega t = a \sin \frac{2\pi}{T} t$$

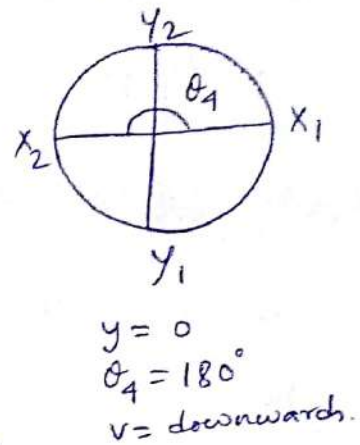
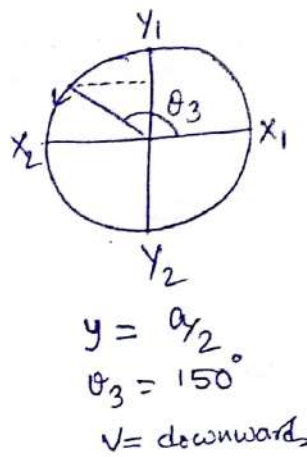
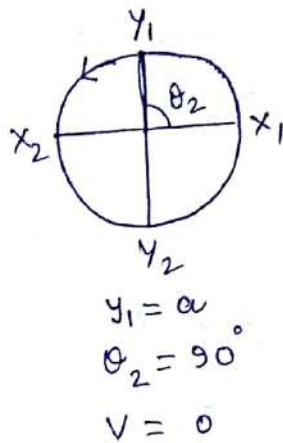
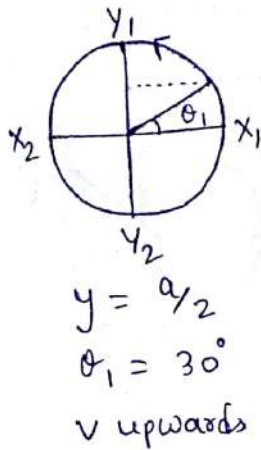
$$v = a\omega \cos \omega t = a\omega \cos \frac{2\pi}{T} t$$

$$f = -a\omega^2 \sin \omega t = -a\omega^2 \sin \frac{2\pi}{T} t$$



Phase

you see, a & ω (angular velocity) are constant.
(amplitude) $\theta = \omega t$ is changing = phase.



phase difference

2 particles.

$$\phi = \theta_1 - \theta_2 = 0 \quad (\text{in phase})$$

$$= 180^\circ \quad (\text{out of phase})$$

Differential form & solution

Homogeneous, 2nd order, CDE with constant coefficient

$$F = -kx \quad \text{or} \quad m\ddot{x} = -kx \quad \text{or} \quad \ddot{x} + \omega^2 x = 0, \quad \omega = \sqrt{\frac{k}{m}}$$

Solution: Multiply by $2\dot{x}$, $2\dot{x}\ddot{x} + 2\omega^2 x\dot{x} = 0$

Integrating $\dot{x}^2 = -\omega^2 x^2 + C$

when displacement is maximum, $x = a$, $\dot{x} = 0 \Rightarrow C = \omega^2 a^2$

$$\therefore v = \dot{x} = \pm \omega \sqrt{a^2 - x^2}$$

$$\text{or } \pm \frac{dx}{\sqrt{a^2 - x^2}} = \omega dt, \quad \text{Integrating } \sin^{-1} \frac{x}{a} = \omega t + \phi$$

$$\text{or } x = a \sin(\omega t + \phi)$$

See, $x = a \cos(\omega t + \phi)$ also satisfy $\ddot{x} + \omega^2 x = 0$.

$$x = a \sin(\omega t + \phi) = a \sin \omega t \cos \phi + a \cos \omega t \sin \phi$$

$$= A \sin \omega t + B \cos \omega t.$$

In operator form, $\frac{d^2 x}{dt^2} = D^2 x, \quad \frac{dx}{dt} = Dx$

$$D^2 x + \omega^2 x = 0 \quad \text{or} \quad D^2 = -\omega^2 \quad \text{or} \quad D = \pm i\omega$$

\therefore General solution $x = A e^{i\omega t} + B e^{-i\omega t}$

For real value of x , $A = B^*$ $A = a+ib$, $B = a-ib$

you can also have $x = ae^{i(\omega t + \phi)}$

Sinusoidal or cosinusoidal.

CW 1. Oscillatory motion of a particle is represented by $x = ae^{i\omega t}$. Establish the motion is SHM. Similarly if $x = a\cos\omega t + b\sin\omega t$ then SHM.

$$x = ae^{i\omega t}, \quad \dot{x} = ai\omega e^{i\omega t}, \quad \ddot{x} = -a\omega^2 e^{i\omega t} = -\omega^2 x \quad (\text{SHM})$$

$$x = a\cos\omega t + b\sin\omega t, \quad \dot{x} = -a\omega\sin\omega t + b\omega\cos\omega t$$

$$\ddot{x} = -a\omega^2\cos\omega t - b\omega^2\sin\omega t = -\omega^2 x \quad (\text{SHM})$$

2. Which periodic motion is not oscillatory?

→ earth around sun or moon around earth.

3. Dimension of force constant of vibrating spring.

$$F = -Kx \quad [K] = \frac{[\text{Force}]}{[\text{displacement}]} = \frac{[\text{Newton}]}{[\text{metre}]}$$

also called
"stiffness"

$$= \frac{[MLT^{-2}]}{[L]} = [MT^{-2}]$$

HW 1. In SHM, displacement is $x = a\sin(\omega t + \phi)$. at $t=0$, $x=x_0$ with velocity v_0 . show that $a = \sqrt{x_0^2 + \frac{v_0^2}{\omega^2}}$ & $\tan\phi = \frac{\omega x_0}{v_0}$.

2. Particle is vibrated at frequency 5 Hz in SHM. Show that when displacement exceeds 10^{-2} metre, the particle loses contact with the vibrator. Given $g = 9.8 \text{ m/s}^2$

3. In SHM, a particle has speed 80 cm/s & 60 cm/s with displacement 3 cm & 4 cm. Calculate amplitude of vibration

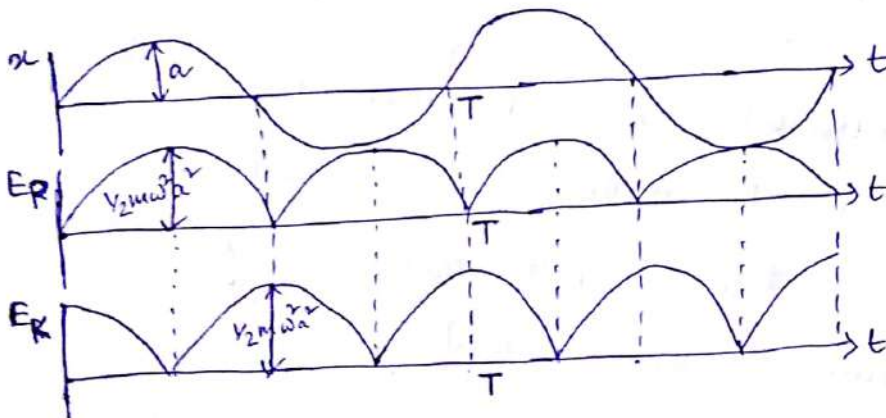
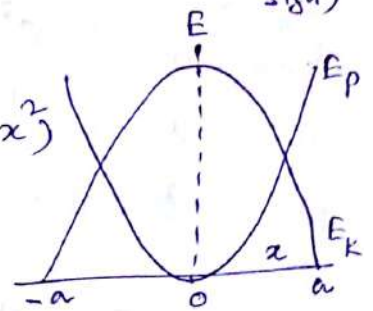
Energy of a particle in SHM

Work is done on particle to displace \rightarrow restoring force. So P.E. in spring stored & motion is K.E. Total energy constant

P.E. $F = mf = -m\omega^2 x \therefore dW = Fdx = m\omega^2 x dx$ (against so no -ive sign)
 $\therefore E_p = \int_0^x m\omega^2 x dx = \frac{1}{2} m\omega^2 x^2$

K.E. $v = \omega \sqrt{a^2 - x^2}$, $E_k = \frac{1}{2} mv^2 = \frac{1}{2} m\omega^2 (a^2 - x^2)$

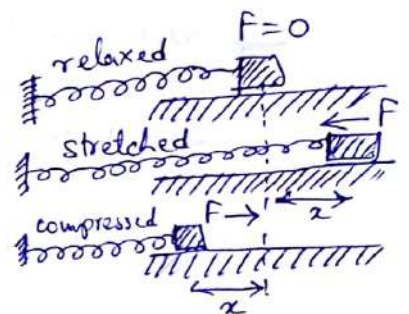
$E_{Tot} = E_k + E_p = \frac{1}{2} m\omega^2 a^2 = \text{constant.}$



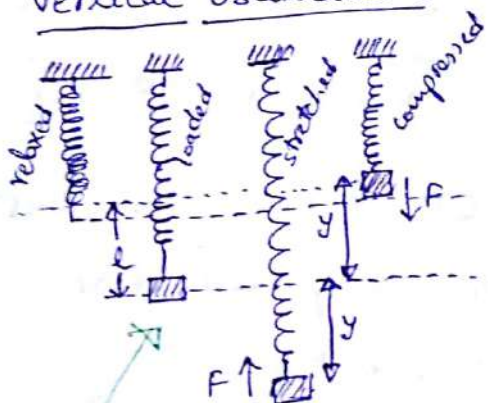
Examples of SHM

Horizontal oscillations

$F = -Kx = m\ddot{x}$
 $\ddot{x} + \omega^2 x = 0$ $\omega = \sqrt{\frac{K}{m}}$
 $x = A \cos(\omega t + \phi)$, $T = 2\pi \sqrt{\frac{m}{K}}$
initial cond. material.



Vertical oscillations



static equilibrium

Tension on spring $F_0 = Kl$

force on mass = mg .

Static eq. $mg = Kl$.

stretched tension on spring = $K(l+y)$

$mg - F = K(l+y) = Kl + Ky$
 $= \cancel{mg} + Ky$

$F = -Ky$.

compressed

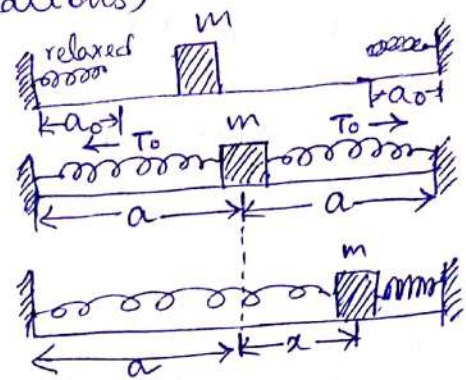
$mg + F = K(l-y) = \cancel{mg} - Ky$

$F = -Ky$.

Two spring system (Longitudinal oscillations)

horizontal frictionless surface,
rigid wall, massless spring,
relaxed length a_0 .

After connection, static equilibrium



$$T_0 = K(a - a_0)$$

x = displacement to right. restoring force by left spring $-K(a + x - a_0)$
force on right spring $K(a - x - a_0)$

$$\therefore F_x = K(a - x - a_0) - K(a + x - a_0) = -2Kx$$

$$m\ddot{x} = -2Kx \quad \text{or} \quad \ddot{x} + \omega^2 x = 0$$

$$\omega = \sqrt{\frac{2K}{m}}$$

$$T_{\text{long}} = 2\pi \sqrt{\frac{m}{2K}}$$

Two spring system (Transverse oscillations)

$$T_0 = K(a - a_0)$$

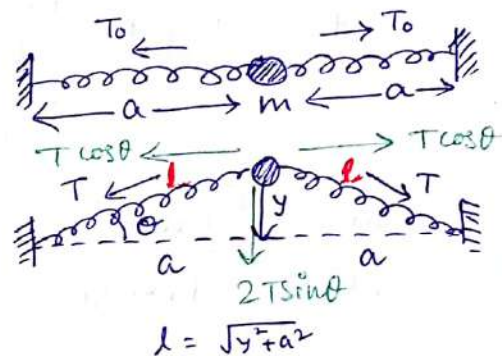
$$T = K(l - a_0)$$

$$F_y = -2T \sin \theta = -2T \frac{y}{l}$$

$$\therefore m\ddot{y} + \frac{2T}{l}y = 0 \quad \text{or} \quad \ddot{y} + \omega^2 y = 0$$

$$\omega^2 = \frac{2T}{ml} = \frac{2K(l - a_0)}{ml}, \quad \text{but } l = f(y).$$

$$\text{So } \ddot{y} + \frac{2K}{m} \left(1 - \frac{a_0}{f(y)}\right) y = 0 \text{ is not a SHM.}$$



$$l = \sqrt{y^2 + a^2}$$

① slinky approximation $a \gg a_0$ or $\frac{a_0}{a} \ll 1$.

$$\omega^2 = \frac{2K}{m} \left(1 - \frac{a_0}{l}\right) = \frac{2K}{m} \left(1 - \frac{a_0}{a} \frac{a}{l}\right) \quad \text{as } l > a.$$

$$= \frac{2K}{m}. \quad \text{Then SHM.}$$

$$\omega = \sqrt{\frac{2K}{m}}$$

$$T_{\text{trans}} = 2\pi \sqrt{\frac{m}{2K}}$$

"large" harmonic oscillations

② small oscillation approximation $a \not\gg a_0$ but $y \ll a$ or l .

$$\therefore l = \sqrt{y^2 + a^2} = a \sqrt{\frac{y^2}{a^2} + 1} \approx a$$

$$\text{Then also } \omega^2 = \frac{2K}{m} \left(1 - \frac{a_0}{a}\right)$$

SHM

$$T_{\text{trans}} = 2\pi \sqrt{\frac{m}{2K \left(1 - \frac{a_0}{a}\right)}}$$

$$\therefore T_{\text{long}} = \sqrt{1 - \frac{a_0}{a}} T_{\text{trans}}$$

So longitudinal is faster than transverse.

Simple pendulum

$$F' = mg \cos \theta$$

(tension in string)

$$F = -mg \sin \theta$$

(restoring force)

$$= -mg \left(\theta - \frac{\theta^3}{3!} + \frac{\theta^5}{5!} - \dots \right) \approx -mg\theta \quad \left[\lim_{\theta \rightarrow 0} \right]$$

$$\text{or, } m\ddot{x} = -mg \frac{x}{l} \quad \text{or} \quad \ddot{x} + \frac{g}{l}x = 0.$$

$$\therefore \omega = \sqrt{\frac{g}{l}}, \quad T = 2\pi \sqrt{\frac{l}{g}}. \quad (\text{mass independent})$$

String tension when pendulum at mean position

$$F' = mg + \frac{mv^2}{l}$$

(centrifugal force)

at A, Energy = KE + PE = 0 + mgh = mgh

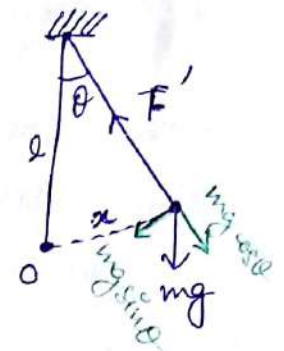
at O, Energy = KE + PE = $\frac{1}{2}mv^2 + 0 = \frac{1}{2}mv^2$

Conservation of energy $\Rightarrow \frac{1}{2}mv^2 = mgh$ or $v^2 = 2gh$.

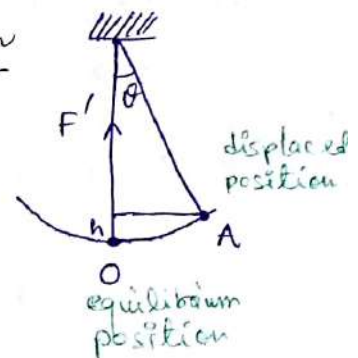
$$\text{or } v^2 = 2g(l - l \cos \theta) = 2gl(1 - \cos \theta) = 2gl \times 2 \sin^2 \frac{\theta}{2}$$

$$\approx 4gl \left(\frac{\theta}{2} \right)^2 = gl\theta^2.$$

$$\therefore F' = mg + \frac{m}{l} gl\theta^2 = mg(1 + \theta^2).$$



$$x = l\theta$$



Compound Pendulum

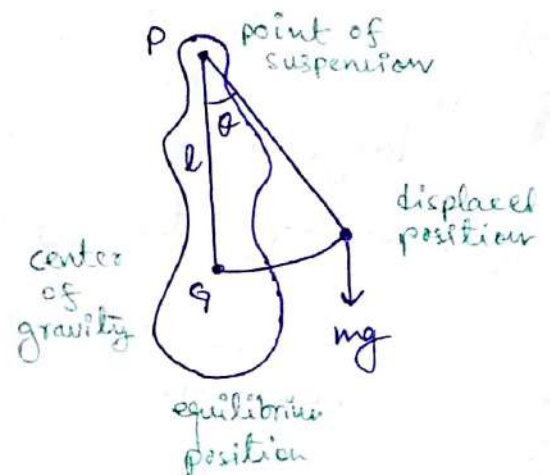
arbitrary shaped rigid body oscillating about a horizontal axis passing through it.

restoring force \leftrightarrow reactive couple or torque

moment of restoring force

$$= -mgl \sin \theta$$

angular acceleration $\alpha = \frac{d^2\theta}{dt^2}$, moment of inertia = I.



$$\tau = I\alpha = I \frac{d^2\theta}{dt^2} = -mgl \sin\theta$$

$$\text{or } \frac{d^2\theta}{dt^2} = -\frac{mgl}{I} \sin\theta \approx -\frac{mgl}{I} \theta \quad \text{or } \frac{d^2\theta}{dt^2} + \omega^2\theta = 0$$

$$T = 2\pi \sqrt{\frac{I}{mgl}}$$

If we consider moment of inertia about a parallel axis through G, K = radius of gyration then using parallel axis theorem,

$$I = mk^2 + ml^2 \quad \therefore T = 2\pi \sqrt{\frac{K^2 + l^2}{g}} = 2\pi \sqrt{\frac{l'}{g}}$$

$$\text{equivalent length of simple pendulum} = \frac{K^2}{l} + l.$$

Torsional Pendulum

twist of shaft \rightarrow torsional oscillations

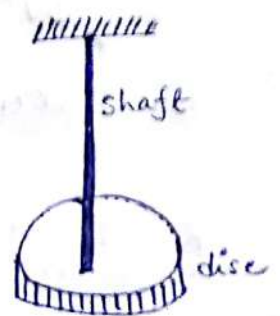
$$\text{torsional couple} = -\tau\theta$$

$$\text{couple due to acceleration} = I \frac{d^2\theta}{dt^2}$$

$$I \frac{d^2\theta}{dt^2} = -\tau\theta \quad , \quad T = 2\pi \sqrt{\frac{I}{\tau}}$$

$$\text{From classical mechanics course, } \tau = \frac{\pi \eta d^4}{32L} = \frac{\pi \eta r^4}{2L}$$

d = shaft diameter, η = modulus of rigidity,
 $= 2\tau$



Electrical Oscillator

Capacitor is charged \Rightarrow electrostatic energy in dielectric media. It discharges through the inductor electrostatic energy \Leftrightarrow magnetic energy. (no dissipation of heat)

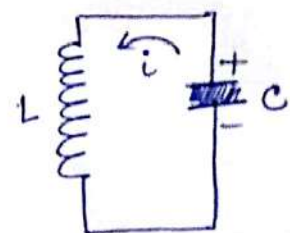
$$\text{voltage across inductor} = -L \frac{di}{dt} = -L \frac{d^2q}{dt^2}$$

$$\text{voltage across capacitor} = \frac{q}{C}$$

$$\text{No e.m.f. circuit, } \frac{q}{C} = -L \frac{d^2q}{dt^2} \quad \text{or } \frac{d^2q}{dt^2} + \frac{q}{LC} = 0$$

$$\omega^2 = \frac{1}{LC} \quad , \quad q = q_0 \sin(\omega t + \phi)$$

charge on capacitor varies harmonically.



$$i = \frac{dq}{dt} = \omega q_0 \cos(\omega t + \phi)$$

$$V = \frac{q}{C} = \frac{q_0}{C} \sin(\omega t + \phi)$$

Total energy = magnetic energy + electric energy

$$= \int iV dt + \frac{1}{2} CV^2 = \int i L \frac{di}{dt} dt + \frac{1}{2} CV^2$$

$$= \int L i di + \frac{1}{2} CV^2 = \frac{1}{2} Li^2 + \frac{1}{2} CV^2 = \frac{1}{2} L \dot{q}^2 + \frac{1}{2} CV^2$$

In mechanical oscillation, Total energy = $\frac{1}{2} m \dot{x}^2 + \frac{1}{2} K x^2$

$$\frac{1}{2} CV^2 = \frac{1}{2} C \left(\frac{q}{C} \right)^2 = \frac{q^2}{2C}$$

In electrical oscillation, Total energy = $\frac{1}{2} L \dot{q}^2 + \frac{1}{2C} q^2$

equivalence