

THERMODYNAMICS (T.D.)

"Entropy Change ΔS "

Isothermal $W_{\text{system}} \rightarrow U_{\text{reservoir}}$

$$\Delta S_{\text{system}} = 0, \Delta S_{\text{surrounding}} = \frac{Q}{T} = \frac{W}{T}$$

$$\Delta S_{\text{universe}} = 0 + \frac{W}{T} = \boxed{\frac{W}{T}}$$

Adiabatic

$$\Delta S_{\text{system}} = \int_{T_i}^{T_f} \frac{dq}{T} = \int_{T_i}^{T_f} \frac{C_p dT}{T} = C_p \ln \frac{T_f}{T_i}$$

$$\Delta S_{\text{surrounding}} = 0, \Delta S_{\text{universe}} = \boxed{C_p \ln \frac{T_f}{T_i}}$$

Free Expansion

$$\Delta S_{\text{system}} = \int_{V_i}^{V_f} \frac{dq}{T} = \int_{V_i}^{V_f} \frac{P dV}{T} = nR \ln \frac{V_f}{V_i}$$

$$\Delta S_{\text{surrounding}} = 0, \Delta S_{\text{universe}} = \boxed{nR \ln \frac{V_f}{V_i}}$$

Conduction

$$\Delta S_{\text{system}} = 0, \Delta S_{\text{reservoir hot}} = -\frac{Q}{T_1},$$

$$\Delta S_{\text{reservoir cold}} = +\frac{Q}{T_2}$$



$$\Delta S_{\text{reservoir}} = \boxed{\frac{Q}{T_2} - \frac{Q}{T_1}} = \Delta S_{\text{universe}}$$

Chemical reaction



Adiabatic wall

$$\Delta S_{\text{system1}} = nR \ln V_f/V_i = \Delta S_{\text{system2}},$$

$$\Delta S_{\text{surrounding}} = 0, \quad \Delta S_{\text{universe}} = 2nR \ln V_f/V_i$$

LEGENDRE TRANSFORMATION IN T.D.

{ Extensive parameter = Independent variable

{ Intensive parameter = Derived

NOT GOOD FOR EXPERIMENTS. XXX

Desired

{ Intensive parameter = Independent variable

{ Extensive parameter = Derived variable
 \rightarrow Fundamental relation in X representation

Given $Y = Y(X_0, X_1, X_2, \dots, X_t)$, we want

$$P_K = \frac{\partial Y}{\partial X^K}$$
 as independent variable

without losing any information \Rightarrow

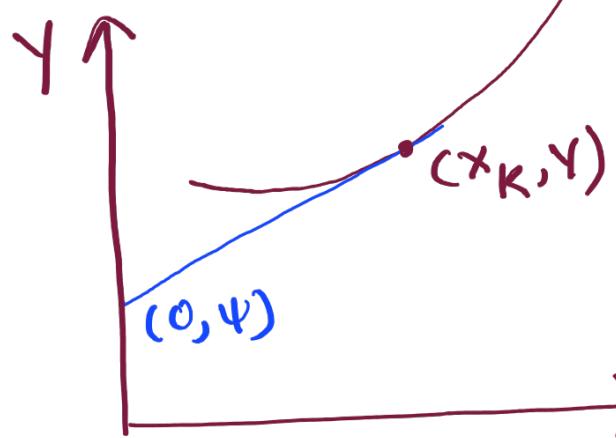
impossible to replace function with derivative! \rightarrow Fundamental relation in P representation

So $\Psi = \Psi(P_0, P_1, P_2, \dots, P_t)$ given

'slope' & 'intercept' (P_K, Ψ) is equivalent

to $Y = Y(X_0, X_1, X_2, \dots, X_t)$

Legendre Transform



$$P_k = \frac{Y - \Psi}{X_k - 0} \Rightarrow \boxed{\Psi = Y - \sum_k P_k X_k}$$

Ψ is the L.T. of Y and its reversible

$$\begin{aligned} d\Psi &= \cancel{dY} - P dx - X dp \\ &= -X dp \quad (\because P = \frac{dY}{dx}) \end{aligned}$$

$$\therefore \boxed{X_k = -\frac{\partial \Psi}{\partial P_k}}$$

$$d\Psi = -\sum_k X_k dP_k$$

1-variable L.T.

$$Y = Y(x)$$

$$P = \frac{dY}{dx}$$

$$\Psi = -Px + Y$$

Eliminating x, Y

$$\Psi = \Psi(P)$$

$$\Psi = \Psi(P)$$

$$-x = \frac{d\Psi}{dP}$$

$$Y = xP + \Psi$$

Eliminating Y, Ψ

$$Y = Y(x)$$

Example (Classical Mechanics):

$$\begin{aligned} \text{Given } \mathcal{L} &= \mathcal{L}(v_1, v_2, \dots, v_r, q_1, q_2, \dots, q_r) \\ \text{Also } \mathcal{L}' &= \mathcal{L}(p_1, p_2, \dots, p_r, q_1, q_2, \dots, q_r) \end{aligned}$$

(Thermodynamics) : $dU < 0, dS > 0$

$$U = U(S, V, N_1, N_2, \dots, N_t)$$

derivative $T, -P, \mu_1, \mu_2 \dots$ intensive parameters

$$f = U - TS \quad \{ \text{d.T. of } U \text{ replaces } S \text{ by } T \\ \text{as independent variable } F = U[T] \}$$

$$U = U(S, V, N_1, N_2, \dots)$$

$$F = F(T, V, N_1, N_2, \dots)$$

$$-P = \frac{\partial U}{\partial V}$$

$$-S = \frac{\partial F}{\partial T}$$

$$F = U - TS$$

$$U = F + TS$$

Eliminating $U \& S$

Eliminating $F \& T$

$$F = F(T, V, N_1, N_2, \dots)$$

$$U = U(S, V, N_1, N_2, \dots)$$

$$H = U + PV \quad \{ \text{d.T. of } U \text{ replaces } V \text{ by } P \}$$

$$\text{as independent variable } H = U[P]$$

$$U = U(S, V, N_1, N_2, \dots)$$

$$H = H(S, P, N_1, N_2, \dots)$$

$$-P = \frac{\partial U}{\partial V}, H = U + PV$$

$$V = \frac{\partial H}{\partial P}, U = H - PV$$

Eliminating $U \& V$

Eliminating $H \& P$

$$H = H(S, P, N_1, N_2, \dots)$$

$$U = U(S, V, N_1, N_2, \dots)$$

$$G = U - TS + PV \quad \{ \text{d.T. of } U$$

replaces S by T and V by $P \quad G = U[T, P] \}$

$$U = U(S, V, N_1, N_2, \dots)$$

$$T = \frac{\partial U}{\partial S}, \quad -P = \frac{\partial U}{\partial V}$$

$$G = U - TS + PV$$

Eliminating U, S, V

$$G = G(T, P, N_1, N_2, \dots)$$

$$G = G(T, P, N_1, N_2, \dots)$$

$$-S = \frac{\partial G}{\partial T}, \quad V = \frac{\partial G}{\partial P}$$

$$U = G + TS - PV$$

Eliminating G, T, P

$$U = U(S, V, N_1, N_2, \dots)$$

Similarly more potentials can be derived

like Grand Canonical potential $U[T, \mu]$

MASSIEU FUNCTIONS

Fundamental relation
 $S = S(U, V, N_1, N_2, \dots)$

$$S\left[\frac{1}{T}\right] = S - \frac{U}{T} = -\frac{F}{T} \quad \{ \text{d.T. of } S$$

replaces U by $\frac{1}{T}$ as independent variable}

$$S\left[\frac{P}{T}\right] = S - \frac{P}{T}V \quad \{ \text{d.T. of } S \text{ replaces } V \text{ by } \frac{P}{T} \text{ as independent variable} \}$$

$$S\left[\frac{1}{T}, \frac{P}{T}\right] = S - \frac{U}{T} - \frac{P}{T}V = -\frac{G}{T} \quad \{ \text{d.T.}$$

of S that replace U by $\frac{P}{T}$ and V by $\frac{P}{T}$
as independent variable }

$$S = S(U, V, N_1, N_2, \dots)$$

$$\frac{P}{T} = \frac{\partial S}{\partial V}$$

$$S\left[\frac{P}{T}\right] = S - \frac{P}{T}V$$

eliminating S & V

$$S\left[\frac{P}{T}\right] = S\left[U, \frac{P}{T}, N_1, N_2, \dots\right]$$

$$S\left[\frac{P}{T}\right] = S\left[U, \frac{P}{T}, N_1, N_2, \dots\right]$$

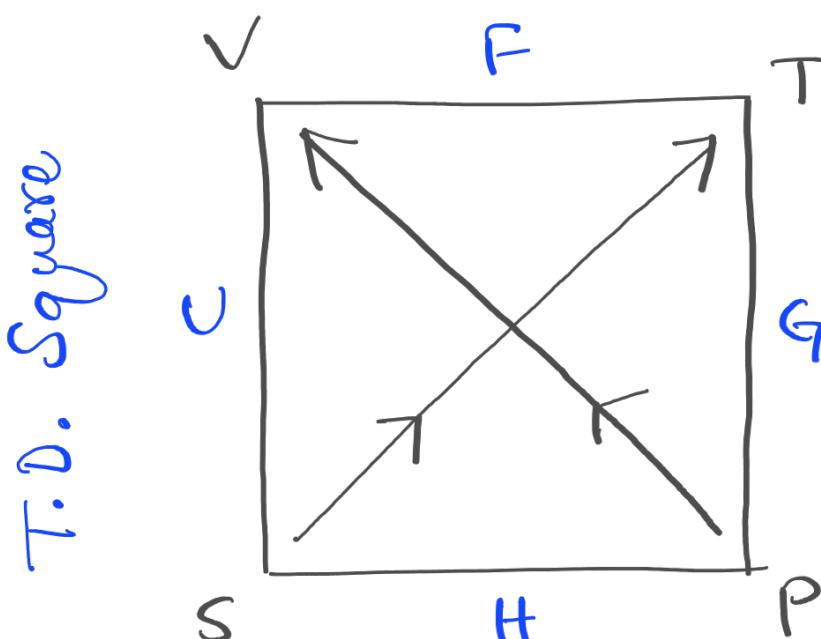
$$-V = \frac{\partial S\left[\frac{P}{T}\right]}{\partial\left[\frac{P}{T}\right]}$$

$$S = S\left[\frac{P}{T}\right] + \frac{P}{T}V$$

eliminating $S\left[\frac{P}{T}\right]$ & $\frac{P}{T}$

$$S = S(U, V, N_1, N_2, \dots)$$

THERMODYNAMIC MNEMONIC DIAGRAM



$$dU = Tds - pdv + \sum_K \mu_K dN_K$$

$$dF = -sdT - pdV + \sum_K \mu_K dN_K$$

$$dG = -sdT + vdP + \sum_K \mu_K dN_K$$

$$dF = Tds + Vdp + \sum_K \mu_K dN_K$$

