
THE KLS CONJECTURE

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§0. Notation

- We refer to measures by greek symbols such as μ and ν and their densities by lowercase alphabets beginning from p .
- B refers to the Euclidean ball of radius 1 in \mathbb{R}^n (the value of n is usually understandable from context).
- Given a measure μ on \mathbb{R}^n and an $(n-1)$ -dimensional surface(?) S in \mathbb{R}^n , $\mu^+(S)$ refers to the “surface area” of the set S , that is,

$$\mu^+(X) = \lim_{\varepsilon \rightarrow 0} \frac{\mu(X + \varepsilon B)}{2\varepsilon}.$$

Alternatively, if $X \subseteq \mathbb{R}^n$ is compact, then

$$\mu^+(\partial X) = \lim_{\varepsilon \rightarrow 0} \frac{\mu(X + \varepsilon B) - \mu(X)}{\varepsilon}.$$

- While needles in [KLS95] refer to one-dimensional segments with a polynomial factor (ℓ^{n-1} where ℓ is linear) in particular, we use them more generally to refer to log-concave measures with a one-dimensional support.

§1. Measure Disintegration

1.1. Introduction

To begin with, let us describe a needle decomposition procedure given in [LV18] to prove the KLS Conjecture. Suppose that we are given a log-concave measure μ with density p with compact convex support K . Let us also fix a subset $E \subseteq K$ of measure $1/2$. We would like to bound $\mu^+(\partial E)$ from below (over all such E).

Now, suppose that we have some hyperplane H that divides space into two half-spaces H_1 and H_2 . Let $K_i = K \cap H_i$ and further assume that $\mu(E \cap H_i) = \frac{1}{2}\mu(K_i)$ for each i . Consider the measures μ_1 and μ_2 with densities given by

$$p_i(x) = \begin{cases} p(x) \frac{\mu(K)}{\mu(K_i)}, & x \in K_i, \\ 0, & \text{otherwise.} \end{cases}$$

Observe that

$$\begin{aligned} p &= p_1 \frac{\mu(K_1)}{\mu(K)} + p_2 \frac{\mu(K_2)}{\mu(K)} \\ \mu &= \mu_1 \frac{\mu(K_1)}{\mu(K)} + \mu_2 \frac{\mu(K_2)}{\mu(K)} \end{aligned} \tag{1.1}$$

More generally, suppose we have some space Ω with a probability measure ν on it such that

$$\mu = \int_{\Omega} \mu_{\omega} d\nu(\omega), \tag{1.2}$$

where the (μ_ω) are log-concave measures on \mathbb{R}^n . In the above example, we can take $\Omega = \{1, 2\}$ and $\nu(\{i\}) = \mu(K_i)/\mu(K)$ for $i \in \Omega$.

Then, given any set E of measure $1/2$, we have

$$\begin{aligned} \mu^+(\partial E) &= \int_{\Omega} \mu_\omega^+(\partial E) d\nu(\omega) \\ &\geq \int_{\Omega} \psi_\omega \mu_\omega(E) (1 - \mu_\omega(E)) d\nu(\omega), \end{aligned} \quad (1.3)$$

where ψ_ω is the isoperimetric constant of μ_ω . If we manage to bound the expression on the right from below by some constant independent of E , then the KLS conjecture follows. It is also worth noting that the decomposition we choose may be dependent on E itself, we only require that the lower bound constant does not depend on this choice of E .

1.2. A proof of the $n^{-1/2}$ bound using needle decompositions

“Needle decomposition” refers to the process of performing the step we used to obtain (1.1) until the bodies K_ω become one-dimensional. We repeatedly split the bodies in a way that the quantity $\mu_\omega(E)$ remains constant at $1/2$. Suppose that we do so and the final limiting set of needles is $(K_\omega)_{\omega \in \Omega}$. Then, we can use one-dimensional isoperimetry to get that for any ω , $\psi_\omega \gtrsim \|A_\omega\|_{\text{op}}^{-1/2}$. We also have that $\mu_\omega(E) = 1/2$, so

$$\mu^+(\partial E) \gtrsim \int_{\Omega} \|A_\omega\|_{\text{op}}^{-1/2} d\nu(\omega). \quad (1.4)$$

We wish to bound the integral on the right from below.

To do so, consider (1.2) (or rather, the similar expression for the density p). Then, we have that

$$\int_{\mathbb{R}^n} p(x) x x^\top dx = \int_{\Omega} \int_{\mathbb{R}^n} p_\omega(x) x x^\top dx d\nu(\omega).$$

Thus,

$$A + b b^\top = \int_{\Omega} A_\omega + b_\omega b_\omega^\top d\nu(\omega), \quad (1.5)$$

where A and b (resp. A_ω and b_ω) refer to the covariance matrix and barycenter of μ (resp. μ_ω) respectively. Assume without loss of generality that $b = 0$. Taking the trace on either side of the above expression,

$$\begin{aligned} \text{Tr}(A) &= \int_{\Omega} \text{Tr}(A_\omega) + \|b_\omega\|^2 d\nu(\omega) \\ &\geq \int_{\Omega} \|A_\omega\|_{\text{op}} d\nu(\omega), \end{aligned}$$

where the inequality follows from the fact that A_ω is a covariance matrix so is positive semi-definite. One can then use Hölder’s inequality to get

$$\left(\int_{\Omega} \|A_\omega\|_{\text{op}} d\nu(\omega) \right) \left(\int_{\Omega} \|A_\omega\|_{\text{op}}^{-1/2} d\nu(\omega) \right)^2 \geq 1$$

and so,

$$\int_{\Omega} \|A_\omega\|_{\text{op}}^{-1/2} d\nu(\omega) \gtrsim \text{Tr}(A)^{-1/2}.$$

Substituting this back in (1.4), we get $\psi_p \gtrsim \text{Tr}(A)^{-1/2}$, that is, $\psi_n \gtrsim n^{-1/2}$.

1.3. An alternate way to look at stochastic localization

Let us return to (1.3). In the above method of needle decomposition, we attempted to exercise control over the quantity $\mu_\omega(E)(1 - \mu_\omega(E))$ for all ω by fixing $\mu_\omega(E)$ at $1/2$.

How does stochastic localization fit into this? Instead of controlling $\mu_\omega(E)$, we try to control ψ_ω by defining a martingale (p_t) whose isoperimetric constant is easily bounded. That is, $\mathbf{E}[p_t] = p$ (this is just an integral of the form of (1.1)) and further, the isoperimetric constant of μ_t is lower bounded by $t^{1/2}$. Then, the problem comes down to estimating

$$\int_{\Omega} \mu_t(E)(1 - \mu_t(E)) d\nu(\omega),$$

which is exactly what papers such as [Che21] do.

1.4. What next?

Going back to needle decompositions again, we wish to show that there exists a needle decomposition conserving $\mu_\omega(E) = 1/2$ such that

$$\int_{\Omega} \|A_\omega\|_{\text{op}}^{-1/2} d\nu(\omega) \gtrsim \|A\|_{\text{op}}^{-1/2}.$$

(1.5) for $b = 0$ gives

$$\|A\|_{\text{op}} = \left\| \int_{\Omega} A_\omega + b_\omega b_\omega^\top d\nu(\omega) \right\|_{\text{op}}.$$

Therefore, it would suffice to show that

$$\int_{\Omega} \|A_\omega\|_{\text{op}}^{-1/2} d\nu(\omega) \gtrsim \left\| \int_{\Omega} A_\omega + b_\omega b_\omega^\top d\nu(\omega) \right\|_{\text{op}}^{-1/2}$$

for some needle decomposition that conserves $\mu_\omega(E)$.

Using Hölder's inequality as we did in the proof of the $n^{-1/2}$ bound, it is seen that it suffices to show

$$\int_{\Omega} \|A_\omega\|_{\text{op}} d\nu(\omega) \lesssim \left\| \int_{\Omega} A_\omega + b_\omega b_\omega^\top d\nu(\omega) \right\|_{\text{op}}$$

for some needle decomposition preserving $\mu_\omega(E)$ (it would in fact be enough to show this with some set $A \subseteq \Omega$ instead of Ω such that $\nu(A)$ is lower-bounded by a constant).

Neglecting the $b_\omega b_\omega^\top$ term, it suffices to show that

$$\int_{\Omega} \|A_\omega\|_{\text{op}} d\nu(\omega) \lesssim \left\| \int_{\Omega} A_\omega d\nu(\omega) \right\|_{\text{op}}. \quad (1.6)$$

The above inequality essentially asks if there exists a needle decomposition where the needles are “nearly aligned”. Indeed, if the segments of the needles are perfectly aligned, then equality holds above. We are allowing a constant factor of leeway. If the direction of the one-dimensional body K_ω is u_ω , then the above is equivalent to

$$\sup_{\|\zeta\| \leq 1} \int_{\Omega} \mathbf{Var}_{x \sim p_\omega}(x) \langle \zeta, u_\omega \rangle^2 d\nu(\omega) \gtrsim \int_{\Omega} \mathbf{Var}_{x \sim p_\omega}(x) d\nu(\omega) \quad (1.7)$$

Another way to prove the $n^{-1/2}$ bound is to take the expectation of the term on the left of the above for ζ being drawn uniformly from S^{n-1} .

§2. More on decompositions

While we have not mentioned it thus far, a useful result to know is that the KLS conjecture on the class of log-concave measures is equivalent to the KLS conjecture on the class of indicator functions on convex bodies. In particular, the class of log-concave functions is the smallest containing the class of indicator functions on convex bodies that is closed under taking marginals and weak limits, as shown in [AGB15].

2.1. Hyperplane bisections

As before, suppose we have a log-concave probability measure μ with density p on the body K , and we fix some $E \subseteq K$ with $\mu(E) = 1/2$. Let us define the function $f_{E,K} : \mathbb{R}^n \setminus \{0\} \rightarrow \mathbb{R}$ by

$$f_{E,K}(x) = \left| \int_{\{z \in \mathbb{R}^n : \langle z, x \rangle \geq \|x\|^2\}} p(y)(\mathbb{1}_E - \mathbb{1}_{\mathbb{R}^n \setminus E}) dy \right|.$$

That is, if H_x is the hyperplane defined by x (orthogonal to x and passing through it) and H_x^+ is either of the resulting halfspaces, the value of the above function at x is equal to $|\mu(E \cap H_x^+) - \mu((\mathbb{R}^n \setminus E) \cap H_x^+)|$.

This serves as a measure of how “imbalanced” the hyperplane corresponding to x is – $f_{E,K}(x) = 0$ iff the hyperplane corresponds to x is a bisecting hyperplane (where bisecting means that $\mu(E \cap K_\omega) = \frac{1}{2}\mu(K_\omega)$, as in needle decompositions).

For nice E , $f_{E,K}$ is continuous.

The primary tool used in [LS93] to prove the localization lemma was that there exists a bisecting hyperplane passing through any $(n-2)$ -dimensional affine space. How would this translate in terms of the above defined function?

Suppose we have an $(n-2)$ -dimensional affine space orthogonal to the subspace spanned by $y, z \in \mathbb{R}^n$ and passing through y .

Suppose that x defines a hyperplane containing this affine space. Then x is orthogonal to the plane, and so orthogonal to the space itself. That is, it must lie in the subspace spanned by y, z . Further, $y - x$ is orthogonal to x . That is, the set of all these x forms a circle passing through 0 contained in the 2-dimensional subspace spanned by y, z .

The conclusion of the localization method is that for any circle S passing through 0, either

- $f_{E,K}(w) = 0$ for some $w \in S \setminus \{0\}$ or
- The limit of $f_{E,K}(w)$ as w goes to 0 along the circle is equal to 0 – this corresponds to a bisecting hyperplane passing through the origin itself. It is not too difficult to check that this is well-defined and that the directional limit along either direction of the circle is the same.

More generally, suppose we have some smooth curve C in \mathbb{R}^n that passes through the origin.¹ Then, as before, either $f_{E,K}(w) = 0$ for some $w \in C \setminus \{0\}$ or one of the directional limits at 0 (along C) is equal to 0.

It may be shown that as a consequence, there exists a surface passing through the origin on which $f_{E,K}$ is equal to 0.

In general, if we start off with E of some measure (not necessarily 1/2) that we want to preserve, we can define $f_{E,K}(x)$ as

$$\left| \int_{\{z \in \mathbb{R}^n : \langle z, x \rangle \geq \|x\|^2\}} p(y)(\mu(\mathbb{R}^n \setminus E) \mathbb{1}_E - \mu(E) \mathbb{1}_{\mathbb{R}^n \setminus E}) dy \right|$$

An interesting question is to generally characterize these functions $f_{E,K}$.

2.2. Aligned 2-dimensional decompositions are always possible

Suppose we have an n -dimensional body K with $n > 2$ along with some direction u in \mathbb{R}^n . We claim that it is possible to decompose this into a set of $(n-1)$ -dimensional bodies $\{K_\omega\}$ such that any of these bodies contains our specified direction u (meaning that a translational shift of $\text{span}(\{u\})$ is contained in the minimal affine space containing any K_ω).

To prove this, assume without loss of generality that $u = e_n$. Consider the set of $(n-2)$ -dimensional affine spaces

$$S = \{x \in \mathbb{R}^n : x_i = q_1, x_j = q_2 : q_1, q_2 \in \mathbb{Q}, 1 \leq i < j \leq n-1\}.$$

This is similar to the argument involved in [LS93] except that we only consider the set of $(n-2)$ -dimensional affine spaces that contain u . As the argument goes there, all the bodies must decompose into at most $(n-1)$ -dimensional bodies in the limiting step – if not, then there exists some affine space in S that intersects the n -dimensional body,

¹Smoothness is not required, only that the curve is differentiable at the origin.

and choosing the corresponding bisecting hyperplane results in a contradiction.

In fact, it turns out that we can decompose it into a set of 2-dimensional bodies that all contain our specified direction!

This is easily done using induction on n . Reducing the n -dimensional body to a set of $(n - 1)$ -dimensional bodies and then each of these smaller bodies to 2-dimensional bodies gets the job done. It should be noted that this argument does not work out if the body under consideration is 2-dimensional, since it does not make sense to have a 0-dimensional affine space containing our direction.

In general, given k directions, it is possible to decompose a body K into $\{K_\omega\}$ such that each of the K_ω is $(k + 1)$ -dimensional and each of them contains these k directions.

A natural next question is: can we give up perfect alignedness in exchange for near alignedness, which is all we really need to show KLS?

2.3. A potential function

Let us fix μ , p , K , and E as usual. Also suppose we have some direction u . We wish to decompose the body into needles in a way that all of them are nearly in the direction of u . Equivalently, the hyperplanes chosen for bisection should all nearly contain u . That is, the set of x corresponding to the hyperplanes $\{H_x\}$ must all be nearly orthogonal to u . So, at each step, the x chosen must be such that $\langle x, u \rangle$ is small – more precisely, $1 - \frac{\langle x, u \rangle^2}{\|x\|^2} \gtrsim 1$.

Also, as seen from (1.3), all we really want is that $\mu_\omega(E)(1 - \mu_\omega(E)) \gtrsim 1$, it might be fine to instead just minimize $f_{E,K}$ instead of ensuring that it is exactly equal to 0. So, one may choose the x corresponding to the bisecting hyperplane at each step by constructing a potential function such as

$$\Phi(x) = (1 + f_{E,K}(x)) \left(1 + \frac{|\langle x, u \rangle|}{\|x\|} \right)$$

and at each step, choosing the x that minimizes Φ . The reason for adding the 1 is that otherwise, the expression would trivially be minimized if the corresponding term is 0 irrespective of the other term.

We might be able to choose our potential function more wisely.

Recall that in our bound of the isoperimetric constant using needle decompositions, we don't really require that $\mu_\omega(E) = 1/2$ for each ω , we only require that $\mu_\omega(E)(1 - \mu_\omega(E)) \gtrsim 1$. Suppose we have a body K and we split it using the hyperplane H_x . Let $K_1 = H_x^+ \cap K$. Then, it is not too difficult to show that

$$\mu_2(E) = \mu(E) + \frac{f_{E,K}(x)}{2\mu(H_x^+)}.$$

In general, if we start off with $K_1 = K$, and we form K_{i+1} by intersecting K_i with the half-space $H_{x_i}^+$ such that the limiting body is K_ω , then

$$\mu_\omega(E) = \mu(E) + \sum_{n \in \mathbb{N}} \frac{f_{E,K_n}(x_n)}{2\mu_n(H_{x_n}^+)}.$$

Our goal should be to minimize the quantity on the right, so a suitable potential might involve something along the lines of $f_{E,K}(x)/2 \min\{\mu(H_x^+), \mu(H_x^-)\}$.

It would be nice to have a similar handy quantity for the alignment part of it as well – if we start off with a body K and bisect it using $H_{x_1}^+, H_{x_2}^+, \dots$ to get a limiting one-dimensional body K_ω , how do we represent the “direction” of K_ω in terms of the x_n (and perhaps some quantities involving K)?

2.4. The Poincaré Inequality

Given a probability measure μ on \mathbb{R}^n with density p , its Poincaré constant is defined by

$$\zeta_p = \inf_{g \text{ smooth}} \frac{\mathbf{E}_{x \sim p} \|\nabla g(x)\|_2^2}{\mathbf{Var}_{x \sim p} g(x)}.$$

We also define the Cheeger constant by

$$h_p = \inf_{g \text{ smooth}} \frac{\mathbf{E}_{x \sim p} \|\nabla g(x)\|_2}{\mathbf{E}_{x \sim p} |g(x) - \mathbf{E}_{x \sim p} g(x)|}.$$

Equation (5.8) in [Led04] shows that for log-concave μ , $h_p^2 \sim \zeta_p$.² Further, more relevant to our interests, $\zeta_p \sim \psi_p^2$ (this is a consequence of Cheeger's inequality – see [Che15, Maz60]).

How is the isoperimetric inequality related to these? Suppose that in the definition of the Cheeger constant, we set $g = \mathbb{1}_E$ for some set E (or rather, a sequence of smooth functions converging to $\mathbb{1}_E$). Then, $\|\nabla g(x)\|$ behaves like a Dirac delta function on ∂E , and we get that $\mathbf{E}_{x \sim p} \|\nabla g(x)\|$ is just $\mu^+(\partial E)$. The denominator on the other hand is the variance of a Bernoulli random variable with parameter $\mu(E)$, which is equal to $\mu(E)(1 - \mu(E))$.

So, the inside expression as a whole becomes $\mu^+(\partial E)/\mu(E)(1 - \mu(E))$, which is precisely the expression involved in the isoperimetric constant!

2.5. Arbitrary cuts through the barycenter tend to be good

Suppose μ is an isotropic log-concave measure with density p . Let $E \subseteq \mathbb{R}^n$ with $\mu(E) = 1/2$. Consider the function $g_0 : S^{n-1} \rightarrow \mathbb{R}$ given by

$$g_0(x) = \mu(H_x^+ \cap E) - \mu(H_x^+ \cap (\mathbb{R}^n \setminus E)),$$

where $H_x^+ = \{z \in \mathbb{R}^n : \langle z, x \rangle \geq 0\}$.

It is not too difficult to show that g_0 is Lipschitz with a Lipschitz constant independent of μ and n .

A consequence of the concentration of measure phenomenon is that given a 1-Lipschitz function f on S^{n-1} with median $\text{med } f$, if σ is the rotation-invariant (uniform) probability measure on S^{n-1} ,

$$\sigma(\{|f - \text{med } f| \geq \varepsilon\}) \leq 2e^{-n\varepsilon^2/2}$$

and $|\text{med } f - \mathbf{E}f| \leq 12n^{-1/2}$.

In our context, $\mathbf{E}g_0 = 0$. Therefore,

$$\sigma(\{|g_0 - \text{med } g_0| \geq \varepsilon\}) \leq 2e^{-nC\varepsilon^2},$$

where C is some constant independent of n and μ . So, for large n , we see that g_0 is almost constant on the entirety of S^{n-1} , so a random hyperplane through the barycenter is good with high probability.

§3. The Slicing Conjecture

Like the KLS Conjecture, the slicing conjecture is one of the main open questions in convex geometry.

Conjecture (Slicing Conjecture). Any convex body $K \subseteq \mathbb{R}^n$ of volume 1 has at least one hyperplane section H such that

$$\text{vol}_{n-1}(K \cap H) \gtrsim 1.$$

[Bal88] shows that this is equivalent to

Conjecture (Slicing Conjecture, Reformulated). For any isotropic log-concave density p on \mathbb{R}^n , the slicing constant $L_p := p(0)^{1/n}$ is $\lesssim 1$.

²Even for μ that is not log-concave, it is true that $h_p^2 \lesssim \zeta_p$.

Bourgain's original paper [Bou86] proves the slicing conjecture for unconditional convex bodies, that is, convex bodies K such that $(x_1, \dots, x_n) \in K$ iff $(|x_1|, \dots, |x_n|) \in K$.

Proposition I.2 in [BM10] shows that there is a link between the slicing conjecture of a distribution and its entropy. Indeed, we have

$$1 + \log \|f\|_\infty^{1/n} \geq \frac{1}{n} h(X) \geq \log \|f\|_\infty^{1/n}, \quad (3.1)$$

where X is a random vector in \mathbb{R}^n with log-concave density f and $\|f\|_\infty = \sup_{x \in \mathbb{R}^n} f(x)$.

The *entropy power* of a random vector $X : \Omega \in \mathbb{R}^n$ is given by

$$\mathcal{N}(X) = \frac{1}{2\pi e} e^{2h(X)/n}.$$

Thus, the slicing conjecture is true if and only if for any random variable X in \mathbb{R}^n with isotropic log-concave density f , $\mathcal{N}(X) \gtrsim 1$.

[MNT18] extensively discusses the slicing conjecture and the related problem of entropy minimization over 1-dimensional symmetric (even) log-concave distributions with a fixed variance.

Over the class of 1-dimensional symmetric log-concave distributions, it shows

$$\mathbf{Var}(X) \geq \mathcal{N}(X) \geq \frac{6}{\pi e} \mathbf{Var}(X),$$

with the upper bound being attained uniquely for the Gaussian distribution and the lower uniquely for the uniform distribution.

The result on the slicing constant directly is that for any even log-concave function f on \mathbb{R} ,

$$f(0)^p \int |x|^p f(x) dx \leq 2^{-p} \Gamma(p+1) \left(\int f(x) dx \right)^{p+1}$$

with equality being attained iff f is a symmetrized exponential distribution.

Let us attempt to generalize this to unconditional convex bodies. Let f be an unconditional log-concave function. We claim that

$$f(0)^p \left(\prod_{1 \leq i \leq n} \int_{\mathbb{R}^n} |x_i|^p f(x_i) dx_i \right) \leq (2^{-p} \Gamma(p+1))^n \left(\int_{\mathbb{R}^n} f(x) dx \right)^{p+1}. \quad (3.2)$$

Assume that $f(0) = 1$. Let $g(x) = \exp(-\sum_{i=1}^n a_i |x_i|)$, where the a_i are chosen such that $\int_{\mathbb{R}^n} f = \int_{\mathbb{R}^n} g$. We claim that for any $i \in \{1, 2, \dots, n\}$,

$$\int_{\mathbb{R}^n} |x_i|^p f(x) dx \leq \int_{\mathbb{R}^n} |x_i|^p g(x) dx.$$

If we let f_i, g_i be the marginal of f, g on the x_i axis, then the above is equivalent to

$$\int_{\mathbb{R}} |x_i|^p f_i(x_i) dx_i \leq \int_{\mathbb{R}} |x_i|^p g_i(x_i) dx_i.$$

Due to the unconditionality of f (and g), f_i and g_i are symmetric log-concave functions. In fact, since g_i is of the form $Ce^{-a_i|x_i|}$, $f_i - g_i$ changes sign either not at all (in which case we take $x_0 = 0$) or exactly twice, at $x_i = \pm x_0$. Thus,

$$\int_{\mathbb{R}} |x_i|^p (f_i(x_i) - g_i(x_i)) dx_i = \int_{\mathbb{R}} (|x_i|^p - |x_0|^p) (f_i(x_i) - g_i(x_i)) dx_i \leq 0$$

since the integrand is always non-positive.

So, it suffices to check the inequality for g , which is easily done.

In particular, for any unconditional isotropic log-concave density p , $L_p = p(0)^{1/n} \leq \sqrt{\Gamma(3)/4} = 1/\sqrt{2}$, where equality holds iff each of its marginals are symmetrized exponentials.

In the above proof, we have n “degrees of freedom” (each of the a_i), but we are only really using one of them to ensure that $\int f = \int g$. Could we generalize to an even broader class of functions, hopefully such that we actually use all the freedom we have?

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