

# HW3

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## 1 Homework: Lagrangian Mechanics

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## 2 Stating the Equations

The lagrangian is defined as  $L \equiv T - U$ , where  $T$  is the kinetic energy for the considered mass, and  $U$  is the potential energy on due to forces on the mass. For the  $j$ th particle, we have  $T_j$  and  $U_j$  as

$$T_j = \frac{1}{2}m\dot{x}_j^2 \quad U_j = \frac{1}{2}kx_j^2$$

We realize that we have  $N$  particles connected by  $N + 1$  springs. Start analyzing by assuming all the particles have shifted from their equilibrium a distance  $x_j$  to the right.

### 2.1 First Atom

Consider the free body diagram for the first atom we find its potential as

$$U_1 = \frac{k}{2}kx_1^2 + \frac{k}{2}(x_2^2 - x_1^2) = \frac{k}{2}(x_2^2 - 2x_1x_2 + 2x_1^2)$$

And the kinetic energy is

$$T_1 = \frac{m}{2}\dot{x}_1^2$$

Then the Lagrangian is

$$L_1 = T_1 - U_1 = \frac{m}{2}\dot{x}_1^2 - \frac{k}{2}(x_2^2 - 2x_1x_2 + 2x_1^2)$$

Since we take the partial derivatives with respect to  $x_1$  it is enough to use  $L_1$ , then the Lagrange equation is

$$\frac{\partial L_1}{\partial x_1} - \frac{d}{dt} \frac{\partial L_1}{\partial \dot{x}_1} = 0$$

So,

$$\begin{aligned} -\frac{k}{2}(-2x_2 + 4x_1) - \frac{d}{dt}m\dot{x}_1 &= 0 \\ -k(-x_2 + 2x_1) - m\ddot{x}_1 &= 0 \\ \ddot{x}_1 + 2\omega_0^2x_1 - \omega_0^2x_2 &= 0 \end{aligned}$$

where,  $\omega_0^2 \equiv k/m$

## 2.2 Last Atom; $N$

Consider the last atom at  $j = N$ , its previous spring stretched by a distance  $x_N - x_{N-1}$ , and its next spring compressed by a distance  $x_N$ . The Lagrangian is

$$L_N = \frac{m}{2} \dot{x}_N^2 - \left( \frac{k}{2} x_N^2 + \frac{k}{2} (x_N - x_{N-1})^2 \right)$$

We can use the last equation and just replace  $x_N$  with  $x_N$  and  $x_{N+1}$  with  $x_{N+1}$ , that is,

$$\ddot{x}_N + 2\omega_0^2 x_N - \omega_0^2 x_{N-1} = 0$$

## 2.3 Middle Atoms

Again, we analyze the free body diagram for a given in between atom to its potential energy. We find that the Lagrangian is,

$$L_j = T_j - U_j = \frac{m}{2} \dot{x}_j^2 - \frac{k}{2} [(x_j - x_{j-1})^2 + (x_{j+1} - x_j)^2]$$

Or

$$L_j = \frac{m}{2} \dot{x}_j^2 - \frac{k}{2} [2x_j^2 - 2x_j x_{j-1} - 2x_j x_{j+1} + x_{j-1}^2 + x_{j+1}^2]$$

Then the equation of motion is,

$$\frac{\partial L}{\partial x_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}_j} = 0$$

Then,

$$\begin{aligned} m\ddot{x} + k[2x_j - x_{j-1} - x_{j+1}] &= 0 \\ \ddot{x} + 2\omega_0^2 x_j - \omega_0^2 x_{j-1} - \omega_0^2 x_{j+1} &= 0 \end{aligned}$$

## 2.4 Generalization of Lagrangian

$$L = \sum_{j=1}^N (T_j - U_j) = \sum_{j=1}^N \left( \frac{m}{2} \dot{x}_j^2 - \frac{k}{2} (x_{j+1}^2 - 2x_j x_{j+1} + 2x_j^2) \right)$$

The Lagrange equation is

$$\frac{\partial L}{\partial x_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}_j} = 0, \quad j = 1, 2, 3, \dots, N$$

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## 3 The Eigenmodes

We write the equation of motion in matrix form as,  $\ddot{X} + \Omega X = 0$

```
[ ]: # import needed packages
import numpy as np

# graph settings
import matplotlib.pyplot as plt
from matplotlib_inline import backend_inline
plt.rcParams['figure.figsize'] = [15, 4]
backend_inline.set_matplotlib_formats('png', 'pdf')
```