## HW10

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# 1 Thermal Physics - Homework 10

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```
[]: from math import comb
  import numpy as np
  import seaborn as sns
  import matplotlib.pyplot as plt
  from IPython.display import Math
  from scipy import constants as con
  from scipy.integrate import quad
  from scipy.optimize import curve_fit
  from matplotlib_inline.backend_inline import set_matplotlib_formats

set_matplotlib_formats('pdf')
  plt.rcParams |= {
    'text.usetex': True,
    'figure.figsize': (10, 4)
  }
  sns.set_theme()
  set_matplotlib_formats('svg', 'pdf')
```

#### 1.1 Question 1

The change in pressure is

$$P_2 - P_1 = \rho g \Delta h$$

And also is given by

$$P_2-P_1=\frac{n_BRT}{V}$$

Combine those two,

$$\rho g \Delta h = \frac{n_B RT}{V}$$

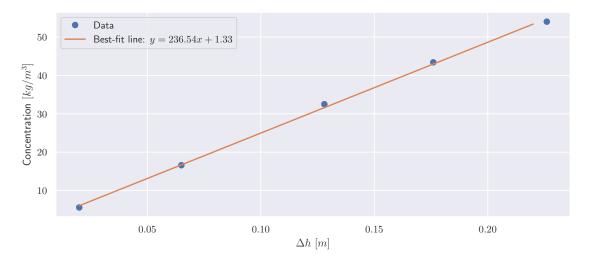
By definition number of solute moles equals the mass of the solute per molecular weight,  $n_B = m/M$ , so

$$\rho g \Delta h = \frac{mRT}{VM}$$

The concentration c equals the mass divided by the volume of the solution,

$$c = \frac{\rho g M}{RT} \delta h$$

Now for the given data we plot them to find the best-fit line's slop.



So we conclude that

$$\frac{\rho gM}{RT} = 236.54 \quad \Rightarrow \quad M = 236.54 \frac{RT}{\rho q} = 236.54 \frac{R \cdot 276.15}{997 * 9.8}$$

So we have

$$\frac{\text{[]}: \rho gM}{RT} = 55.59$$

#### 1.2 Question 3

#### 1.2.1 (a)

For simplification we define  $v_0 = \sqrt{2kT/m}$ , where T = 300 K and atomic mass of nitrogen is 28u. So

```
[]: m = 28 * con.atomic_mass
v0 = np.sqrt(2 * 300 * con.k / m)
Math(r'v_0 = v_{max}= \sqrt{\frac{2kT}{m}} = %.2f' % v0)
```

[]:  $v_0 = v_{max} = \sqrt{\frac{2kT}{m}} = 422.10$ 

#### 1.2.2 (b)

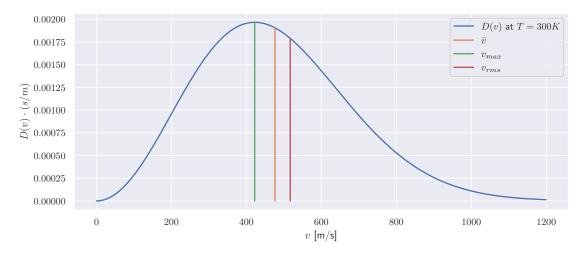
```
[]: # Define maxwell speed distribution

def maxwell_dist(v, t=1, v0=v0):
    """ 't' is multiple of 300K"""

return (4 / np.sqrt(np.pi) * v**2 / v0**3
    * t**(-3/2) * np.exp(- v**2 / v0**2 / t))
```

```
[]: vel_range = np.arange(0, 1200)
    v_max = v0
    v_avg = np.sqrt(8 * con.k * 300 / np.pi / m)
    v_rms = np.sqrt(3 * 300 * con.k / m)

plt.plot(vel_range, maxwell_dist(vel_range, 1), label='$D(v)$ at $T=300K$')
    plt.plot([v_avg]*2,[0, maxwell_dist(v_avg)], label=r'$\bar{v}$')
    plt.plot([v_max]*2,[0, maxwell_dist(v_max)], label=r'$v_{max}$')
    plt.plot([v_rms]*2,[0, maxwell_dist(v_rms)], label=r'$v_{rms}$')
    plt.ylabel(r'$D(v) \cdot (s/m)$')
    plt.xlabel(f'$v$ [m/s]')
    plt.legend()
    plt.show()
```



#### 1.2.3 (c)

We I use the Python quad to integrate numerically the maxwell\_dist function.

[]: print('Probability of finding a nitrogen molecule with speed' f' between 300 and 572 is {quad(maxwell\_dist, 1, 572)[0]:.2f}')

Probability of finding a nitrogen molecule with speed between 300 and 572 is 0.70

#### 1.3 Question 4

#### 1.3.1 (b)

We find first the second derivative of Z with respect to  $\beta$ 

$$\frac{\partial^2}{\partial \beta^2} = \frac{\partial^2}{\partial \beta^2} \sum_s e^{-\beta E(s)} = \sum_s \frac{\partial}{\partial \beta} [-E(s)e^{\beta E(s)}] = \sum_s [E(s)]^2 e^{\beta E(s)}$$

$$= Z \cdot \sum_{s} [E(s)]^2 \frac{e^{-\beta E(s)}}{Z} = Z \cdot \bar{E^2}$$

Now using result from Problem 6.16 twice:

$$\bar{E^2} = \frac{1}{Z}\frac{\partial}{\partial\beta}\left(\frac{\partial Z}{\partial\beta}\right) = \frac{1}{Z}\frac{\partial}{\partial\beta}(-Z\bar{E}) = -\frac{1}{Z}\left(\frac{\partial\bar{E}}{\partial\beta}Z + \bar{E}\frac{\partial Z}{\partial\beta}\right) = \frac{\partial\bar{E}}{\partial\beta} + (\bar{E})^2$$

#### 1.3.2 (a)

Or

$$\bar{E^2} - (\bar{E})^2 = -\frac{\partial \bar{E}}{\partial \beta} = -\frac{\partial T}{\partial \beta} \frac{\partial \bar{E}}{\partial \beta}$$

But  $\partial T/\partial \beta=(\partial \beta/\partial T)^{-2}=-kT^2$ , while  $\partial \bar{E}/\partial T$  is just the heat capacity at constant volume,  $C_V$ , then

$$\sigma_E^2=\bar{E^2}-(\bar{E})^2=kT^2C_V$$

#### 1.3.3 (c)

or

$$\sigma_E = kT\sqrt{\frac{C_V}{k}}$$

#### 1.4 Question 5

Suppose we have an isolated Einstein solid composed of 200 oscillators with a total of 200 units of energy. We take 1 oscillator as the system and the rest as the reservoir, with  $\epsilon = 1$ 

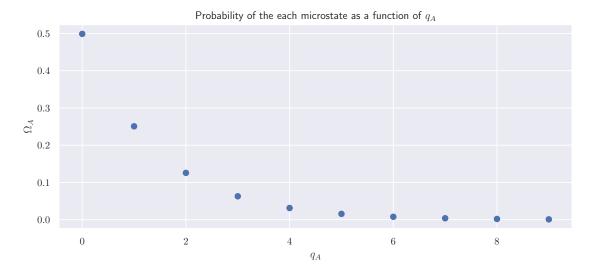
#### 1.4.1 (a)

```
[]: N_A = 1; N_B = 199
N = N_A + N_B
q_total = 200

q_A = np.arange(0, q_total+1)
q_B = q_A[::-1]

omega_A = np.array([float(comb(q + N_A - 1, q)) for q in q_A])
omega_B = np.array([float(comb(q + N_B - 1, q)) for q in q_B])
omega_total = omega_A * omega_B
prob = omega_total / np.sum(omega_total)

plt.title('Probability of the each microstate as a function of $q_A$')
plt.xlabel('$q_A$')
plt.ylabel('$\mathref{Omega_A}$')
plt.plot(q_A[:10], prob[:10], 'o')
plt.show()
```



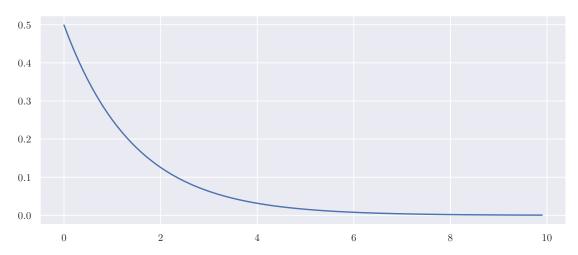
#### 1.4.2 (b)

We use curve\_fit to fit the general exponential form

$$y = A e^{Bx}$$

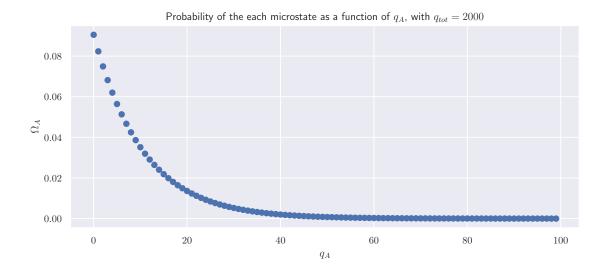
```
(A, B), error = curve_fit(gen_exp, x, y)
x = np.arange(0, 10, 0.1)

plt.plot(x, gen_exp(x, A, B))
plt.show()
```



### 1.4.3 (d)

Repeat (a) through (b) but with 2,000 energy units, plotting up to 100 on x-axis.



```
[]: x, y = q_A[:100], prob[:100]

(A, B), error = curve_fit(gen_exp, x, y, p0=(0, prob[0]))
x = np.arange(0, 100, 0.1)

plt.plot(q_A[:100], prob[:100], 'o', label='Data', alpha=0.5)
plt.plot(x, gen_exp(x, A, B), label='Fit curve')
plt.legend()
plt.show()
```

