

Derivation of axisymmetric equilibria

Ideal MHD $\frac{\partial}{\partial t}$ eqns are

$$\nabla \cdot (\sigma v) = 0 \quad (1)$$

$$\sigma v \cdot \nabla v - \sigma \times B = -\nabla p \quad (2)$$

$$v \cdot \nabla S = 0 \quad (3)$$

$$E + v \times B = 0 \quad (4)$$

$$\nabla \times B = \mu_0 J \quad (5)$$

$$\nabla \times E = 0 \quad (6)$$

As $\nabla \cdot B = 0$ \downarrow scalar

$$\Rightarrow B = B_\phi e_\phi + \frac{1}{r} \nabla \psi \times e_\phi \quad (7)$$

As $\nabla \cdot J = 0$ (from 5)

We have \downarrow scalar

$$J = J_\phi e_\phi + \frac{1}{r} \nabla f \times e_\phi \quad (8)$$

Thus (5) is -

$$\mu_0 J = \nabla \times B = -\frac{\Delta^* \psi}{r} e_\phi + \frac{1}{r} \nabla (\gamma B_\phi) \times e_\phi$$

$$\Rightarrow \mu_0 J_\phi = -\frac{\Delta^* \psi}{r} \quad (9)$$

$$\text{or } \mu_0 f = \gamma B_\phi \quad (10)$$

Case I: No flow

$$\mathbf{J} \times \mathbf{B} = \nabla p \quad (11)$$

$$\Rightarrow \mathbf{B} \cdot (\mathbf{J} \times \mathbf{B}) = \mathbf{B} \cdot \nabla p$$

$$\text{or } 0 = \frac{1}{\mu} (\nabla p \times \nabla \psi) \cdot \epsilon \phi \quad (\text{Identity 52})$$

$$\Rightarrow p = p(x, z) = P(\psi) \quad \psi = \psi(x, z) \quad (12)$$

$$\text{also } \mathbf{J} \cdot (\mathbf{J} \times \mathbf{B}) = \mathbf{J} \cdot \nabla p$$

$$\Rightarrow 0 = \frac{1}{\mu} \nabla p \times \nabla (\sigma B \phi)$$

$$\Rightarrow 0 = \nabla \psi \times \nabla (\sigma B \phi)$$

$$\Rightarrow \sigma B \phi = \mu_0 f(x, z) = F(\psi) \quad (13)$$

Now, (11) can be written as:

$$\frac{\partial \phi}{\partial x} \nabla \psi - \frac{B \phi}{\mu_0} \nabla (\sigma B \phi) + \frac{1}{\mu_0} \nabla (\sigma B \phi) \times \nabla \psi = P' \nabla \psi$$

$$\Rightarrow \left(\frac{\partial \phi}{\partial x} - \frac{E E'}{\mu_0 \sigma^2} \right) \nabla \psi = P' \nabla \psi$$

$$\Rightarrow \frac{\partial \phi}{\partial x} = \frac{\sigma P' + E E'}{\mu_0 \sigma^2} \quad (14)$$

Thus using in (9) gives

$$\begin{aligned} \Delta^* \psi &= -\sigma \mu_0 \frac{\partial \phi}{\partial x} \\ &= -\mu_0 \sigma^2 P' - E E' \end{aligned} \quad \} \quad (15)$$

No flow eqn

Case II :  Flow

Eqns (4) & (6) give

$$\nabla \times (\mathbf{v} \times \mathbf{B}) = 0$$

$$\Rightarrow \mathbf{v} \times \mathbf{B} = \nabla g \quad (16)$$

\downarrow scalar

Eq (1) gives

$$\mathbf{v} = \nabla \phi \mathbf{e}_\phi + \frac{1}{\sigma \epsilon_0} \nabla \psi \times \mathbf{e}_\phi \quad (17)$$

Thus

(identity 53)

$$\frac{\nabla \phi}{\sigma} \nabla \psi - \frac{B\phi}{\sigma \epsilon_0} \nabla \psi_\perp + \frac{1}{\sigma^2 \epsilon_0} \nabla \psi_\perp \times \nabla \psi = \nabla g$$

$$\Rightarrow (\nabla \psi_\perp \times \nabla \psi) \cdot \mathbf{e}_\phi = 0$$

$$\text{Thus } \psi_\perp = \bar{\Psi}(\psi)$$

& also:

$$\frac{\nabla \phi}{\sigma} \nabla \psi - \frac{B\phi}{\sigma \epsilon_0} \bar{\Psi}' \nabla \psi = \nabla g$$

$$\Rightarrow \left(\frac{\nabla \phi}{\sigma} - \frac{B\phi}{\sigma \epsilon_0} \bar{\Psi}' \right) \nabla \psi = \nabla g.$$

$$\Rightarrow g = g(r, z) = \bar{G}(\psi) \quad (17b)$$

$$\Rightarrow \frac{\nabla \phi}{\sigma} - \frac{B\phi}{\sigma \epsilon_0} \bar{\Psi}' = \bar{G}'(\psi) = G(\psi) \quad (18)$$

$$\text{or } \frac{\nabla \phi}{\sigma} - \frac{B\phi}{\sigma} \bar{\Psi}' = \sigma G(\psi) \quad (19)$$

(4)

Use identity

$$\mathbf{v} \cdot \nabla v + v \times \nabla \times v = \nabla \left(\frac{v^2}{2} \right)$$

in eq (2) to give

$$-\mathbf{v} \times \boldsymbol{\omega} - \frac{\mathbf{J} \times \mathbf{B}}{\sigma} + \nabla \left(\frac{v^2}{2} \right) = -\frac{\nabla p}{\sigma} \quad (20)$$

$$\leftarrow \boldsymbol{\omega} = \nabla \times \mathbf{v}$$

(21)

$$\Rightarrow \boldsymbol{\omega} = \omega_\phi \mathbf{e}_\phi + \frac{1}{\sigma} \nabla \lambda \times \mathbf{e}_\phi$$

(identity 48)

$$= -\frac{\Delta_S^* \Psi_v}{\sigma} \mathbf{e}_\phi + \frac{1}{\sigma} \nabla(\sigma v_\phi) \times \mathbf{e}_\phi$$

where

$$\Delta_S^* \Psi_v \equiv \frac{\partial}{\partial z} \left(\frac{1}{\sigma} \frac{\partial \Psi_v}{\partial z} \right) + \frac{\partial}{\partial \sigma} \left(\frac{1}{\sigma} \frac{\partial \Psi_v}{\partial \sigma} \right) \quad (22)$$

$$\text{Thus } \omega_\phi = -\frac{\Delta_S^* \Psi_v}{\sigma} \quad (23)$$

$$\lambda = \sigma v_\phi \quad (24)$$

Take \mathbf{e}_ϕ component of (20)

$$\Rightarrow (\mathbf{v} \times \boldsymbol{\omega} + \frac{\mathbf{J} \times \mathbf{B}}{\sigma}) \cdot \mathbf{e}_\phi = 0$$

use identity 53 & 51 to get

$$\left[\frac{1}{\sigma^2} \nabla \Psi_v \times \nabla(\sigma v_\phi) + \frac{1}{\sigma^2} \nabla(\sigma B_\phi) \times \nabla \Psi \right] \cdot \mathbf{e}_\phi = 0$$

$$\Rightarrow [\nabla(\sigma B_\phi) - \bar{\Psi}' \nabla(\sigma v_\phi)] \times \nabla \Psi \cdot \mathbf{e}_\phi = 0$$

$$\Rightarrow \nabla(\sigma B_\phi) - \bar{\Psi}' \nabla(\sigma v_\phi) \parallel \nabla \Psi$$

$$\Rightarrow \sigma B_\phi - \bar{\Psi}' \sigma v_\phi = K(\Psi) \quad (25)$$

Now from thermodynamics (see thermo notes Eq 3)

$$\frac{dp}{s} = dh - Tds$$

thus momentum Eq (20) becomes :

$$-\mathbf{v} \times \omega - \frac{\mathbf{J} \times \mathbf{B}}{s} + \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) = T \nabla s \quad (26)$$

Take dot product with \mathbf{v} :

$$-\mathbf{v} \cdot \frac{\mathbf{J} \times \mathbf{B}}{s} + \mathbf{v} \cdot \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) = + v \frac{T s}{s} \text{ to (eq 3)}$$

Now

$$-\mathbf{v} \cdot \frac{\mathbf{J} \times \mathbf{B}}{s} = \mathbf{v} \cdot \frac{\mathbf{B} \times \mathbf{J}}{s} = \mathbf{v} \times \mathbf{B} \cdot \frac{\mathbf{J}}{s}$$

$$= \nabla g \cdot \frac{\mathbf{J}}{s} = G(\psi) \nabla \psi \cdot \frac{\mathbf{J}}{s} \quad (\text{see eq 17b})$$

$$= G(\psi) \frac{\nabla \psi}{s} \cdot \left[\mathbf{e}_\phi \mathbf{J} \phi + \frac{1}{r} \nabla(r \mathbf{B} \phi) \times \mathbf{e}_\phi \right]$$

$$= G(\psi) \frac{\nabla \psi}{s} \cdot \nabla(r \mathbf{B} \phi) \times \mathbf{e}_\phi$$

$$= G(\psi) \frac{\nabla \psi \times \nabla(r \mathbf{B} \phi)}{s} \cdot \mathbf{e}_\phi = -G(\psi) \frac{\mathbf{J}(r \mathbf{B} \phi) \times \nabla \psi}{s} \cdot \mathbf{e}_\phi$$

Also $\mathbf{v} \cdot \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) = \frac{1}{r s} \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) \times \nabla \psi \cdot \mathbf{v} \quad (\text{identity } 54)$

$$= \frac{\bar{\Psi}'}{s} \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) \times \nabla \psi$$

Thus $\left[-G(\psi) \nabla(r \mathbf{B} \phi) + \bar{\Psi}' \nabla \left(\frac{\mathbf{v}^2}{2} + h \right) \right] \times \nabla \psi \cdot \mathbf{e}_\phi = 0$

$\propto \left[-\nabla(Gr \mathbf{B} \phi) + \nabla \left(\bar{\Psi}' \left\{ \frac{\mathbf{v}^2}{2} + h \right\} \right) \right] \times \nabla \psi \cdot \mathbf{e}_\phi$
(verify by expansion)

(6)

$$\frac{\partial \phi}{\psi} = \pi v \phi + K$$

Thus $\nabla \left[-G \tau \phi + \bar{\Psi}' \left(\frac{v^2}{2} + h \right) \right] \times \nabla \psi \cdot e \phi = 0$

or $-G \tau \phi + \bar{\Psi}' \left(\frac{v^2}{2} + h \right) = \bar{H}(\psi) \quad (27)$

$$\Rightarrow -G \frac{\partial \phi}{\psi} + \frac{v^2}{2} + h = \frac{\bar{H}(\psi)}{\bar{\Psi}'}$$

$$\Rightarrow -G (K(\psi) + \pi v \phi) + \frac{v^2}{2} + h = \frac{\bar{H}(\psi)}{\bar{\Psi}'}$$

or finally:

$$h + \frac{v^2}{2} - \pi v \phi G = H(\psi) \quad (28)$$

\uparrow "total" enthalpy

Now take dot product with $\nabla \psi$ (7)

(A)

(B)

(C)

$$-\mathbf{v} \times \omega \cdot \nabla \psi - \frac{\mathbf{J} \times \mathbf{B}}{s} \cdot \nabla \psi + \nabla \left(\frac{v^2}{2} + h \right) \cdot \nabla \psi = \tau s' |\nabla \psi|^2$$

From (A)

$$\mathbf{v} \times \omega \cdot \nabla \psi$$

$$= \left(\frac{\omega_\phi}{r} \nabla(\sigma v_\phi) - \frac{\omega_\phi}{rs} \nabla \psi_r + \frac{1}{r^2} \nabla \psi \times \nabla(\sigma v_\phi) \right) \cdot \nabla \psi$$

$$= \frac{v_\phi}{r} \nabla(\sigma v_\phi) \cdot \nabla \psi - \frac{\omega_\phi \bar{\Psi}' |\nabla \psi|^2}{rs}$$

From (B)

$$\frac{\mathbf{J} \times \mathbf{B}}{s} \cdot \nabla \psi$$

$$= \frac{1}{s} \left[\frac{\mathbf{J}_\phi}{r} \nabla \psi - \frac{\mathbf{B}_\phi}{r} \nabla(\sigma B_\phi) + \frac{1}{r^2} \nabla(\sigma B_\phi) \times \nabla \psi \right] \cdot \nabla \psi$$

$$= \frac{1}{sr} \mathbf{J}_\phi |\nabla \psi|^2 - \frac{B_\phi}{rs} \nabla(\sigma B_\phi) \cdot \nabla \psi$$

From (C)

$$\nabla \left(\frac{v^2}{2} + h \right) \cdot \nabla \psi = \tau \left[\sigma v_\phi G + H \right] \cdot \nabla \psi$$

$$= G \nabla(\sigma v_\phi) \cdot \nabla \psi + \sigma v_\phi G' |\nabla \psi|^2 + H' |\nabla \psi|^2$$

$$\text{Thus } -A - B + C = \tau s' |\nabla \psi|^2$$

\Rightarrow

$$-\frac{v_\phi}{r} \nabla(\sigma v_\phi) \cdot \nabla \psi + \frac{\omega_\phi}{rs} \bar{\Psi}' |\nabla \psi|^2 \quad (I)$$

$$- \frac{\mathbf{J}_\phi}{rs} \cdot \nabla \psi + \frac{B_\phi}{rs} \nabla(\sigma B_\phi) \cdot \nabla \psi \quad (II)$$

$$+ G \nabla(\sigma v_\phi) \cdot \nabla \psi + \sigma v_\phi G' |\nabla \psi|^2 + H' |\nabla \psi|^2$$

$$= \tau s' |\nabla \psi|^2$$

(8)

Now $-I + II + III$

$$= \left[-\frac{V\phi}{\sigma} \nabla(\sigma V\phi) + \frac{B\phi}{\sigma S} \nabla(\sigma B\phi) + G \nabla(\sigma V\phi) \right] \nabla \psi$$

$$\text{Now } \nabla(\sigma B\phi) = \nabla [\bar{\Psi}' \tau V\phi + K] \quad (\text{from Eq 25})$$

$$= \bar{\Psi}' \nabla(\tau V\phi) + \sigma V\phi \bar{\Psi}'' \nabla \psi + K' \nabla \psi$$

thus

$$\rightsquigarrow \left[-\frac{V\phi}{\sigma} \nabla(\sigma V\phi) + \frac{B\phi}{\sigma S} \bar{\Psi}' \nabla(\sigma V\phi) + \frac{B\phi}{\sigma S} \tau V\phi \bar{\Psi}'' \nabla \psi + \frac{B\phi}{\sigma S} K' \nabla \psi \right. \\ \left. + G \nabla(\sigma V\phi) \right] \nabla \psi$$

$$\Rightarrow \left(G - \frac{V\phi}{\sigma} + \frac{B\phi}{\sigma S} \bar{\Psi}' \right) \nabla(\sigma V\phi) \nabla \psi + \frac{B\phi}{\sigma S} \tau V\phi \bar{\Psi}'' |\nabla \psi|^2 + \frac{B\phi}{\sigma S} K' |\nabla \psi|^2$$

[Eq 18]

Thus

$$\frac{B\phi V\phi}{S} \bar{\Psi}'' + \frac{B\phi}{\sigma S} K' + \frac{\omega\phi}{\sigma S} \bar{\Psi}' - \frac{J\phi}{\sigma S} + \sigma V\phi G' + H' \\ = T \tilde{s}'$$

or

$$\frac{\bar{\Psi}'}{\sigma S} \omega\phi - \frac{J\phi}{\sigma S} = T \tilde{s}' - H' - \sigma V\phi G' - \frac{B\phi}{\sigma S} K' - \frac{B\phi V\phi}{\sigma S} \bar{\Psi}''$$

(29)

is the needed 4F equation.

$$\left| \tau^2 \nabla \cdot \left(\frac{\nabla \psi}{\tau^2} \right) \right. \equiv \Delta^* \psi$$

(A1)

Appendix

Let $\nabla f = K(r, z) \nabla \psi$

thus $\frac{\partial f}{\partial r} = K \frac{\partial \psi}{\partial r}$ (A1)

$$\frac{\partial f}{\partial z} = K \frac{\partial \psi}{\partial z} \quad (\text{A2})$$

axisymmetric system

Take z derivative of A1:

$$\frac{\partial}{\partial z} \left(\frac{\partial f}{\partial z} \right) = \frac{\partial K}{\partial z} \frac{\partial \psi}{\partial r} + K \frac{\partial^2 \psi}{\partial z \partial r}$$

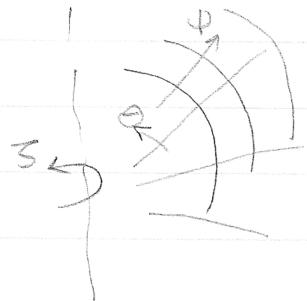
$$\Rightarrow \frac{\partial}{\partial r} \left(K \frac{\partial \psi}{\partial z} \right) = \frac{\partial K}{\partial z} \frac{\partial \psi}{\partial r} + K \frac{\partial^2 \psi}{\partial z \partial r}$$

$$\Rightarrow \frac{\partial K}{\partial r} \frac{\partial \psi}{\partial z} - \frac{\partial K}{\partial z} \frac{\partial \psi}{\partial r} = 0 \quad (\text{A3})$$

Now: $K = K(r, z) = \bar{K}(\psi, \theta)$

$$\Rightarrow \frac{\partial K}{\partial r} = \frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial r} + \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial r}$$

$$\frac{\partial K}{\partial z} = \frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial z} + \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial z}$$



and in A3:

$$\left(\frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial r} + \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial r} \right) \frac{\partial \psi}{\partial z} - \left(\frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial z} + \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial z} \right) \frac{\partial \psi}{\partial r} = 0$$

$$\Rightarrow \cancel{\frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial r}} \frac{\partial \psi}{\partial z} - \cancel{\frac{\partial \bar{K}}{\partial \psi} \frac{\partial \psi}{\partial z}} \frac{\partial \psi}{\partial r}$$

$$+ \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial r} \frac{\partial \psi}{\partial z} - \frac{\partial \bar{K}}{\partial \theta} \frac{\partial \theta}{\partial z} \frac{\partial \psi}{\partial r} = 0$$

$$\Rightarrow \frac{\partial \bar{K}}{\partial \theta} \left(\frac{\partial \theta}{\partial r} \frac{\partial \psi}{\partial z} - \frac{\partial \theta}{\partial z} \frac{\partial \psi}{\partial r} \right) = 0$$

(note $\nabla \theta$)

$$\Rightarrow \frac{\partial \bar{K}}{\partial \theta} (\nabla \theta \times \nabla \psi) \cdot \hat{S} = 0 \Rightarrow \frac{\partial \bar{K}}{\partial \theta} = 0 \quad \begin{matrix} \text{is never} \\ \parallel \text{to } \nabla \psi \end{matrix}$$

Thus

$$\boxed{\bar{K} = \bar{K}(\psi)}$$

 \Rightarrow if $\nabla f = K(r, z) \nabla \psi$
 then $f = F(\psi)$