

December 15, 2025 (Numerical prototypes)

First Numerical Prototypes from $H_3(\mathbb{O})$

Simple vacua, concrete spectra

Key Insight. The octonionic model is not only an abstract Playground of exceptional algebras. Already for very simple vacuum configurations $\langle H \rangle$ in the Albert algebra $H_3(\mathbb{O})$ one can compute explicit eigenvalue multiplets of the mass map $\Pi(\langle H \rangle)$. These multiplets form “numerical prototypes” for fermion mass spectra: banded, hierarchical, and organised in patterns that resemble quark and lepton families.

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From symbols to numbers

So far, the exceptional structures have been presented in a largely symbolic way: octonions, the Albert algebra $H_3(\mathbb{O})$, the mass map $\Pi(H)$ and its relation to vacuum configurations $\langle H \rangle$. At some point, however, the model must leave the realm of pure algebra and produce numbers that can be compared, at least qualitatively, to particle physics.

The good news is that this can already be done with remarkably simple choices of $\langle H \rangle$. By freezing most degrees of freedom and retaining only a few dominant internal directions, one obtains toy vacua that are simple enough to diagonalise exactly, but rich enough to display hierarchical spectra.

Diagonal prototype vacua

The cleanest starting point is a diagonal configuration in $H_3(\mathbb{O})$:

$$\langle H \rangle = \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix}, \quad \lambda_i \in \mathbb{R}.$$

Such a vacuum preserves a maximal amount of internal symmetry: off-diagonal octonionic entries are set to zero, and only three real parameters remain. Despite this simplicity, the induced mass map

$$\Pi(\langle H \rangle) : \mathcal{H}_{\text{int}} \longrightarrow \mathcal{H}_{\text{int}}$$

already exhibits a nontrivial eigenvalue structure.

In the eigenvalue problem

$$\Pi(\langle H \rangle) \Psi = m \Psi,$$

the internal Hilbert space \mathcal{H}_{int} decomposes into sectors that correspond to quark-like and lepton-like states, coloured and colourless states, left- and right-handed components, and so on. Each sector contributes a set of eigenvalues m .

For suitable choices of $(\lambda_1, \lambda_2, \lambda_3)$ one observes:

- Eigenvalues cluster into a few *bands*, separated by gaps.

- Within each band, degeneracies and small splittings appear, induced by the residual internal symmetry.
- The overall pattern can be read as “light, medium, heavy” families, evocative of the observed hierarchy between generations.

Off-diagonal perturbations and splitting patterns

To go beyond the simplest diagonal picture, one can add controlled off-diagonal octonionic entries to $\langle H \rangle$:

$$\langle H \rangle = \begin{pmatrix} \lambda_1 & x & 0 \\ \bar{x} & \lambda_2 & y \\ 0 & \bar{y} & \lambda_3 \end{pmatrix}, \quad x, y \in \mathbb{O}.$$

Even when x and y are chosen along a single imaginary octonion unit, the eigenvalue structure of $\Pi(\langle H \rangle)$ changes qualitatively:

- Degenerate bands split into sub-bands.
- Some states are pushed up or down in mass, mimicking the effect of mixing between different internal directions.
- Eigenvectors acquire nontrivial support across different triality-related components, foreshadowing the mixing phenomena associated with CKM and PMNS matrices.

These controlled deformations are the “numerical workhorses” of the model: they allow us to scan how the spectrum responds when the vacuum is moved inside $H_3(\mathbb{O})$.

Prototype spectra and their interpretation

The resulting eigenvalue sets are not yet to be interpreted as precise predictions for Standard Model fermion masses. Instead, they serve as *prototypes* that illustrate what the exceptional machinery tends to produce generically:

1. **Banded hierarchies:** eigenvalues group into bands with ratios that are naturally large or small, depending on how many internal directions contribute coherently.
2. **Sector-dependent patterns:** quark-like and lepton-like sectors display different splitting patterns, reflecting their different embedding in the internal algebra.
3. **Robustness under deformation:** certain qualitative hierarchies persist under a wide range of small changes in $(\lambda_1, \lambda_2, \lambda_3, x, y)$, suggesting that they are structural rather than fine-tuned.

From this perspective, the complicated observed mass spectrum is seen as one point in a structured space of possible spectra generated by the octonionic geometry.

The role of the Ersatzbank

Within the larger Advent project, these numerical prototypes form an “Ersatzbank” of ready-made examples:

- They can be used to replace more philosophical or qualitative discussion days by concrete numerical sheets.
- They demonstrate to the reader that the abstract constructions genuinely lead to computable spectra with clear patterns.
- They provide anchor points for later, more refined fits where the parameters of $\langle H \rangle$ and the definition of $\Pi(H)$ are tuned to approximate real-world data.

Even very simple vacua in $H_3(\mathbb{O})$ already yield banded, hierarchical spectra: numerical prototypes of a world built from exceptional geometry.

In teaching or outreach contexts, such an Ersatzbank is particularly valuable: it allows one to show explicit tables of “eigenvalues from an exceptional universe” without requiring the audience to follow every technical step.

An invitation to explore

Technically, numerical exploration of $\Pi(\langle H \rangle)$ is straightforward: once a concrete basis for $H_3(\mathbb{O})$ and the relevant operators is chosen, the problem reduces to diagonalising large but structured matrices. Modern linear algebra libraries can handle this comfortably.

Conceptually, the challenge is to interpret the resulting patterns in a way that respects both the algebraic symmetries and the phenomenological constraints. The prototypes presented here are first steps in that direction: they show that the exceptional machinery is capable of producing spectra that look qualitatively like what one expects from quark and lepton masses, without any artificial “by hand” hierarchies inserted into the Lagrangian.

References

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