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## Spectral geometry: bridge to gravity

When a Dirac spectrum encodes both spacetime and octonions

**Key Insight.** The central object of this sheet is a single operator, the Dirac operator  $D$ . In spectral geometry one defines a “spectral action”

$$S_{\text{spec}} = \text{Tr } f(D^2/\Lambda^2),$$

and shows that its expansion produces terms of Einstein–Hilbert type, cosmological-constant type and gauge-kinetic type. This offers a concrete way to connect an internal Dirac matrix  $D_F$  from the octonionic construction with an effective spacetime geometry, without adding a separate gravitational action by hand.

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THIS sheet shifts the focus from the detailed octonionic construction back to a single unifying object: the Dirac operator. We briefly recall how a geometric Dirac operator encodes the metric of spacetime, how the spectral action principle turns its spectrum into an effective action, and how a finite internal matrix can be coupled to it. In this way, spectral geometry provides a precise framework in which internal structures can be tested against gravitational and gauge dynamics.

### Dirac operators in ordinary geometry

In the standard (continuum) description of gravity one starts from a metric  $g_{\mu\nu}$  on a four-dimensional manifold. From  $g_{\mu\nu}$  one builds a Levi-Civita connection, curvature tensors and finally the Einstein–Hilbert action. Fermions enter through a *Dirac operator*  $D_M$ , constructed from  $g_{\mu\nu}$  and the spin connection.

Concretely,  $D_M$  is a first-order differential operator of the schematic form

$$D_M = i \gamma^\mu \nabla_\mu,$$

where  $\nabla_\mu$  is the covariant derivative determined by the metric and  $\gamma^\mu$  are gamma matrices adapted to  $g_{\mu\nu}$ . The metric enters  $D_M$  twice: through the choice of  $\gamma^\mu$  and through the connection in  $\nabla_\mu$ . Thus,  $D_M$  can be viewed as a compact encoding of the metric data.

Spectral geometry turns this observation around: instead of viewing  $D_M$  as derived from  $g_{\mu\nu}$ , it asks how much of  $g_{\mu\nu}$  and its curvature can be *recovered* from the spectrum of  $D_M$ . The answer is: a lot. Certain standard coefficients  $a_n(D_M^2)$ , the Seeley–DeWitt coefficients, carry integrated information about volume and curvature.

For our purposes we only need three facts:

- $a_0(D_M^2)$  is proportional to the total volume,
- $a_2(D_M^2)$  contains the scalar curvature  $R$ ,
- $a_4(D_M^2)$  contains combinations of  $R^2$  and, if gauge fields are included, terms like  $F_{\mu\nu}F^{\mu\nu}$ .

### The spectral action and its structure

The spectral action principle takes the Dirac operator  $D$  as the primary input and defines

$$S_{\text{spec}} = \text{Tr } f(D^2/\Lambda^2),$$

where  $\Lambda$  is a large mass scale and  $f$  is a fixed, positive test function. The trace is taken over the Hilbert space on which  $D$  acts.

If  $D$  is of the usual geometric form built from  $g_{\mu\nu}$ , one can express  $S_{\text{spec}}$  as an integral over spacetime. In four dimensions this yields schematically

$$S_{\text{spec}} \sim \int d^4x \sqrt{|g|} \left( \underbrace{\Lambda^4 a_0(D^2)}_{\text{vacuum / cosmological}} + \underbrace{\Lambda^2 a_2(D^2)}_{\text{Einstein–Hilbert type}} + \underbrace{a_4(D^2)}_{\text{curvature}^2 \text{ and gauge}} + \dots \right).$$

Several points are important:

- the *same* operator  $D$  controls all three sectors: cosmological term, Einstein–Hilbert term and gauge kinetic terms,
- the organisation into these sectors is automatic once  $D$  is given,
- no separate gravitational Lagrangian is added by hand; it appears as part of the expansion of  $S_{\text{spec}}$ .

The precise numerical coefficients depend on  $f$  and on the detailed form of  $D$ , but the qualitative structure is robust.

### Adding a finite internal space

So far  $D$  only knew about spacetime. To include internal quantum numbers one enlarges the Hilbert space from pure spinors to

$$\mathcal{H} = L^2(\text{spinors on } M) \otimes \mathcal{H}_F,$$

where  $\mathcal{H}_F$  is a finite-dimensional Hilbert space. All internal labels (gauge charges, generation indices, etc.) live in  $\mathcal{H}_F$ .

The Dirac operator is then modified to

$$D = D_M \otimes \mathbf{1} + \gamma^5 \otimes D_F.$$

Here:

- $D_M$  is the geometric Dirac operator built from  $g_{\mu\nu}$ ,
- $D_F$  is a finite matrix acting on  $\mathcal{H}_F$ ,
- $\gamma^5$  is the chirality matrix: it distinguishes left-handed from right-handed spinor components.

The factor  $\gamma^5$  appears for a concrete reason. In realistic models the internal symmetries act differently on left- and right-handed fermions. Placing  $D_F$  next to  $\gamma^5$  ensures that this chiral pattern is correctly reflected when  $D$  is squared and inserted into  $S_{\text{spec}}$ .

For the Standard Model there is a specific choice of  $(\mathcal{H}_F, D_F)$  for which the leading terms in the expansion of  $S_{\text{spec}}$  reproduce the known gauge group and fermion content. This shows that a *finite* internal matrix  $D_F$  can already carry realistic particle-physics information once it is coupled to  $D_M$  in this way.

## Octonionic data as input for an internal Dirac operator

In the octonionic part of this calendar we have already built a rich internal structure:

- an exceptional internal space based on  $H_3(\mathbb{O})$ ,
- rotor and compressor operators acting on the states of a generation,
- mass maps such as  $\Pi(\langle H \rangle)$ .

Later sheets introduce additional invariants (for example universe parameters  $(p_S, p_G, p_T)$ ) that classify sectors of this internal algebra.

From the viewpoint of spectral geometry, these objects are natural candidates for the blocks of a finite internal Dirac operator  $D_F$ :

- combinations of rotors and compressors suggest gauge-covariant parts of  $D_F$ ,
- mass maps provide off-diagonal blocks that mix left- and right-handed components in  $D_F$ ,
- sector invariants indicate which combinations of internal operators could appear in the coefficients  $a_n(D^2)$  and thus influence effective coupling strengths in  $S_{\text{spec}}$ .

*The Dirac operator links spacetime and internal physics. Spectral geometry uses a combined Dirac operator to automatically generate an action that includes both gravity and gauge forces. Our octonionic model offers the internal part; spectral action shows how its properties shape the universe.*

This sheet does not attempt to write down a final  $D_F$ . Its aim is more modest and more precise:

- given any explicit choice of  $D_F$  built from the octonionic operator toolbox, spectral geometry provides a clear recipe for constructing a combined operator

$$D = D_M \otimes \mathbf{1} + \gamma^5 \otimes D_F$$

and for computing the corresponding spectral action,

- the resulting  $S_{\text{spec}}$  automatically organises terms into cosmological, Einstein–Hilbert and gauge-kinetic parts.

The octonionic model supplies the internal matrix; the spectral framework supplies the rule “turn  $D$  into an action”.

## Summary of the bridge

The bridge built here has three steps:

1. A Dirac operator  $D_M$  compactly encodes a space-time metric  $g_{\mu\nu}$  and its curvature.
2. Enlarging  $D_M$  to  $D = D_M \otimes \mathbf{1} + \gamma^5 \otimes D_F$  allows a finite internal matrix  $D_F$  to carry gauge and mass information.
3. The spectral action  $S_{\text{spec}} = \text{Tr} f(D^2/\Lambda^2)$  turns the spectrum of this combined Dirac operator into a unified expression whose expansion contains both gravitational and gauge-kinetic terms.

In this sense, the spectral-geometry framework gives a precise target: if the octonionic construction can produce a concrete internal Dirac matrix  $D_F$ , then there is a mathematically well-defined way to test how its spectrum feeds into an effective action that already knows about gravity.

## References

- [1] A. Connes and A. H. Chamseddine, “The spectral action principle,” *Commun. Math. Phys.* **186**, 731–750 (1997).
- [2] A. H. Chamseddine, A. Connes and M. Marcolli, “Gravity and the standard model with neutrino mixing,” *Adv. Theor. Math. Phys.* **11**, 991–1089 (2007).