Faster than Hermitian Quantum Mechanics by Carl M Bender et al

Ana Fabela Hinojosa



Supervisors: Jesper Levinsen Meera Parish

Outline

Introduction

Time evolution

Brachistochrone problem

Conclusion

Contents

Introduction

Time evolution

Brachistochrone problem

Conclusion

Hilbert spaces

A Hilbert space is a vector space that can be infinite dimensional.

Vector spaces are equipped with an inner product.

Inner product:
Define a distance function.

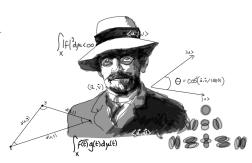


Fig.1: Hilbert space stuff

Hamiltonians as Observables

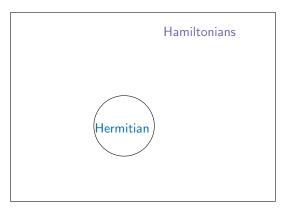


Fig.1: The set of all possible Hamiltonians.

- 1. Observables are **self-adjoint** operators.
- 2. A state ψ of the quantum system is a unit vector of \hat{H} .
- 3. \hat{H} has a real energy spectrum with defined lowest energy.
- 4. Unitarity.

Hamiltonians as Observables

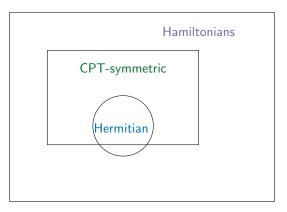


Fig.1: The set of all possible Hamiltonians.

- 1. Observables are **self-adjoint** operators.
- 2. A state ψ of the quantum system is a unit vector of \hat{H} .
- 3. \hat{H} has a real energy spectrum with defined lowest energy.
- 4. Unitarity.

Hamiltonians as Observables

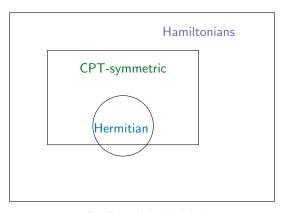


Fig.1: The set of all possible Hamiltonians.

- 1. Observables are self-adjoint operators.
- 2. A state ψ of the quantum system is a unit vector of \hat{H} .
- 3. \hat{H} has a real energy spectrum with defined lowest energy.
- 4. Unitarity.

 $\mathcal{C} \to \mathrm{charge}\,\mathrm{conjugation},$

 $\mathcal{P} \to \mathrm{spatial} \ \mathrm{inversion},$

 $\mathcal{T} \to \text{complex conjugation and time reversal.}$

Hilbert space $\rightarrow \hat{H} = \hat{H}^{CPT}$

 $\mathcal{C} \to \text{charge conjugation},$

 $\mathcal{P} \to \mathrm{spatial} \ \mathrm{inversion},$

 $\mathcal{T} \to \text{complex conjugation and time reversal.}$

Hilbert space $\rightarrow \hat{H} = \hat{H}^{CPT}$

Inner product is determined dynamically in terms of the Hamiltonian.

 $\mathcal{C} \to \mathrm{charge}\,\mathrm{conjugation},$

 $\mathcal{P} \to \mathrm{spatial} \ \mathrm{inversion},$

 $\mathcal{T} \to \text{complex conjugation and time reversal.}$

Hilbert space $\rightarrow \hat{H} = \hat{H}^{C\mathcal{PT}}$

Inner product is determined dynamically in terms of the Hamiltonian.

If the eigenfunctions of a \mathcal{PT} -symmetric Hamiltonian are **not** also eigenfunctions of the \mathcal{PT} operator we say the Hamiltonian possesess **broken** \mathcal{PT} -symmetry.

 $\mathcal{C} \to \mathrm{charge}\,\mathrm{conjugation},$

 $\mathcal{P} \to \mathrm{spatial} \ \mathrm{inversion},$

 $\mathcal{T} \to \text{complex conjugation and time reversal.}$

Hilbert space $\rightarrow \hat{H} = \hat{H}^{CPT}$

Inner product is determined dynamically in terms of the Hamiltonian.

If the eigenfunctions of a \mathcal{PT} -symmetric Hamiltonian are **not** also eigenfunctions of the \mathcal{PT} operator we say the Hamiltonian possesess **broken** \mathcal{PT} -symmetry.

Interesting physical phenomena occurs in the broken symmetry region

Contents

Introduction

Time evolution

Brachistochrone problem

Conclusion

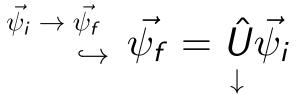
$$\vec{\psi_i}
ightarrow \vec{\psi_f}$$

$$ec{\psi_i} \stackrel{
ightarrow}{
ightarrow} ec{\psi_f} \stackrel{
ightarrow}{\psi_f} = \hat{U} ec{\psi_i}$$

$$\vec{\psi_i} \stackrel{\vec{\psi_f}}{\hookrightarrow} \vec{\psi_f} = \hat{U}\vec{\psi_i}$$

1. Hermitian quantum mechanics:

$$\hat{U} = e^{-i\hat{H}t/\hbar}$$
, where $t > 0$.



1. Hermitian quantum mechanics:

$$\hat{U}=e^{-i\hat{H}t/\hbar}$$
, where $t>0$.

2. CPT quantum mechanics: ?



Contents

Introduction

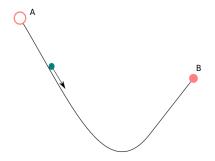
Time evolution

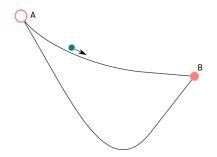
Brachistochrone problem

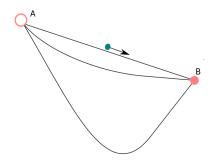
Conclusion

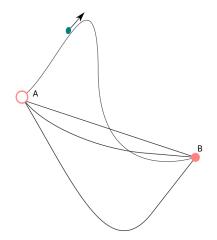


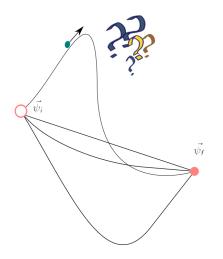












βράχιστος χρόνος brákhistos khrónos: "shortest time"

Fig.2: A particle travels from left to right in time t, Can we make this trip nearly instantaneous?

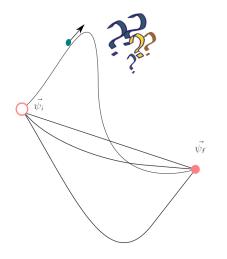


Fig.2: A particle travels from left to right in time t, Can we make this trip nearly instantaneous?

βράχιστος χρόνος brákhistos khrónos: "shortest time"

How fast can we evolve a state?

This space is spanned by $\vec{\psi_i}$ and $\vec{\psi_f}$.

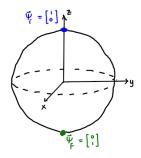


Fig.3: Bloch sphere with initial and final states

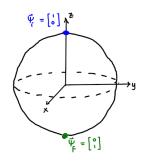


Fig.3: Bloch sphere with initial and final states

This space is spanned by $\vec{\psi_i}$ and $\vec{\psi_f}$.

We want the **fastest** time evolution possible, without violating the time-energy uncertainty principle.

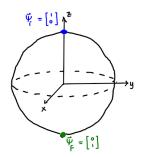


Fig.3: Bloch sphere with initial and final states

This space is spanned by $\vec{\psi_i}$ and $\vec{\psi_f}$.

We want the **fastest** time evolution possible, without violating the time-energy uncertainty principle.

The chosen Hamiltonian must satisfy:

Energy constraint:

$$\omega = \textit{E}_{\rm max} - \textit{E}_{\rm min}$$

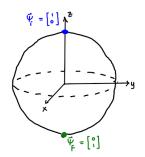


Fig.3: Bloch sphere with initial and final states

This space is spanned by $\vec{\psi_i}$ and $\vec{\psi_f}$.

We want the **fastest** time evolution possible, without violating the time-energy uncertainty principle.

The chosen Hamiltonian must satisfy:

Energy constraint: $\omega = E_{\rm max} - E_{\rm min}$

Will a complex non-Hermitian Hamiltonian give time-optimal evolution?

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} \mathbf{s} & r \mathbf{e}^{-i\theta} \\ r \mathbf{e}^{i\theta} & u \end{pmatrix}, \{r, \mathbf{s}, \theta, u\} \in \mathbb{R},$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} s & re^{-i\theta} \\ re^{i\theta} & u \end{pmatrix}, \{r, s, \theta, u\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = (s-u)^2 + 4r^2,$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} s & re^{-i\theta} \\ re^{i\theta} & u \end{pmatrix}, \{r, s, \theta, u\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = (s-u)^2 + 4r^2,$$

Time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} \mathbf{s} & r \mathbf{e}^{-i\theta} \\ r \mathbf{e}^{i\theta} & u \end{pmatrix}, \{r, \mathbf{s}, \theta, u\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = (s-u)^2 + 4r^2,$$

Time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = \mathrm{e}^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

CPT-symmetric case

$$\tilde{H} = \begin{pmatrix} re^{i\theta} & s \\ s & re^{-i\theta} \end{pmatrix}, \quad \{r, s, \theta\} \in \mathbb{R},$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} s & re^{-i\theta} \\ re^{i\theta} & u \end{pmatrix}, \{r, s, \theta, u\} \in \mathbb{R},$$

Energy constraint

$$\omega^2=(s-u)^2+4r^2,$$

Time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = \mathrm{e}^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

CPT-symmetric case

$$ilde{H} = egin{pmatrix} r e^{i heta} & s \ s & r e^{-i heta} \end{pmatrix}, \quad \{r, s, heta\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = 4s^2 - 4r^2\sin^2\theta > 0,$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

Hermitian case

$$\hat{H} = \begin{pmatrix} s & r e^{-i\theta} \\ r e^{i\theta} & u \end{pmatrix}, \{r, s, \theta, u\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = (s - u)^2 + 4r^2,$$

Time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

CPT-symmetric case

$$\tilde{H} = \begin{pmatrix} re^{i\theta} & s \\ s & re^{-i\theta} \end{pmatrix}, \quad \{r, s, \theta\} \in \mathbb{R},$$

Energy constraint

$$\omega^2 = 4s^2 - 4r^2 \sin^2 \theta > 0,$$

Time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = \frac{e^{\frac{-itr\cos\theta}{\hbar}}}{\cos\alpha} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar} - \alpha) \\ -i\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

where $\sin(\alpha) = \frac{r}{s}\sin(\theta)$.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$ec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad ec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\Rightarrow |b| = \frac{2r}{\omega}\sin\left(\frac{\omega t}{2\hbar}\right),$$

$$ec{\psi_i} = egin{pmatrix} 1 \ 0 \end{pmatrix}, \quad ec{\psi_f} = egin{pmatrix} a \ b \end{pmatrix} = egin{pmatrix} 0 \ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{2\hbar} \sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\Rightarrow |b| = \frac{2r}{\omega}\sin\left(\frac{\omega t}{2\hbar}\right),$$

$$\Rightarrow t = \frac{2\hbar}{\omega}\arcsin\left(\frac{\omega|b|}{2r}\right),$$

$$ec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad ec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{split} \begin{pmatrix} a \\ b \end{pmatrix} &= e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{2\hbar}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}. \\ \Rightarrow |b| &= \frac{2r}{\omega}\sin\left(\frac{\omega t}{2\hbar}\right), \\ \Rightarrow t &= \frac{2\hbar}{\omega}\arcsin\left(\frac{\omega|b|}{2r}\right), \end{split}$$

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω .

$$\therefore \tau = \frac{2\hbar}{\omega} \arcsin(|b|) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

$$ec{\psi_i} = egin{pmatrix} 1 \ 0 \end{pmatrix}, \quad ec{\psi_f} = egin{pmatrix} a \ b \end{pmatrix} = egin{pmatrix} 0 \ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{pmatrix} a \\ b \end{pmatrix} = e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\Rightarrow |b| = \frac{2r}{\omega}\sin\left(\frac{\omega t}{2\hbar}\right),$$

$$\Rightarrow t = \frac{2\hbar}{\omega}\arcsin\left(\frac{\omega|b|}{2r}\right),$$

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω .

$$\therefore \tau = \frac{2\hbar}{\omega} \arcsin(|b|) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

$$\begin{pmatrix} a \\ b \end{pmatrix} = \frac{e^{\frac{-itr\cos\theta}{\hbar}}}{\cos\alpha} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar} - \alpha) \\ -i\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$ec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad ec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{split} \begin{pmatrix} a \\ b \end{pmatrix} &= e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega}\sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega}e^{i\theta}\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}. \\ \\ \Rightarrow |b| &= \frac{2r}{\omega}\sin\left(\frac{\omega t}{2\hbar}\right), \\ \\ \Rightarrow t &= \frac{2\hbar}{\omega}\arcsin\left(\frac{\omega|b|}{2r}\right), \end{split}$$

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω .

$$\therefore \tau = \frac{2\hbar}{\omega} \arcsin(|b|) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

$$\begin{pmatrix} a \\ b \end{pmatrix} = \frac{e^{\frac{-itr\cos\theta}{\hbar}}}{\cos\alpha} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar} - \alpha) \\ -i\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\Rightarrow |a| = \frac{\cos\frac{\omega t}{2\hbar} - \alpha}{\cos\alpha} \sin\left(\frac{\omega t}{2\hbar}\right),$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{split} \begin{pmatrix} a \\ b \end{pmatrix} &= e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega} \sin(\frac{\omega t}{2\hbar}) \\ -i\frac{2r}{\omega} e^{i\theta} \sin(\frac{\omega t}{2\hbar}) \end{pmatrix}. \\ \Rightarrow |b| &= \frac{2r}{\omega} \sin\left(\frac{\omega t}{2\hbar}\right), \\ \Rightarrow t &= \frac{2\hbar}{\omega} \arcsin\left(\frac{\omega|b|}{2r}\right), \end{split}$$

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω

$$\therefore \tau = \frac{2\hbar}{\omega} \arcsin(|b|) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

$$\begin{pmatrix} a \\ b \end{pmatrix} = \frac{e^{\frac{-itr\cos\theta}{\hbar}}}{\cos\alpha} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar} - \alpha) \\ -i\sin(\frac{\omega t}{2\hbar}) \end{pmatrix}.$$

$$\Rightarrow |a| = \frac{\cos\frac{\omega t}{2\hbar} - \alpha}{\cos\alpha} \sin\left(\frac{\omega t}{2\hbar}\right),$$

$$\Rightarrow t = \frac{\hbar}{\omega} (2\alpha + \pi),$$

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω .

$$\therefore \tau = \frac{2\hbar}{\omega}\arcsin\left(|b|\right) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Hermitian time evolution

$$\begin{split} \begin{pmatrix} a \\ b \end{pmatrix} &= e^{\frac{-i(s+u)t}{2\hbar}} \begin{pmatrix} \cos(\frac{\omega t}{2\hbar}) - i\frac{s-u}{\omega} \sin(\frac{\omega t}{2\hbar}) \\ -i\frac{s^2}{\omega} e^{i\theta} \sin(\frac{\omega t}{2\hbar}) \end{pmatrix}. \\ \\ \Rightarrow |b| &= \frac{2r}{\omega} \sin\left(\frac{\omega t}{2\hbar}\right), \\ \\ \Rightarrow t &= \frac{2\hbar}{\omega} \arcsin\left(\frac{\omega|b|}{2r}\right), \end{split}$$

For all r > 0 subject to $\omega^2 = (s - u)^2 + 4r^2$ for a fixed ω .

$$\therefore \tau = \frac{2\hbar}{\omega}\arcsin\left(|b|\right) = \frac{\pi\hbar}{\omega}$$

is the minimum passage time.

CPT time evolution

 $\rightarrow \omega^2 = 4s^2 \cos^2(\alpha)$

Then $t \to 0$ when $\alpha \to -\frac{\pi}{2}$. but we also require s, r >> 1.

 $\mathcal{C}\mathcal{P}\mathcal{T}$ inner products are defined depending on the Hamiltonian used.

 $\mathcal{C}\mathcal{P}\mathcal{T}$ inner products are defined depending on the Hamiltonian used.

$$ec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad ec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

1. orthogonal \rightarrow Hermitian inner product

 $\mathcal{C}\mathcal{P}\mathcal{T}$ inner products are defined depending on the Hamiltonian used.

$$ec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad ec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

1. orthogonal \rightarrow Hermitian inner product

 \mathcal{CPT} inner products are defined depending on the Hamiltonian used.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

1. orthogonal \rightarrow Hermitian inner product vector separation: $\delta = 2\arccos |\langle \psi_f | \psi_i \rangle| = \pi$

 \mathcal{CPT} inner products are defined depending on the Hamiltonian used.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- 1. orthogonal \rightarrow Hermitian inner product vector separation: $\delta = 2\arccos |\langle \psi_f | \psi_i \rangle| = \pi$
- 2. not orthogonal \rightarrow CPT inner product

 $\mathcal{C}\mathcal{P}\mathcal{T}$ inner products are defined depending on the Hamiltonian used.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- 1. orthogonal \rightarrow Hermitian inner product vector separation: $\delta = 2\arccos |\langle \psi_f | \psi_i \rangle| = \pi$
- 2. not orthogonal \rightarrow CPT inner product vector separation: $\delta = \pi 2|\alpha|$

 \mathcal{CPT} inner products are defined depending on the Hamiltonian used.

$$\vec{\psi_i} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \vec{\psi_f} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- 1. orthogonal \rightarrow Hermitian inner product vector separation: $\delta = 2\arccos |\langle \psi_f | \psi_i \rangle| = \pi$
- 2. not orthogonal \rightarrow CPT inner product vector separation: $\delta = \pi 2|\alpha|$

We can choose α to create a "wormhole" effect in Hilbert space.

Contents

Introduction

Time evolution

Brachistochrone problem

Conclusion

Conclusion

- 1. This paper demonstrates the theoretical differences in both \mathcal{CPT} -symmetric and conventional Hermitian quantum mechanics.
- 2. The "wormhole" effect could be of importance in design and implementation of fast quantum computation and comunication algorithms.
- 3. There could be quantum protection mechanisms limiting the applicability of Hilbert-space "wormholes".

References



C. M. Bender, D. C. Brody, H. F. Jones, and B. K. Meister. Faster than hermitian quantum mechanics.

Phys. Rev. Lett. 98, 040403 (2007). DOI: 10.1103/PhysRevLett.98.040403.