## Simple finite element methods in Python

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### 1 Heat equation

Let  $\Omega \subset \mathbb{R}^d$ , d=2,3 be the computational domain. We suppose to have a disjoined partition of its boundary:  $\partial\Omega = \Gamma_D \cup \Gamma_N \cup \Gamma_R$ . We consider the parabolic equation for the temperature T, heat flux  $\vec{q}$  and heat release  $\dot{q}$ 

$$\begin{cases} \vec{q} = -k\nabla T \\ \rho C_p \frac{\partial T}{\partial t} + \operatorname{div}(\vec{v}T) + \operatorname{div}\vec{q} = \dot{q} & \text{in } \Omega \\ T = T^D & \text{in } \Gamma_D \\ k \frac{\partial T}{\partial n} = q^N & \text{in } \Gamma_N \\ c_R T + k \frac{\partial T}{\partial n} = q^R & \text{in } \Gamma_R \end{cases} \tag{1.1}$$

Let  $H^1_f:=\Big\{u\in H^1(\Omega)\ \Big|\ T_{|_{\Gamma_D}}=f\Big\}.$  The standard weak formulation looks for  $T\in H^1_{T^D}$  such that for all  $\varphi\in H^1_0(\Omega)$ 

$$\int_{\Omega} \rho C_{p} \frac{\partial T}{\partial t} \varphi - \int_{\Omega} \vec{v} T \cdot \nabla \varphi + \int_{\Omega} k \nabla T \cdot \nabla \varphi + \int_{\Gamma_{R}} c_{R} T \varphi + \int_{\Gamma_{R} \cup \Gamma_{N}} \vec{v}_{n} T \varphi = \int_{\Omega} \dot{q} \varphi + \int_{\Gamma_{R}} q^{R} \varphi \quad (1.2)$$

We can derive (1.2) from (1.1) by the divergence theorem

$$\int_{\Omega} \operatorname{div} \vec{F} = \int_{\partial \Omega} \vec{F}_n \quad \stackrel{F \to F \varphi}{\Longleftrightarrow} \quad \int_{\Omega} (\operatorname{div} \vec{F}) \varphi = - \int_{\Omega} \vec{F} \cdot \nabla \varphi + \int_{\partial \Omega} \vec{F}_n \varphi,$$

which gives with  $\vec{F} = \vec{v} + \vec{q}$ 

$$\int_{\Omega} \operatorname{div} \left( \vec{v} + \vec{q} \right) \varphi = - \int_{\Omega} \vec{v} \cdot \nabla \varphi + \int_{\Omega} k \nabla T \cdot \nabla \varphi + \int_{\partial \Omega} \vec{F}_n \varphi.$$

Using that  $\phi$  vanishes on  $\Gamma_D$  we have

$$\int_{\partial\Omega} \vec{F}_n \varphi = \! \int_{\Gamma_N \cup \Gamma_R} \vec{F}_n \varphi = \! \int_{\Gamma_N \cup \Gamma_R} \vec{v}_n \varphi + \! \int_{\Gamma_N \cup \Gamma_R} \vec{q}_n \varphi,$$

and then with the different boundary conditions, we find

$$\int_{\Gamma_{N}\cup\Gamma_{R}}\vec{q}_{n}\varphi=\int_{\Gamma_{D}}q^{N}\varphi+\int_{\Gamma_{R}}\left(\mathbb{R}-c_{R}T\right)\varphi$$

# Appendices

### A Finite elements on simplices

### A.1 Simplices

We consider an arbitrary non-degenerate simplex  $K = (x_0, x_1, \dots, x_d)$ . The (signed) volume of K is given by

$$|K| = \frac{1}{d!} \det(x_1 - x_0, \dots, x_d - x_0) = \frac{1}{d!} \det(1, x_0, x_1, \dots, x_d) \quad 1 = (1, \dots, 1)^\mathsf{T}.$$
 (A.1)

The d+1 sides  $S_k$  (co-dimension one, d-1-simplices or facets) are defined by  $S_k = (x_0, \dots, x_K, \dots, x_d)$ . The height is  $d_k = |P_{S_k}x_k - x_k|$ , where  $P_S$  is the orthogonal projection on the hyperplane associated to  $S_k$ . We have

$$d_k = d \frac{|K|}{|S_k|}$$
 (and for  $d = 3 |S_k| = \frac{1}{2} |u \times v|$ )

### A.2 Integration on simplices

Any polynomial in the barycentric coordinates can be integrated exactly.

$$\int_{K} \prod_{i=1}^{d+1} \lambda_{i}^{n_{i}} d\nu = d! |K| \frac{\prod_{i=1}^{d+1} n_{i}!}{\left(\sum_{i=1}^{d+1} n_{i} + d\right)!}$$
(A.2)

see [EisenbergMalvern73], [VermolenSegal18].

#### A.3 Finite elements

The d+1 basis functions of the  $P^1$  (Courant) element are the barycentric coordinates  $\lambda_i$  defined as being affine with respect to the coordinates and  $\lambda_i(x_j) = \delta_{ij}$ . Their constant gradient is given by

$$\nabla \lambda_{i} = -\frac{1}{d_{i}} \vec{n_{i}}.$$